

Cite this: *Nanoscale Adv.*, 2021, 3, 1047

Solution processed transparent anatase TiO₂ nanoparticles/MoO₃ nanostructures heterojunction: high performance self-powered UV detector for low-power and low-light applications†

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Ultraviolet (UV) photodetectors are considered as the major players in energy saving technology of the future. Efforts are needed to further develop such devices, which are capable of operating efficiently at low driving potential as well as with weak illumination. Herein, we report an all-oxide, highly transparent TiO₂/MoO₃ bilayer film, with nanoparticulate anatase TiO₂ as the platform, fabricated by a simple solution based method and demonstrate its use in UV photodetection. Photoconductivity measurement with 352 nm light reveals the self-powered UV detection capability of the device due to the built-in potential at the bilayer interface. The device exhibits a high photoresponsivity (46.05 A W⁻¹), detectivity (2.84 × 10¹² Jones) and EQE (16 223%) even with a weak illumination of 76 μW cm⁻², at a low bias of only -1 V. The self-powered performance of the bilayer device is comparable to that of commercial Si photodetectors as well as other such UV detectors reported based on metal oxide heterojunctions. The improved and faster photoresponse shown by the device is due to the formation of an effective heterojunction, as evidenced by XPS, electrochemical and *I*-*V* studies. It can be further attributed to the better charge transport through the densely aligned nanostructures, reduced recombination and the better mobility of anatase TiO₂ nanoparticles. The performance is best-in-class and proves the potential of the transparent heterojunction to be used in highly responsive, self-powered UV detectors for low bias, low light applications.

Received 20th September 2020

Accepted 17th December 2020

DOI: 10.1039/d0na00780c

rsc.li/nanoscale-advances

Introduction

Wide bandgap metal oxide semiconductors have been widely accepted for their potential use in a variety of applications due to their low cost, less toxicity and earth abundant nature.¹ Enormous efforts have been made towards the evolution of transparent electronic devices employing wide bandgap

transparent semiconductors. Transparent diodes,² transparent transistors,³ sensors,⁴ displays, solar cells,⁵ UV LEDs⁶ and photodetectors^{7,8} have been demonstrated. Properties such as optical transparency, environmental-friendly nature, thermal and chemical stability can make them more attractive among semiconductor-based devices.^{2,9}

TiO₂ is a popular n-type wide bandgap semiconductor serving as a prospective candidate in energy conversion applications because of its chemical and thermal stability, low cost and environmental-friendly characteristics. Due to these remarkable properties, TiO₂ is used in many applications including photodetectors, photocatalysis, photovoltaics, electrochromic devices and sensors.¹⁰⁻¹² Even though both ZnO and TiO₂ have been explored for most optoelectronic applications, the photo and chemical stability of TiO₂ are known to be higher than those of ZnO.¹²⁻¹⁵ However, it was observed that the charge separation efficiency of TiO₂ in applications such as photocatalytic activity gets strongly suppressed due to the recombination of photogenerated carriers.^{16,17} Fabricating a heterojunction with TiO₂ is one of the important techniques for reducing the recombination, thereby increasing the lifetime of the carriers, enhancing the charge separation efficiency and improving the charge transfer efficiency. The heterojunction of

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† Electronic supplementary information (ESI) available: XRD patterns of TiO₂, MoO₃ and TiO₂/MoO₃ films; FE-SEM cross-sectional image of the bilayer film; UV absorption spectra; *I*-*V* and *I*-*t* characteristics of the TiO₂ single layer device; *I*-*V* characteristics of the MoO₃ single layer device, log *I*-log *V* plot of the bilayer device; equivalent circuit model for EIS; table comparing the resistance values of TiO₂, MoO₃ and TiO₂/MoO₃ films as calculated using the EIS equivalent circuit model; calculated diode parameters of the bilayer device; table comparing the performance of TiO₂ based UV detectors; table comparing the performance of TiO₂/MoO₃ based UV detectors; statistical data on the performance of the bilayer device based on the 34 devices made. See DOI: 10.1039/d0na00780c

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TiO₂ with numerous semiconductors including CdTe, CdS, CuI, Cu₂O, MoS₂, and ZnO was studied where an enhancement in the photoelectrochemical properties, photovoltaic or photocatalytic performance was observed.^{5,18–20}

α -MoO₃ is one of the well-known layered structured materials widely used in electrodes of energy storage devices. Owing to its high hole mobility and work function, it is used in energy conversion devices as well. The electronic characteristics and the layered structure are the key factors making MoO₃ a potential semiconductor in various applications including sensors, photocatalysts, and smart windows.^{21,22} It is well known that the charge separation efficiency in case of a type II heterojunction is much greater than that of a type I heterojunction.^{23,24} As the energy bands of TiO₂ align well and form a type II heterojunction with MoO₃, the TiO₂/MoO₃ heterojunction structures have also been explored for various applications including photocatalysis,²⁵ supercapacitors²⁶ and ion storage devices.²⁷ For example, a TiO₂/MoO₃ heterojunction fabricated using TiO₂ nanofibers as the template showed a significant increase in the adsorption of dye and improvement in the photocatalytic activity as compared to bare TiO₂ nanofibres.²⁵ The increase in the areal and specific capacitances was observed for the supercapacitor using a TiO₂ nanotube/MoO₃ nanoplate heterostructure.²⁶ S. H. Paek *et al.* observed an enhancement in the lithium ion storage capacity and reversible capacity in the porous TiO₂ pillared MoO₃.²⁷ N. K. Shrestha *et al.* investigated the electrochromic properties of TiO₂/MoO₃ heterostructures and observed a 7 fold increase in the charge storage capacity compared to TiO₂ nanotubes.²⁸

UV light detection is essential for diverse applications including environmental monitoring, remote control, imaging and opto-electronic circuits.²⁹ The key features required for the UV detectors are the strong light absorption and efficient separation of photogenerated carriers.³⁰ Self-powered UV detectors are of great interest due to their zero energy consumption, sustainable nature and minimal maintenance.³¹ Such detectors are possible with efficient photoelectrochemical interfaces, Schottky junctions, or any other heterojunctions.³² The self-powered performance arises due to the built-in potential at the interface or junction, which drives the separation of photogenerated carriers *via* a photovoltaic effect.^{31,33} TiO₂ based heterojunctions are also evolving as promising UV detectors and show efficient photodetection ability with an ultrafast response and high photosensitivity.^{34,35} Although there are reports on TiO₂ based heterojunctions with organic layers, NiO, SnO₂, CuZnS, and ZnS for self-powered photodetection application,^{31,32,36–39} such studies on TiO₂/MoO₃ were rarely reported. A recent report on a TiO₂/MoO₃ UV detector employed hydrothermally synthesised 1D TiO₂ (rutile) nanowires with MoO₃ nanostructures, and used a bias voltage of 2.2 V, with no observation of self-powered performance.⁴⁰

