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Exotic physical properties of 2D materials modulated by moiré superlattices

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van der Waals heterostructures of two-dimensional materials are naturally endowed with the nanoscale moiré pattern, which has become a versatile platform for studying novel quantum phenomena during the past few decades. Here, we discussed the moiré superlattice combining graphene and hexagonal boron nitride (h-BN) which led to the observation of electronic minibands, superlattice Dirac points and anomalous physical phenomena such as the quantum Hall effect and the Hofstadter butterfly pattern. The twist bilayer graphene (TBG) system has witnessed the appearance of angle-dependent van Hove singularities. Moreover, in the magic-angle twisted bilayer graphene (MATBG), there is a superconductor–insulator transition due to the enhancement of van der Waals forces by interlayer resonant conditions. Besides, owing to the moiré superlattice constant being comparable to the Bohr radius of excitons, exciton physics dominated by the moiré potential for the heterojunction based on transition metal dichalcogenide (TMD) monolayers has attracted extensive attention. This review mainly covers the exotic physical properties, ranging from electronic structures to excitonic properties of graphene-based and TMD-based van der Waals-coupled twistrionic materials, and discusses the remaining unsolved strongly correlated questions and possible expected semiconductor device applications in the quantum matter world.

1. Introduction

Two-dimensional (2D) materials have become the most popular material system since 2004 when graphene was mechanically exfoliated from bulk graphite. Galvanized by this breakthrough,

researchers have isolated and studied dozens of 2D crystals and predicted the existence of many more. These atomically layered materials that exhibit intriguing properties distinct from those of their bulk counterparts can be fabricated as building blocks for achieving advanced material properties. A so-called van der Waals (vdW)¹ material is characterized by covalent bonding within the layer and weak vdW forces between the 2D plane. This has led many researchers to study the flatland regime of 2D materials ranging from fundamental physics^{2,3} to device applications such as high-performance photonic and optoelectronic

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devices.^{4–8} When two single layers are stacked with a rotation, they produce a moiré pattern of much larger length scale than the periodicity of each layer. The pattern is more often than not an incommensurate structure and exhibits twist-dependent emergent electronic behavior,^{9–12} whereas an overarching periodicity emerges in the local atomic registry of the constituent crystal structures, which is known as a moiré superlattice.^{13–18}

The moiré superlattice induced by the lattice mismatch or crystal orientation misalignment can significantly modify the electronic and optical properties of heterostructures.¹⁹ Graphene/hexagonal boron nitride (h-BN), a typical vdW heterostructure, exhibits greatly improved properties for device applications by stacking graphene on top of h-BN. In the graphene/h-BN heterostructure, the h-BN can reduce ripples and suppress charge inhomogeneities to improve the carrier mobility of graphene,^{20,21} as well as provide unique opportunities to open the band gap of graphene at the Dirac point by the periodic superlattice potential,^{13,22–27} which leads to various anomalous quantum phenomena including the quantum Hall effect and the Hofstadter butterfly pattern.^{13–15,28} Twisted bilayer graphene (TBG) is simply fabricated by changing the relative orientation between the layers, which induces low-energy Van Hove singularities (VHSs).²⁹ This may provide intriguing prospects for VHS engineering of electronic phases. In some special cases, the two graphene layers are twisted relative to each other at a small angle which is called the magic angle, and there is a correlated insulator and unconventional superconductor transition due to the localization of electrons within the moiré superlattice.^{30,31}

Although the distinctive linear dispersion of graphene yields unique electronic and optical properties, the zero-band gap greatly hinders its applicability in semiconductor based elements and devices. The TMD monolayer exfoliated from a bulk crystal possesses direct energy gaps in the range of 1–2 eV and opened up a new window for optoelectronics. The reduced dimensionality and the associated reduced dielectric screening result in a greatly enhanced exciton binding energy in monolayer TMDs. The strong excitonic effects lead to room temperature excitonic devices that have promising prospects for overcoming the dilemma of response time and integration in electron or/and photon based

devices.^{32–34} Heterobilayers of TMDs are particularly attractive for low-dimensional semiconductor optics because they host interlayer excitons – with electrons and holes localized in different layers – and thereby possess various novel and appealing properties, especially in a typical type II band alignment, where electrons and holes find their energy minima in different layers.^{35–37} Besides, owing to the high sensitivity of TMDs to crystal structure symmetry and interlayer coupling, the bandgaps of TMD heterobilayers composed of various 2D materials are of big difference. Furthermore, the moiré potential determined by interlayer coupling in TMD heterobilayers can localize interlayer excitons and modulate the multiple interlayer exciton resonances, suggesting that the feasibility of engineering artificial excitonic crystals can provide a unique platform for studying novel quantum phases of matter.^{38–44}

In this review, we mainly aim to describe exotic electronic and excitonic properties modulated by the nanoscale moiré pattern in graphene-based and TMD-based heterostructures. Firstly, we overview how the electronic structure of graphene is modified by the moiré superlattice and the anomalous quantum phenomena of graphene/h-BN. Then, we discuss the transport properties of twisted bilayer graphene. Moreover, the interplay of correlated insulating states and superconductivity of twisted bilayer graphene is also described. Finally, we review excitonic phenomena determined by the moiré potential in twisted TMD heterobilayers and come up with several possible perspectives to inspire the research enthusiasm for semiconductor quantum light-emitting devices based on moiré superlattices.

2. Quantum phenomena in the graphene moiré superlattice

Graphene has attracted great attention owing to its remarkable properties as the first 2D vdW material.^{45–48} The recent development of experimental techniques for combining graphene with other atomically thin vdW crystals has enabled the exploration of the properties of these so-called vdW heterostructures, which provides an exciting approach to design quantum materials that can harness and extend the already fascinating properties of the constituents. In this section, we first discuss the exotic experimental phenomena in the emergent graphene/h-BN moiré superlattice, such as superlattice Dirac points and Hofstadter butterfly spectrum. Then we illustrate the intriguing properties in twisted bilayer graphene (TBG) due to a stronger interlayer coupling than that in the graphene/h-BN moiré superlattice. Finally, we describe two correlated quantum states of correlated insulating states and unconventional superconductivity caused by a strong electron–electron Coulomb interaction in the magic-angle twisted bilayer moiré superlattice of graphene.

2.1 Graphene/h-BN moiré superlattice

Heterostructures of 2D vdW materials offer a platform for engineering novel optoelectronic properties. h-BN is another well-known vdW material that is commonly used. It can provide



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high and homogeneous dielectric conditions for graphene, enabling the observation of the quantum Hall effect (QHE) and fractional quantum Hall effect (FQHE) in graphene/hBN heterostructures.^{49–51} Additionally, h-BN can modify the optoelectronic properties of graphene when aligned with graphene under certain circumstances, for example, the emergence of superlattice Dirac points, accompanied by the appearance of the Hofstadter butterfly spectrum.^{13–18,28} The graphene/h-BN heterostructure provides an opportunity for engineering the properties that can be achieved by tuning each of these degrees of freedom, thereby facilitating the design and realization of graphene-based devices.

2.1.1 Quantum Hall effect in graphene. Quantum-mechanically enhanced transport phenomena such as the quantum Hall effect can be observed when electrons are confined in two-dimensional materials. Fig. 1a shows a schematic of a Hall bar device made of graphene and h-BN used to measure the quantum Hall effect.⁵² Highly degenerate Landau-levels (LLs) can be formed when graphene is applied with a perpendicular magnetic field.^{51,53} The energy dispersion of the LLs is given by

$$E_n = \text{sgn}(n)\sqrt{2e\hbar v_F^2|n|B}, \quad (1)$$

where e is the electron charge, \hbar is Planck's reduced constant, the integer n represents an electron-like ($n > 0$) or a hole-like ($n < 0$) LL index, v_F is the Fermi velocity and B is the magnitude of the applied magnetic field. Unlike the traditional two-dimensional electron gas, the LLs formed in graphene under a magnetic field are not equally spaced in energy, as shown in Fig. 1b. When the Fermi level (E_F) of graphene crosses a LL by adjusting the magnetic field, the Hall conductance σ_{xy} exhibits a quantized plateau:

$$\sigma_{xy} = \nu \frac{e^2}{h} = g \left(n + \frac{1}{2} \right) \frac{e^2}{h}, \quad (2)$$

where $\frac{1}{2}$ comes from the Berry phase, ν is the quantized filling factor of LLs and E_F jumps by an amount of $g e^2/h$. The longitudinal resistance R_{xx} is zero, while the Hall resistance R_{xy} exhibits a quantum Hall plateau, indicating that the system enters the quantum Hall regime (Fig. 1c). Note that $\sigma_{xy} = -R_{xy}^{-1}$ and $\sigma_{xx} = -R_{xx}^{-1}$. This Landau quantization is a result of the linear band dispersion of graphene. Additionally, the valley degree of freedom in graphene adds an extra twofold degeneracy, but, in most materials, the degeneracy of each LL is doubled due to the spin degree of freedom. Considering the valley and spin degrees of freedom of graphene, there are half-integer multiple quantum Hall plateaus, that is, the fractional quantum Hall effect. Fig. 1d shows an expanded view in the $N=0$ LL between $\nu = 0$ and $\nu = 1$. Two-flux composite fermion states (centered around $\nu = 1/2$) and four-flux composite fermion states (centered around $\nu = 1/4$) are observed.⁵⁴

