



Cite this: *Nanoscale Adv.*, 2021, **3**, 1047

## Solution processed transparent anatase $\text{TiO}_2$ nanoparticles/ $\text{MoO}_3$ nanostructures heterojunction: high performance self-powered UV detector for low-power and low-light applications†

Bhuvaneshwari Ezhilmaran,<sup>ab</sup> M. Dhanasekar,<sup>†ab</sup> and S. Venkataprasad Bhat  <sup>ab</sup>

Ultraviolet (UV) photodetectors are considered as the major players in energy saving technology of the future. Efforts are needed to further develop such devices, which are capable of operating efficiently at low driving potential as well as with weak illumination. Herein, we report an all-oxide, highly transparent  $\text{TiO}_2/\text{MoO}_3$  bilayer film, with nanoparticulate anatase  $\text{TiO}_2$  as the platform, fabricated by a simple solution based method and demonstrate its use in UV photodetection. Photoconductivity measurement with 352 nm light reveals the self-powered UV detection capability of the device due to the built-in potential at the bilayer interface. The device exhibits a high photoresponsivity ( $46.05 \text{ A W}^{-1}$ ), detectivity ( $2.84 \times 10^{12} \text{ Jones}$ ) and EQE (16.223%) even with a weak illumination of  $76 \mu\text{W cm}^{-2}$ , at a low bias of only  $-1 \text{ V}$ . The self-powered performance of the bilayer device is comparable to that of commercial Si photodetectors as well as other such UV detectors reported based on metal oxide heterojunctions. The improved and faster photoresponse shown by the device is due to the formation of an effective heterojunction, as evidenced by XPS, electrochemical and  $I-V$  studies. It can be further attributed to the better charge transport through the densely aligned nanostructures, reduced recombination and the better mobility of anatase  $\text{TiO}_2$  nanoparticles. The performance is best-in-class and proves the potential of the transparent heterojunction to be used in highly responsive, self-powered UV detectors for low bias, low light applications.

Received 20th September 2020

Accepted 17th December 2020

DOI: 10.1039/d0na00780c

rsc.li/nanoscale-advances

## Introduction

Wide bandgap metal oxide semiconductors have been widely accepted for their potential use in a variety of applications due to their low cost, less toxicity and earth abundant nature.<sup>1</sup> Enormous efforts have been made towards the evolution of transparent electronic devices employing wide bandgap

transparent semiconductors. Transparent diodes,<sup>2</sup> transparent transistors,<sup>3</sup> sensors,<sup>4</sup> displays, solar cells,<sup>5</sup> UV LEDs<sup>6</sup> and photodetectors<sup>7,8</sup> have been demonstrated. Properties such as optical transparency, environmental-friendly nature, thermal and chemical stability can make them more attractive among semiconductor-based devices.<sup>2,9</sup>

$\text{TiO}_2$  is a popular n-type wide bandgap semiconductor serving as a prospective candidate in energy conversion applications because of its chemical and thermal stability, low cost and environmental-friendly characteristics. Due to these remarkable properties,  $\text{TiO}_2$  is used in many applications including photodetectors, photocatalysis, photovoltaics, electrochromic devices and sensors.<sup>10-12</sup> Even though both  $\text{ZnO}$  and  $\text{TiO}_2$  have been explored for most optoelectronic applications, the photo and chemical stability of  $\text{TiO}_2$  are known to be higher than those of  $\text{ZnO}$ .<sup>12-15</sup> However, it was observed that the charge separation efficiency of  $\text{TiO}_2$  in applications such as photocatalytic activity gets strongly suppressed due to the recombination of photogenerated carriers.<sup>16,17</sup> Fabricating a heterojunction with  $\text{TiO}_2$  is one of the important techniques for reducing the recombination, thereby increasing the lifetime of the carriers, enhancing the charge separation efficiency and improving the charge transfer efficiency. The heterojunction of

<sup>a</sup>SRM Research Institute, SRM Institute of Science and Technology, Kattankulathur, Kancheepuram-603203, India. E-mail: venkatab@srmist.edu.in

<sup>b</sup>Department of Physics and Nanotechnology, SRM Institute of Science and Technology, Kattankulathur, Kancheepuram-603203, India

† Electronic supplementary information (ESI) available: XRD patterns of  $\text{TiO}_2$ ,  $\text{MoO}_3$  and  $\text{TiO}_2/\text{MoO}_3$  films; FE-SEM cross-sectional image of the bilayer film; UV absorption spectra;  $I-V$  and  $I-t$  characteristics of the  $\text{TiO}_2$  single layer device;  $I-V$  characteristics of the  $\text{MoO}_3$  single layer device, log  $I$ -log  $V$  plot of the bilayer device; equivalent circuit model for EIS; table comparing the resistance values of  $\text{TiO}_2$ ,  $\text{MoO}_3$  and  $\text{TiO}_2/\text{MoO}_3$  films as calculated using the EIS equivalent circuit model; calculated diode parameters of the bilayer device; table comparing the performance of  $\text{TiO}_2$  based UV detectors; table comparing the performance of  $\text{TiO}_2/\text{MoO}_3$  based UV detectors; statistical data on the performance of the bilayer device based on the 34 devices made. See DOI: 10.1039/d0na00780c

‡ Current address: Department of Physics, SRM Institute of Science and Technology, Ramapuram, Chennai-600089.



$\text{TiO}_2$  with numerous semiconductors including  $\text{CdTe}$ ,  $\text{CdS}$ ,  $\text{CuI}$ ,  $\text{Cu}_2\text{O}$ ,  $\text{MoS}_2$ , and  $\text{ZnO}$  was studied where an enhancement in the photoelectrochemical properties, photovoltaic or photocatalytic performance was observed.<sup>5,18–20</sup>

$\alpha\text{-MoO}_3$  is one of the well-known layered structured materials widely used in electrodes of energy storage devices. Owing to its high hole mobility and work function, it is used in energy conversion devices as well. The electronic characteristics and the layered structure are the key factors making  $\text{MoO}_3$  a potential semiconductor in various applications including sensors, photocatalysts, and smart windows.<sup>21,22</sup> It is well known that the charge separation efficiency in case of a type II heterojunction is much greater than that of a type I heterojunction.<sup>23,24</sup> As the energy bands of  $\text{TiO}_2$  align well and form a type II heterojunction with  $\text{MoO}_3$ , the  $\text{TiO}_2/\text{MoO}_3$  heterojunction structures have also been explored for various applications including photocatalysis,<sup>25</sup> supercapacitors<sup>26</sup> and ion storage devices.<sup>27</sup> For example, a  $\text{TiO}_2/\text{MoO}_3$  heterojunction fabricated using  $\text{TiO}_2$  nanofibers as the template showed a significant increase in the adsorption of dye and improvement in the photocatalytic activity as compared to bare  $\text{TiO}_2$  nanofibres.<sup>25</sup> The increase in the areal and specific capacitances was observed for the supercapacitor using a  $\text{TiO}_2$  nanotube/ $\text{MoO}_3$  nanoplate heterostructure.<sup>26</sup> S. H. Paek *et al.* observed an enhancement in the lithium ion storage capacity and reversible capacity in the porous  $\text{TiO}_2$  pillared  $\text{MoO}_3$ .<sup>27</sup> N. K. Shrestha *et al.* investigated the electrochromic properties of  $\text{TiO}_2/\text{MoO}_3$  heterostructures and observed a 7 fold increase in the charge storage capacity compared to  $\text{TiO}_2$  nanotubes.<sup>28</sup>