According to the literature, 1D TiO₂ nanostructures have been extensively used in TiO₂ based UV detectors as well as in TiO₂/MoO₃ based devices for various other energy harvesting applications.^{21–24} Although 1D TiO₂ nanostructures offer high specific area, defined geometry and efficient unidirectional carrier transport, the high surface recombination due to their

high surface to volume ratio is considered as one of the potential limiting factors for their use in energy harvesting systems.⁴¹ The surface recombination is known to be less in materials with a lower aspect ratio, where the parasitic surface area is lower.⁴² Thus, the TiO₂/MoO₃ bilayer heterojunction using TiO₂ nanoparticles would not have this limitation and also would be more transparent, unlike the one reported based on hydrothermally grown TiO₂ nanowires.^{40,43} As already mentioned, transparent devices are considered as the major players in future energy saving technology due to their unique optical and electrical properties. It's also notable that previous reports did not explore TiO₂ with anatase crystal structure, for which the electron mobility is known to be relatively higher.¹¹ Thus, it's interesting to study a solution processed TiO₂/MoO₃ heterojunction device with anatase TiO₂ nanoparticle based bottom layer, owing to its transparent nature, superior mobility, lower aspect ratio, and also the simpler cost-effective preparation methods.^{11,12} Based on the above considerations, herein we attempted to fabricate a transparent anatase TiO₂ nanoparticle α -MoO₃ nanostructure based heterojunction device with a simple solution based approach. The charge transport properties of the heterojunction were investigated by analysing the *I*–*V* characteristics in detail. Also, the effect of the morphology and junction formation on the charge transfer and the self-powered UV photodetection capability has been analysed.

Experimental details

All the chemicals were purchased from Sisco Research Laboratories Pvt. Ltd., India and used without further purification.

A transparent solution for making TiO₂ film was prepared by the hydrolysis reaction of titanium tetraisopropoxide (TTIP) in HNO₃ aqueous solution.⁴⁴ First, an alkoxide solution was prepared by dissolving 26.5 M TTIP in ethanol and an acid solution was prepared by mixing 0.3 M nitric acid in water under stirring for 2 hours. The volume ratio of water to ethanol was maintained at 1 : 1. The acid solution was added dropwise into the alkoxide solution under vigorous stirring. The resultant mixture was kept at 45 °C for 36 hours to get a transparent solution. The solution was then spin coated on cleaned FTO/glass substrates at 3000 rpm for 30 s followed by drying at 100 °C for 10 min. This process was repeated seven times to attain the desired thickness of the TiO₂ and finally the film was annealed at 500 °C for 1 hour on a hotplate.

For MoO₃ coating, 1 mmol of ammonium heptamolybdate hexahydrate was dissolved in DI water and kept at 80 °C for 1 h, to obtain a colorless transparent solution.⁴⁵ The solution was spin coated on the FTO/glass substrate at 2000 rpm for 20 s followed by drying at 130 °C for 15 min. This process was repeated seven times and finally, the film was annealed at 380 °C for 1 hour on a hotplate. TiO₂/MoO₃ film was prepared by following the procedure explained above to sequentially deposit TiO₂ and MoO₃ films as a bilayer on the FTO/glass substrate. Finally, the devices for studying the photodetection performance were made by thermally evaporating Au top contacts on these single layer and bilayer films (Fig. 1).



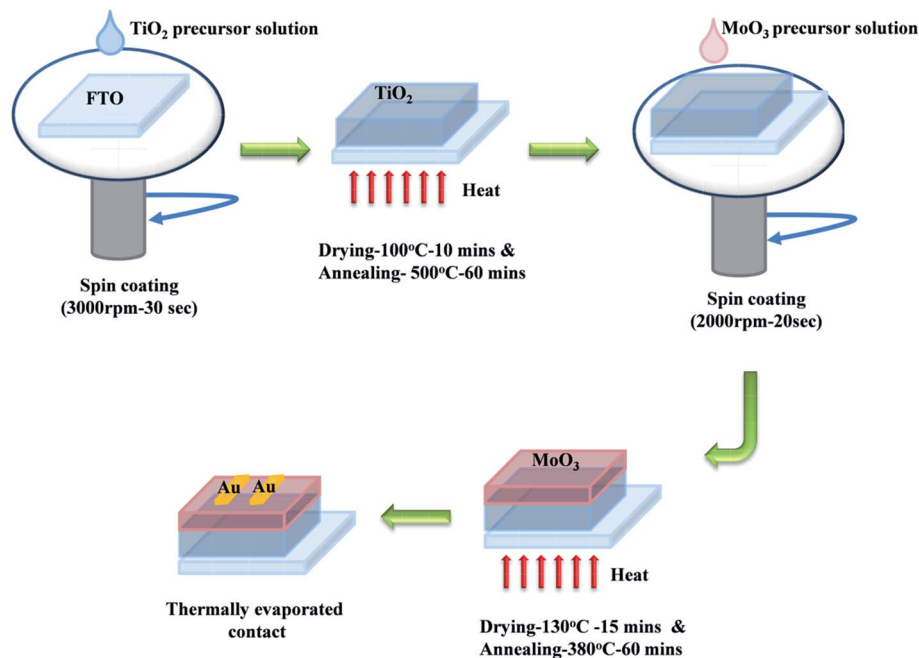


Fig. 1 Schematic illustrating the fabrication of a bilayer UV detector device.

X-ray diffraction was carried out using Cu K α radiation (Powder X-ray Diffractometer, PANalytical, Xpert) with the diffraction angle from 10 to 80°. UV-visible absorption and transmittance studies were carried out using an Agilent Cary UV-visible absorption spectrometer. FE-SEM analysis was carried out with a Quanta FEG 2000. X-ray photoelectron spectroscopy was carried out using a PHI VersaProbe III scanning XPS microprobe, Physical Electronics, USA. Electrochemical impedance spectroscopy with a two-electrode system was carried out using a Zahner Zennium electrochemical workstation in the frequency range from 10 mHz to 100 kHz with an amplitude of 10 mV. I - V and I - t characteristics of the device (active area of 0.24 cm²) were measured using a 2450 Keithley source meter unit. I - V characteristics were measured in the dark, under 1 sun illumination and illumination by a 352 nm UV lamp with an intensity of 76 μ W cm⁻². I - t characteristics were measured in the dark and under illumination by a 352 nm UV lamp with an intensity of 76 μ W cm⁻². The spectral response measurements were carried out using a Bentham PVE300 photovoltaic characterization setup with a TMC 300 monochromator.

Results and discussion

XRD analysis was carried out to confirm the crystalline nature of the samples. Fig. S1† shows the XRD patterns of TiO₂, MoO₃ and TiO₂/MoO₃ films coated on FTO/glass. The evident peaks of all the films could be matched well with the standard data of anatase-TiO₂ (ICSD-98-00-5225) and α -MoO₃ (ICSD-98-001-7296) after excluding the peaks of FTO. The grain size was calculated using the Scherrer formula and the mean grain size was 48.14 nm for TiO₂, 17.12 nm for MoO₃ and 15.17 nm for TiO₂/MoO₃ film.

