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Shock wave and modelling study of the unimolecular dissociation of $\text{Si}(\text{CH}_3)_2\text{F}_2$: an access to spectroscopic and kinetic properties of SiF_2^\dagger

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The thermal dissociation of $\text{Si}(\text{CH}_3)_2\text{F}_2$ was studied in shock waves between 1400 and 1900 K. UV absorption-time profiles of its dissociation products SiF_2 and CH_3 were monitored. The reaction proceeds as a unimolecular process not far from the high-pressure limit. Comparing modelled and experimental results, an asymmetric representation of the falloff curves was shown to be most realistic. Modelled limiting high-pressure rate constants agreed well with the experimental data. The UV absorption spectrum of SiF_2 was shown to be quasi-continuous, with a maximum near 222 nm and a wavelength-integrated absorption cross section of $4.3 (\pm 1) \times 10^{-23} \text{ cm}^2$ (between 195 and 255 nm, base e), the latter being consistent with radiative lifetimes from the literature. Experiments over the range 1900–3200 K showed that SiF_2 was not consumed by a simple bond fission $\text{SiF}_2 \rightarrow \text{SiF} + \text{F}$, but by a bimolecular reaction $\text{SiF}_2 + \text{SiF}_2 \rightarrow \text{SiF} + \text{SiF}_3$ (rate constant in the range 10^{11} – $10^{12} \text{ cm}^3 \text{ mol}^{-1} \text{ s}^{-1}$), followed by the unimolecular dissociation $\text{SiF}_3 \rightarrow \text{SiF}_2 + \text{F}$ such that the reaction becomes catalyzed by the reactant SiF_2 . The analogy to a pathway $\text{CF}_2 + \text{CF}_2 \rightarrow \text{CF} + \text{CF}_3$, followed by $\text{CF}_3 \rightarrow \text{CF}_2 + \text{F}$, in high-temperature fluorocarbon chemistry is stressed. Besides the high-temperature absorption cross sections of SiF_2 , analogous data for SiF are also reported.

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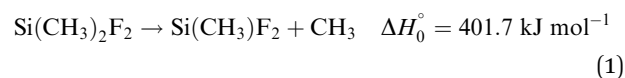
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Introduction

The role of SiF_2 in the etching of silicon by fluorine atoms continues to be under debate (see, for example, the review in ref. 1). Part of the problem lies in the scarcity of quantitative information on the properties of this intermediate in chemical or plasma-assisted etching processes employing fluorine-containing compounds. A limited amount of kinetic data is available for room temperature conditions (e.g., ref. 2–6), while less is known for elevated temperatures. The present article intends to improve this situation by investigating kinetic properties of SiF_2 under high-temperature conditions in shock waves.

First, a suitable source for SiF_2 had to be selected. It has been shown that the thermal dissociation of Si_2F_6 , in a process $\text{Si}_2\text{F}_6 \rightarrow \text{SiF}_2 + \text{SiF}_4$ and with a rate constant $10^{12.41} \exp(-193.5 \text{ kJ mol}^{-1}/RT) \text{ s}^{-1}$, directly produces SiF_2 .^{7,8} The thermal dissociation of SiF_4 , on the other hand, in a

sequence $\text{SiF}_4 \rightarrow \text{SiF}_3 \rightarrow \text{SiF}_2$ also leads to SiF_2 , but, because of the large thermal stability of SiF_4 , requires considerably higher temperatures than the dissociation of Si_2F_6 . In the present work, instead of Si_2F_6 or SiF_4 , it appeared more suitable to use $\text{Si}(\text{CH}_3)_2\text{F}_2$ as the precursor for SiF_2 . This compound is easy to handle in shock wave experiments and, at comparably low temperatures, it forms SiF_2 in a sequence of the two steps



(the given reaction enthalpies at 0 K, ΔH_0° , were determined by quantum-chemical calculations as described in the ESI†).

Next, a detection method for SiF_2 had to be chosen. As the UV absorption of CF_2 has been found useful to study high-temperature fluorocarbon chemistry,^{9–11} one may try to employ the analogous spectrum of SiF_2 to investigate reactions of the latter. At room temperature, SiF_2 has a band spectrum which is similarly structured and intense as that of CF_2 .¹³ One may expect that this spectrum at high temperatures becomes similarly quasi-continuous as that of CF_2 . One of the goals of the present work, therefore, was the characterization of the UV absorption spectrum of SiF_2 at elevated temperatures and to

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determine its absorption cross section as a function of temperature and wavelength in comparison to quantum-chemical calculations of its oscillator strength (as described in the ESI†).

