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## High-performance wide bandgap perovskite solar cells fabricated in ambient high-humidity conditions†

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Lead-halide perovskite solar cells (PSCs) are currently the most promising emergent thin-film photovoltaic technology, having already reached power conversion efficiency (PCE) levels of state-of-the-art wafer-based silicon cells. The class of wide bandgap PSCs has also demonstrated high PCE values, thus becoming highly attractive for top sub-cells in tandem devices constructed with silicon or other types of bottom sub-cells. In this study, wide bandgap double-halide ( $\text{Cs}_{0.17}\text{FA}_{0.83}\text{PbI}_{3-x}\text{Br}_x$ ) perovskite absorbers were developed with different bromine content, aiming to obtain bandgap values between 1.66 to 1.74 eV, by a glovebox-free (ambient) procedure. Low-cost inorganic materials, *i.e.*  $\text{TiO}_2$  and  $\text{CuSCN}$ , were used for the electron and hole transport layers, respectively. The 1.70 eV bandgap perovskite resulted in the highest reproducibility and stability (>80% initial PCE after 3500 hours) properties of the PSCs, remarkably attaining 16.4% PCE even with ambient and high humidity (~70%) fabrication conditions.

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## 1. Introduction

Inorganic–organic metal-halide perovskite solar cells (PSCs) have become the most attractive class of thin-film photovoltaic (PV) technology in the last decade, due to the astonishingly fast rise of their power conversion efficiency (PCE) values from 3.8 to 25.5% within 10 years only.<sup>1,2</sup> Besides, this emergent PV technology has shown enormous potential for cost savings, due to the non-vacuum low-temperature deposition and crystallization<sup>3</sup> of the cell layers, enabling also application on inexpensive flexible substrates.<sup>4–8</sup>

The inorganic–organic hybrid perovskites have the general formula  $\text{ABX}_3$ , where A is a monovalent cation composed of methylammonium ( $\text{MA}$ :  $\text{CH}_3\text{NH}_3^+$ ), formamidinium ( $\text{FA}$ :  $\text{CH}_3(\text{NH}_2)_2^+$ ), and/or cesium ( $\text{Cs}^+$ ),<sup>9</sup> B is a divalent metal cation (*e.g.*  $\text{Pb}^{2+}$ ,  $\text{Ge}^{2+}$ ,  $\text{Sn}^{2+}$ ) and X is a halide anion ( $\text{I}^-$ ,  $\text{Br}^-$ , and/or  $\text{Cl}^-$ ).<sup>10</sup> The highest PCEs for single-junction devices have been attained with perovskite absorbers with ~1.55 eV bandgap. Nonetheless, their easy bandgap tuning favours PSC application in tandem (double-junction) architectures, not only with wide-bandgap PSC/low-bandgap PSC configurations<sup>11,12</sup> but also with wide-bandgap

PSCs<sup>13,14</sup> coupled to distinct types of bottom cells based on silicon<sup>15–17</sup> or CIGS materials.<sup>18,19</sup> In this work we fabricated wide-bandgap perovskite layers with different bandgap values by changing the bromine (Br) to iodine (I) ratio at the halide anion site, while a mixture of FA and Cs was used for the A cations.

Albeit PSCs are already competitive with other PV technologies in terms of PCE values, the commercialisation of perovskite-based PV modules is still in suspense; not only because of the well-known stability issues but also due to the presently high fabrication costs. However, techno-economic assessments indicate that PSCs hold the promise of becoming one of the most affordable PV technologies with low levelized cost of energy (LCOE) in the near future.<sup>20,21</sup> Yet, state-of-the-art solution-processed PSCs are commonly fabricated in temperature and humidity controlled glovebox systems, mainly due to the highly environmentally-sensitive organic materials composing the structure.<sup>22,23</sup> This adds critical processing costs and strong limitations in the scalability of this PV technology, thus hindering its industrial exploitation. Therefore, there is a strong motivation to develop the production of high-efficient and stable PSCs in ambient atmosphere conditions, as achieved in this work.

Previous studies in the literature<sup>24–28</sup> explored this possibility, but mainly employing MA-based perovskites with conventional (~1.55 eV) bandgap for single-junction cells. These and many other works have analysed the influence of the ambient humidity on the deposition and annealing of the perovskite layers. Even though it has been generally believed that humidity

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is an adverse detrimental factor in the PSCs' processing, some contributions have challenged this view. For instance, Pathak *et al.* reported that the humidity can improve the perovskite crystallinity, and thus enhance the PSCs' efficiency.<sup>29</sup> Trughton *et al.* fabricated PSCs in high relative humidity (RH) conditions, up to 75%, by applying solvent engineering using ethyl acetate as an antisolvent to absorb moisture during the perovskite layer deposition.<sup>30</sup> In 2019, Wang *et al.* published a detailed study concerning the humidity impacts on MAPbI<sub>3</sub> PSCs, presenting methods for deposition of the perovskite layers in high-mid and low humidity conditions.<sup>31</sup> The authors claimed that, by preheating the substrates before the perovskite layer deposition, the humidity effect on the layer quality and properties can be minimized, thus naming this a humidity-insensitive antisolvent method. Preheating the substrate before the perovskite layer deposition in high humidity levels has also been investigated by Sveinbjörnsson *et al.* for (FAPbI<sub>3</sub>)<sub>1-x</sub>(MAPbBr<sub>3</sub>)<sub>x</sub> perovskite films, showing that preheating has a strong influence on the thickness and the uniformity of the perovskite layer.<sup>32</sup> Cheng *et al.* fabricated MAPbI<sub>3</sub> PSCs with inverted layer structure *via* a two-step deposition method, observing that moisture and oxygen can be removed by preheating the substrate after the first step of PbI<sub>2</sub> deposition.<sup>33</sup>

Nevertheless, despite all the great developments thus far, at present there is still no agreement in the literature regarding the favorable PSC processing conditions in air/humid conditions. Besides, up to the authors' knowledge, the present study is the first report demonstrating highly efficient wide-bandgap PSCs deposited in ambient air conditions with high humidity.