UV light detection is essential for diverse applications including environmental monitoring, remote control, imaging and opto-electronic circuits.<sup>29</sup> The key features required for the UV detectors are the strong light absorption and efficient separation of photogenerated carriers.<sup>30</sup> Self-powered UV detectors are of great interest due to their zero energy consumption, sustainable nature and minimal maintenance.<sup>31</sup> Such detectors are possible with efficient photoelectrochemical interfaces, Schottky junctions, or any other heterojunctions.<sup>32</sup> The self-powered performance arises due to the built-in potential at the interface or junction, which drives the separation of photogenerated carriers *via* a photovoltaic effect.<sup>31,33</sup>  $\text{TiO}_2$  based heterojunctions are also evolving as promising UV detectors and show efficient photodetection ability with an ultrafast response and high photosensitivity.<sup>34,35</sup> Although there are reports on  $\text{TiO}_2$  based heterojunctions with organic layers,  $\text{NiO}$ ,  $\text{SnO}_2$ ,  $\text{CuZnS}$ , and  $\text{ZnS}$  for self-powered photodetection application,<sup>31,32,36–39</sup> such studies on  $\text{TiO}_2/\text{MoO}_3$  were rarely reported. A recent report on a  $\text{TiO}_2/\text{MoO}_3$  UV detector employed hydrothermally synthesised 1D  $\text{TiO}_2$  (rutile) nanowires with  $\text{MoO}_3$  nanostructures, and used a bias voltage of 2.2 V, with no observation of self-powered performance.<sup>40</sup>

According to the literature, 1D  $\text{TiO}_2$  nanostructures have been extensively used in  $\text{TiO}_2$  based UV detectors as well as in  $\text{TiO}_2/\text{MoO}_3$  based devices for various other energy harvesting applications.<sup>21–24</sup> Although 1D  $\text{TiO}_2$  nanostructures offer high specific area, defined geometry and efficient unidirectional carrier transport, the high surface recombination due to their

high surface to volume ratio is considered as one of the potential limiting factors for their use in energy harvesting systems.<sup>41</sup> The surface recombination is known to be less in materials with a lower aspect ratio, where the parasitic surface area is lower.<sup>42</sup> Thus, the  $\text{TiO}_2/\text{MoO}_3$  bilayer heterojunction using  $\text{TiO}_2$  nanoparticles would not have this limitation and also would be more transparent, unlike the one reported based on hydrothermally grown  $\text{TiO}_2$  nanowires.<sup>40,43</sup> As already mentioned, transparent devices are considered as the major players in future energy saving technology due to their unique optical and electrical properties. It's also notable that previous reports did not explore  $\text{TiO}_2$  with anatase crystal structure, for which the electron mobility is known to be relatively higher.<sup>11</sup> Thus, it's interesting to study a solution processed  $\text{TiO}_2/\text{MoO}_3$  heterojunction device with anatase  $\text{TiO}_2$  nanoparticle based bottom layer, owing to its transparent nature, superior mobility, lower aspect ratio, and also the simpler cost-effective preparation methods.<sup>11,12</sup> Based on the above considerations, herein we attempted to fabricate a transparent anatase  $\text{TiO}_2$  nanoparticle  $\alpha\text{-MoO}_3$  nanostructure based heterojunction device with a simple solution based approach. The charge transport properties of the heterojunction were investigated by analysing the *I*–*V* characteristics in detail. Also, the effect of the morphology and junction formation on the charge transfer and the self-powered UV photodetection capability has been analysed.

## Experimental details

All the chemicals were purchased from Sisco Research Laboratories Pvt. Ltd., India and used without further purification.

A transparent solution for making  $\text{TiO}_2$  film was prepared by the hydrolysis reaction of titanium tetraisopropoxide (TTIP) in  $\text{HNO}_3$  aqueous solution.<sup>44</sup> First, an alkoxide solution was prepared by dissolving 26.5 M TTIP in ethanol and an acid solution was prepared by mixing 0.3 M nitric acid in water under stirring for 2 hours. The volume ratio of water to ethanol was maintained at 1 : 1. The acid solution was added dropwise into the alkoxide solution under vigorous stirring. The resultant mixture was kept at 45 °C for 36 hours to get a transparent solution. The solution was then spin coated on cleaned FTO/glass substrates at 3000 rpm for 30 s followed by drying at 100 °C for 10 min. This process was repeated seven times to attain the desired thickness of the  $\text{TiO}_2$  and finally the film was annealed at 500 °C for 1 hour on a hotplate.

For  $\text{MoO}_3$  coating, 1 mmol of ammonium heptamolybdate hexahydrate was dissolved in DI water and kept at 80 °C for 1 h, to obtain a colorless transparent solution.<sup>45</sup> The solution was spin coated on the FTO/glass substrate at 2000 rpm for 20 s followed by drying at 130 °C for 15 min. This process was repeated seven times and finally, the film was annealed at 380 °C for 1 hour on a hotplate.  $\text{TiO}_2/\text{MoO}_3$  film was prepared by following the procedure explained above to sequentially deposit  $\text{TiO}_2$  and  $\text{MoO}_3$  films as a bilayer on the FTO/glass substrate. Finally, the devices for studying the photodetection performance were made by thermally evaporating Au top contacts on these single layer and bilayer films (Fig. 1).



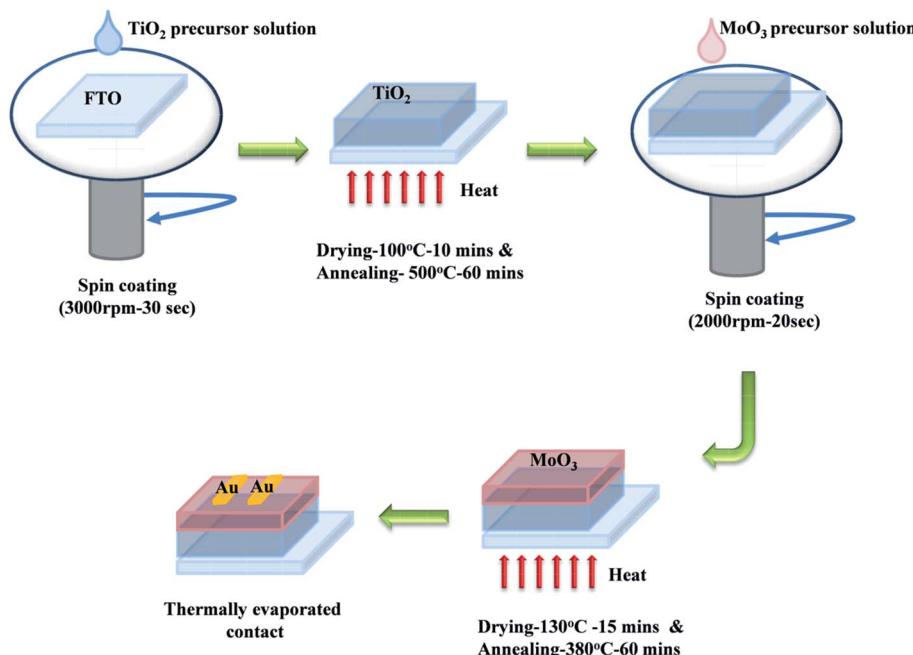


Fig. 1 Schematic illustrating the fabrication of a bilayer UV detector device.

X-ray diffraction was carried out using  $\text{Cu K}\alpha$  radiation (Powder X-ray Diffractometer, PANalytical, Xpert) with the diffraction angle from  $10$  to  $80^\circ$ . UV-visible absorption and transmittance studies were carried out using an Agilent Cary UV-visible absorption spectrometer. FE-SEM analysis was carried out with a Quanta FEG 2000. X-ray photoelectron spectroscopy was carried out using a PHI VersaProbe III scanning XPS microprobe, Physical Electronics, USA. Electrochemical impedance spectroscopy with a two-electrode system was carried out using a Zahner Zennium electrochemical workstation in the frequency range from  $10$  mHz to  $100$  kHz with an amplitude of  $10$  mV.  $I$ - $V$  and  $I$ - $t$  characteristics of the device (active area of  $0.24\text{ cm}^2$ ) were measured using a 2450 Keithley source meter unit.  $I$ - $V$  characteristics were measured in the dark, under 1 sun illumination and illumination by a  $352$  nm UV lamp with an intensity of  $76\text{ }\mu\text{W cm}^{-2}$ .  $I$ - $t$  characteristics were measured in the dark and under illumination by a  $352$  nm UV lamp with an intensity of  $76\text{ }\mu\text{W cm}^{-2}$ . The spectral response measurements were carried out using a Bentham PVE300 photovoltaic characterization setup with a TMC 300 monochromator.