2.1.2 Quantum phenomena in the graphene/h-BN moiré superlattice. In the graphene/h-BN heterostructure wherein rotational mismatch is small, a long-period geometric interference pattern can be produced. Because the electronic and

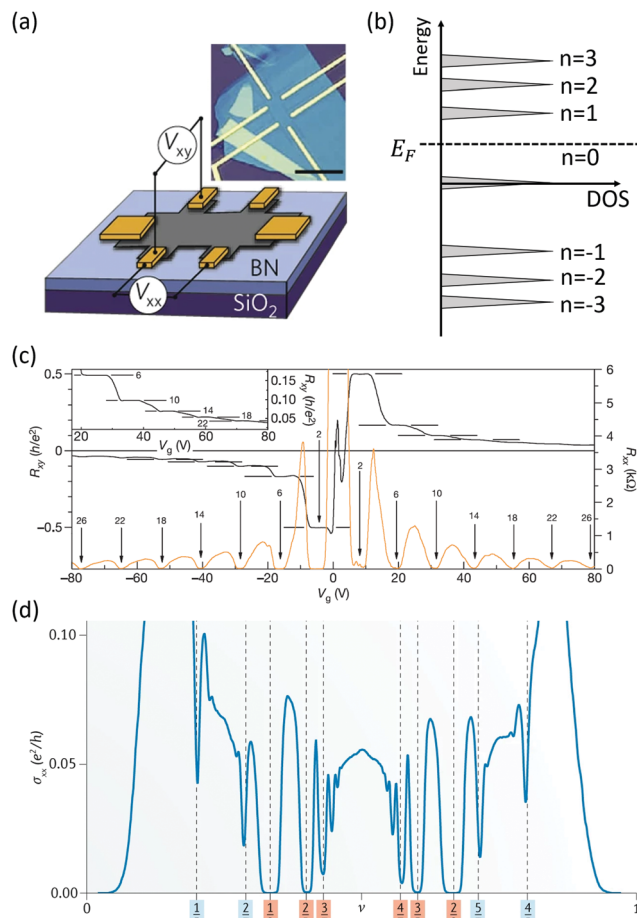


Fig. 1 (a) Schematic of a Hall bar device made of graphene and h-BN used to measure the quantum Hall effect.⁵² Copyright 2011, Springer Nature Limited. (b) A schematic diagram of the DOS of LLs formed by quantization of Dirac fermions under an external perpendicular magnetic field.⁵¹ Copyright 2005, Springer Nature Limited. (c) Hall resistance (black) and magnetoresistance (orange) as a function of gate voltage under a fixed magnetic field and a low temperature. The upper inset is a diagram of a high-filling-factor plateau measured at different temperatures.⁵¹ Copyright 2005, Springer Nature Limited. (d) Device conductance (σ_{xx}) at $B = 15$ T between $\nu = 0$ and $\nu = 1$. The red labels indicate the 2-flux fractional quantum Hall effect and blue labels indicate the 4-flux sequence.⁵⁵ Copyright 2019, Springer Nature Limited.

optical properties of the heterostructure can be significantly modified by the moiré pattern, the unique low-temperature transport phenomena of the graphene/h-BN moiré superlattice have been studied extensively. Fig. 2a shows the STM image of the graphene/h-BN moiré pattern. The period of this moiré pattern is very sensitive to lattice mismatch δ between graphene and h-BN and their relative rotational orientation. This relationship can be described as¹⁶

$$\lambda = \frac{(1 + \delta)a}{\sqrt{2(1 + \delta)(1 - \cos \theta) + \delta^2}} \quad (3)$$

where λ is the moiré periodicity, δ is the lattice mismatch between h-BN and graphene which is less than 2%, a is the lattice constant of graphene, and θ is the relative rotation angle



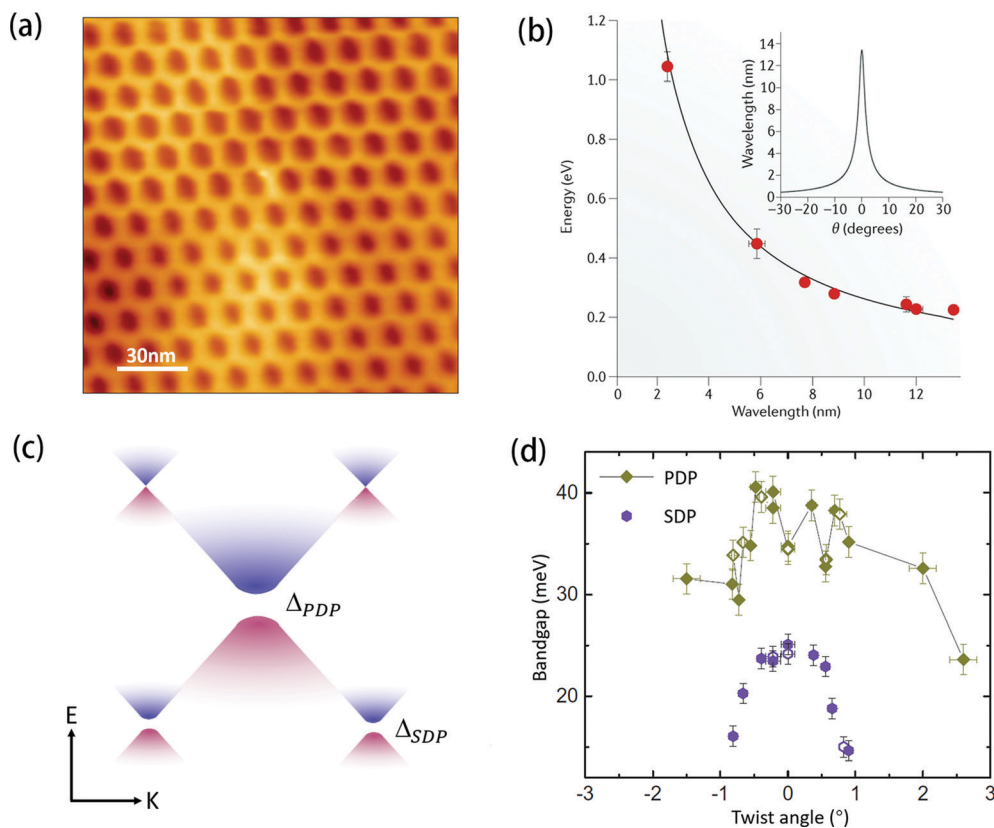


Fig. 2 (a) STM image of the moiré pattern observed in the graphene/h-BN heterostructure owing to the lattice mismatch between them. (b) The energy of the superlattice Dirac points as a function of moiré periodicity. The inset displays the wavelength of the moiré superlattice as a function of the twist angle (θ) between graphene and the h-BN substrate. (c) Schematic band structure of the graphene/h-BN heterostructure with a small twist angle. E , energy; k , momentum. (d) Measured bandgaps at the primary and superlattice Dirac points as a function of twist angle within a single graphene/h-BN device. Figure adapted with permission from: (a) ref. 25 Copyright 2014, Springer Nature Limited; (b) ref. 55 Copyright 2019, Springer Nature Limited; (c and d) ref. 59 Copyright 2018, AAAS.

between the two lattices. The inset of Fig. 2b shows that the moiré periodicity increases when the relative rotation angle decreases. STM measurements reveal that the moiré superlattice has a fluctuation of 20 pm in space through the epitaxial growth of graphene/h-BN samples. The strength of van der Waals interaction between carbon atoms and the h-BN substrate is inversely proportional to their distance. Therefore, the moiré pattern effectively acts as a periodic superlattice electrical potential for graphene.¹⁷ The resulting periodic potential enables scattering processes along the directions of its reciprocal lattice vectors which are otherwise forbidden in graphene. The so-called superlattice Dirac points (SDPs) are created in the band structure where the periodic potential connects the k and $-k$ bands. The appearance of the SDPs originates from the moiré superlattice attached to graphene and is located at the Brillouin zone boundaries of the superlattice in the reciprocal space. Indeed, researchers have systematically studied the SDP. Scanning tunneling spectroscopy (STS) revealed the DOS of these SDPs to be symmetric with respect to the primary Dirac point (PDP) of the graphene/h-BN moiré superlattice.¹⁶ The energy dispersion of these SDPs is extracted as a function of the moiré periodicity (Fig. 2b). It is clear that the energy of the SDP decreases as the moiré

periodicity increases. The observation of SDPs confirms the modulation of the band structure of graphene by moiré periodic potential.^{13,16,17} Moreover, various experimental measurements demonstrate sizeable bandgaps at the PDP and SDPs in the valence band.^{13,25,56–61} A schematic band structure of the graphene/h-BN heterostructure with a small twist angle is shown in Fig. 2c. These bandgaps are understood to result from sublattice symmetry breaking, which is due to the imbalance of carbon-on-boron and carbon-on-nitrogen stacking within the graphene. The measured bandgaps at the primary and superlattice Dirac points are extracted as a function of twist angle (Fig. 2d). It is clear that the gap at the SDP is typically smaller than that at the PDP and vanishes at $>1^\circ$ misalignment.^{18,59} Note that the bandgap can be further broadened by many-body interactions.^{62,63}