Another critical limitation to the commercialization of PSCs concerns the material costs of the present highly efficient devices. In PSCs, the perovskite absorber is sandwiched between selective contacts for the photo-generated carriers extraction: the electron (ETL)<sup>34</sup> and hole (HTL)<sup>35</sup> transport layers, which block the opposite charges to avoid recombination and current loss. TiO<sub>2</sub> has been the most commonly used ETL

material, benefiting from a favorable band alignment with the perovskite, either in compact form for planar cells,<sup>36</sup> in mesoporous films with pristine TiO<sub>2</sub> crystalline layers<sup>37</sup> as employed in this work, or even with further ETL improvements.<sup>38–40</sup> Being an inexpensive abundant material, TiO<sub>2</sub> does not significantly affect the PSCs' material costs, but the same cannot be said for the conventional HTL materials. Namely, the common organic compounds (Spiro-MeOTAD, PTAA or P3HT) used to synthesize the HTLs are quite expensive. One of the most costly is the typically employed Spiro-MeOTAD, whose price is currently above 300 € per g mainly due to the high number of synthesis and purification steps, as well as reaction steps down to –78 °C required for its production.<sup>37,41–43</sup> This also strongly boosts the CAPEX (capital expenditure) and LCOE of PSC technology.

As such, in view of enabling the market deployment of perovskite-based PV, it is extremely desirable to replace such expensive organic-based HTL materials by inexpensive compounds. There is a recent strong interest towards alternative inorganic-based HTLs, such as CuSCN (copper(i) thiocyanate) or NiO<sub>x</sub>, not only due to their much lower material cost (*e.g.* CuSCN costs ~300 times less than Spiro-MeOTAD),<sup>35,44</sup> but also because they are much more chemically and thermally robust, thus enabling improved PSC stability under environmental conditions. CuSCN revealed to be one the most efficient and less expensive material for the conventional n–i–p PSC structure, while NiO<sub>x</sub> has been mainly used in inverted (p–i–n) PSCs.<sup>45–47</sup>

Here, we present wide-bandgap (WBG) PSCs using low cost, inorganic CuSCN, as the hole transport material, and each step of the synthesis, depositions (except back contact evaporation) and analysis are performed in ambient environment with high humidity conditions (>70%). Our best performing cell exhibits 16.4% PCE with  $V_{oc} = 1.08$  V,  $J_{sc} = 21.0$  mA cm<sup>–2</sup>, FF = 0.72. The PSCs also exhibit long term stability, maintaining >80% of the initial PCE for 3500 hours without any encapsulation. Fig. 1 shows the schematic of the solar cells' layers and energy band diagram of the devices' structure.

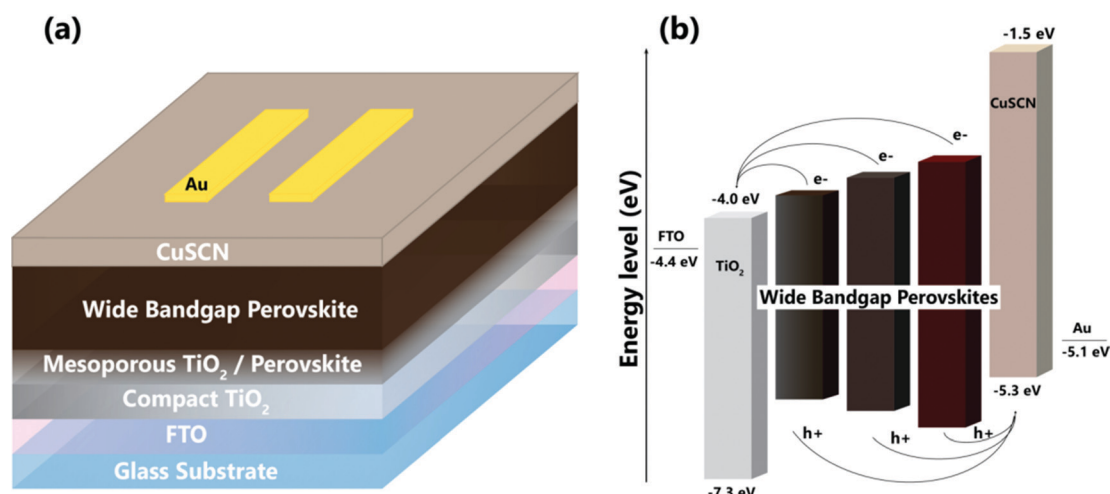


Fig. 1 Architecture (a) and energy band diagram (b) of the wide-bandgap (WBG) perovskite solar cells (PSCs) fabricated in this work, with a superstrate (n–i–p) layer structure, under ambient environment with high humidity conditions (>70%).



## 2. Results and discussion

### 2.1 Analysis of wide-bandgap (WBG) perovskite layers

WBG perovskite films of the type  $\text{Cs}_{0.17}\text{FA}_{0.83}\text{PbI}_{3-x}\text{Br}_x$  were first deposited on FTO-coated glass substrates for optical, structural and morphological characterization. The thicknesses of the layers measured by profilometry are approximately  $550 \pm 50$  nm. It was observed that the change in Br content influences the layers' morphology, causing surfaces wrinkles with roughness up to 45 nm as presented further below, but it does not have a considerable impact on the average thickness of the perovskite layers.