## Results and discussion

XRD analysis was carried out to confirm the crystalline nature of the samples. Fig. S1† shows the XRD patterns of  $\text{TiO}_2$ ,  $\text{MoO}_3$  and  $\text{TiO}_2/\text{MoO}_3$  films coated on FTO/glass. The evident peaks of all the films could be matched well with the standard data of anatase- $\text{TiO}_2$  (ICSD-98-00-5225) and  $\alpha$ - $\text{MoO}_3$  (ICSD-98-001-7296) after excluding the peaks of FTO. The grain size was calculated using the Scherrer formula and the mean grain size was  $48.14\text{ nm}$  for  $\text{TiO}_2$ ,  $17.12\text{ nm}$  for  $\text{MoO}_3$  and  $15.17\text{ nm}$  for  $\text{TiO}_2/\text{MoO}_3$  film.

The changes observed in the grain size can be correlated with the changes in the oxygen vacancy concentration as indicated by the variation in the intensity of the secondary peaks.<sup>46,47</sup>

FE-SEM analysis was carried out to investigate the surface morphology and the obtained images are shown in Fig. 2a–c. The surface of the  $\text{TiO}_2$  film was observed to have agglomerated irregular spherical nanoparticles. The  $\text{MoO}_3$  film had a nano-spike-like structure on FTO while the structure shown by the  $\text{TiO}_2/\text{MoO}_3$  film was not exactly similar. A decrease in the length of nanospikes of  $\text{MoO}_3$  and densely packed alignment was observed in the case of  $\text{MoO}_3$  film deposited on  $\text{TiO}_2$ . This may be due to the substrate dependent growth of  $\text{MoO}_3$  structures. For the bare  $\text{MoO}_3$  film, the growth was on the smooth film of flat FTO/glass and for the  $\text{TiO}_2/\text{MoO}_3$  film, the growth of  $\text{MoO}_3$  was on the film of agglomerated  $\text{TiO}_2$  nanoparticles.<sup>41,48</sup> The change can also be related to the change in the oxygen vacancy concentration that affects the particle size and shape.<sup>49</sup> It's known in the literature that denser  $\text{ZnO}$  nanostructures show a higher photoresponse<sup>50</sup> and analogously, we could predict that the reduction in the length and dense packing of  $\text{MoO}_3$  nanostructures could lead to an improved photoresponse. The FESEM cross-sectional image (Fig. S2†) of the bilayer film confirms the successive growth of  $\text{TiO}_2$  ( $\sim 300\text{ nm}$  thick) and  $\text{MoO}_3$  ( $\sim 260\text{ nm}$  thick) layers.

UV-visible transmission spectra of the films recorded in the wavelength range of  $280$ – $800$  nm are shown in Fig. 2d. The spectra confirmed that both  $\text{TiO}_2$  and  $\text{TiO}_2/\text{MoO}_3$  films absorb light in the UV region while exhibiting high transparency to the visible light (above  $400$  nm). The average transmittance of  $85\%$  and  $74\%$  were observed in the visible region ( $400$ – $800$  nm) for the  $\text{TiO}_2$  and  $\text{TiO}_2/\text{MoO}_3$  films, respectively. It should also be noted that the absorption of light by the bilayer heterojunction



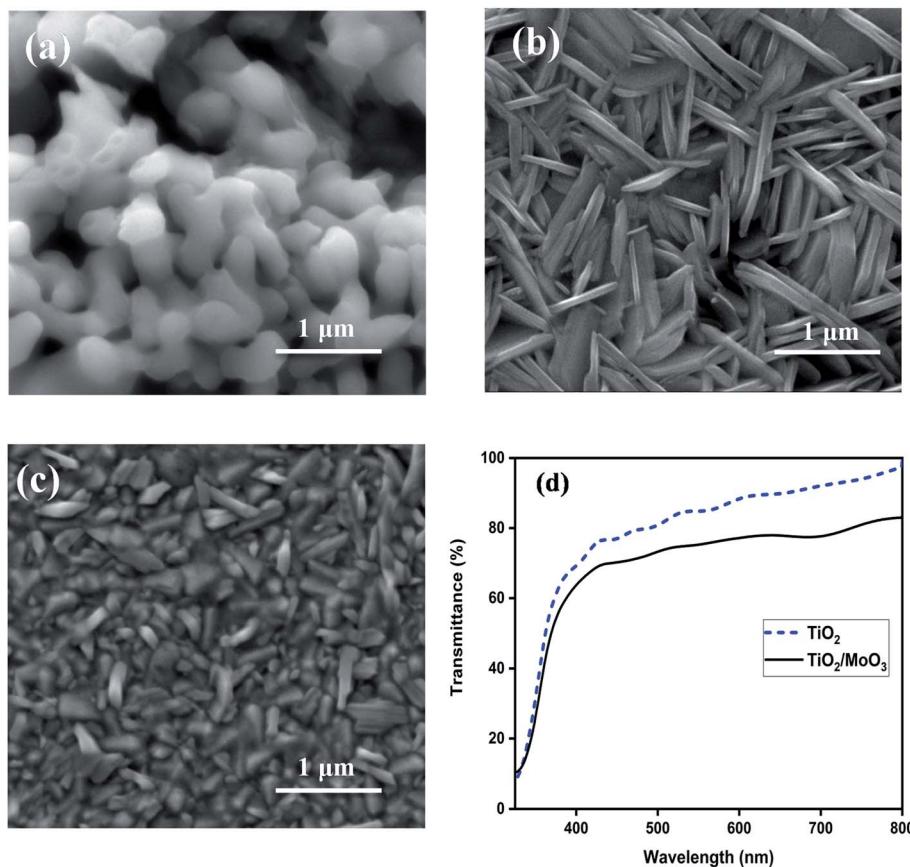


Fig. 2 FESEM images of (a)  $\text{TiO}_2$ , (b)  $\text{MoO}_3$ , and (c)  $\text{TiO}_2/\text{MoO}_3$  films, and (d) UV-visible transmission spectra of  $\text{TiO}_2$  and  $\text{TiO}_2/\text{MoO}_3$  films.

film was significantly higher as compared to the  $\text{TiO}_2$  film (Fig. S3†).