The changes observed in the grain size can be correlated with the changes in the oxygen vacancy concentration as indicated by the variation in the intensity of the secondary peaks.^{46,47}

FE-SEM analysis was carried out to investigate the surface morphology and the obtained images are shown in Fig. 2a-c. The surface of the TiO₂ film was observed to have agglomerated irregular spherical nanoparticles. The MoO₃ film had a nanospike-like structure on FTO while the structure shown by the TiO₂/MoO₃ film was not exactly similar. A decrease in the length of nanospikes of MoO₃ and densely packed alignment was observed in the case of MoO₃ film deposited on TiO₂. This may be due to the substrate dependent growth of MoO₃ structures. For the bare MoO₃ film, the growth was on the smooth film of flat FTO/glass and for the TiO₂/MoO₃ film, the growth of MoO₃ was on the film of agglomerated TiO₂ nanoparticles.^{41,48} The change can also be related to the change in the oxygen vacancy concentration that affects the particle size and shape.⁴⁹ It's known in the literature that denser ZnO nanostructures show a higher photoresponse⁵⁰ and analogously, we could predict that the reduction in the length and dense packing of MoO₃ nanostructures could lead to an improved photoresponse. The FESEM cross-sectional image (Fig. S2†) of the bilayer film confirms the successive growth of TiO₂ (~300 nm thick) and MoO₃ (~260 nm thick) layers.

UV-visible transmission spectra of the films recorded in the wavelength range of 280–800 nm are shown in Fig. 2d. The spectra confirmed that both TiO₂ and TiO₂/MoO₃ films absorb light in the UV region while exhibiting high transparency to the visible light (above 400 nm). The average transmittance of 85% and 74% were observed in the visible region (400–800 nm) for the TiO₂ and TiO₂/MoO₃ films, respectively. It should also be noted that the absorption of light by the bilayer heterojunction



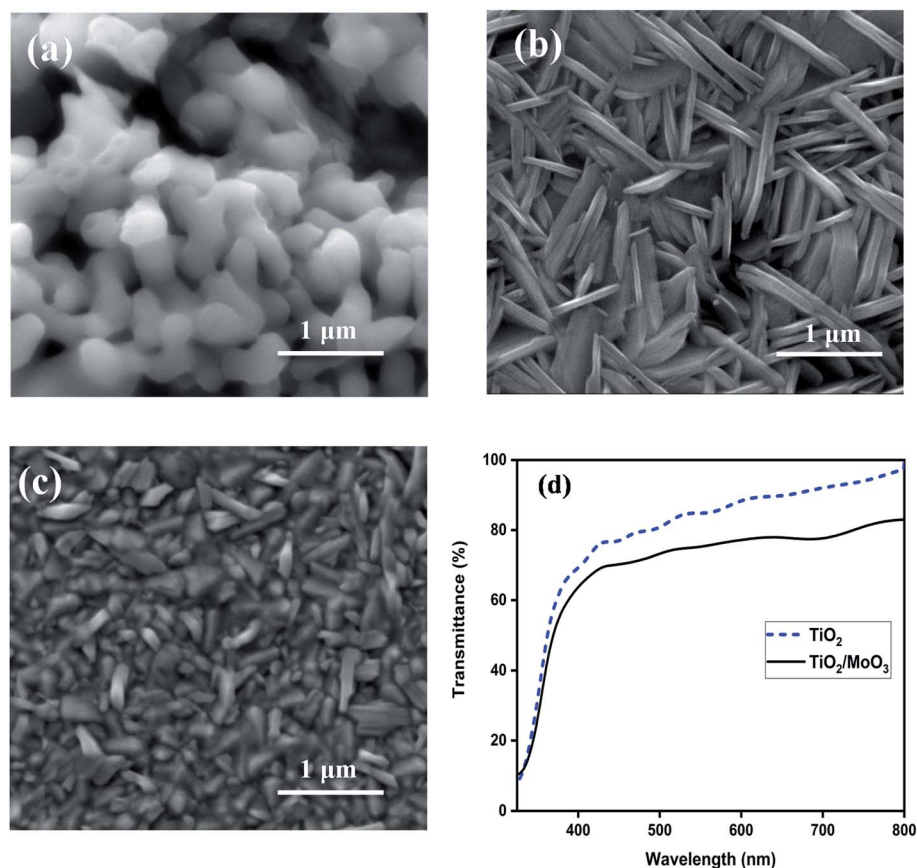


Fig. 2 FESEM images of (a) TiO_2 , (b) MoO_3 , and (c) $\text{TiO}_2/\text{MoO}_3$ films, and (d) UV-visible transmission spectra of TiO_2 and $\text{TiO}_2/\text{MoO}_3$ films.

film was significantly higher as compared to the TiO_2 film (Fig. S3[†]).

XPS analysis was carried out to confirm the chemical states, as well as the elemental composition of the samples and the results are shown in Fig. 3. The binding energies of Mo $3d_{3/2}$ and Mo $3d_{5/2}$ for the MoO_3 film were located at 236.0 eV and 232.9 eV, respectively, corresponding to the Mo^{6+} oxidation state.⁵¹ These binding energies were shifted to 235.8 eV and 232.7 eV for the Mo 3d in the $\text{TiO}_2/\text{MoO}_3$ film. Though there was

a slight shift in Mo 3d binding energies, the oxidation state of Mo^{6+} remained the same. The binding energy of the O 1s peak in the $\text{TiO}_2/\text{MoO}_3$ heterojunction was located at 530.8 eV which was also shifted downwards by 0.2 eV compared to that of the O 1s peak for MoO_3 . Both Mo 3d and O 1s peaks of the $\text{TiO}_2/\text{MoO}_3$ film were shifted to lower energies than those of the MoO_3 film and this can be attributed to the electron transfer between TiO_2 and MoO_3 in the heterojunction.²⁵

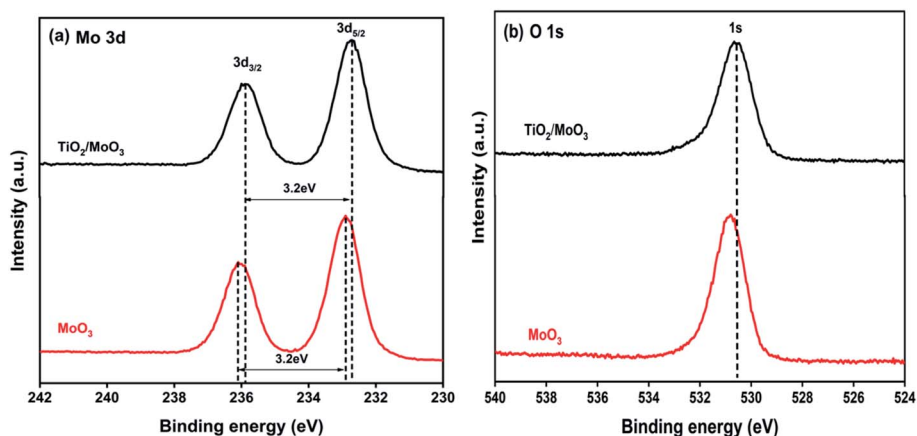


Fig. 3 XPS spectra of (a) Mo 3d and (b) O 1s of MoO_3 & $\text{TiO}_2/\text{MoO}_3$ films.