The long-period moiré pattern in the graphene/h-BN heterostructure provides a unique system to explore the interplay between electrical potentials and magnetic fields over similar length scales. The transport properties of graphene/h-BN devices indicate that the SDPs manifest themselves as additional peaks in the resistance symmetrically flanking the primary Dirac point.^{13–15} The charge carrier density n -dependent conductance measurements show that the longitudinal



resistivity ρ_{xx} reaches a maximum at both sides of PDP and SDP (Fig. 3a), whereas the Hall resistivity ρ_{xy} drops suddenly at one side of the PDP and SDP and increases at the other side (Fig. 3b), reversing in sign. To present the detailed electronic structures of PDP and SDPs, the R_{xx} and R_{xy} are plotted as a function of gate voltage and magnetic field, which is called a Landau fan. The Landau fan diagram shows that the degeneracy of Landau levels formed by the two-dimensional electron gas increases and the system enters the quantum Hall state gradually as the magnetic field increases. R_{xx} reaches a minimum when the Fermi level crosses a LL and R_{xy} exhibits a quantum Hall plateau at the corresponding position. Similarly to the Dirac point, these SDPs break up into sequences of LLs when applying a magnetic field. Transport measurements show that these LLs emerge from each of these points, including PDP and SDPs, at low magnetic fields and intersect as the magnetic field increases, resulting in a series of energy levels, which is known as a Hofstadter butterfly spectrum (Fig. 3c and d).

The Hofstadter butterfly was predicted in 1976, and it was pointed out that the 2D electron gas would form a fractal energy spectrum if it is simultaneously under periodic electric field and magnetic field.⁶⁴ By tuning the magnetic field, a fractal of

the energy band is completed when a unit cell is filled with one magnetic flux quantum h/e . However, the small period (~ 1 nm) of electrical potentials in an ordinary crystal lattice leads to one magnetic flux quantum filling needing a magnetic field of more than 1000 T, which is difficult to achieve in experiments. Intriguingly, the graphene/h-BN moiré superlattice with a moiré period of about 10 nm can reduce this magnetic field to approximately 30 T, making the Hofstadter butterfly spectrum visible in laboratory-scale magnetic fields. Moreover, in the graphene/h-BN moiré superlattice system, these LLs intersect when^{13–15}

$$\frac{\phi_0}{\phi} = q \quad (4)$$

where ϕ_0 is the magnetic flux quantum through the superlattice unit cell, q is the magnetic flux and q is a positive integer. Fig. 3e displays a Wannier diagram showing the quantum Hall effect states as a function of charge carrier density, highlighting the Hofstadter states experimentally observed in graphene/h-BN. Particularly, black, blue and red lines indicate the fractal integer quantum Hall, conventional fractional quantum Hall and anomalous quantum Hall states, respectively. Although

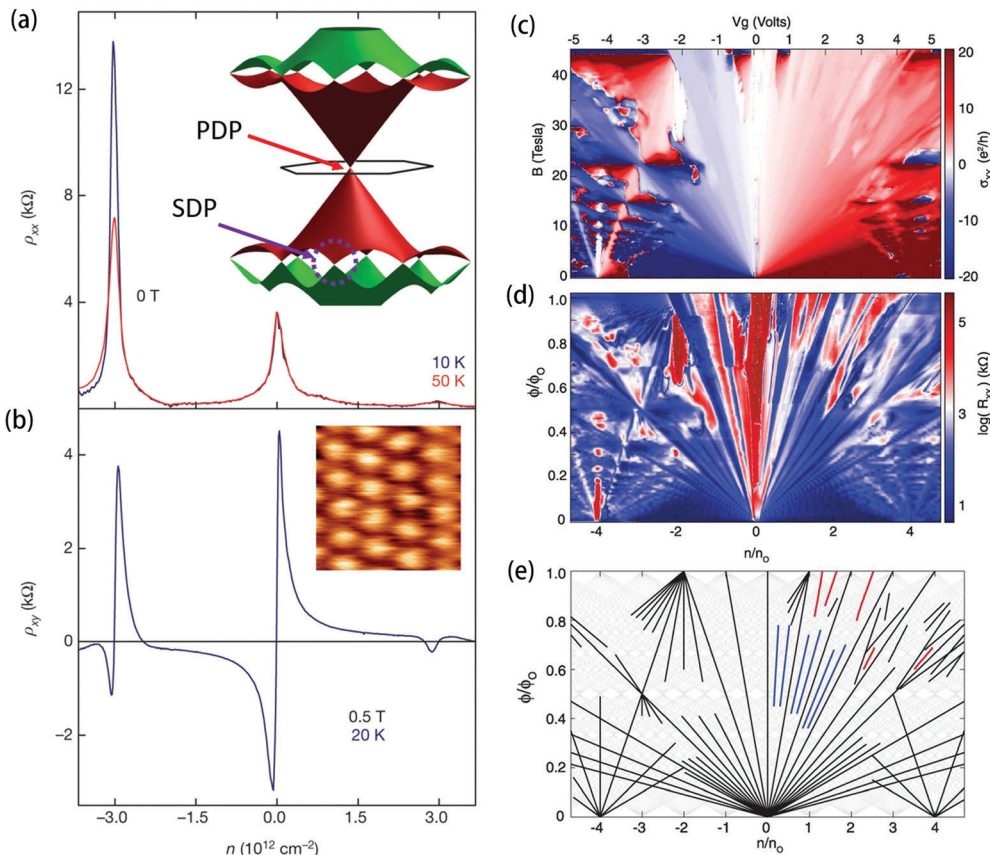


Fig. 3 Emergence of anomalous quantum Hall states. Longitudinal resistivity ρ_{xx} (a) and Hall resistivity ρ_{xy} (b) plotted as a function of carrier density n . Positive values of n indicate electrons, and negative values indicate holes. Hall conductivity σ_{xy} plotted versus magnetic field B and gate bias V_g (c) and longitudinal resistance R_{xx} as a function of the magnetic field flux ϕ/ϕ_0 and charge density n/n_0 (d) in a graphene/h-BN device. (e) A simplified Wannier diagram in which the positions of the most prominent QHE states are plotted as a function of magnetic flux and charge carrier density. Figure adapted with permission from: (a and b) ref. 15 Copyright 2013, Springer Nature Limited; (c–e) ref. 18 Copyright 2015, AAAS.



there are many studies about the Hofstadter spectrum, further studies are needed to understand the rich physics in the Hofstadter butterfly.

2.2 Strongly correlated quantum states in twisted bilayer graphene

Among existing 2D materials, graphene is the most stable material with the most mature and successful preparation technology as well as the most extensive and in-depth research into its physical properties. In previous reports, the properties of graphene can be understood very well based on a single electron approximation, without considering the electron–electron interaction. However, there is a wealth of strongly correlated quantum states in magic-angle twisted bilayer graphene.^{30,31} Jarillo-Herrero's group observed two correlated quantum states of correlated insulating states and unconventional superconductivity caused by a strong electron–electron interaction.

Note that there are three methods that can experimentally induce an electron–electron interaction in graphene. The most commonly used one is to generate Landau levels in graphene by applying a perpendicular external magnetic field. The fractional quantum Hall effect is a kind of strongly correlated quantum state of matter induced by a strong magnetic field. The second method is to modulate the Fermi level of monolayer graphene located on the Van Hove peak, which may make graphene exhibit anomalous magnetic or superconducting properties.^{65–70} However, the gap between Van Hove peaks of monolayer graphene and the neutral point, which is the Dirac point, is about 3 eV, and the current experimental technology is unable to realize such a modulation. Therefore, inducing magnetism and superconductivity into monolayer graphene through an electron–electron interaction has always been studied theoretically. The third method proposes that a strongly correlated quantum state can be induced in twisted bilayer graphene. The intuitive physical scenario is as follows: there are twist-angle-dependent low-energy Van Hove singularities (VHSs) in the TBG and the energy gap between VHSs and Dirac points of the system is generally less than a few hundred millielectron volts, which make it possible to tune the Fermi level to VHS accurately in the experiment.^{71–73} The twisted bilayer graphene system provides an opportunity to experimentally realize the strongly correlated electronic state.

2.2.1 Van Hove singularities in TBG. Prior to the experimental discovery of graphene/h-BN superlattices, the related experimental research on twisted bilayer graphene has attracted intense interest since 2009. The most important difference between the two kinds of moiré superlattice is that the band gap of h-BN is about 5.9 eV, which makes the low energy band of graphene unable to hybridize with h-BN, while in the twisted bilayer graphene moiré superlattice the upper and lower graphene energy bands are coupled with each other. In graphene/h-BN structures, the presence of a moiré superlattice can lead to the observation of electronic minibands, whereas in twisted graphene bilayers its effects are enhanced by interlayer resonant conditions. The interlayer coupling strength of electrons is periodically distributed in the TBG moiré

superlattice system, which may strongly modify the electronic properties of individual graphene.