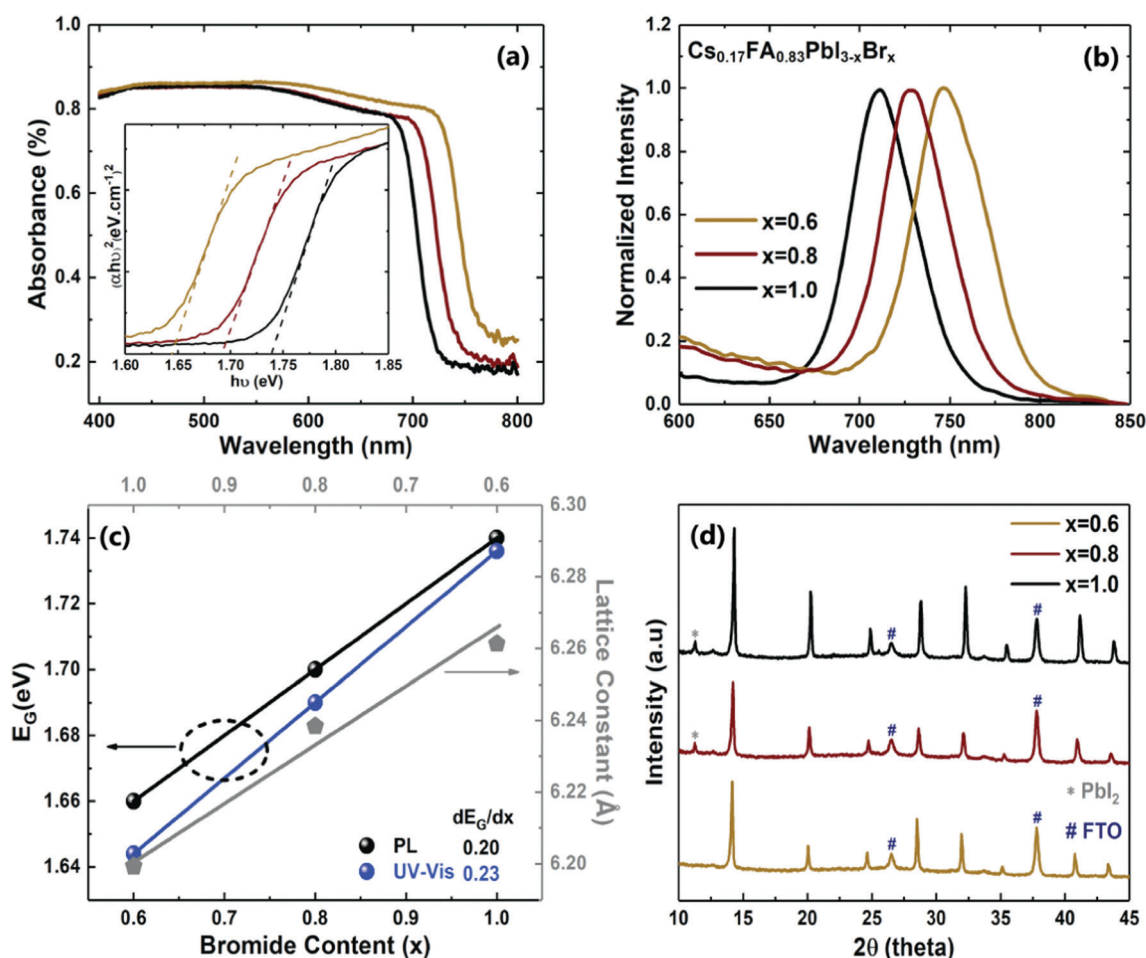
Fig. 2a shows that the absorbance of the perovskite layers on the FTO substrates is more than 80% for the visible region and, depending on the bandgap value (determined by Br content,  $x$ ), the absorbance onset changes in the near-infrared wavelength range. The energy bandgap ( $E_G$ ) values displayed in Fig. 2c were determined *via* the steady-state photoluminescence (PL) peaks (see Fig. 2b) and the Tauc plot method (see inset of Fig. 2a) for direct bandgap semiconductors. It can be clearly seen that the  $E_G$  values increase linearly with higher Br content ratio, as

expected. The slope of the linear regression displaying this characteristic is  $dE_G/dx = 0.2\text{--}0.23$  eV, which proves a facile tuning of the perovskite layer bandgap. The  $E_G$  values obtained from absorbance and PL concur well, with slight differences that reduce as the Br content increases. Besides, the observed stronger PL intensities with increasing Br content (see Fig. S1 in ESI†) is in agreement with the crystallinity measurements described below.<sup>48</sup>

The crystallinity of the perovskite layer is a key factor for the PSCs' efficiency. To scrutinize the crystalline properties of the perovskite absorbers depending on Br content, we calculated the Urbach energies ( $E_U$ ) of the perovskite layers from the absorption spectra *via* eqn (1):

$$\alpha = \alpha_0 + e^{h\nu/E_U} \quad (1)$$

where  $\alpha$  is the absorption coefficient,  $\alpha_0$  is a fitting constant and  $E_U$  is the Urbach energy. The calculated  $E_U$  values (see Table 1) are inversely proportional to the Br content, indicating that the crystallinity properties were enhanced with Br ratio ( $x$ ).<sup>49</sup> As  $x$  increases the Urbach tail energy (attributed to lattice disorder)



**Fig. 2** Optical and structural results of  $\text{Cs}_{0.17}\text{FA}_{0.83}\text{PbI}_{3-x}\text{Br}_x$  perovskite layers on FTO-coated glass substrates, with different Br content ( $x = 0.6, 0.8, 1.0$ ). (a) Absorbance spectra with Tauc plots  $[(\alpha h\nu)^2 \text{ vs. } h\nu]$  in inset. (b) Normalized steady-state PL measurements. (c) Relation between the perovskite bandgap,  $E_G$ , and the lattice constant, as a function of Br content. (d) X-Ray diffraction plots of the WBG perovskite layers, whose peaks are shown in further detail in Fig. S2 of the ESI.†





**Table 1** Urbach energy, surface roughness and average grain size values of the perovskite layers with different Br content in the  $\text{Cs}_{0.17}\text{FA}_{0.83}\text{PbI}_{3-x}\text{Br}_x$  composition

$x$ (Br content ratio)	0.6	0.8	1.0
Urbach energy ( $E_U$ , meV)	63	57	53
RMS roughness (nm)	29.68	42.29	45.02
Mean grain size (nm)	137.3	190.1	205.6

decreases,<sup>49</sup> thus lower values of  $E_U$  correspond to higher order in the crystal structure. The Urbach values attained for our ambient-deposited perovskite films are slightly higher than those of state-of-the-art perovskites deposited under glovebox controlled environment,<sup>17</sup> and this leads to relatively higher open-circuit voltage deficit as analysed in ESI† (Fig. S4).