XPS analysis was carried out to confirm the chemical states, as well as the elemental composition of the samples and the results are shown in Fig. 3. The binding energies of Mo  $3d_{3/2}$  and Mo  $3d_{5/2}$  for the  $\text{MoO}_3$  film were located at 236.0 eV and 232.9 eV, respectively, corresponding to the  $\text{Mo}^{6+}$  oxidation state.<sup>51</sup> These binding energies were shifted to 235.8 eV and 232.7 eV for the Mo 3d in the  $\text{TiO}_2/\text{MoO}_3$  film. Though there was

a slight shift in Mo 3d binding energies, the oxidation state of  $\text{Mo}^{6+}$  remained the same. The binding energy of the O 1s peak in the  $\text{TiO}_2/\text{MoO}_3$  heterojunction was located at 530.8 eV which was also shifted downwards by 0.2 eV compared to that of the O 1s peak for  $\text{MoO}_3$ . Both Mo 3d and O 1s peaks of the  $\text{TiO}_2/\text{MoO}_3$  film were shifted to lower energies than those of the  $\text{MoO}_3$  film and this can be attributed to the electron transfer between  $\text{TiO}_2$  and  $\text{MoO}_3$  in the heterojunction.<sup>25</sup>

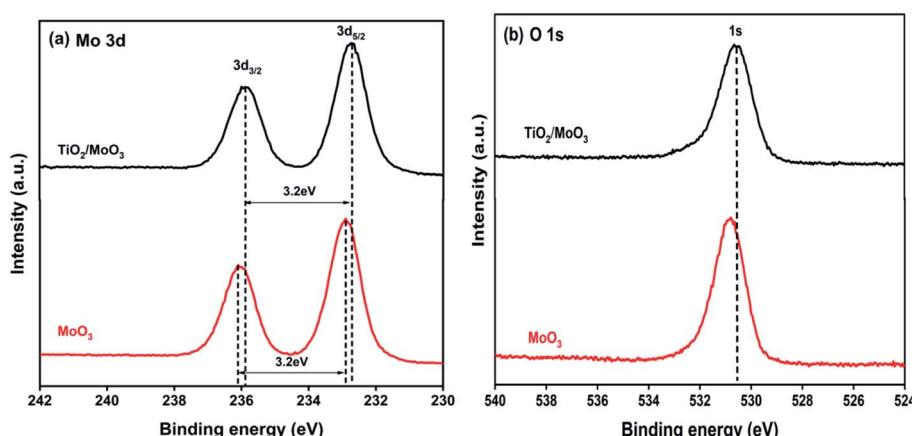


Fig. 3 XPS spectra of (a) Mo 3d and (b) O 1s of  $\text{MoO}_3$  and  $\text{TiO}_2/\text{MoO}_3$  films.



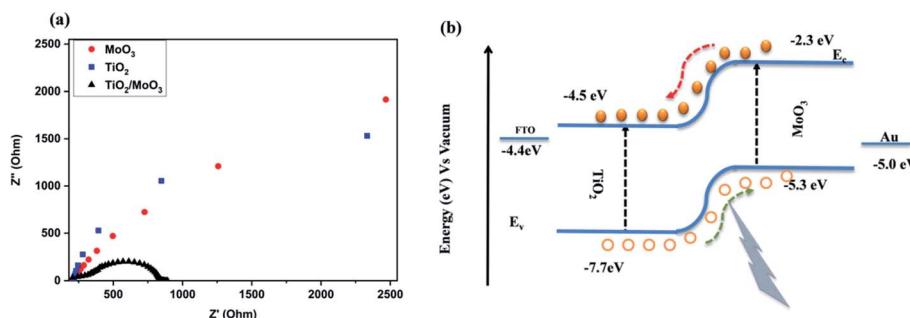


Fig. 4 (a) Electrochemical impedance spectra of MoO<sub>3</sub>, TiO<sub>2</sub>, and TiO<sub>2</sub>/MoO<sub>3</sub> films, and (b) energy level diagram of the TiO<sub>2</sub>/MoO<sub>3</sub> heterojunction.

Electrochemical impedance spectroscopy was carried out to analyse the charge transfer resistance and separation efficiency of the heterojunction. The spectra were recorded with the back contact FTO and the front contact Au in a 2 electrode configuration at room temperature under dark conditions. The obtained Nyquist plots are shown in Fig. 4a. It was evident from the spectra that the radius of the semicircle for the bilayer heterojunction device was much smaller than that of the single layer device. The small semi-circular arc for the TiO<sub>2</sub>/MoO<sub>3</sub> heterojunction indicated a smaller charge transfer resistance which can be related to the faster charge transfer across the interface. This also corresponds to the higher efficiency of the heterojunction for the separation of carriers.<sup>26</sup> The equivalent circuit model used for fitting and the resistance values are given in Fig. S9 and Table S1.†

To understand the electrical behaviour of the TiO<sub>2</sub>/MoO<sub>3</sub> bilayer, two probe *I*-*V* measurements were done in the dark at room temperature with Au and FTO as a front and back contact, respectively, and the results are shown in Fig. 5a. Before the fabrication of the heterojunction, TiO<sub>2</sub> and MoO<sub>3</sub> were studied separately. The *I*-*V* curve of the TiO<sub>2</sub> based device exhibited a Schottky behaviour (Fig. S4a†), while the MoO<sub>3</sub> based device showed ohmic conduction (Fig. S5†). However, the *I*-*V* curve of the bilayer device exhibited a diode-like behaviour. The diode parameters were analyzed using the ideal diode equation<sup>52</sup>

$$I = I_s \left[ \exp\left(\frac{qV}{nkT}\right) - 1 \right] \quad (1)$$

$$I_s = A A^* T^2 \exp\left(\frac{-q\phi_b}{kT}\right) \quad (2)$$

$$\phi_b = \frac{kT}{q} \ln\left(\frac{AA^*T^2}{I_s}\right) \quad (3)$$

where  $I_s$  is the reverse saturation current,  $q$  is the electron charge,  $n$  is the ideality factor,  $k$  is the Boltzmann constant,  $T$  is the temperature,  $A$  is the active area,  $A^*$  is the effective Richardson constant and  $\phi_b$  is the barrier height. The ideality factor was calculated from the slope of the semi-logarithmic plot and reverse saturation current was calculated from the current intercept at  $V = 0$ . The Richardson constant was taken based on the literature ( $1200 \text{ A cm}^{-2} \text{ K}^{-2}$  for TiO<sub>2</sub>).<sup>53</sup> The obtained parameters are listed in Table S2.†

The reverse current started increasing slightly at  $-0.4 \text{ V}$ , but the corresponding forward current was one order of magnitude higher which is comparable to the earlier observations with other heterojunction devices.<sup>40,54,55</sup> The ideality factor which was greater than unity (2.98) indicates that the charge transport mechanism deviated from the thermionic emission theory applicable for an ideal diode. To further investigate the charge transport mechanism, power-law compliance was considered and the forward *I*-*V* data were plotted in a log *I*-log *V* scale<sup>56</sup>

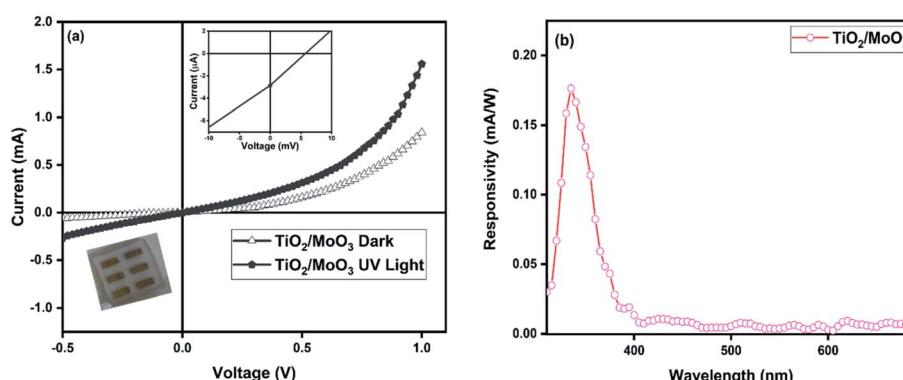


Fig. 5 Device characteristics of the bilayer device: (a) *I*-*V* characteristics in the dark and under UV illumination, insets: *I*-*V* curve under UV light, enlarged around 0 V & a photograph of the device, and (b) spectral responsivity.



(Fig. S6†). The plot was fitted well with the power-law expression mentioned below with two different slopes.

$$I = KV^m$$

where  $K$  is a constant and  $m$  is the exponent derived from the slope of the curve.