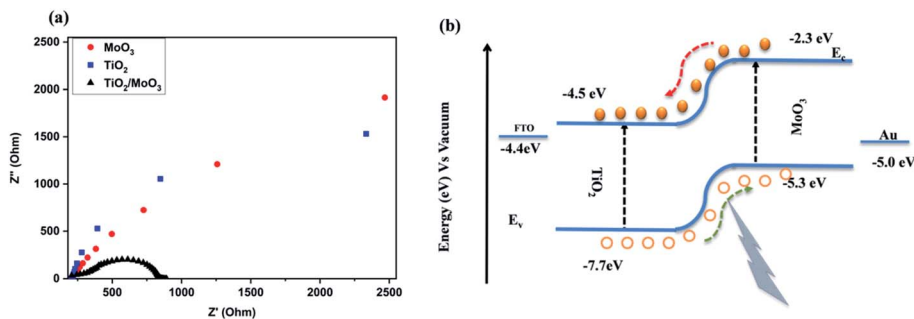


Fig. 4 (a) Electrochemical impedance spectra of MoO₃, TiO₂, and TiO₂/MoO₃ films, and (b) energy level diagram of the TiO₂/MoO₃ heterojunction.

Electrochemical impedance spectroscopy was carried out to analyse the charge transfer resistance and separation efficiency of the heterojunction. The spectra were recorded with the back contact FTO and the front contact Au in a 2 electrode configuration at room temperature under dark conditions. The obtained Nyquist plots are shown in Fig. 4a. It was evident from the spectra that the radius of the semicircle for the bilayer heterojunction device was much smaller than that of the single layer device. The small semi-circular arc for the TiO₂/MoO₃ heterojunction indicated a smaller charge transfer resistance which can be related to the faster charge transfer across the interface. This also corresponds to the higher efficiency of the heterojunction for the separation of carriers.²⁶ The equivalent circuit model used for fitting and the resistance values are given in Fig. S9 and Table S1.†

To understand the electrical behaviour of the TiO₂/MoO₃ bilayer, two probe *I*-*V* measurements were done in the dark at room temperature with Au and FTO as a front and back contact, respectively, and the results are shown in Fig. 5a. Before the fabrication of the heterojunction, TiO₂ and MoO₃ were studied separately. The *I*-*V* curve of the TiO₂ based device exhibited a Schottky behaviour (Fig. S4a†), while the MoO₃ based device showed ohmic conduction (Fig. S5†). However, the *I*-*V* curve of the bilayer device exhibited a diode-like behaviour. The diode parameters were analyzed using the ideal diode equation⁵²

$$I = I_s \left[\exp\left(\frac{qV}{nkT}\right) - 1 \right] \quad (1)$$

$$I_s = AA^*T^2 \exp\left(\frac{-q\phi_b}{kT}\right) \quad (2)$$

$$\phi_b = \frac{kT}{q} \ln\left(\frac{AA^*T^2}{I_s}\right) \quad (3)$$

where I_s is the reverse saturation current, q is the electron charge, n is the ideality factor, k is the Boltzmann constant, T is the temperature, A is the active area, A^* is the effective Richardson constant and ϕ_b is the barrier height. The ideality factor was calculated from the slope of the semi-logarithmic plot and reverse saturation current was calculated from the current intercept at $V = 0$. The Richardson constant was taken based on the literature ($1200 \text{ A cm}^{-2} \text{ K}^{-2}$ for TiO₂).⁵³ The obtained parameters are listed in Table S2.†

The reverse current started increasing slightly at -0.4 V , but the corresponding forward current was one order of magnitude higher which is comparable to the earlier observations with other heterojunction devices.^{40,54,55} The ideality factor which was greater than unity (2.98) indicates that the charge transport mechanism deviated from the thermionic emission theory applicable for an ideal diode. To further investigate the charge transport mechanism, power-law compliance was considered and the forward *I*-*V* data were plotted in a log *I*-log *V* scale⁵⁶

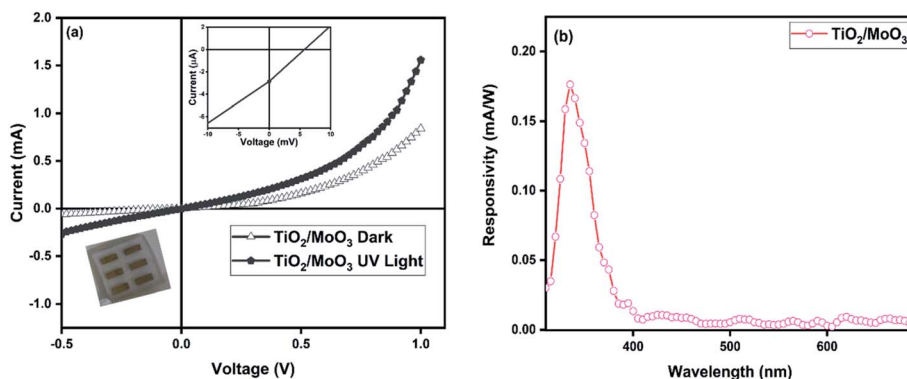


Fig. 5 Device characteristics of the bilayer device: (a) *I*-*V* characteristics in the dark and under UV illumination, insets: *I*-*V* curve under UV light, enlarged around 0 V & a photograph of the device, and (b) spectral responsivity.



(Fig. S6†). The plot was fitted well with the power-law expression mentioned below with two different slopes.

$$I = KV^m$$

where K is a constant and m is the exponent derived from the slope of the curve.

At lower voltages (0.02–0.28 V), the exponent determined was close to unity with a correlation coefficient of 0.99645. This indicates that the current varied linearly with the voltage and the charge transport was governed by Ohm's law. Hence, the thermally generated free electrons were responsible for the current produced at lower bias. Furthermore, an enhancement in the current was observed with the increase in the voltage, indicating that the injected carriers contribute to the rise in current.^{54,57,58} At higher voltages (0.32–1.0 V), the exponent was determined to be 2.18 with a correlation coefficient of 0.99589. This indicates that the space charge limited charge transport became pre-dominant, where the number of injected carriers becomes comparable to the thermally generated carriers. The value of the exponent is slightly greater than 2 and this suggests that the injected carriers have just entered into the space charge region.⁵⁹