The aforementioned impact of Br content in the perovskite crystallization is confirmed by X-ray diffractometry (XRD) measurements shown in Fig. 2d. All samples display the perovskite peak around  $2\theta = 14.2^\circ$ , referring to the (100) perovskite plane. Additionally, the peaks at  $20^\circ$ ,  $24.8^\circ$ ,  $28^\circ$  and  $41.5^\circ$  are also found in wide-bandgap perovskite layers.<sup>50</sup> In general, with the increase of Br content the XRD peak intensities increase, indicating improved crystallinity. On the other hand, the perovskite layer with the lowest Br content ( $x = 0.6$ ) does not exhibit any photoinactive  $\text{PbI}_2$  peaks ( $\sim 12.6^\circ$ ), which can be attributed to the lower precursor concentration, and thus better dissolubility and homogeneity of the perovskite solution. This sample also has a pronouncedly lower surface roughness determined from AFM measurements, as presented in Table 1 and discussed further below.

The XRD results show that, when increasing the Br:I ratio, a slight shift to higher angle values occurs, as presented in further detail in Fig. S2 of ESI†. It is known that higher Br content causes the lattice to shrink, due to the smaller size of the bromide ion when compared with the iodide ion,<sup>17,51</sup> which transforms the tetragonal phase of  $\text{MAPbI}_3$  into the cubic phase of pure  $\text{MAPbBr}_3$ . The same behaviour occurs in FA and Cs mixed-cation perovskites.<sup>52</sup> The lattice constant calculations (see Fig. 2c) for the main perovskite peak of  $2\theta = 14.2^\circ$  confirm

such lattice contraction, and the linear relation between the lattice constant and Br content matches Vegard's law.<sup>17</sup>

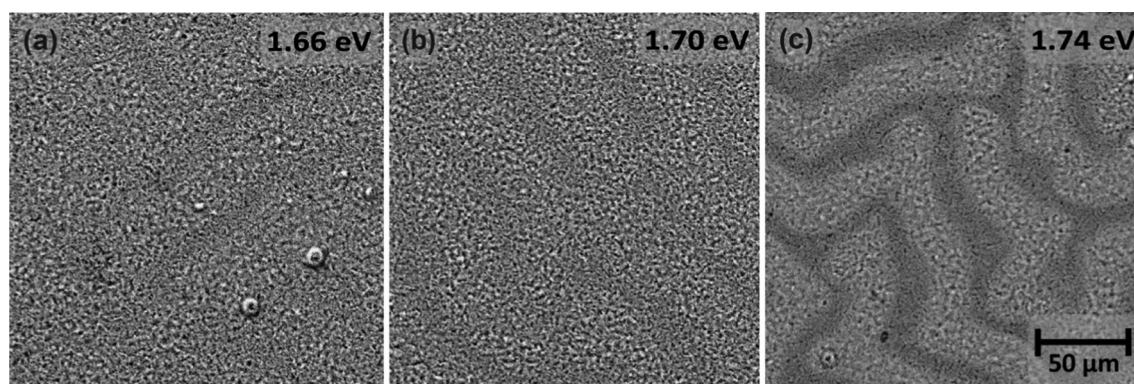
The AFM images (shown in Fig. S3, ESI†) indicate that when increasing  $x$ , the RMS surface roughness also increases (Table 1). The reduced roughness of the perovskites with lowest  $x$  should be caused by the homogeneity of the precursor solutions and Br:I ratio in the layers. Also, the change in the roughness with Br content is attributed to the wrinkled morphology of the layers, as visualized by scanning electron microscopy (SEM, see Fig. 3). As suggested in the literature, the wrinkled-shaped structures are a consequence of compact stress during the film formation, depending on the precursor solution, antisolvent washing and annealing step.<sup>53–55</sup>

Although the wrinkled surface morphology does not significantly affect the optical properties of the films, the device performance was reduced mainly due to the less conformal deposition of the hole transport layer.<sup>49,56</sup> As shown in ESI† (Fig. S3), higher  $x$  leads to relatively deeper cracks in the films, which causes higher density of recombination centers, short-circuits and overall loss in the PSCs' performance.<sup>57</sup> Therefore, even though the perovskite layer with the highest Br content has better crystallinity properties and lower Urbach energy values, the resulting solar cell parameters were worse than the other cells, as shown in the next section.

Besides, as shown in the SEMs of Fig. 4 with higher magnification and in Table 1, with the increment of bromide content the average grain size of the perovskite films increases, mainly due to the appearance of large ( $>300$  nm) grains causing a higher positive skewness (*i.e.* lower homogeneity) of the size distribution.

## 2.2 PSCs' optoelectronic assessment and optimization

**2.2.1 Solar cell characterization.** The current density–voltage ( $JV$ ) measurements of the solar cells composed of the WBG perovskites with the three bandgap values are shown in Fig. 5a. The cells were fabricated with the layer structure: FTO/ $\text{TiO}_2$ / $\text{Cs}_{0.17}\text{FA}_{0.83}\text{PbI}_{3-x}\text{Br}_x$ /CuSCN/Au, as described in the experimental methods. All the preparations and layer depositions were



**Fig. 3** Low magnification SEM images of the WBG perovskite layers on FTO-coated glass substrates, composed with different Br content: (a)  $x = 0.6$ , (b)  $x = 0.8$ , (c)  $x = 1.0$ . It can be seen that the overall homogeneity is better for the perovskite with the lower Br content (and lower  $E_G$ , displayed in top-right corner), while increased wrinkling is observed with higher  $x$ .