At lower voltages (0.02–0.28 V), the exponent determined was close to unity with a correlation coefficient of 0.99645. This indicates that the current varied linearly with the voltage and the charge transport was governed by Ohm's law. Hence, the thermally generated free electrons were responsible for the current produced at lower bias. Furthermore, an enhancement in the current was observed with the increase in the voltage, indicating that the injected carriers contribute to the rise in current.<sup>54,57,58</sup> At higher voltages (0.32–1.0 V), the exponent was determined to be 2.18 with a correlation coefficient of 0.99589. This indicates that the space charge limited charge transport became pre-dominant, where the number of injected carriers becomes comparable to the thermally generated carriers. The value of the exponent is slightly greater than 2 and this suggests that the injected carriers have just entered into the space charge region.<sup>59</sup>

$I$ – $V$  characteristics measured under UV light illumination (352 nm) clearly showed the photoconductive nature of the devices. The photocurrent (the difference between the current under UV illumination and dark) of the bilayer device was observed as 1.94  $\mu$ A at 0 V and it reached a value of 840  $\mu$ A at a reverse bias of –1 V. The photocurrent without bias was 40 times higher and with bias, it was 3 times higher when compared to that of the single layer  $\text{TiO}_2$  nanoparticle based device for which the photocurrent was only 42 nA at 0 V and 228  $\mu$ A at –1 V (Fig. S4†). The devices exhibited photocurrent even under zero bias conditions, due to the photovoltaic effect under UV illumination. The  $I$ – $V$  curves enlarged around 0 V showed a photovoltaic behaviour upon UV illumination for both single and bilayer devices (inset in Fig. 4a and S4a†). The  $\text{TiO}_2/\text{MoO}_3$  bilayer device showed an open-circuit voltage ( $V_{oc}$ ) of 5.7 mV and a short-circuit current ( $I_{sc}$ ) of 2.856  $\mu$ A, while for the  $\text{TiO}_2$  single layer, the values of  $V_{oc}$  and  $I_{sc}$  were 4.4 mV and 0.953  $\mu$ A, respectively. It can be attributed to the built-in potential at the FTO/ $\text{TiO}_2$  and  $\text{TiO}_2/\text{MoO}_3$  interfaces arising due to alignment of energy bands, which act to separate the photogenerated carriers efficiently.<sup>33,60</sup>

Based on the  $I$ – $V$  characteristics, the photoresponsivity ( $R$ ), detectivity ( $D$ ) and external quantum efficiency (EQE) were calculated using the formulae mentioned below<sup>30</sup>

$$R = \frac{I_{\text{light}} - I_{\text{dark}}}{P_{\text{ill}}} \quad (5)$$

$$D = \frac{R}{(2eJ_d)^{1/2}} \quad (6)$$

$$\text{EQE}_\lambda \% = \frac{1240R}{\lambda} 100$$

where  $P_{\text{ill}} = I_{\text{ill}} \times A_0$ ,  $I_{\text{ill}}$  is the illumination power density,  $A_0$  is the active area,  $e$  is the charge of an electron,  $J_d$  is the dark current density,  $R$  is the photoresponsivity, and  $\lambda$  is the excitation wavelength (352 nm). For the  $\text{TiO}_2/\text{MoO}_3$  device,  $R$ ,  $D$ , and EQE were estimated to be 108 mA W<sup>–1</sup>,  $2.26 \times 10^{10}$  Jones and 38.04% without bias, while at –1 V bias the values were found to be 46.05 A W<sup>–1</sup>,  $2.84 \times 10^{12}$  Jones and  $1.62 \times 10^4\%$ , respectively. When compared to the single layer device, the value of  $R$  was 40 times higher without bias and 3 times higher with bias.  $D$  and EQE were observed to be 1 order higher in magnitude with and without bias.

When the  $I$ – $V$  measurements were carried out under 1 sun white light illumination, a negligible photocurrent was observed (Fig. S7†). This confirms that the response was significant only with illumination by UV light of the wavelength corresponding to the bandgap of the film, concurrent with the absorption spectrum. The spectral responsivity measured with 0 V bias (Fig. 5b) has again proven the excellent UV-light detection selectivity of the bilayer heterojunction film, with the highest response shown in the spectral range between 300 and 385 nm. The results confirmed that the device has the ability to work as a visible blind UV detector.

To further understand the self-powering ability,  $I$ – $t$  measurements were performed under 0 V bias with an interval of 30 s between UV illuminations (Fig. 6). The photocurrent, rise time and fall time for our bilayer film, as estimated from the  $I$ – $t$  characteristics, were found to be 1.94  $\mu$ A, 1.82 s and 1.42 s, respectively. The high photocurrent upon UV light illumination can be attributed to the photo-desorption of oxygen species on the surface. The dynamics associated with the photo-adsorption/desorption process results in the increase of free carrier density and consequently the photocurrent. The apparent overshoot phenomenon could be related to the charge accumulation at the metal/semiconductor interface.<sup>61</sup> It is noteworthy that the time taken to reach the steady state was less than 2 s and the response remained unaltered even after 10 cycles. To understand the photoresponse under external bias,  $I$ –

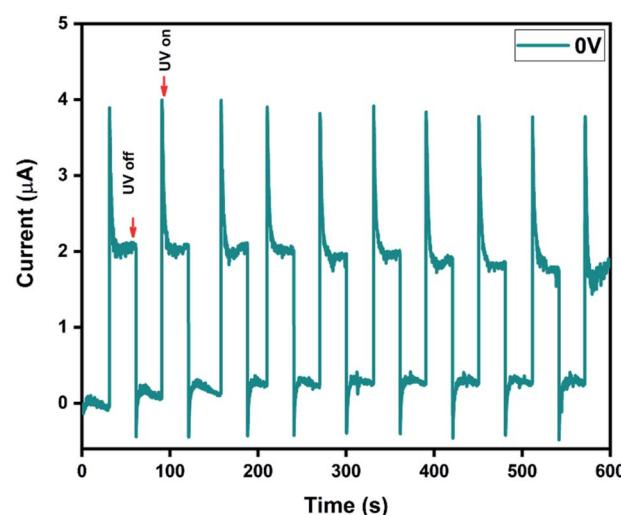


Fig. 6  $I$ – $t$  characteristics of the self-powered bilayer device.



Table 1 Performance of  $\text{TiO}_2$  based self-powered UV photodetectors

Device	Wavelength (nm)	Intensity of the source ( $\text{mW cm}^{-2}$ )	Bias (V)	Photoresponsivity ( $\text{mA W}^{-1}$ )	Detectivity (Jones)	Rise time	Fall time	Ref.
$\text{TiO}_2/\text{Ag}$	350	1.2	0	1.1	—	2 ms	47 ms	64
$\text{TiO}_2/\text{Ag}$	350	0.038	0	32.5	$6 \times 10^9$	44 ns	1.85 $\mu\text{s}$	65
$\text{NiO}/\text{TiO}_2$	370	0.2	0	860	$1.6 \times 10^{13}$	3 ms	60 ms	60
$\text{NiO}/\text{TiO}_2/\text{TiO}_x$	380	0.5	0	5.66	$2.5 \times 10^{12}$	<0.1 s	<0.1 s	39
$\text{SnS}_x/\text{TiO}_2$	365	2	0	5.85	$6.54 \times 10^{10}$	3 ms	25 ms	37
$\text{PANI}/\text{TiO}_2$	320	0.87	0	3.6	$3.9 \times 10^{11}$	1.2 ms	22.8 ms	66
$\text{SnO}_2/\text{TiO}_2$	365	40	0	600	—	0.02 s	0.004 s	38
$\text{Au}/\text{P3HT}/\text{TiO}_2$	350	1	0	0.25	$2.9 \times 10^{10}$	0.72 s	0.5 s	31
$\text{ZnO}/\text{Ga}_2\text{O}_3$	266	1.6	-0.1	9.7	$2.58 \times 10^{12}$	100 $\mu\text{s}$	900 $\mu\text{s}$	67
$\text{TiO}_2$	352	<b>0.076</b>	<b>0</b>	<b>2.68</b>	$3.34 \times 10^9$	—	—	<b>This work</b>
$\text{TiO}_2/\text{MoO}_3$	352	<b>0.076</b>	<b>0</b>	<b>108</b>	$2.26 \times 10^{10}$	<b>1.82 s</b>	<b>1.48 s</b>	

*t* measurements were also performed at  $-1\text{ V}$  (Fig. S8†). The photocurrent, rise time and fall time were found to be  $840\text{ }\mu\text{A}$ ,  $1.22\text{ s}$  and  $1.4\text{ s}$ , respectively. The cycling stability of the device with an applied voltage of  $-1\text{ V}$  was observed to be slightly inferior, and the dark current was found to increase after 5 cycles. This may be attributed to the redistributed electric field caused by the applied bias, which can induce a decrease in voltage dependent carrier separation and transport.<sup>36,62,63</sup>

However, it's noteworthy that the stability loss observed here was only 5% and there was almost no loss when the device was operated under self-powered mode.