The TBG moiré superlattice is a homostructure, and thus the lattice mismatch δ is zero. Therefore, formula (3) can be simplified as

$$\lambda = \frac{a}{\sqrt{2(1 - \cos \theta)}} \quad (5)$$

In the energy band structure of TBG, Dirac points corresponding to the upper and lower layers are separated with the rotation angle, forming K and K_θ points (Fig. 4a). Along the direction of separation of the Dirac points, a saddle point is formed in the middle of the two Dirac points, which is called a Van Hove singularity,⁷⁴ corresponding to the two maxima of the density of states. The twisted bilayer graphene with different rotation angles was studied by using STM and scanning tunneling spectroscopy.^{29,75,76} Because the STS spectra reflect the local electronic density of states of the sample, there are two characteristic peaks corresponding to Van Hove peaks, to the two saddle points in the band structure of TBG (Fig. 4b). The left panel shows the spectra of two VHSs and the right panel shows two saddle points formed at $k = 0$ between the two Dirac cones, K and K_θ , with a separation of $\Delta K = 2K \sin(\theta/2)$. Many intriguing phenomena caused by the VHSs in the TBG, including enhanced optical absorption, Raman G-band resonance and enhanced chemical reactivity, were studied in the last few decades.^{77–85} Moreover, for the TBG moiré superlattice, it has been demonstrated that the energy interval of Van Hove peaks, ΔE_{VHS} , is twist angle-dependent, and the relationship approximately satisfies^{75–77,86,87}

$$\Delta E_{\text{VHS}} = \hbar v_{\text{F}} \Delta K - 2t_\theta \quad (6)$$

where v_{F} is the renormalized Fermi velocity and t_θ is the interlayer coupling coefficient. Assuming that the coupling coefficient between layers is constant, $\sin \theta \approx \theta$ for small rotation angles ($\theta < 6^\circ$), the gap between the Van Hove peaks is proportional to the twist angle. E_{VHS} increases as the twist angle increases, which is consistent with the experimental results (Fig. 4c). Moreover, the photocurrent can be significantly enhanced when the incident photon energy matches the energy interval between VHSs of the conduction and valence bands in the TBG photodetection device, which suggests the existence of strong light–matter interactions.⁸⁸ Therefore, the emergence of twist angle-dependent VHSs in TBG can selectively enhance photocurrent generation, providing valuable insights for designing graphene-based photodetectors with variable wavelengths.

Additionally, the experiments also proved that the low-energy region of the TBG has an angle-dependent Fermi velocity.^{89,90} When the twist angle is not very small ($\theta > 3^\circ$), the Fermi velocity can be renormalized, and the relationship is as follows:^{71–73,91}

$$\frac{\tilde{v}_{\text{F}}}{v_{\text{F}}} = 1 - 9 \left(\frac{t_\theta}{\hbar v_{\text{F}} \Delta K} \right)^2 \quad (7)$$

Obviously, the Fermi velocity in the low-energy region of the TBG continues to decrease as the twist angle becomes smaller



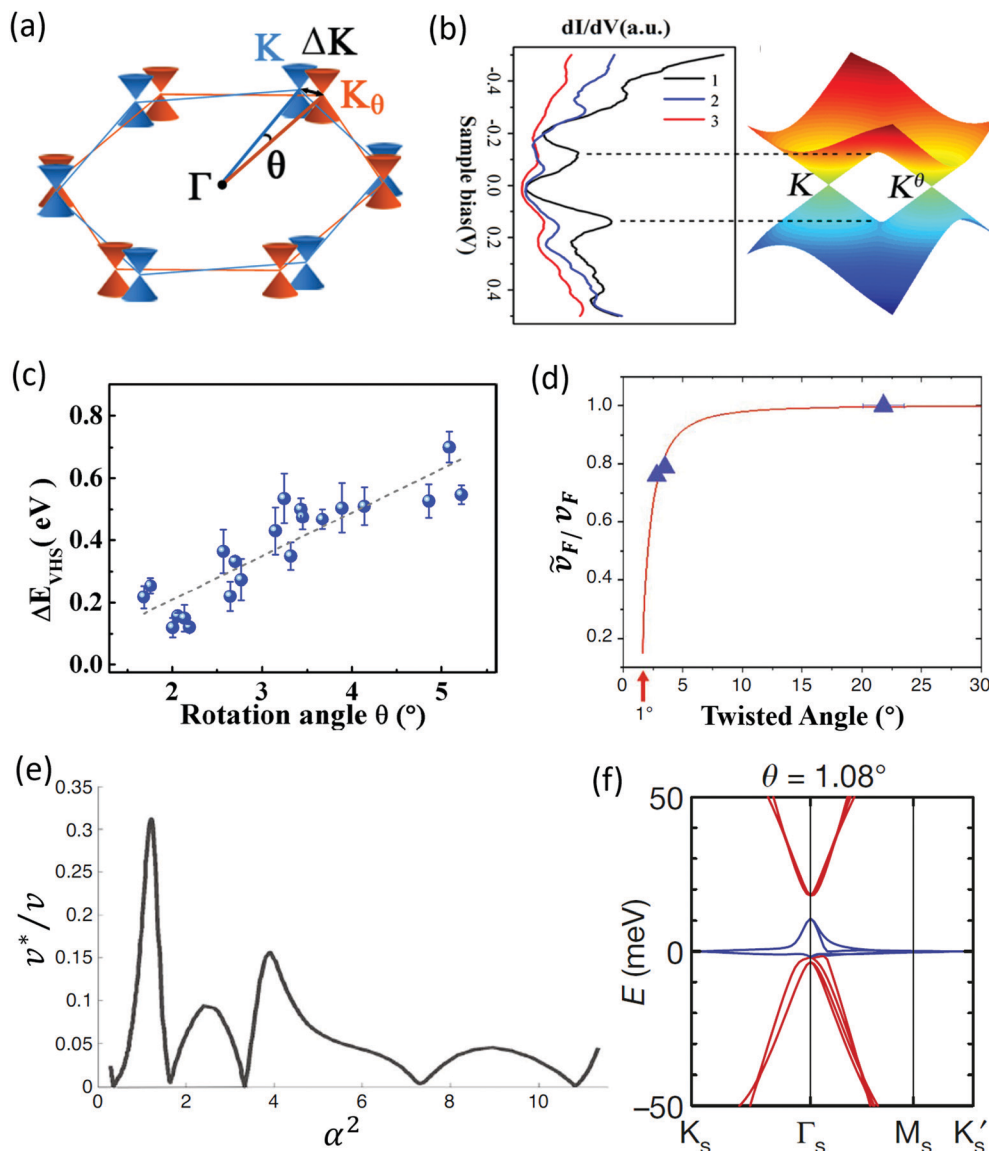


Fig. 4 (a) Schematic of the Dirac cones, K and K_{θ} , of two misoriented graphene layers in the reciprocal space. (b) Left: Tunneling spectra recorded at different positions of a moiré pattern. Right: Electronic band structure of the TBG with a finite interlayer coupling. (c) Angle-dependent VHSs in twisted graphene bilayers. Solid blue circles: average experimental data measured in different twisted graphene bilayers. The gray dotted line: linear fit of the experimental data. (d) Renormalized Fermi velocity as a function of twist angle obtained from Landau-level spectroscopy (symbols) compared to a theoretical simulation (solid line). (e) Renormalized Fermi velocity, $v^* = v_F(1 - 3\alpha^2)/(1 + 6\alpha^2)$, as a function of twist angle in the range of $0.18^\circ < \theta < 1.2^\circ$ shows a sequence of magic twist angles at which the Fermi velocity vanishes. (f) The band structure of MATBG shows the formation of flat bands (blue lines) when approaching the magic angle. Figure adapted with permission from: (a and c) ref. 86 Copyright 2014, American Physical Society; (b) ref. 75 Copyright 2012, American Physical Society; (d) ref. 89 Copyright 2011, American Physical Society; (e) ref. 73 Copyright 2011, National Academy of Sciences; (f) ref. 30 Copyright 2018, Springer Nature Limited.

and drops to zero at the first magic angle (Fig. 4d and e), leading to the appearance of two flat bands in the low-energy band structure of the system (Fig. 4f).⁷³ Fig. 4e shows the twist angle expressed in terms of the dimensionless parameter $\alpha = w/(\hbar v_F \Delta K)$, where w is the interlayer tunnelling strength and $\Delta K \approx 2K \sin(\theta/2)$ as defined in the above text. The calculation results show that the bandwidth of flat bands in magic-angle TBG is approximately 12 meV ($E > 0$) or 2 meV ($E < 0$). The flat band of the MATBG was measured experimentally for the first time in 2015, and the corresponding magic angle was about 1.11° .⁹²