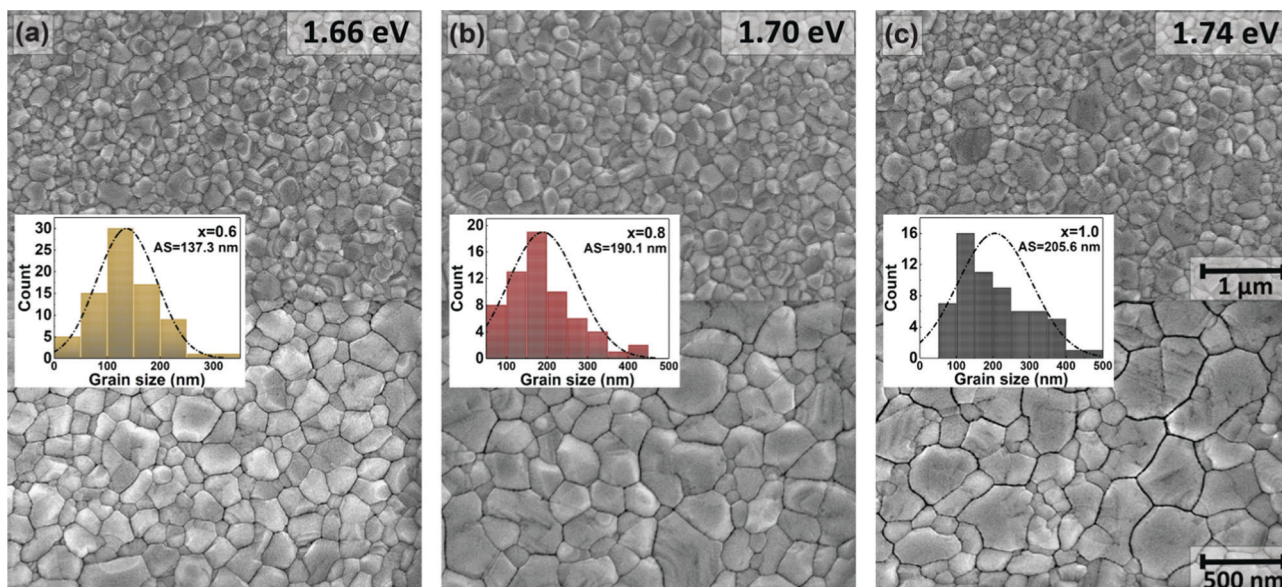


Fig. 4 Top-view SEM images of the WBG perovskite layers on FTO-coated glass substrates, composed with different Br content: (a)  $x = 0.6$ , (b)  $x = 0.8$ , (c)  $x = 1.0$ . The images reveal the multi-crystalline grain structure of the films at two different magnifications. The inset histograms show the grain size distribution and average grain size (AS, also indicated in Table 1).

performed in ambient air with relative humidity,  $RH > 70\%$ . The PV quantities of interest are shown in Fig. 5b and summarized in Table 2.

As expected, there is a reduction in the PCE with increasing  $E_G$  of the perovskite absorber. However, an opposite trend should have been observed for the open-circuit voltage ( $V_{oc}$ ), but here that is not the case since the  $V_{oc}$  deficit ( $\Delta V_{oc} = \frac{E_G}{q} - V_{oc}$ ) increases with the bandgap widening (see Fig. S4, ESI†). This effect is attributed to inevitable phase segregation in the perovskite layer, which becomes more pronounced when the fabrication occurs in a humid environment.<sup>58,59</sup> Also, the aforementioned non-ideal interface properties, between the perovskite and charge transport layers, cause  $V_{oc}$  reduction and higher series resistance. As for the

Table 2 Main solar cell parameters of the WBG PSCs with different  $E_G$  values of the perovskite layer

$E_g$ (eV)	$V_{oc}$ (V)	$J_{sc}$ (mA cm <sup>-2</sup> )	FF	PCE (%)	$R_{series}$ (Ω)	$R_{shunt}$ (Ω)
1.66	1.00	20.96	0.71	15.1	62.8	25 517
1.70	1.03	20.43	0.68	14.4	86.6	32 335
1.74	0.99	17.61	0.64	11.1	120	13 734

shunt resistance ( $R_{shunt}$ ), it does not present a linear trend since it is higher for the mid-bandgap solar cell relative to the others.

Regarding the short-circuit current density ( $J_{sc}$ ), it decreases with higher bandgap of the perovskite layer. This is expected optically since the absorbance onset is blue-shifted, as seen in Fig. 2a.<sup>60</sup> In addition, the increased roughness with higher Br

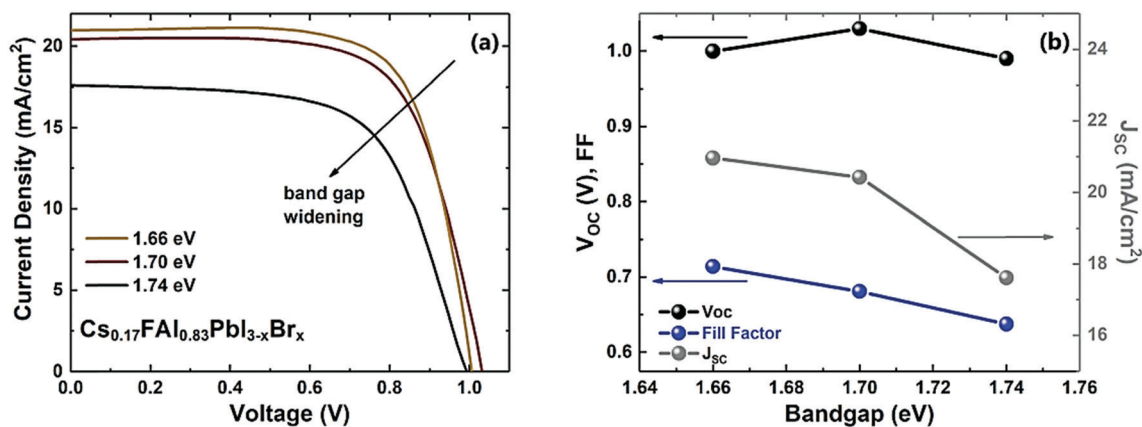
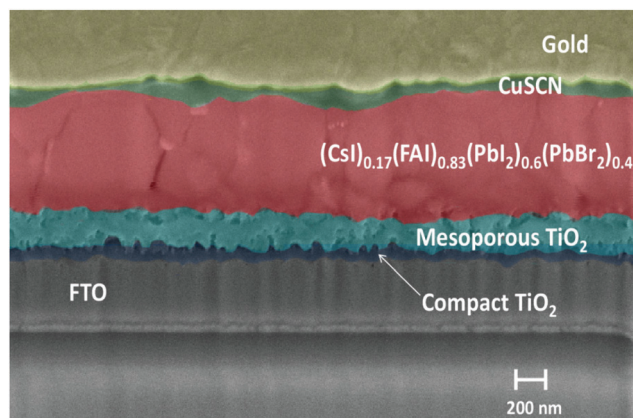


Fig. 5  $J$ - $V$  curves (a) and variation of characteristic parameters (b) of the WBG PSCs with the  $E_G$  values.