As compared to the literature, our single layer device with the configuration of  $\text{FTO}/\text{TiO}_2/\text{Au}$  itself showed better performance, under bias conditions lower than the reported ones (Table S3†). The detection capability of the  $\text{TiO}_2/\text{MoO}_3$  bilayer device was found to be remarkable in terms of the bias conditions as well as the responsivity when compared with the other reported anatase  $\text{TiO}_2$  based UV detectors. Also, the rise time is 3 times less and the fall time is 12 times less (Table S4†) when compared to the values reported earlier for the  $\text{TiO}_2/\text{MoO}_3$  UV detector.<sup>40</sup> This faster response may be attributed to the effective collection of the photogenerated carriers and the consequent recombination loss.<sup>68</sup> This corresponds to the effective junction formation between  $\text{TiO}_2$  and  $\text{MoO}_3$  layers and the efficient separation of photogenerated carriers.<sup>69</sup>

It can be seen from Table 1 that the performance of our devices was better when compared with many such reports on  $\text{TiO}_2$  based self-powered UV detectors. The bilayer device had higher responsivity as well as detectivity at zero bias, even under the relatively weakest illumination. It stands out as the best so far among the devices studied under low light conditions. It is also noteworthy that the photoresponsivity of our self-powered bilayer device at zero bias itself is comparable to that of commercial Si photodetectors ( $0.1\text{--}0.2\text{ A W}^{-1}$ )<sup>70</sup> and with a bias of  $-1\text{ V}$ , it is much higher. However, the slow response speed of the device under  $0\text{ V}$  bias was related to the oxygen vacancies.<sup>71</sup> The oxygen vacancies produce subgap states and can act as recombination centers. The subgap states hinder carrier excitation upon UV illumination<sup>72</sup> and they also act as trap centres to the carriers upon the shut off of UV irradiation, resulting in a slow response.<sup>73</sup>

The observed photoresponse is also better than that of most self-powered photodetectors based on various metal oxides such as  $\text{ZnO}$ ,  $\text{Ga}_2\text{O}_3$ ,  $\text{NiO}$ , and  $\text{SnO}_2$ , which perform under illumination conditions much stronger than ours.<sup>60,65,74-77</sup> Overall, the superior performance of the bilayer device, as evident from the discussion above, can be ascribed to the formation of an effective heterojunction between  $\text{TiO}_2$  and  $\text{MoO}_3$ . The built-in potential at the interface of  $\text{TiO}_2/\text{MoO}_3$  leads to an enhanced separation efficiency as well as an effective charge transfer, and enables the device to be self powered.<sup>60</sup> The improved and faster response can be related to the reduced recombination. As already mentioned, these improvements could also be due to the morphological factor of having densely aligned nanostructures. It is well known that the adsorption/photodesorption of  $\text{O}_2$  species upon UV illumination<sup>78</sup> followed by the transport of photogenerated carriers<sup>79</sup> plays a vital role in UV detection performance. The depletion region would shrink upon application of the reverse bias due to the hole-trapping at the surface. This increases the electron tunneling, and the electron migration in large numbers is related to the EQE over 100%.<sup>80,81</sup> It can be proposed that the densely aligned nanostructures could expose multiple junctions sensitive to the UV light and could also provide efficient pathways for the transport of carriers, leading to the enhancement in the photocurrent and EQE.<sup>50,82</sup> When compared to previous reports, the improvement can also be related to the mobility of the anatase  $\text{TiO}_2$  used in this work which is known to be better than that of rutile  $\text{TiO}_2$ .<sup>83</sup> Though the densely aligned morphology offers better performance in terms of faster response and high photocurrent, the photocurrent stability with time can be further improved by tuning the surface area.<sup>84</sup> There is room for achieving better performance by more carefully optimizing the particle size, morphology, crystallinity, and thickness of the films.

## Conclusion

In summary, we have fabricated and studied a fully solution processed transparent  $\text{TiO}_2/\text{MoO}_3$  bilayer thin film device with an anatase  $\text{TiO}_2$  nanoparticle based bottom layer for the first time, for self-powered UV detector application. XPS,



electrochemical measurements and current–voltage characteristics of the bilayer have confirmed the formation of an effective heterojunction diode, where the interaction of the energy bands could lead to an efficient charge transfer. The charge transport mechanism of the bilayer device was governed by Ohm's law at low voltages and space charge limited conduction at high voltages. The bilayer device exhibited a very good photoresponsivity and detectivity under 352 nm UV light illumination of very low intensity. The performance was on par with other  $\text{TiO}_2$  based heterojunctions reported as self-powered UV detectors, showing a photoresponse even under zero bias conditions which was further enhanced at a low bias of  $-1$  V. The self-powered performance of the device was attributed to the built-in potential of the  $\text{TiO}_2/\text{MoO}_3$  interface that led to the separation of photogenerated carriers without bias. The improved detection capability exhibited by the  $\text{TiO}_2/\text{MoO}_3$  bilayer device in terms of the low bias conditions, high responsivity, high EQE and faster response can be mainly attributed to the better carrier collection as well as the reduced recombination across the heterojunction. The morphological factors of nanostructures and the better mobility of anatase  $\text{TiO}_2$  nanoparticles could also be the reasons behind the observed enhancement in the performance. The results suggest that the bilayer heterojunction, where the  $\text{TiO}_2$  nanoparticle film is the platform with the  $\alpha\text{-MoO}_3$  film on the top, is a promising candidate for self-powered UV detector application with its high transparency as well as the best-in-class photoresponse under low bias and weak illumination conditions.

## Author contributions

BE carried out all the experiments including synthesis, device fabrication and characterization with initial help from MD. SVB conceptualized the project and supervised the research work. BE and SVB discussed the results, prepared the manuscript and carried out the required revisions.

## Conflicts of interest

There are no conflicts of interest to declare.

## Acknowledgements

The authors acknowledge the financial support of SRMIST towards this research work and all the in-campus characterization facilities such as the UV-visible spectrometer at the Department of Physics & Nanotechnology, XPS and FE-SEM facility at NRC and the XRD facility with the support from MNRE (Project No. 31/03/2014-15/PVSE-R&D), Govt. of India. We thank Dr S. Harinipriya and coworkers for the help with the access to electrochemical impedance measurements. We also thank Dr P. Malar and coworkers for the help with the use of a thermal evaporator. We thank Prof. B. Neppolian and coworkers for the access to the EQE facility supported by Department of Science and Technology – Solar Energy Research Initiative (DST; SERB), File No. DST/SERI/FR/2016/S170. We would also like to acknowledge the research grants from DST–

SERB [File. No. (ECR/2015/000513)] and MNRE [File. No. 31/03/2014-15/PVSE-R&D], Govt. of India.