I - V characteristics measured under UV light illumination (352 nm) clearly showed the photoconductive nature of the devices. The photocurrent (the difference between the current under UV illumination and dark) of the bilayer device was observed as 1.94 μA at 0 V and it reached a value of 840 μA at a reverse bias of -1 V. The photocurrent without bias was 40 times higher and with bias, it was 3 times higher when compared to that of the single layer TiO_2 nanoparticle based device for which the photocurrent was only 42 nA at 0 V and 228 μA at -1 V (Fig. S4†). The devices exhibited photocurrent even under zero bias conditions, due to the photovoltaic effect under UV illumination. The I - V curves enlarged around 0 V showed a photovoltaic behaviour upon UV illumination for both single and bilayer devices (inset in Fig. 4a and S4a†). The $\text{TiO}_2/\text{MoO}_3$ bilayer device showed an open-circuit voltage (V_{oc}) of 5.7 mV and a short-circuit current (I_{sc}) of 2.856 μA , while for the TiO_2 single layer, the values of V_{oc} and I_{sc} were 4.4 mV and 0.953 μA , respectively. It can be attributed to the built-in potential at the FTO/ TiO_2 and $\text{TiO}_2/\text{MoO}_3$ interfaces arising due to alignment of energy bands, which act to separate the photogenerated carriers efficiently.^{33,60}

Based on the I - V characteristics, the photoresponsivity (R), detectivity (D) and external quantum efficiency (EQE) were calculated using the formulae mentioned below³⁰

$$R = \frac{I_{\text{light}} - I_{\text{dark}}}{P_{\text{ill}}} \quad (5)$$

$$D = \frac{R}{(2eJ_{\text{d}})^{1/2}} \quad (6)$$

$$\text{EQE}_{\lambda}\% = \frac{1240R}{\lambda} 100$$

where $P_{\text{ill}} = I_{\text{ill}} \times A_0$, I_{ill} is the illumination power density, A_0 is the active area, e is the charge of an electron, J_{d} is the dark current density, R is the photoresponsivity, and λ is the excitation wavelength (352 nm). For the $\text{TiO}_2/\text{MoO}_3$ device, R , D , and EQE were estimated to be 108 mA W^{-1} , 2.26×10^{10} Jones and 38.04% without bias, while at -1 V bias the values were found to be 46.05 A W^{-1} , 2.84×10^{12} Jones and 1.62 $\times 10^4\%$, respectively. When compared to the single layer device, the value of R was 40 times higher without bias and 3 times higher with bias. D and EQE were observed to be 1 order higher in magnitude with and without bias.

When the I - V measurements were carried out under 1 sun white light illumination, a negligible photocurrent was observed (Fig. S7†). This confirms that the response was significant only with illumination by UV light of the wavelength corresponding to the bandgap of the film, concurrent with the absorption spectrum. The spectral responsivity measured with 0 V bias (Fig. 5b) has again proven the excellent UV-light detection selectivity of the bilayer heterojunction film, with the highest response shown in the spectral range between 300 and 385 nm. The results confirmed that the device has the ability to work as a visible blind UV detector.

To further understand the self-powering ability, I - t measurements were performed under 0 V bias with an interval of 30 s between UV illuminations (Fig. 6). The photocurrent, rise time and fall time for our bilayer film, as estimated from the I - t characteristics, were found to be 1.94 μA , 1.82 s and 1.42 s, respectively. The high photocurrent upon UV light illumination can be attributed to the photo-desorption of oxygen species on the surface. The dynamics associated with the photo-adsorption/desorption process results in the increase of free carrier density and consequently the photocurrent. The apparent overshoot phenomenon could be related to the charge accumulation at the metal/semiconductor interface.⁶¹ It is noteworthy that the time taken to reach the steady state was less than 2 s and the response remained unaltered even after 10 cycles. To understand the photoresponse under external bias, I -

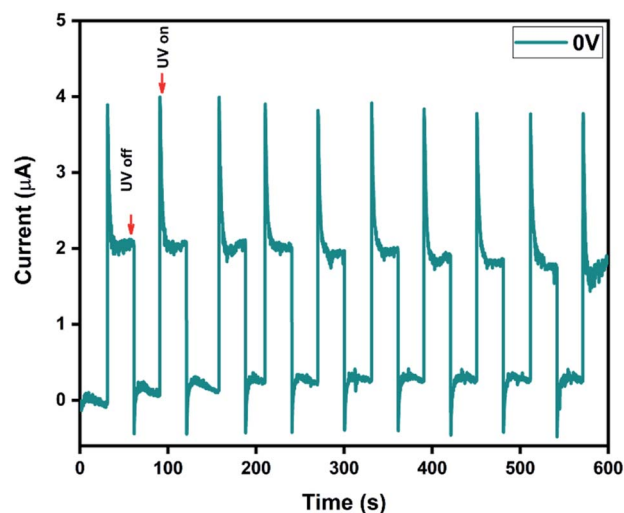


Fig. 6 I - t characteristics of the self-powered bilayer device.



Table 1 Performance of TiO₂ based self-powered UV photodetectors

Device	Wavelength (nm)	Intensity of the source (mW cm ⁻²)	Bias (V)	Photoresponsivity (mA W ⁻¹)	Detectivity (Jones)	Rise time	Fall time	Ref.
TiO ₂ /Ag	350	1.2	0	1.1	—	2 ms	47 ms	64
TiO ₂ /Ag	350	0.038	0	32.5	6 × 10 ⁹	44 ns	1.85 μs	65
NiO/TiO ₂	370	0.2	0	860	1.6 × 10 ¹³	3 ms	60 ms	60
NiO/TiO ₂ /TiO _x	380	0.5	0	5.66	2.5 × 10 ¹²	<0.1 s	<0.1 s	39
SnS _x /TiO ₂	365	2	0	5.85	6.54 × 10 ¹⁰	3 ms	25 ms	37
PANI/TiO ₂	320	0.87	0	3.6	3.9 × 10 ¹¹	1.2 ms	22.8 ms	66
SnO ₂ /TiO ₂	365	40	0	600	—	0.02 s	0.004 s	38
Au/P3HT/TiO ₂	350	1	0	0.25	2.9 × 10 ¹⁰	0.72 s	0.5 s	31
ZnO/Ga ₂ O ₃	266	1.6	-0.1	9.7	2.58 × 10 ¹²	100 μs	900 μs	67
TiO ₂	352	0.076	0	2.68	3.34 × 10 ⁹	—	—	This work
TiO ₂ /MoO ₃	352	0.076	0	108	2.26 × 10 ¹⁰	1.82 s	1.48 s	

t measurements were also performed at -1 V (Fig. S8†). The photocurrent, rise time and fall time were found to be 840 μA, 1.22 s and 1.4 s, respectively. The cycling stability of the device with an applied voltage of -1 V was observed to be slightly inferior, and the dark current was found to increase after 5 cycles. This may be attributed to the redistributed electric field caused by the applied bias, which can induce a decrease in voltage dependent carrier separation and transport.^{36,62,63}

However, it's noteworthy that the stability loss observed here was only 5% and there was almost no loss when the device was operated under self-powered mode.