**Table 3** Solar cell parameters of the best performing WBG (1.70 eV) PSC

Best results:	$E_g$ (eV)	$V_{oc}$ (V)	$J_{sc}$ (mA cm <sup>-2</sup> )	FF	PCE (%)	$R_{series}$ ( $\Omega$ )	$R_{shunt}$ ( $\Omega$ )
	1.70	1.08	21.0	0.72	16.4	73.2	25 244

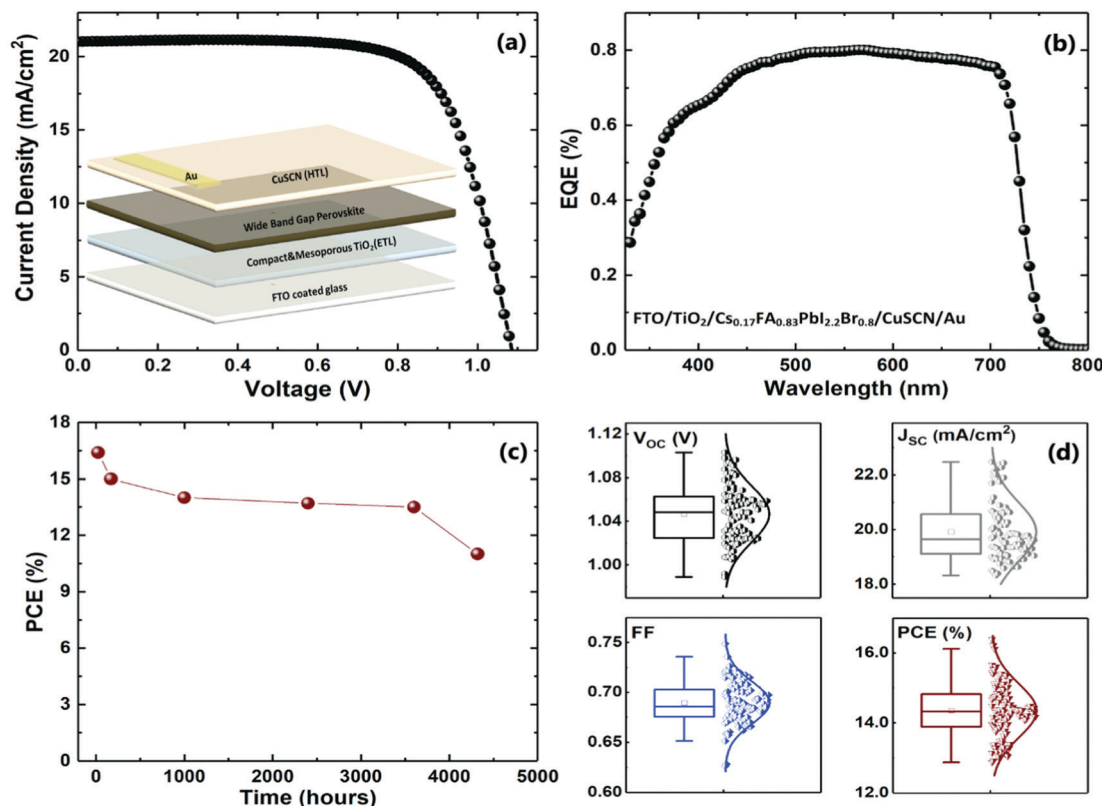
**Fig. 6** Cross-sectional SEM-FIB image of FTO/TiO<sub>2</sub>/(CsI)<sub>0.17</sub>(FAI)<sub>0.83</sub>(PbI<sub>2</sub>)<sub>0.6</sub>(PbBr<sub>2</sub>)<sub>0.4</sub>/CuSCN/Au solar cell.

content causes more charge recombination at the interfaces, which also justifies the decrease of the fill factor (FF) values.<sup>61</sup>

Although the highest PCE of this set was attained with the 1.66 eV bandgap perovskite, the higher  $V_{oc}$  and  $R_{shunt}$  with the 1.70 eV perovskite motivated a subsequent study focused in the improvement of the 1.70 eV PSCs, in view of their promising application for tandem solar cells due to the suitable bandgap value.

**2.2.2 Optimization of WBG perovskite solar cells with  $E_g = 1.7$  eV absorber.** We proceed with further optimization of the wide-bandgap (1.70 eV) PSCs with the FTO/TiO<sub>2</sub>/Cs<sub>0.17</sub>FA<sub>0.83</sub>PbI<sub>3-x</sub>Br<sub>x</sub>/CuSCN/Au configuration, where  $x = 0.8$ . The device fabrication procedure was fine tuned mainly in the UV treatment time applied on the ETL, substrate preheating prior to perovskite deposition, antisolvent washing and application parameters (volume, spin-coating speed) during perovskite crystallization, as well as perovskite layer annealing conditions. Consequently, we attained the highest solar cell efficiency (PCE = 16.4%) of this study, presented in Table 3, employing the optimized fabrication parameters indicated in Table S1 (ESI<sup>†</sup>). As shown further below, such optimized process also allowed a noteworthy reproducibility and stability without any encapsulation.