## References

- 1 K. D. Leedy, M. Higashiwaki, R. L. Peterson and S. James, *APL Mater.*, 2019, **7**, 022201.
- 2 H. Ohta, M. Hirano, K. Nakahara, H. Maruta, T. Tanabe, M. Kamiya, T. Kamiya and H. Hosono, *Appl. Phys. Lett.*, 2003, **83**, 1029–1031.
- 3 C. C Shih, W. Y. Lee, Y. C. Chiu, H. W. Hsu, H. C. Chang, C. L. Liu and W. C. Chen, *Sci. Rep.*, 2016, **6**, 20129.
- 4 Z. Q. Zheng, J. D. Yao, B. Wang and G. W. Yang, *Sci. Rep.*, 2015, **5**, 11070.
- 5 M. Pavan, S. Rühle, A. Ginsburg, D. A. Keller, H. N. Barad, P. M. Sberna, D. Nunes, R. Martins, A. Y. Anderson, A. Zaban and E. Fortunato, *Sol. Energy Mater. Sol. Cells*, 2015, **132**, 549–556.
- 6 L. Ahtapodov, B. Fimland, K. Kishino and H. Weman, *Nano Lett.*, 2019, **19**, 1649–1658.
- 7 S. Abbas, D. Ban and J. Kim, *Sens. Actuators, A*, 2019, **293**, 215–221.
- 8 J. Chen, Z. Li, F. Ni, W. Ouyang and X. Fang, *Mater. Horiz.*, 2020, **7**, 1828–1833.
- 9 C. C. Forin, M. Purica, E. Budianu and P. Schiopu, *Proc. Int. Semicond. Conf. CAS*, 2012, **1**, 131–134.
- 10 Z. Li, M. K. Joshi, J. Chen, Z. Zhang, Z. Li and X. Fang, *Adv. Funct. Mater.*, 2020, **30**, 2005291.
- 11 G. S. Mital and T. Manoj, *Chin. Sci. Bull.*, 2011, **56**, 1639–1657.
- 12 X. Chen and S. S. Mao, *Chem. Rev.*, 2007, **107**, 2891–2959.
- 13 M. Quintana, T. Edvinsson, A. Hagfeldt and G. Boschloo, *J. Phys. Chem. C*, 2007, **111**, 1035–1041.
- 14 J. Sobuś, G. Burdziński, J. Karolczak, J. Idigoras, J. A. Anta and M. Ziółek, *Langmuir*, 2014, **30**, 2505–2512.
- 15 T. Ling, J. G. Song, X. Y. Chen, J. Yang, S. Z. Qiao and X. W. Du, *J. Alloys Compd.*, 2013, **546**, 307–313.
- 16 D. Friedmann, C. Mendive and D. Bahnemann, *Appl. Catal., B*, 2010, **99**, 398–406.
- 17 B. Liu, X. Zhao, J. Yu, I. P. Parkin and A. Fujishima, *J. Photochem. Photobiol., C*, 2019, **39**, 1–57.
- 18 S. Mishra, P. Yogi, P. R. Sagdeo and R. Kumar, *ACS Appl. Energy Mater.*, 2018, **1**, 790–798.
- 19 Y. Li, C. Cai, Y. Gu, W. Cheng, W. Xiong and C. Zhao, *Appl. Surf. Sci.*, 2017, **414**, 34–40.
- 20 M. Sun, J. Hu, C. Zhai, M. Zhu and J. Pan, *Electrochim. Acta*, 2017, **245**, 863–871.
- 21 I. A. De Castro, R. S. Datta, J. Z. Ou and A. Castellanos-gomez, *Adv. Mater.*, 2017, **29**, 1701619.
- 22 D. Kwak, M. Wang, K. J. Koski, L. Zhang, H. Sokol, R. Maric and Y. Lei, *ACS Appl. Mater. Interfaces*, 2019, **11**, 10697–10706.
- 23 S. Sonu, V. Dutta, S. Sharma, P. Raizada, A. Hosseini-Bandegharaei, V. Kumar Gupta and P. Singh, *J. Saudi Chem. Soc.*, 2019, **23**, 1119–1136.
- 24 W. Yang, K. Hu, F. Teng, J. Weng, Y. Zhang and X. Fang, *Nano Lett.*, 2018, **18**, 4697–4703.