2.2.2 Strongly correlated quantum states in MATBG. Due to the lack of mature and accurate twisted graphene preparation technology, it was until 2018 when Cao *et al.* successfully measured two correlated quantum states of insulating and unconventional superconductivity caused by a strong electron–electron interaction in the magic-angle twisted bilayer moiré superlattice of graphene at extremely low-temperatures.^{30,31} In the Hubbard model, the correlation effect of electrons is determined by U/W , in which U and W represent the mutual Coulomb repulsion between electrons and the kinetic energy of electrons,



respectively. The correlation effect of electrons is significant when the mutual Coulomb repulsion between electrons dominates over their kinetic energy, $U/W \gg 1$.^{93–97} In the TBG superlattice, the Coulomb interaction U satisfies

$$U = \frac{e^2}{4\pi\epsilon_0\epsilon_r d} \quad (8)$$

where d is the distance between electrons and ϵ_0 and ϵ_r are the vacuum permittivity and relative permittivity, respectively. Let us

take $d = \lambda/2 = 6.5$ nm ($\theta = 1.1^\circ$, $\lambda = 13$ nm) and ϵ_r is set to 4, then $U \approx 55.4$ meV. Here, 1.1° is the so-called first magic angle. As mentioned above, the Fermi velocity could drop to zero at the magic angle in TBG, leading to the emergence of flat bands, which presents weak dispersion in the momentum space, with a bandwidth of about 10 meV. The upper limit of electron kinetic energy W is the flat band bandwidth when the electron concentration is only tuned to fill the flat band near zero energy (Fig. 5a). Therefore, strongly correlated electronic phases are

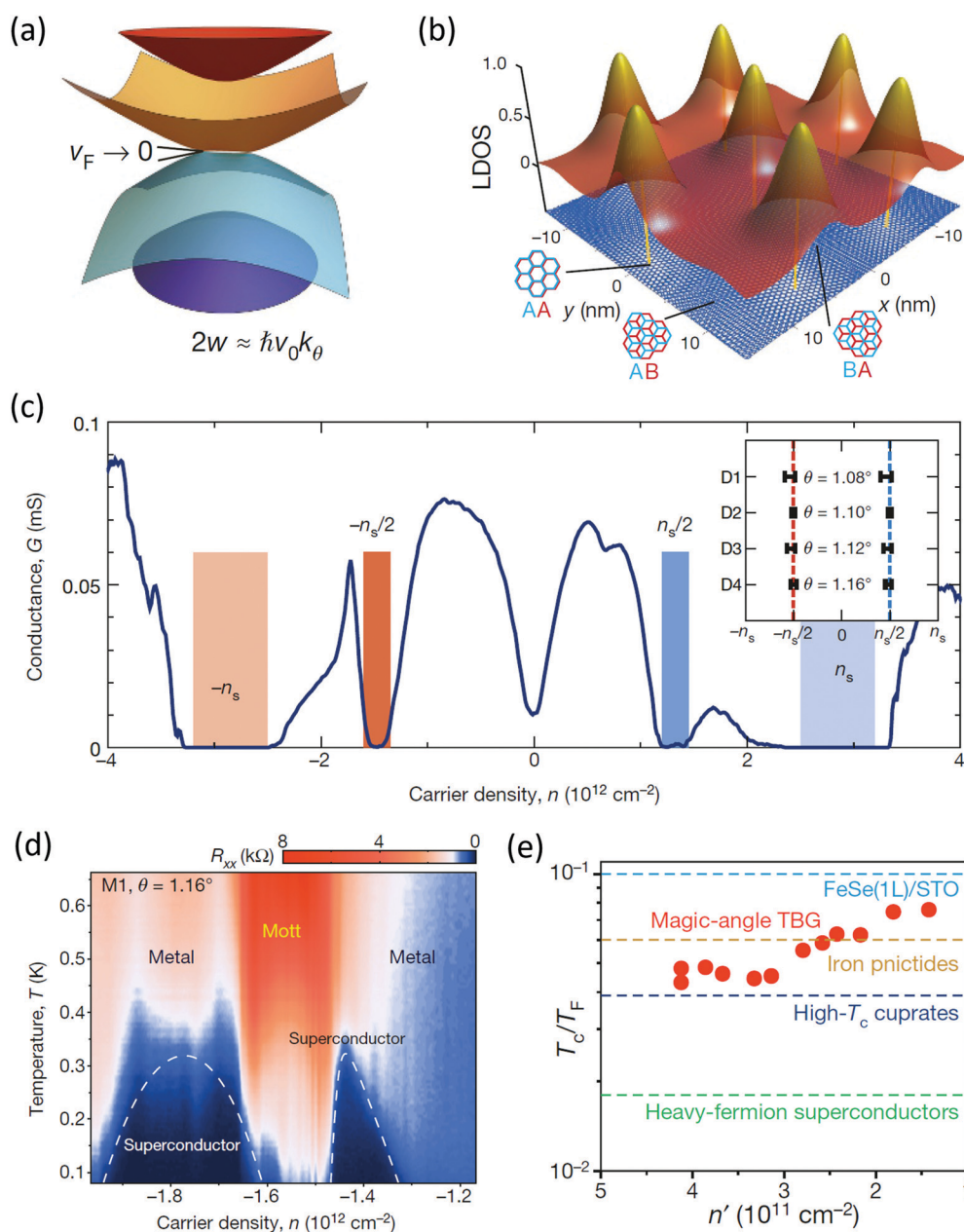


Fig. 5 (a) Illustration of the effect of interlayer hybridization for $2W \approx \hbar v_0 k$. (b) Spatial distribution diagram of local density of states (LDOS) of electrons calculated for flat bands with $E > 0$ in the MATBG. The electron density is strongly concentrated in regions with AA stacking order, but it is mostly depleted in AB- and BA-stacked regions. (c) Carrier dependence of conductance in MATBG. (d) Phase diagram showing the carrier density-dependent superconducting critical temperature in the partially filled flat band. There are two dome-shaped superconducting regions (blue) flanking an insulator (red). (e) Density dependence of T_c/T_F as a function of doping for MATBG (red filled circles) compared to other strongly correlated systems. T_c : superconducting temperature, T_F : Fermi energy. Figure adapted with permission from: (a–c), ref. 30, Springer Nature Ltd; (d and e), ref. 31, Springer Nature Ltd.



expected to be generated in MATBG because the condition $U/W \gg 1$ is satisfied. The flat-band DOS peak in MATBG revealed by STM/STS measurements illustrates that the electron is highly localized in space with the inversion symmetry breaking of the electronic state distribution when the Fermi level lies within the flat bands in the low-energy region, which was attributed to the strong electron–electron interaction in the system (Fig. 5b).

The strong correlation effect of electrons is often accompanied by the Mott-like insulator state transition.⁹⁸ The sample is expected to behave like a metal according to single-particle band energy theory when the band is not fully filled with electrons. However, it actually behaves like an insulator owing to the strong correlation effect of electrons when the energy band is half-filled, which is usually observed in cuprate materials.^{99–102} Similarly, low-temperature (~ 1 K) electrical transport measurements on MATBG revealed interaction-induced correlated insulating states when the flat band is half-filled with electrons, and thus the state near $\pm n_s/2$ is called a half-filling insulating state (Fig. 5c). The MATBG system exhibits two pronounced superconducting domes on each side of the half-filling correlated insulating state when using gate voltage to continuously tune the carrier concentration of the system, with the currently obtained highest superconducting transition temperature of 1.7 K (Fig. 5d), whereas the MATBG system exhibits a quantum anomalous Hall effect (QAHE) at 3/4 filling of the conduction miniband. The QAHE combines topology and magnetism to produce precisely quantized Hall resistance at zero magnetic field. The effect is driven by the spontaneous valley polarization in valley resolved moiré bands with an opposite Chern number.^{103–106} In recent papers, the researchers observed emergent ferromagnetic hysteresis, reflecting a topological state with Chern number $C = 1$ that is an orbital ferromagnet. Moreover, the magnetic order of chiral edge states of opposite polarization can be reversed by applying a small direct current, which may have potential applications such as an electrically rewritable magnetic memory.^{107,108}

In MATBG, the ratio of the superconducting transition temperature to Fermi temperature, $T_c/T_F \approx 0.1$, exceeds the range in which the weak coupling theory of superconductivity can be used, and is similar to T_c/T_F in other materials (Fig. 5e) that exhibit superconductivity close to metal–insulator phase transitions. Very recently, nematicity has been observed across the entire superconducting dome of TBG, which indicates that nematic fluctuations might play an important role in the low-temperature phases of MATBG.¹⁰⁹ The phenomenology of MATBG is very similar to the phase-diagram of cuprate and pnictide^{110,111} superconductors in which superconducting domes flank Mott-like insulating phases. Moreover, it is much easier to characterize and control the MATBG than a high-temperature cuprate superconducting system. Therefore, in-depth research on MATBG is expected to help solve the high-temperature superconducting mechanism that has been troubling physicists for more than thirty years.