As compared to the literature, our single layer device with the configuration of FTO/TiO₂/Au itself showed better performance, under bias conditions lower than the reported ones (Table S3†). The detection capability of the TiO₂/MoO₃ bilayer device was found to be remarkable in terms of the bias conditions as well as the responsivity when compared with the other reported anatase TiO₂ based UV detectors. Also, the rise time is 3 times less and the fall time is 12 times less (Table S4†) when compared to the values reported earlier for the TiO₂/MoO₃ UV detector.⁴⁰ This faster response may be attributed to the effective collection of the photogenerated carriers and the consequent recombination loss.⁶⁸ This corresponds to the effective junction formation between TiO₂ and MoO₃ layers and the efficient separation of photogenerated carriers.⁶⁹

It can be seen from Table 1 that the performance of our devices was better when compared with many such reports on TiO₂ based self-powered UV detectors. The bilayer device had higher responsivity as well as detectivity at zero bias, even under the relatively weakest illumination. It stands out as the best so far among the devices studied under low light conditions. It is also noteworthy that the photoresponsivity of our self-powered bilayer device at zero bias itself is comparable to that of commercial Si photodetectors (0.1–0.2 A W⁻¹)⁷⁰ and with a bias of -1 V, it is much higher. However, the slow response speed of the device under 0 V bias was related to the oxygen vacancies.⁷¹ The oxygen vacancies produce subgap states and can act as recombination centers. The subgap states hinder carrier excitation upon UV illumination⁷² and they also act as trap centres to the carriers upon the shut off of UV irradiation, resulting in a slow response.⁷³

The observed photoresponse is also better than that of most self-powered photodetectors based on various metal oxides such as ZnO, Ga₂O₃, NiO, and SnO₂, which perform under illumination conditions much stronger than ours.^{60,65,74–77} Overall, the superior performance of the bilayer device, as evident from the discussion above, can be ascribed to the formation of an effective heterojunction between TiO₂ and MoO₃. The built-in potential at the interface of TiO₂/MoO₃ leads to an enhanced separation efficiency as well as an effective charge transfer, and enables the device to be self powered.⁶⁰ The improved and faster response can be related to the reduced recombination. As already mentioned, these improvements could also be due to the morphological factor of having densely aligned nanostructures. It is well known that the adsorption/photodesorption of O₂ species upon UV illumination⁷⁸ followed by the transport of photogenerated carriers⁷⁹ plays a vital role in UV detection performance. The depletion region would shrink upon application of the reverse bias due to the hole-trapping at the surface. This increases the electron tunneling, and the electron migration in large numbers is related to the EQE over 100%.^{80,81} It can be proposed that the densely aligned nanostructures could expose multiple junctions sensitive to the UV light and could also provide efficient pathways for the transport of carriers, leading to the enhancement in the photocurrent and EQE.^{50,82} When compared to previous reports, the improvement can also be related to the mobility of the anatase TiO₂ used in this work which is known to be better than that of rutile TiO₂.⁸³ Though the densely aligned morphology offers better performance in terms of faster response and high photocurrent, the photocurrent stability with time can be further improved by tuning the surface area.⁸⁴ There is room for achieving better performance by more carefully optimizing the particle size, morphology, crystallinity, and thickness of the films.

Conclusion

In summary, we have fabricated and studied a fully solution processed transparent TiO₂/MoO₃ bilayer thin film device with an anatase TiO₂ nanoparticle based bottom layer for the first time, for self-powered UV detector application. XPS,



electrochemical measurements and current–voltage characteristics of the bilayer have confirmed the formation of an effective heterojunction diode, where the interaction of the energy bands could lead to an efficient charge transfer. The charge transport mechanism of the bilayer device was governed by Ohm's law at low voltages and space charge limited conduction at high voltages. The bilayer device exhibited a very good photoresponsivity and detectivity under 352 nm UV light illumination of very low intensity. The performance was on par with other TiO₂ based heterojunctions reported as self-powered UV detectors, showing a photoresponse even under zero bias conditions which was further enhanced at a low bias of -1 V. The self-powered performance of the device was attributed to the built-in potential of the TiO₂/MoO₃ interface that led to the separation of photogenerated carriers without bias. The improved detection capability exhibited by the TiO₂/MoO₃ bilayer device in terms of the low bias conditions, high responsivity, high EQE and faster response can be mainly attributed to the better carrier collection as well as the reduced recombination across the heterojunction. The morphological factors of nanostructures and the better mobility of anatase TiO₂ nanoparticles could also be the reasons behind the observed enhancement in the performance. The results suggest that the bilayer heterojunction, where the TiO₂ nanoparticle film is the platform with the α -MoO₃ film on the top, is a promising candidate for self-powered UV detector application with its high transparency as well as the best-in-class photoresponse under low bias and weak illumination conditions.

Author contributions

BE carried out all the experiments including synthesis, device fabrication and characterization with initial help from MD. SVB conceptualized the project and supervised the research work. BE and SVB discussed the results, prepared the manuscript and carried out the required revisions.

Conflicts of interest

There are no conflicts of interest to declare.

Acknowledgements

The authors acknowledge the financial support of SRMIST towards this research work and all the in-campus characterization facilities such as the UV-visible spectrometer at the Department of Physics & Nanotechnology, XPS and FE-SEM facility at NRC and the XRD facility with the support from MNRE (Project No. 31/03/2014-15/PVSE-R&D), Govt. of India. We thank Dr S. Harinipriya and coworkers for the help with the access to electrochemical impedance measurements. We also thank Dr P. Malar and coworkers for the help with the use of a thermal evaporator. We thank Prof. B. Neppolian and coworkers for the access to the EQE facility supported by Department of Science and Technology – Solar Energy Research Initiative (DST; SERB), File No. DST/SERI/FR/2016/S170. We would also like to acknowledge the research grants from DST-

SERB [File. No. (ECR/2015/000513)] and MNRE [File. No. 31/03/2014-15/PVSE-R&D], Govt. of India.