In Fig. 6, the cross-sectional SEM shows that the perovskite layer thickness is in the 560–600 nm range, exhibiting some defects in the layer bulk and at both interfaces with the ETL and HTL. This reveals that there is still room for improvement in order to further increase the performance, and reach that of

**Fig. 7**  $J$ - $V$  characteristic under AM 1.5 illumination (a), and EQE spectrum (b) of the best performing WBG (1.70 eV) PSC, which presents remarkable stability up to 4500 hours (c) even without any cell encapsulation. (d) Distribution and statistical analysis of the solar cell parameters of over 100 fabricated cells.

state-of-the-art cells produced in a controlled glovebox environment. Moreover, the dark  $I$ - $V$  characteristic (Fig. S5, ESI†) shows a high leakage current in reverse bias, causing a worse diode character, and thus reduced efficiency, compared to the record literature values with similar WBG-PSC structures. This also contributes to the commonly observed hysteresis effect of PSCs, shown in Fig. S6 of ESI†, attributed to the mobile ions/ion-vacancies in the perovskite material which cause a pronounced dependence of the cell  $JV$  characteristic with the scan direction.

As referred before, Wang *et al.*<sup>31</sup> described a “turbid point (TP)” observed during the perovskite layer deposition when the transparent perovskite film becomes turbid during the spin coating, and this point is dependent on the humidity levels of the ambient air. In our work, the most efficient PSC was achieved by preheating the substrates and applying antisolvent (chlorobenzene, CB) dripping 1–2 seconds before the turbid point, as well as by using warm antisolvent to help the solvent evaporation process and removal of oxygen and water. When the perovskite layers were deposited in  $\leq 50\%$  RH, compatible with Wang’s observations, the turbid point appears sooner, and as a result, a thinner layer grows. We also applied the CB washing placing the pipette tip very close (less than 5 mm) to the substrate, and noticed a significant enhancement on the perovskite layer formation. This impact can be attributed to speeding up the evaporation process, hence promoting better layer crystallization. Fig. S7 in ESI† shows the effects of key fabrication steps in the resulting  $JV$  curves of the PSCs, whose optimization essentially allows minimizing the defect density in the perovskite layer and, thereby, improving the devices’ performance. All in all, this study reinforces Wang’s humidity insensitive antisolvent method, providing further understanding of the process of PSCs fabrication in humid ambient air.

To confirm the PSCs reproducibility, we fabricated more than 100 solar cells with WBG perovskite layers in the FTO/compact-TiO<sub>2</sub>/meso-TiO<sub>2</sub>/(CsI)<sub>0.17</sub>(FAI)<sub>0.83</sub>(PbI<sub>2</sub>)<sub>0.6</sub>(PbBr<sub>2</sub>)<sub>0.4</sub>/CuSCN/Au configuration, and under  $\sim 70\%$  RH ambient conditions. The  $J$ - $V$  characteristics and external quantum efficiency (EQE) measurements of the best performing cell are shown in Fig. 7a and b. The integrated current density values derived from the EQE spectra match that of the  $J$ - $V$  characteristics acquired with the solar simulator. Remarkably, without any encapsulation the best performing solar cell kept more than 80% of its initial efficiency for more than 3500 hours, and almost 70% after 4500 hours, which is shown in Fig. 7c.

In Fig. 7d, the solar cell parameters of the  $E_G = 1.7$  eV PSCs are displayed over 100 fabricated devices. The main noteworthy aspect is that, although the humidity and temperature conditions in the laboratory are not controlled, the optimized process developed in this work allows highly reproducible cells with high efficiency values. The average of the  $V_{oc}$ ,  $J_{sc}$ , FF and PCE values are 1.05 V, 19.9 mA cm<sup>-2</sup>, 0.70 and 14.4%, respectively. These results are among the best attained up to now in the PSCs’ field, namely in comparison with the previous reported advances attained with un-encapsulated WBG PSCs fabricated in ambient humidity conditions.<sup>62,63</sup>

### 3. Conclusions

We present a low-cost glovebox-free procedure to contrive WBG PSCs, whose optical bandgap values were tuned between 1.66 to 1.74 eV by changing the Br content in the perovskite structure. Here, the wrinkled morphology of the surface deteriorated the interface properties, and thus caused relatively lower efficiency for higher Br:I ratio. The perovskite layer with 1.70 eV bandgap was selected for further optimization due to its potential application in tandem devices. A large set of mixed-cation mixed-halide WBG PSCs with FTO/compact-TiO<sub>2</sub>/meso-TiO<sub>2</sub>/(CsI)<sub>0.17</sub>(FAI)<sub>0.83</sub>(PbI<sub>2</sub>)<sub>0.6</sub>(PbBr<sub>2</sub>)<sub>0.4</sub>/CuSCN/Au configuration were fabricated, in which a record cell is reported with the best efficiency (16.4%) attained thus far for ambient-processed WBG PSCs. In addition, this cell presents a remarkable stability, maintaining 80% of the initial efficiency for 3500 hours, and 70% up to 4500 hours. These outcomes reveal that it is possible to reproducibly fabricate high-performing and stable WBG PSCs without a controlled glovebox environment, and employing inexpensive charge transport materials such as the one used here (CuSCN) which is much cheaper than the conventional organic-based HTLs. This was attained *via* a humidity-insensitive antisolvent method, in which it was found crucial to preheat the substrate before the perovskite deposition, as well as to preheat the antisolvent and apply the antisolvent washing from a very close tip-sample distance, which enhances the evaporation and crystallization process of the perovskite film by reducing the moisture effects.

In addition to the set of morphological, structural and optoelectronic characterization techniques employed in this work, further insight into the role of the resulting bulk and interfaces’ properties on the PSCs’ transport mechanisms can be gained *via* electrochemical impedance/admittance analysis, which can also probe and provide a better understanding of the charge distribution and transfer within the PSC structure, including ion migration and, thus, hysteresis effects, as discussed in related works.<sup>38,64–68</sup>

Lastly, the main avenues that should be explored for further advances are the extended optimization of the absorber layer with wider (up to  $\sim 2$  eV) bandgap values, in view of application in tandem devices with silicon and low bandgap PSCs. Additionally, the implementation of photonic structures in the PSCs should be pursued, not only as a means to increase the photocurrent generation (and thereby the PCE) *via* light capturing and trapping,<sup>69</sup> but also to improve the devices’ stability against UV<sup>70,71</sup> and water<sup>72</sup> penetration as demonstrated in recent contributions.