However, there are still a lot of efforts required to confirm the resemblance between the phenomenology of MATBG and high-temperature superconductors. For example, the parent material of the high-temperature cuprate superconducting

system is an anti-ferromagnetic Mott insulator. Current experiments have confirmed that the MATBG is an insulator when the flat band is half-filled, but there is no definite evidence of whether it is a Mott insulator. Moreover, it is needed to be further studied whether the MATBG has anti-ferromagnetism if Mott-like insulator states are confirmed in it. Besides, we still do not know the superconducting pairing mechanism in the MATBG. Apart from this, it is unclear whether there is a pseudogap. On the other hand, there are many problems in the MATBG superconductivity needing to be settled. For example, it is still not clear how the superconductivity in MATBG changes along with angles near the magic angle under the same filling conditions. It is also indefinite whether there is a smaller magic-angle at which similar superconductivity phenomena can occur. In each moiré pattern of the TBG, the stacking order undergoes a change from AB stacking to BA stacking and there are topologically protected boundary states on the AB–BA stacking domain boundary,^{112,113} which can intrigue an interesting question of how do the superconducting and topological boundary state networks in the MATBG interact with and affect each other. Nevertheless, we hold optimistic opinions about more “magic” properties in the MATBG in the foreseeable future.

3. Moiré excitons in TMD heterobilayers

Transition metal dichalcogenides are another class of materials that attract people’s attention apart from graphene because they are natural semiconductor materials with considerable band gaps. Similar to graphene, strong intralayer covalent bonding and weak interlayer vdW interactions make it possible to exfoliate the thin 2D material from its bulk counterpart. Moreover, TMDs display unique semiconducting properties which render them optically active. The anomalous phenomena in graphene-based moiré superlattices attract the interest of researchers to investigate twisted bilayer TMD superlattices. A TMD moiré superlattice is produced when two layers with lattice mismatch are grown on top of one another, naturally endowing them with a nanoscale moiré pattern, which can significantly modify the electronic and optical properties of TMD heterobilayers. In this section, we mainly present the contemporary experimental and theoretical understanding of the unique effects of a moiré superlattice on the interlayer exciton potential landscape and optical properties.

3.1 Theory of interlayer exciton moiré potential energy

Various TMD heterobilayers have a type II band alignment where electrons (holes) have lower energies in different layers.^{37,114,115} The excitons composed of electrons and holes located at different layers through strong Coulombic interactions are referred to as interlayer excitons, shown as violet circles in Fig. 6a,^{116–118} whereas intralayer excitons are shown as blue and red circles, which were formed by electrons and holes coming from individual layers. Fig. 6b shows the excitation, charge



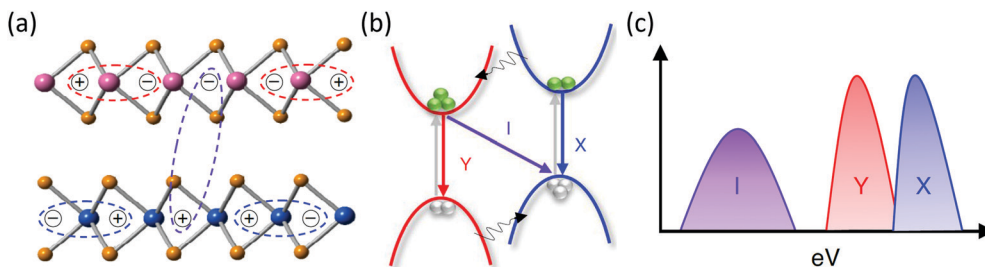


Fig. 6 Type II alignment of 2D heterostructures with interlayer interaction strength. (a) Populations of intralayer and interlayer excitons. (b) Excitation, charge separation and emission processes in simplified band pictures. (c) Photoluminescence spectra expected from the heterostructure. Figure adapted with permission from: (a–c), ref. 116 Copyright 2020, Springer Nature Limited.

separation and emission processes of interlayer and intralayer excitons. The photoluminescence spectra expected from the heterostructure are shown in Fig. 6c. The blue (X) and red (Y) photoluminescence peaks correspond to pristine intralayer excitons, whereas the violet (I) peaks are from interlayer excitons.

The moiré pattern in TMD superlattices induced by the lattice mismatch or crystal orientation misalignment can cause modulations in the electronic band structure and optical properties, such as a spatially dependent band gap¹¹⁹ or twist-controlled interlayer excitons.^{41–43,120} Consider a moiré superlattice of the common chalcogen TMD bilayer WX_2/MoX_2 ($X = Se, S$) with a small twist angle θ and an in-plane displacement \mathbf{d} , which has two distinct stacking orders of AA stacking (Fig. 7a) and AB stacking (Fig. 7d). This heterojunction has a small lattice mismatch ($\sim 0.1\%$) and the metal heteroatoms of adjacent layers overlap each other only for periodic moiré periodicities, and thus the wavelength of the moiré superlattice a_M can be simplified as $a_M = a_0/\theta$, in which a_0 is the lattice constant of a monolayer TMD. a_M (~ 18.3 nm) is much larger than the exciton Bohr radius (~ 1.3 nm) when the twist angle is smaller than 1° . Therefore, the exciton energy can be approximated by following the local band gap.

In this system, the moiré potential modulated interlayer exciton energy is $E_g - E_b$, where E_g is the band gap between the conduction band minimum (CBM) of MoX_2 and the valence band maximum (VBM) of WX_2 at the \mathbf{K} point, and E_b is the exciton binding energy. The variation of E_b in the moiré pattern is usually neglected, since it is typically smaller than the variation of E_g . According to the *ab initio* continuum model, the electronic properties calculated from the untwisted reference stackings can be leveraged to approximate the twisted bilayer with a smoothly varying field over the moiré pattern, which can accurately capture atomic relaxations and electronic phenomena.^{38,39,63,121–124} The procedures for the calculation of excitons' moiré potential energy in TMD heterobilayers by using the *ab initio* continuum model are as follows.³⁸

Firstly, calculate the density functional theory (DFT) band structure of untwisted reference cells. Obtain E_g as a periodic function of relative displacement \mathbf{d} at the zero-twist angle from the DFT band structure (Fig. 7b and e). The bandgap E_g is defined as the energy difference between the conduction band minimum (CBM) of MoX_2 and the valence band maximum (VBM) of WX_2 at the \mathbf{K} point.

Secondly, obtain the band structure at different local stackings. Since the potential is a smooth periodic function of \mathbf{d} in both AA and AB stackings, it can be approximated as $\Delta(\mathbf{d}) \equiv E_g(\mathbf{d}) - \langle E_g \rangle \approx \sum_{j=1}^6 V_j \exp(i\mathbf{G}_j \cdot \mathbf{d})$, in which $\langle E_g \rangle$ is the average of E_g over \mathbf{d} , \mathbf{G}_j is one of the first-shell reciprocal lattice vectors (Fig. 7i), and V_j is the moiré potential parameter, denoting the coupling strength between MoX_2 and WX_2 in TMD heterobilayers. Moreover, the existence of threefold rotational symmetry (\hat{C}_3) in each TMD layer leads to $V_1 = V_3 = V_5$, $V_2 = V_4 = V_6$ (Fig. 7h). Therefore, $\Delta(\mathbf{d})$ is always real.

Finally, approximate the interlayer Bloch-wave coupling over a single moiré supercell. Assume that, in a moiré supercell, the variation of E_g locally follows $\Delta(\mathbf{r}) \approx \Delta[\mathbf{d}(\mathbf{r})] \approx \sum_{j=1}^6 V_j \exp[i\mathbf{G}_j \cdot \mathbf{d}(\mathbf{r})] \approx \sum_{j=1}^6 V_j \exp[i\mathbf{b}_j \cdot \mathbf{r}]$ (Fig. 7c and f), where the local configuration, \mathbf{d} , is a function of space, $\mathbf{d}(\mathbf{r})$, and $\mathbf{b}_j = \theta\mathbf{G}_j \times \hat{z}$ is the reciprocal lattice vector of the moiré pattern (Fig. 7g).