References

- 1 K. D. Leedy, M. Higashiwaki, R. L. Peterson and S. James, *APL Mater.*, 2019, **7**, 022201.
- 2 H. Ohta, M. Hirano, K. Nakahara, H. Maruta, T. Tanabe, M. Kamiya, T. Kamiya and H. Hosono, *Appl. Phys. Lett.*, 2003, **83**, 1029–1031.
- 3 C. C. Shih, W. Y. Lee, Y. C. Chiu, H. W. Hsu, H. C. Chang, C. L. Liu and W. C. Chen, *Sci. Rep.*, 2016, **6**, 20129.
- 4 Z. Q. Zheng, J. D. Yao, B. Wang and G. W. Yang, *Sci. Rep.*, 2015, **5**, 11070.
- 5 M. Pavan, S. Rühle, A. Ginsburg, D. A. Keller, H. N. Barad, P. M. Sberna, D. Nunes, R. Martins, A. Y. Anderson, A. Zaban and E. Fortunato, *Sol. Energy Mater. Sol. Cells*, 2015, **132**, 549–556.
- 6 L. Ahtapodov, B. Fimland, K. Kishino and H. Weman, *Nano Lett.*, 2019, **19**, 1649–1658.
- 7 S. Abbas, D. Ban and J. Kim, *Sens. Actuators, A*, 2019, **293**, 215–221.
- 8 J. Chen, Z. Li, F. Ni, W. Ouyang and X. Fang, *Mater. Horiz.*, 2020, **7**, 1828–1833.
- 9 C. C. Forin, M. Purica, E. Budianu and P. Schiopu, *Proc. Int. Semicond. Conf. CAS*, 2012, **1**, 131–134.
- 10 Z. Li, M. K. Joshi, J. Chen, Z. Zhang, Z. Li and X. Fang, *Adv. Funct. Mater.*, 2020, **30**, 2005291.
- 11 G. S. Mital and T. Manoj, *Chin. Sci. Bull.*, 2011, **56**, 1639–1657.
- 12 X. Chen and S. S. Mao, *Chem. Rev.*, 2007, **107**, 2891–2959.
- 13 M. Quintana, T. Edvinsson, A. Hagfeldt and G. Boschloo, *J. Phys. Chem. C*, 2007, **111**, 1035–1041.
- 14 J. Sobuś, G. Burdziński, J. Karolczak, J. Idígoras, J. A. Anta and M. Ziółek, *Langmuir*, 2014, **30**, 2505–2512.
- 15 T. Ling, J. G. Song, X. Y. Chen, J. Yang, S. Z. Qiao and X. W. Du, *J. Alloys Compd.*, 2013, **546**, 307–313.
- 16 D. Friedmann, C. Mendive and D. Bahnemann, *Appl. Catal., B*, 2010, **99**, 398–406.
- 17 B. Liu, X. Zhao, J. Yu, I. P. Parkin and A. Fujishima, *J. Photochem. Photobiol., C*, 2019, **39**, 1–57.
- 18 S. Mishra, P. Yogi, P. R. Sagdeo and R. Kumar, *ACS Appl. Energy Mater.*, 2018, **1**, 790–798.
- 19 Y. Li, C. Cai, Y. Gu, W. Cheng, W. Xiong and C. Zhao, *Appl. Surf. Sci.*, 2017, **414**, 34–40.
- 20 M. Sun, J. Hu, C. Zhai, M. Zhu and J. Pan, *Electrochim. Acta*, 2017, **245**, 863–871.
- 21 I. A. De Castro, R. S. Datta, J. Z. Ou and A. Castellanos-gomez, *Adv. Mater.*, 2017, **29**, 1701619.
- 22 D. Kwak, M. Wang, K. J. Koski, L. Zhang, H. Sokol, R. Maric and Y. Lei, *ACS Appl. Mater. Interfaces*, 2019, **11**, 10697–10706.
- 23 S. Sonu, V. Dutta, S. Sharma, P. Raizada, A. Hosseini-Bandegharai, V. Kumar Gupta and P. Singh, *J. Saudi Chem. Soc.*, 2019, **23**, 1119–1136.
- 24 W. Yang, K. Hu, F. Teng, J. Weng, Y. Zhang and X. Fang, *Nano Lett.*, 2018, **18**, 4697–4703.