- 25 M. Lu, C. Shao, K. Wang, N. Lu, X. Zhang, P. Zhang, M. Zhang, X. Li and Y. Liu, *ACS Appl. Mater. Interfaces*, 2014, **6**, 9004–9012.
- 26 S. Sun, X. Liao, X. Pu, Y. Sun, G. Yin, Y. Yao, Z. Huang and X. Pu, *RSC Adv.*, 2017, **7**, 22983–22989.
- 27 S. Paek, J. Kang, H. Jung, S. Hwang and J. Choy, *Chem. Commun.*, 2009, **48**, 7536–7538.
- 28 N. K. Shrestha, Y. Nah and P. Schmuki, *Chem. Commun.*, 2009, **15**, 2008–2010.
- 29 Q. Chen, Y. Zhang, T. Zheng, Z. Liu, L. Wu, Z. Wang and J. Li, *Nanoscale Adv.*, 2020, **2**, 2705–2712.
- 30 M. Moun, M. Kumar, M. Garg, R. Pathak and R. Singh, *Sci. Rep.*, 2018, **1**, 1–10.
- 31 L. Zheng, X. Deng, Y. Wang, J. Chen, X. Fang, L. Wang, X. Shi and H. Zheng, *Adv. Funct. Mater.*, 2020, **30**, 2001604.
- 32 X. Li, S. Gao, G. Wang, Z. Xu, S. Jiao, D. Wang, Y. Huang, D. Sang, J. Wang and Y. Zhang, *J. Mater. Chem. C*, 2020, **8**, 1353–1358.
- 33 Y. Shen, X. Yan, Z. Bai, X. Zheng, Y. Sun, Y. Liu, P. Lin, X. Chen and Y. Zhang, *RSC Adv.*, 2015, **5**, 5976–5981.
- 34 J. Xing, H. Wei, E. J. Guo and F. Yang, *J. Phys. D: Appl. Phys.*, 2011, **44**, 375104.
- 35 A. M. Selman and Z. Hassan, *Superlattices Microstruct.*, 2015, **83**, 549–564.
- 36 D. Zhang, C. Liu, B. Yin, R. Xu, J. Zhou, X. Zhang and S. Ruan, *Nanoscale*, 2017, **9**, 9095–9103.
- 37 J. Chen, J. Xu, S. Shi, R. Cao and D. Liu, *ACS Appl. Mater. Interfaces*, 2020, **12**, 23145–23154.
- 38 C. Gao, X. Li, X. Zhu, L. Chen, Y. Wang, F. Teng, Z. Zhang, H. Duan and E. Xie, *J. Alloys Compd.*, 2014, **616**, 510–515.
- 39 Y. Gao, J. Xu, S. Shi, H. Dong, Y. Cheng, C. Wei, X. Zhang, S. Yin and L. Li, *ACS Appl. Mater. Interfaces*, 2018, **10**, 11269–11279.
- 40 B. Yin, Y. Zhang, K. Li, J. Zhou, C. Liu, M. Zhang and S. Ruan, *Nanotechnology*, 2019, **30**, 465501.
- 41 T. Zhai and J. Yao, *One-dimensional nanostructures: principles and applications*, Wiley&Sons, United States, 2012.
- 42 P. Gao, H. Wang, Z. Sun, W. Han, J. Li and J. Ye, *Appl. Phys. Lett.*, 2013, **103**, 213505.
- 43 J. Szułcik, J. Majewski, A. Buczkowski, J. Radojewski, L. Jędrzejak and E. B. Radojewska, *Sol. Energy Mater.*, 1989, **18**, 241–252.
- 44 M. Sasani Ghamsari and A. R. Bahramian, *Mater. Lett.*, 2008, **62**, 361–364.
- 45 S. Murase and Y. Yang, *Adv. Mater.*, 2012, **24**, 2459–2462.
- 46 R. Parra, C. M. Aldao, J. A. Varela and M. S. Castro, *J. Electroceram.*, 2005, **14**, 149–156.
- 47 S. Aškrabić, Z. Dohčević-Mitrović, A. Kremenović, N. Lazarević, V. Kahlenberg and Z. V. Popović, *J. Raman Spectrosc.*, 2012, **43**, 76–81.
- 48 J. Song and S. Lim, *J. Phys. Chem. C*, 2007, **111**, 596–600.
- 49 R. Mehmood, X. Wang, P. Koshy, J. L. Yang and C. C. Sorrell, *CrystEngComm*, 2018, **20**, 1536–1545.
- 50 V. Postica, M. Hoppe, J. Grottrup, P. Hayes, V. Robisch, D. Smazna, R. Adelung, B. Viana, P. Aschhehoug, T. Pauporté and O. Lupan, *Solid State Sci.*, 2017, **71**, 75–86.
- 51 S. Xie, D. Cao, Y. She, H. Wang, J. W. Shi, M. K. Leung and C. Niu, *Chem. Commun.*, 2018, **54**, 7782–7785.
- 52 D. K. Schroder, *Semiconductor Material and Device Characterization*, Wiley&Sons, United States, 2015.
- 53 H. Altuntas, A. Bengi, U. Aydemir, T. Asar, S. S. Cetin, I. Kars, S. Altindal and S. Ozcelik, *Mater. Sci. Semicond. Process.*, 2009, **12**, 224–232.
- 54 P. C. Ramamurthy, W. R. Harrell, R. V. Gregory, B. Sadanandan and A. M. Rao, *Polym. Eng. Sci.*, 2004, **44**, 28–33.
- 55 P. Singh, P. K. Rout, H. Pandey and A. Dogra, *J. Mater. Sci.*, 2018, **53**, 4806–4813.
- 56 H. Elangovan, A. V. Kesavan, K. Chattopadhyay and P. C. Ramamurthy, *J. Appl. Phys.*, 2019, **125**, 245302.
- 57 A. Wahl and J. Augustynski, *J. Phys. Chem. B*, 1998, **5647**, 7820–7828.
- 58 M. A. Lampert and R. B. Schilling, in *Semiconductors and Semimetals*, ed. R. K. Williardson and A. Beer, Elsevier, Amsterdam, 1970, vol. 6, ch. 1, pp. 1–96.
- 59 X. Qiao, J. Chen and D. Ma, *J. Phys. D: Appl. Phys.*, 2010, **43**, 215402.
- 60 R. Cao, J. P. Xu, S. Shi, J. Chen, D. Liu, Y. Bu, X. S. Zhang, S. Yin and L. Li, *J. Mater. Chem. C*, 2020, **8**, 9646–9654.
- 61 Y. Fang, Y. Hou, Y. Hu and F. Teng, *Materials*, 2015, **8**, 4050–4060.
- 62 W. P. Dumke, *Appl. Phys. Lett.*, 1967, **11**, 314–316.
- 63 S. Luryi in *Hot Electrons in Semiconductors*, ed. N. Balkan, Clarendon Press, Oxford, 1998, vol. 6, ch. 17, pp. 389–395.
- 64 J. Xu, W. Yang, H. Chen, L. Zheng, M. Hu, Y. Li and X. Fang, *J. Mater. Chem. C*, 2018, **6**, 3334–3340.
- 65 H. Fang, C. Zheng, L. Wu, Y. Li, J. Cai, M. Hu, X. Fang, R. Ma, Q. Wang and H. Wang, *Adv. Funct. Mater.*, 2019, **29**, 1809013.
- 66 L. Zheng, P. Yu, K. Hu, F. Teng, H. Chen and X. Fang, *ACS Appl. Mater. Interfaces*, 2016, **8**, 33924–33932.
- 67 B. Zhao, W. Fei, C. Hongyu, Z. Lingxia, S. Longxing, Z. Dongxu and X. Fang, *Adv. Funct. Mater.*, 2017, **27**, 1700264.
- 68 S. Li, Y. Zhang, W. Yang, H. Liu and X. Fang, *Adv. Mater.*, 2020, **32**, 1905443.
- 69 M. Razeghi and A. Rogalski, *J. Appl. Phys.*, 1996, **79**, 7433–7473.
- 70 W. Zheng, X. Li, C. Dong, X. Yan and G. He, *RSC Adv.*, 2014, **4**, 44868–44871.
- 71 P. Ivanoff Reyes, C. J. Ku, Z. Duan, Y. Xu, E. Garfunkel and Y. Lu, *Appl. Phys. Lett.*, 2012, **101**, 2010–2014.
- 72 K. Y. Chen, S. P. Chang and C. hung Lin, *RSC Adv.*, 2019, **9**, 87–90.
- 73 Q. Chen, H. Ding, Y. Wu, M. Sui, W. Lu, B. Wang, W. Su, Z. Cui and L. Chen, *Nanoscale*, 2013, **5**, 4162–4165.
- 74 J. Chen, W. Ouyang, W. Yang, J. H. He and X. Fang, *Adv. Funct. Mater.*, 2020, **30**, 1909909.
- 75 Y. Wang, C. Wu, D. Guo, P. Li, S. Wang, A. Liu, C. Li, F. Wu and W. Tang, *ACS Appl. Electron. Mater.*, 2020, **2**, 2032–2038.
- 76 Y. Wang, L. Li, H. Wang, L. Su, H. Chen, W. Bian, J. Ma, B. Li, Z. Liu and A. Shen, *Nanoscale*, 2020, **12**, 1406–1413.
- 77 B. Deka Boruah, *Nanoscale Adv.*, 2019, **1**, 2059–2085.



- 78 D. Gedamu, I. Paulowicz, S. Kaps, O. Lupan, S. Wille, G. Haidarschin, Y. K. Mishra and R. Adelung, *Adv. Mater.*, 2014, **26**, 1541–1550.
- 79 Y. K. Mishra, G. Modi, V. Cretu, V. Postica, O. Lupan, T. Reimer, I. Paulowicz, V. Hrkac, W. Benecke, L. Kienle and R. Adelung, *ACS Appl. Mater. Interfaces*, 2015, **7**, 14303–14316.
- 80 S. Dhar, T. Majumder and S. P. Mondal, *ACS Appl. Mater. Interfaces*, 2016, **8**, 31822; W. Yin, J. Yang, K. Zhao, A. Cui, J. Zhou, W. Tian, W. Li, Z. Hu and J. Chu, *ACS Appl. Mater. Interfaces*, 2020, **12**, 11797–11805.
- 81 W. Yin, J. Yang, K. Zhao, A. Cui, J. Zhou, W. Tian, W. Li, Z. Hu and J. Chu, *ACS Appl. Mater. Interfaces*, 2020, **12**, 11797–11805.
- 82 V. Postica, I. Paulowicz, O. Lupan, F. Schütt, N. Wol and A. Cojocaru, *Vacuum*, 2019, **166**, 393–398.
- 83 Y. Li, W. Xie, X. Hu, G. Shen, X. Zhou, Y. Xiang, X. Zhao and P. Fang, *Langmuir*, 2010, **26**, 591–597.
- 84 R. Gottesman, A. Song, I. Levine, M. Krause, A. T. M. N. Islam, D. Abou-Ras, T. Dittrich, R. Van de Krol and A. C. Chemseddine, *Adv. Funct. Mater.*, 2020, **30**, 1910832.