Therefore, the moiré pattern formed in the twisted TMD heterobilayer can lead to a position-dependent potential for excitons across a single moiré supercell, with spatial modulation up to ~ 150 meV for interlayer excitons in the MoS_2/WSe_2 heterobilayer with rotational alignment and lattice mismatch, which has been experimentally identified by STM/STS (Fig. 8a and b).¹¹⁹ Additionally, TMDs are highly sensitive to crystal structure symmetry and interlayer coupling, which makes them highly susceptible to stacking orders. Atoms stacked in different configurations are known as atomic registries, which can cause valence and conduction bands to possess different energies, giving additional control over the electronic and optical properties in these systems.^{125–129}

For interlayer excitons in a moiré pattern, the theory predicts that the moiré potential generally results in localized exciton states near the potential minimum positions. Fig. 8c shows that the depth of moiré potential for 0° is much larger than that for 60° . It implies that the moiré pattern leads to twist-angle-dependent periodic potential landscapes that govern the transport and relaxation of interlayer excitons, which has been experimentally confirmed.¹³⁰ The existence of moiré potential in twisted TMD heterobilayers suggests a vdW



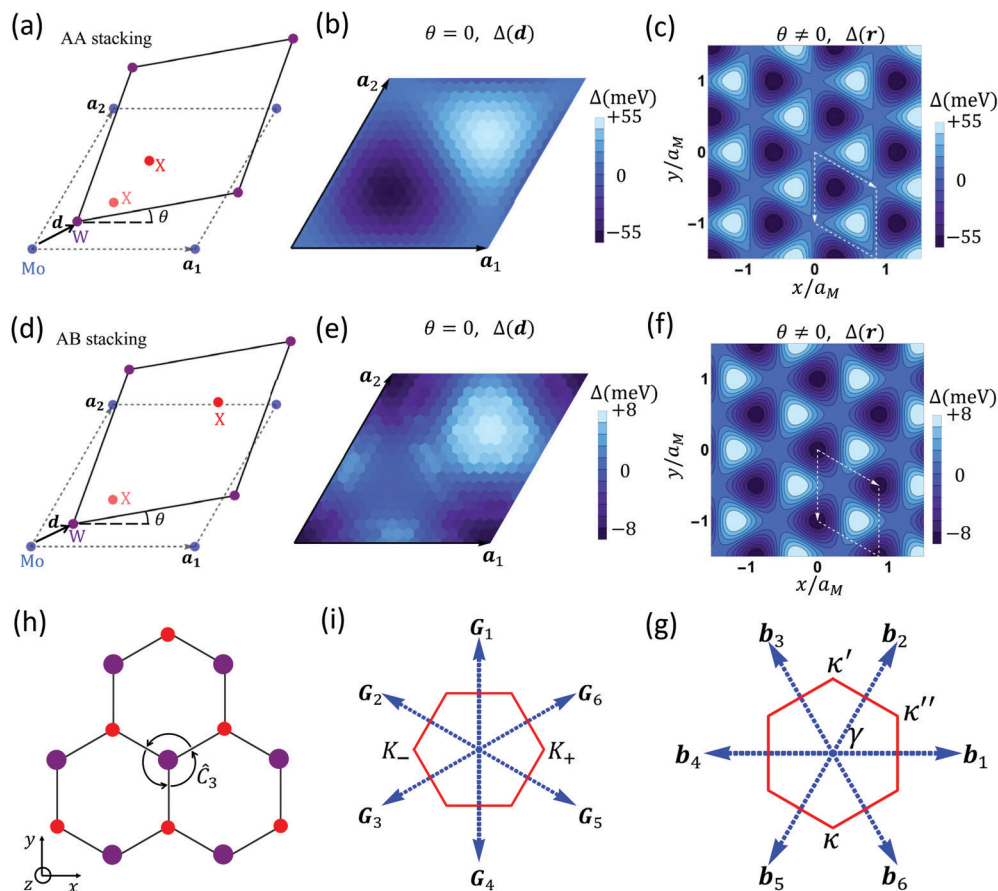


Fig. 7 (a) Illustration of an AA stacked WS_2/MoS_2 heterobilayer with a small twist angle θ and an in-plane displacement \mathbf{d} . The variation as a function of displacement \mathbf{d} (b) and position \mathbf{r} (c) of the interlayer band gap (Δ) within a WS_2/MoS_2 heterobilayer for AA stacking. (d–f) The same plots as those in (a–c) but for the WS_2/MoS_2 heterobilayer in AB stacking. (h) \hat{C}_3 transformation of the TMD lattice. Black arrows denote the phase change by \hat{C}_3 . (i) First-shell reciprocal lattice vectors \mathbf{G}_j of a monolayer TMD triangular lattice and the corresponding Brillouin zone (red hexagon). (g) Moiré reciprocal lattice vectors \mathbf{b}_j and the corresponding Brillouin zone. Figure adapted with permission from: (a, d, i and g) ref. 39, Copyright 2017, American Physical Society; (b, c, e and f) ref. 38, Copyright 2018, American Physical Society.

heterostructure based strategy for realizing two-dimensional nanoscale arrays of uniform quantum dots, which has a great prospect for application in quantum photonic devices. Yu *et al.*⁴⁰ studied interlayer excitons in the MoX_2/WX_2 heterobilayer and found that they have a unique spin-dependent hopping, which can be switched off when applying a perpendicular electric field according to DFT calculations, giving the possibility to turn the superlattices into identical arrays of nanodots that act as uniform quantum emitters. Besides, the spin optical selection rules can be flipped by switching the positions of emitters in the moiré due to the electric field. These properties of moiré excitons in TMD heterobilayers suggest the feasibility of engineering artificial excitonic crystals for nanophotonics and quantum information applications.¹³¹

3.2 Twist-dependent moiré excitons in TMD heterobilayers

Although the theory of exciton potential energy in twisted TMD heterobilayers has predicted notable effects of the moiré potential on interlayer excitons years before,^{38,40,120} the impacts of the moiré potential on light emission and absorption in

heterobilayers of TMDs have not been detected experimentally until 2019. It was found that interlayer excitons have a valley degree of freedom, and different interlayer excitons at a given valley can absorb light with different circular polarization; besides, the excitons interact with light of a specific polarization that depends on the local symmetry of the atoms arranged periodically in the heterobilayer.^{41–44}

Recently, the WSe_2/WS_2 moiré superlattice system with different twist angles was studied experimentally. Jin *et al.*⁴² measured the PL of the WSe_2/WS_2 heterobilayer and found that there is no monolayer PL in the heterostructure, which indicates the existence of strong interlayer coupling. Three peaks emerge around the original A exciton resonance of WSe_2 in the reflection spectrum in the small twisted WSe_2/WS_2 moiré superlattice, corresponding to distinct moiré exciton states, whereas, in the superlattice with a large twist angle, it shows only a single resonance in energy near the A exciton of WSe_2 (Fig. 9a). These phenomena can be explained with a theoretical model in which the moiré potential is much stronger than the kinetic energy of exciton, generating multiple flat exciton minibands.



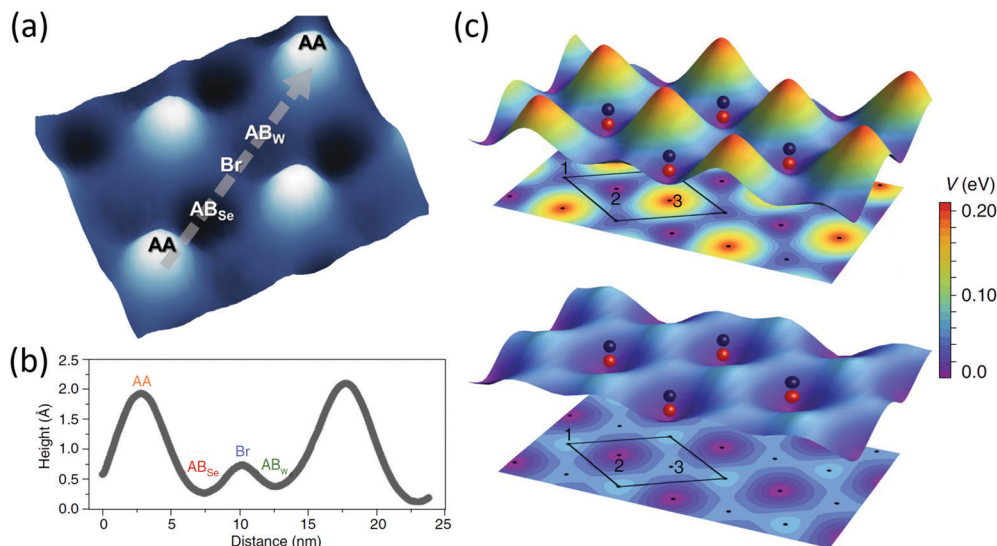


Fig. 8 (a) STM measurement of the topographic height across a single moiré supercell of the MoS₂/WSe₂ moiré pattern. (b) A height profile along the diagonal line from AA to AA in a moiré supercell (arrow in a). (c) Illustration of the 2D K valley moiré potentials in both 3D graphs and 2D projections that can trap long-lived interlayer excitons (red and black spheres) in the local minima as they move from one location to another. 1, 2 and 3 indicate high symmetry points in a TMD heterobilayer of WS₂/WSe₂ with a twist angle of 0° (top) and 60° (bottom). Figure adapted with permission from: (a and b) ref. 119 Copyright 2017, AAAS; (c) ref. 130 Copyright 2020, Springer Nature Limited.