- 25 M. Lu, C. Shao, K. Wang, N. Lu, X. Zhang, P. Zhang, M. Zhang, X. Li and Y. Liu, *ACS Appl. Mater. Interfaces*, 2014, **6**, 9004–9012.
- 26 S. Sun, X. Liao, X. Pu, Y. Sun, G. Yin, Y. Yao, Z. Huang and X. Pu, *RSC Adv.*, 2017, **7**, 22983–22989.
- 27 S. Paek, J. Kang, H. Jung, S. Hwang and J. Choy, *Chem. Commun.*, 2009, **48**, 7536–7538.
- 28 N. K. Shrestha, Y. Nah and P. Schmuki, *Chem. Commun.*, 2009, **15**, 2008–2010.
- 29 Q. Chen, Y. Zhang, T. Zheng, Z. Liu, L. Wu, Z. Wang and J. Li, *Nanoscale Adv.*, 2020, **2**, 2705–2712.
- 30 M. Moun, M. Kumar, M. Garg, R. Pathak and R. Singh, *Sci. Rep.*, 2018, **1**, 1–10.
- 31 L. Zheng, X. Deng, Y. Wang, J. Chen, X. Fang, L. Wang, X. Shi and H. Zheng, *Adv. Funct. Mater.*, 2020, **30**, 2001604.
- 32 X. Li, S. Gao, G. Wang, Z. Xu, S. Jiao, D. Wang, Y. Huang, D. Sang, J. Wang and Y. Zhang, *J. Mater. Chem. C*, 2020, **8**, 1353–1358.
- 33 Y. Shen, X. Yan, Z. Bai, X. Zheng, Y. Sun, Y. Liu, P. Lin, X. Chen and Y. Zhang, *RSC Adv.*, 2015, **5**, 5976–5981.
- 34 J. Xing, H. Wei, E. J. Guo and F. Yang, *J. Phys. D: Appl. Phys.*, 2011, **44**, 375104.
- 35 A. M. Selman and Z. Hassan, *Superlattices Microstruct.*, 2015, **83**, 549–564.
- 36 D. Zhang, C. Liu, B. Yin, R. Xu, J. Zhou, X. Zhang and S. Ruan, *Nanoscale*, 2017, **9**, 9095–9103.
- 37 J. Chen, J. Xu, S. Shi, R. Cao and D. Liu, *ACS Appl. Mater. Interfaces*, 2020, **12**, 23145–23154.
- 38 C. Gao, X. Li, X. Zhu, L. Chen, Y. Wang, F. Teng, Z. Zhang, H. Duan and E. Xie, *J. Alloys Compd.*, 2014, **616**, 510–515.
- 39 Y. Gao, J. Xu, S. Shi, H. Dong, Y. Cheng, C. Wei, X. Zhang, S. Yin and L. Li, *ACS Appl. Mater. Interfaces*, 2018, **10**, 11269–11279.
- 40 B. Yin, Y. Zhang, K. Li, J. Zhou, C. Liu, M. Zhang and S. Ruan, *Nanotechnology*, 2019, **30**, 465501.
- 41 T. Zhai and J. Yao, *One-dimensional nanostructures: principles and applications*, Wiley&Sons, United States, 2012.
- 42 P. Gao, H. Wang, Z. Sun, W. Han, J. Li and J. Ye, *Appl. Phys. Lett.*, 2013, **103**, 213505.
- 43 J. Szlufcik, J. Majewski, A. Buczkowski, J. Radojewski, L. Jędral and E. B. Radojewska, *Sol. Energy Mater.*, 1989, **18**, 241–252.
- 44 M. Sasani Ghamsari and A. R. Bahramian, *Mater. Lett.*, 2008, **62**, 361–364.
- 45 S. Murase and Y. Yang, *Adv. Mater.*, 2012, **24**, 2459–2462.
- 46 R. Parra, C. M. Aldao, J. A. Varela and M. S. Castro, *J. Electroceram.*, 2005, **14**, 149–156.
- 47 S. Aškrabič, Z. Dohčevič-Mitrovič, A. Kremenovič, N. Lazarevič, V. Kahlenberg and Z. V. Popovič, *J. Raman Spectrosc.*, 2012, **43**, 76–81.
- 48 J. Song and S. Lim, *J. Phys. Chem. C*, 2007, **111**, 596–600.
- 49 R. Mehmood, X. Wang, P. Koshy, J. L. Yang and C. C. Sorrell, *CrystEngComm*, 2018, **20**, 1536–1545.
- 50 V. Postica, M. Hoppe, J. Grottrup, P. Hayes, V. Robisch, D. Smazna, R. Adelung, B. Viana, P. Aschehoug, T. Pauporté and O. Lupan, *Solid State Sci.*, 2017, **71**, 75–86.
- 51 S. Xie, D. Cao, Y. She, H. Wang, J. W. Shi, M. K. Leung and C. Niu, *Chem. Commun.*, 2018, **54**, 7782–7785.
- 52 D. K. Schroder, *Semiconductor Material and Device Characterization*, Wiley&Sons, United States, 2015.
- 53 H. A. Altuntas, A. Bengi, U. Aydemir, T. Asar, S. S. Cetin, I. Kars, S. Altindal and S. Ozcelik, *Mater. Sci. Semicond. Process.*, 2009, **12**, 224–232.
- 54 P. C. Ramamurthy, W. R. Harrell, R. V. Gregory, B. Sadanadan and A. M. Rao, *Polym. Eng. Sci.*, 2004, **44**, 28–33.
- 55 P. Singh, P. K. Rout, H. Pandey and A. Dogra, *J. Mater. Sci.*, 2018, **53**, 4806–4813.
- 56 H. Elangovan, A. V. Kesavan, K. Chattopadhyay and P. C. Ramamurthy, *J. Appl. Phys.*, 2019, **125**, 245302.
- 57 A. Wahl and J. Augustynski, *J. Phys. Chem. B*, 1998, **5647**, 7820–7828.
- 58 M. A. Lampert and R. B. Schilling, in *Semiconductors and Semimetals*, ed. R. K. Williardson and A. Beer, Elsevier, Amsterdam, 1970, vol. 6, ch. 1, pp. 1–96.
- 59 X. Qiao, J. Chen and D. Ma, *J. Phys. D: Appl. Phys.*, 2010, **43**, 215402.
- 60 R. Cao, J. P. Xu, S. Shi, J. Chen, D. Liu, Y. Bu, X. S. Zhang, S. Yin and L. Li, *J. Mater. Chem. C*, 2020, **8**, 9646–9654.
- 61 Y. Fang, Y. Hou, Y. Hu and F. Teng, *Materials*, 2015, **8**, 4050–4060.
- 62 W. P. Dumke, *Appl. Phys. Lett.*, 1967, **11**, 314–316.
- 63 S. Luryi in *Hot Electrons in Semiconductors*, ed. N. Balkan, Clarendon Press, Oxford, 1998, vol. 6, ch. 17, pp. 389–395.
- 64 J. Xu, W. Yang, H. Chen, L. Zheng, M. Hu, Y. Li and X. Fang, *J. Mater. Chem. C*, 2018, **6**, 3334–3340.
- 65 H. Fang, C. Zheng, L. Wu, Y. Li, J. Cai, M. Hu, X. Fang, R. Ma, Q. Wang and H. Wang, *Adv. Funct. Mater.*, 2019, **29**, 1809013.
- 66 L. Zheng, P. Yu, K. Hu, F. Teng, H. Chen and X. Fang, *ACS Appl. Mater. Interfaces*, 2016, **8**, 33924–33932.
- 67 B. Zhao, W. Fei, C. Hongyu, Z. Lingxia, S. Longxing, Z. Dongxu and X. Fang, *Adv. Funct. Mater.*, 2017, **27**, 1700264.
- 68 S. Li, Y. Zhang, W. Yang, H. Liu and X. Fang, *Adv. Mater.*, 2020, **32**, 1905443.
- 69 M. Razeghi and A. Rogalski, *J. Appl. Phys.*, 1996, **79**, 7433–7473.
- 70 W. Zheng, X. Li, C. Dong, X. Yan and G. He, *RSC Adv.*, 2014, **4**, 44868–44871.
- 71 P. Ivanoff Reyes, C. J. Ku, Z. Duan, Y. Xu, E. Garfunkel and Y. Lu, *Appl. Phys. Lett.*, 2012, **101**, 2010–2014.
- 72 K. Y. Chen, S. P. Chang and C. hung Lin, *RSC Adv.*, 2019, **9**, 87–90.
- 73 Q. Chen, H. Ding, Y. Wu, M. Sui, W. Lu, B. Wang, W. Su, Z. Cui and L. Chen, *Nanoscale*, 2013, **5**, 4162–4165.
- 74 J. Chen, W. Ouyang, W. Yang, J. H. He and X. Fang, *Adv. Funct. Mater.*, 2020, **30**, 1909909.
- 75 Y. Wang, C. Wu, D. Guo, P. Li, S. Wang, A. Liu, C. Li, F. Wu and W. Tang, *ACS Appl. Electron. Mater.*, 2020, **2**, 2032–2038.
- 76 Y. Wang, L. Li, H. Wang, L. Su, H. Chen, W. Bian, J. Ma, B. Li, Z. Liu and A. Shen, *Nanoscale*, 2020, **12**, 1406–1413.
- 77 B. Deka Boruah, *Nanoscale Adv.*, 2019, **1**, 2059–2085.



- 78 D. Gedamu, I. Paulowicz, S. Kaps, O. Lupan, S. Wille, G. Haidarschin, Y. K. Mishra and R. Adelung, *Adv. Mater.*, 2014, **26**, 1541–1550.
- 79 Y. K. Mishra, G. Modi, V. Cretu, V. Postica, O. Lupan, T. Reimer, I. Paulowicz, V. Hrkac, W. Benecke, L. Kienle and R. Adelung, *ACS Appl. Mater. Interfaces*, 2015, **7**, 14303–14316.
- 80 S. Dhar, T. Majumder and S. P. Mondal, *ACS Appl. Mater. Interfaces*, 2016, **8**, 31822; W. Yin, J. Yang, K. Zhao, A. Cui, J. Zhou, W. Tian, W. Li, Z. Hu and J. Chu, *ACS Appl. Mater. Interfaces*, 2020, **12**, 11797–11805.
- 81 W. Yin, J. Yang, K. Zhao, A. Cui, J. Zhou, W. Tian, W. Li, Z. Hu and J. Chu, *ACS Appl. Mater. Interfaces*, 2020, **12**, 11797–11805.
- 82 V. Postica, I. Paulowicz, O. Lupan, F. Schütt, N. Wol and A. Cojocaru, *Vacuum*, 2019, **166**, 393–398.
- 83 Y. Li, W. Xie, X. Hu, G. Shen, X. Zhou, Y. Xiang, X. Zhao and P. Fang, *Langmuir*, 2010, **26**, 591–597.
- 84 R. Gottesman, A. Song, I. Levine, M. Krause, A. T. M. N. Islam, D. Abou-Ras, T. Dittrich, R. Van de Krol and A. C. Chemseddine, *Adv. Funct. Mater.*, 2020, **30**, 1910832.