### Experimental methods

#### Materials

All chemical reagents used in this study were purchased from commercial vendors. Formamidinium iodide (FAI) (Sigma Aldrich), lead iodide (PbI<sub>2</sub>) (Sigma Aldrich), cesium iodide (CsI) (Acros Organics), lead bromide (PbBr<sub>2</sub>) (Sigma Aldrich), titanium IV-isopropoxide (TTIP) (Sigma Aldrich), titanium



dioxide (TiO<sub>2</sub>) paste (Sigma Aldrich), copper(i) thiocyanate (CuSCN) (Alfa Aesar), diethyl sulfide (DES) (Sigma Aldrich), dimethylformamide (DMF) (Fisher Chemical), dimethyl sulfoxide (DMSO) (Panreac), chlorobenzene (CB) (Alfa Aesar) were used without further purification.

### Substrate preparation and ETL layer deposition

The wide-bandgap perovskite solar cells were fabricated with mesoscopic structure: Glass/FTO/TiO<sub>2</sub>/Cs<sub>0.17</sub>FA<sub>0.83</sub>PbI<sub>3-x</sub>Br<sub>x</sub>/CuSCN/Au. All the preparations and layer depositions were performed in ambient air with RH > 70%.

FTO-coated glasses (13 Ω sq<sup>-1</sup>, 82–84.5% of transmittance) were cut into 2 × 2 cm, then a 2 mm stripe of the FTO layer was etched with zinc powder and HCl solution, and afterwards cleaned by sonication with Hellmanex II, water, acetone, isopropanol alcohol and an additional step of UV-O<sub>3</sub> treatment to remove the organic compounds right before the deposition of the electron transport layer (ETL).

The first deposited layer of the solar cell is a very thin compact TiO<sub>2</sub> film. The solution is prepared by a two-step method. Firstly, 180 μl of TTIP was added to 1.25 ml of absolute ethanol and placed onto a stirrer, and then 1.25 ml of absolute ethanol with 18 μl of HCL was added to the main solution. The compact TiO<sub>2</sub> layers were deposited by spin coating for 35 seconds at 4000 rpm. After the spin coating, the layers were dried at 100 °C for 10 minutes, and then sintered in a furnace at 500 °C in ambient air. After the samples cool down to room temperature, another UV-O<sub>3</sub> treatment was applied for 15 minutes before the mesoporous TiO<sub>2</sub> layer deposition. For this layer, 120 mg of titania paste was dissolved in 1 ml absolute ethanol and deposited by spin coating for 20 seconds at 4000 rpm, followed by a drying step at 120 °C for 10 minutes and sintering at 450 °C for 30 minutes in ambient air.

### Perovskite solutions and deposition

The perovskite precursor solutions were prepared with the composition of (CsI)<sub>0.17</sub>(FAI)<sub>0.83</sub>(PbI<sub>2</sub>)<sub>1-0.5x</sub>(PbBr<sub>2</sub>)<sub>0.5x</sub> using different bromide content by dissolving CsI (88.9 mg), FAI (285 mg) and PbI<sub>2</sub>, PbBr<sub>2</sub> in 1 ml mixture of DMF and DMSO (with a volume ratio of 9:1). For the bandgap variation, *x* was chosen as 0.6, 0.8 and 1 during the precursor solution preparation. The precursor solutions were stirred at 70 °C until they reach a bright yellow colour. The solutions were filtered by a PTFE filter and the substrates coated with the ETL were heated up to 70 °C right before the deposition. The perovskite layer was deposited *via* a single-step deposition method with two different spin rates: 1000 rpm for 10 seconds and then increases to followed by an increase to 5000 rpm for a 20 seconds spin. Chlorobenzene was poured during the spinning at the higher spin rate to remove the solvents of the perovskite solution and improve crystallization. Soon after the deposition of the perovskite layer, the samples are annealed at 110 °C for 20 minutes on a hot plate covered by a Petri dish. As soon as the annealing process starts the layers turn to dark brown color from a light yellow one.

### Hole transport layer and metal contact deposition

Following the perovskite layer deposition, CuSCN (copper(i) thiocyanate) was used as the hole transport material. Dissolving in diethyl sulfide, the solution was prepared as 0.3 M, filtered by a 0.45 μm filter and deposited by spin coating for 20 s at 4000 rpm, followed by 5 minute annealing applied at 90 °C to evaporate the HTL solvent to avoid the deterioration of the perovskite layer. The gold back contacts were deposited by e-beam evaporation under high vacuum at a low deposition rate. The cells were produced with an active area of 0.10 cm<sup>2</sup>.

### Characterization techniques

The optical transmission measurements were performed by UV-Vis-NIR spectrophotometry (PerkinElmer Lambda 950) in 300–800 nm wavelength range. Steady-state photoluminescence (PL) was measured by a PerkinElmer LS55, with an excitation at 400 nm. Scanning Electron Microscopy (SEM) surface images were acquired by a TM3030 Plus Hitachi for lower and a Hitachi Regulus 8220 for higher magnification observations, and the corresponding grain size analysis was performed *via* ImageJ software. For the cross-sectional imaging, a Carl Zeiss AURIGA Cross Beam SEM work station was used. Surface topography Atomic Force Microscopy (AFM) images were obtained by an Asylum Research MFP-3D Standalone system using commercially available silicon AFM probes (Olympus AC160TS, *f*<sub>0</sub> = 300 kHz, *k* = 26 N m<sup>-1</sup>). The X-ray diffraction (XRD) patterns were collected on a PANalytical XPert Pro X-ray powder diffractometer using Cu-Kα radiation (*λ* = 1.54050 Å). The *J*–*V* measurements were performed with Newport's Oriel® VeraSol LSH-7520 LED-based Solar Simulator (100 mW cm<sup>-2</sup>) under AM 1.5 conditions, calibrated by a crystalline silicon reference cell (from VLSI Standards Inc.), in ambient conditions. A Keithley 2500 power source was used to apply voltages and collect current from the devices. The Quantum Efficiency measurements of the solar cells were carried out by Newport's QUANTX-300 system.

### Conflicts of interest

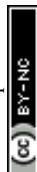
There are no conflicts to declare.

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