Seyler *et al.* studied individual interlayer excitons trapped in the moiré potential in the MoSe₂/WSe₂ moiré superlattice with a small twist angle.⁴³ They found that there are several narrow peaks near the free interlayer exciton energy at extremely low

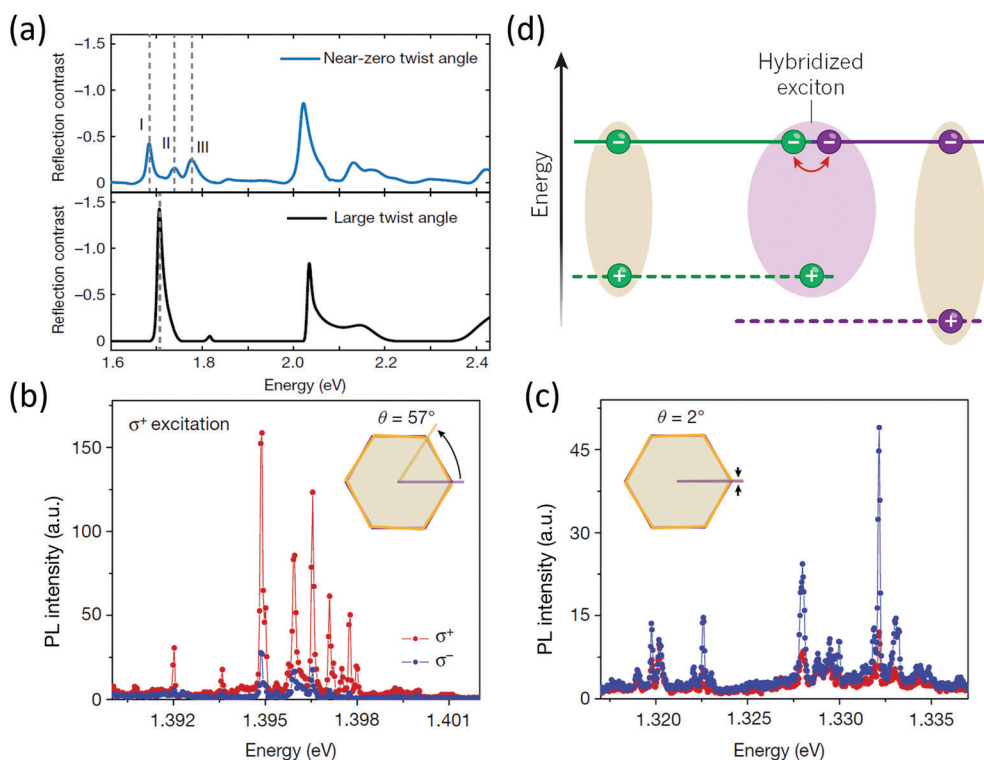


Fig. 9 (a) Reflection contrast spectrum of a small twisted WSe₂/WS₂ heterostructure (light blue, top) compared to that of a WSe₂/WS₂ heterostructure with a large twist angle (black, bottom). Helicity-resolved PL spectra of trapped interlayer excitons of MoSe₂/WSe₂ heterobilayers with twist angles of 57° (b) and 2° (c). The insets illustrate the twist angles of the two samples. (d) The band diagram showing intralayer excitons can be hybridized with interlayer excitons in heterobilayers of MoSe₂/WS₂ (indicated in purple). Figure adapted with permission from: (a) ref. 42 Copyright 2019, Springer Nature Limited; (b and c) ref. 43 Copyright 2019, Springer Nature Limited; (d) ref. 132 Copyright 2019, Springer Nature Limited.



temperatures when using a low excitation power. The linewidth of peaks is about 110 μeV , making it comparable with quantum emitters. However, the excitons can delocalize with broadened linewidth emissions under high power excitation. Besides, the trapped interlayer excitons exhibit strong circular polarization of the same helicity for a given twist angle (Fig. 9b and c). The samples are excited with σ^+ -polarized light at 1.72 eV. The σ^+ and σ^- components of the photoluminescence are shown in red and blue, respectively. The PL from the heterobilayer with a twist angle of 57° is co-circularly polarized, but that from the heterobilayer with a twist angle of 2° is cross-circularly polarized. This phenomenon suggests that the moiré-defined trapping sites preserve \hat{C}_3 symmetry. Therefore, the observation of narrow interlayer PL lines verified the existence of moiré potential with inherited valley-contrasting physics, trapping the interlayer excitons.

The experimental result of Tran *et al.* is the observation of multiple interlayer exciton resonances with circularly polarized emission in the same moiré superlattice system, a $\text{MoSe}_2/\text{WSe}_2$ heterobilayer with a small twist angle.⁴¹ They can locate the emission peaks of interlayer excitons and determine the extent of energy separation between peaks dependent on twist angle in the moiré superlattice. This result provides an exciting opportunity for high-performance semiconductor lasers, quantum emitter arrays and a platform for studying exotic quantum phases.

Additionally, the research also reveals that intralayer excitons can hybridize with interlayer excitons in $\text{MoSe}_2/\text{WS}_2$ owing to the conduction band delocalized over both layers, which leads to a resonant enhancement of moiré superlattice effects.⁴⁴ The essence of hybridization between intralayer and interlayer excitons is that the conduction band edges of MoSe_2 and WS_2 are near-degeneracy (Fig. 9d), which can amplify the impact of the moiré potential on the optical features of these heterobilayers. The energy of hybridized excitons relies on the interlayer twist angle because hybridized excitons are formed by holes that reside in MoSe_2 binding to a twist-dependent superposition of electron states in the adjacent monolayers, which gives an additional degree of freedom to modulate the optical properties of TMD heterojunctions.

The above four studies all studied the impact of the moiré superlattice on the optical properties of moiré excitons in TMD heterobilayers with different twist angles experimentally and found many novel phenomena. The origin of the observed effects is interlayer excitons trapped in a smooth moiré potential with inherited valley-contrasting physics. However, these studies are only the first footstep toward complete understanding of and control over individual moiré potentials. In-depth knowledge of the scale and depth of potentials is expected to provide control over single exciton traps, which provides opportunities for obtaining quantum emitter arrays of versatile tunability in twisted TMD heterostructures.

4. Conclusion and perspective

In this paper, we have mainly reviewed quantum phenomena in two-dimensional heterostructures. The electronic and optical

properties of graphene-based and TMD-based heterojunctions have been summarized, and the physical mechanism of quantum phenomena and the moiré potential in moiré superlattices have been elaborated. The moiré superlattice modulated properties in 2D materials have been studied at a remarkable level for a wide variety of materials, among graphene/h-BN, (magic-angle) twisted graphene bilayers and TMD heterobilayers. In graphene/h-BN heterostructures, the moiré superlattice can strongly modify the Dirac spectrum and assist optical spectroscopy in the fractional quantum Hall regime, which results in superlattice minibands and leads to Hofstadter-butterfly spectra under a strong magnetic field, whereas in twisted graphene bilayers its effects are enhanced by interlayer resonant conditions, resulting in a superconductor-insulator transition at magic twist angles. Moreover, in TMD heterobilayers, the appearance of the moiré superlattice is accompanied by the twist angle-dependent periodic potential landscape that can trap interlayer excitons near the potential minima, resulting in anomalous optical phenomena. Although lots of intriguing physical phenomena have been observed experimentally, exploiting the distinct advantages of the angle-dependent moiré superlattice in 2D materials and their superior devices is still full of challenges. In perspective, we propose several research directions for 2D material-based vdW-coupled twistrionic materials of great importance, which are worth being discussed thoroughly in the future.

(1) Precise control of twist angle

vdW moiré heterostructures can be fabricated by using a variety of methods, including CVD epitaxial growth, monolayer folding and tear-and-stack methods. Each method has its advantages and disadvantages, but the issue of common concern is how to precisely control the twist angle, which is also the most important obstacle to immediate progress in understanding the properties of 2D twistrionic materials. Uncontrolled strains related to details of how mechanical contact is established between the twisted layers can strongly influence the physical properties, such as basic features of the phase diagram that express gapped states, the appearance or absence of superconductivity and orbital magnetism in the MATBG.

(2) Physical mechanism of superconductivity in the MATBG

Two strongly correlated quantum states have been observed in the MATBG due to the enhanced interlayer coupling compared with graphene/h-BN heterostructures. However, there are some important problems that remain to be solved. For example, there is no definite evidence whether the MATBG is a Mott insulator when the flat band is half-filled; moreover, it is needed to be further studied whether the MATBG has anti-ferromagnetism if a Mott-like insulator state is confirmed in it. Besides, we are not yet clear on the superconducting pairing mechanism in the MATBG. The resemblance between the phenomenology of high-temperature superconductors and MATBG is striking, but whether these similarities reflect a deep connection remains to be seen.



(3) Development of multi-functional twistrionic devices

TMD twistrionic materials are a big family with rapid development and innovation. In twisted TMD heterobilayers, the moiré potential generally results in localized exciton states near the potential minimum position, which suggests a vdW heterostructure based strategy for realizing 2D periodic arrays that act as tunable uniform quantum light emitters. Furthermore, the number of trapped excitons could go down to one by carefully tuning the size and depth of the moiré potential, offering a source of single-photon emitter arrays. In addition, interlayer excitons at a given valley can absorb light with different circular polarization, which provides the possibility to fabricate high-performance semiconductor lasers and entangled photon sources as well as artificial excitonic crystals.

In conclusion, the moiré superlattice in 2D materials provides a platform of versatile tunability for studying exotic quantum phases of matter with imaging possibilities. Progress is being made in this direction, and we can expect that studies of moiré superlattices can advance twistrionics physics and help us understand other condensed matter systems.

Conflicts of interest

The authors declare no conflicts of interest.

Acknowledgements

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