There are more aspects of the present work. Monitoring SiF_2 formation in reactions (1) and (2) enables one to follow the unimolecular dissociation reaction of $\text{Si}(\text{CH}_3)_2\text{F}_2$. This molecule is sufficiently large to dissociate (under typical shock wave conditions) not far from the high-pressure limit of the unimolecular reaction. Therefore, a study of the pressure dependence of the dissociation rate constant k appears suitable to analyze its approach to the high-pressure rate constant k_∞ . This is an issue in standard unimolecular rate theory. The various versions of the latter propose different approaches of k to k_∞ .^{14,15} By comparing experimental and modelled rate constants, the present study provides an opportunity to address this problem in particular detail. In addition, in high-temperature fluorocarbon chemistry an autocatalytic pathway for CF_2 decomposition of the type $\text{CF}_2 + \text{CF}_2 \rightarrow \text{CF} + \text{CF}_3$, followed by $\text{CF}_3 \rightarrow \text{CF}_2 + \text{F}$, was observed.¹¹ It appears of interest to search for an analogous pathway $\text{SiF}_2 + \text{SiF}_2 \rightarrow \text{SiF} + \text{SiF}_3$, followed by $\text{SiF}_3 \rightarrow \text{SiF}_2 + \text{F}$, in high-temperature fluorosilicon chemistry. In both cases, the very endothermic direct dissociation of CF_2 or SiF_2 , respectively, then can be circumvented by a faster mechanism which also leads to dissociation.

Experimental technique and results

The present experiments have been performed by heating mixtures of $\text{Si}(\text{CH}_3)_2\text{F}_2$ and Ar in shock waves. $\text{Si}(\text{CH}_3)_2\text{F}_2$ (from abcr with a purity of 99%) could be used without further purification, because it was highly diluted (down to about 30 ppm) in the bath gas Ar (from Air Liquide with a purity of 99.9999%). The shock tube, as well as the UV lamp – quartz monochromator – photomultiplier – data acquisition equipment for recording absorption-time profiles, have, *e.g.*, been detailed in ref. 11 and 16 such that no further description is given here. In the first part of the present experiments, absorption-time profiles of shock-heated mixtures of about 100 ppm of $\text{Si}(\text{CH}_3)_2\text{F}_2$ in Ar were recorded at the wavelength of 222 nm, *i.e.* near to the maximum of the room-temperature absorption of SiF_2 .¹² Fig. 1 shows an example for a temperature of 1660 K behind the reflected shock wave. Directly behind the Schlieren peaks of the incident and reflected shock waves, no absorption signal is observed. This indicates that the absorption continuum of the parent molecule $\text{Si}(\text{CH}_3)_2\text{F}_2$ (having a room-temperature maximum near 155 nm¹⁷) with increasing temperature does not broaden to such an extent that it would influence absorption measurements at 222 nm. The absorption signal of Fig. 1 then can directly be related to the formation of SiF_2 (as reaction (1) is by far more endothermic than reaction (2), $\text{Si}(\text{CH}_3)_2\text{F}_2$ should rapidly dissociate to $\text{CH}_3 + \text{SiF}_2$). One may also look for an absorption signal from CH_3 . It is known that there is an absorption band of CH_3 not far from that of SiF_2 . However, its maximum is located at shorter wavelengths (near to 215 nm) and its maximum absorption cross section is much smaller than that of SiF_2 .¹⁸ Because the absorption

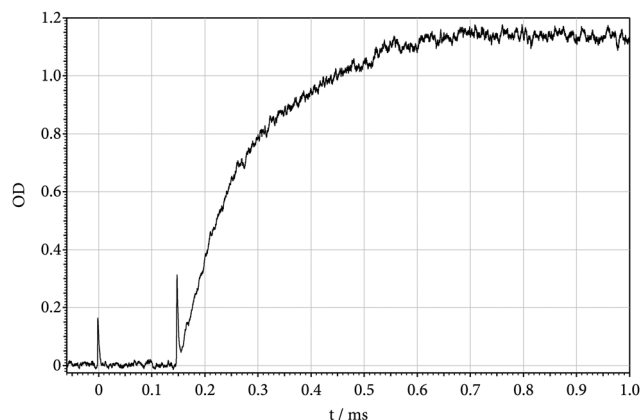


Fig. 1 Absorption-time profile at 222 nm of SiF_2 forming by unimolecular dissociation of $\text{Si}(\text{CH}_3)_2\text{F}_2$ behind reflected shock wave ($T = 1660$ K, $[\text{Ar}] = 8.6 \times 10^{-5} \text{ mol cm}^{-3}$, 100 ppm of $\text{Si}(\text{CH}_3)_2\text{F}_2$ in Ar; $\text{OD} = \sigma[\text{SiF}_2]$ with $l = 9.4$ cm).

cross section of SiF_2 decreases with decreasing wavelength, a signal from CH_3 could, nevertheless, be detected together with that from SiF_2 . Fig. 2 shows a signal recorded at 200 nm and employing larger reactant concentrations than in Fig. 1. Its magnitude corresponds to the marked decrease of the SiF_2 absorption cross section with decreasing wavelength as analyzed below and to the absorption cross section of CH_3 as reported in ref. 18. Fig. 2 indicates that $[\text{SiF}_2]$ and $[\text{CH}_3]$ have different time dependences. While $[\text{CH}_3]$ first increases and then decreases, $[\text{SiF}_2]$ like in Fig. 1 reaches a stationary final level. The decay of the CH_3 signal corresponds to the dimerization $2\text{CH}_3 \rightarrow \text{C}_2\text{H}_6$, whose rate, under the present conditions, is known.¹⁸ In contrast to $[\text{CH}_3]$, $[\text{SiF}_2]$ finally remains constant such as shown in Fig. 1. Apparently, the reverse of reactions (1) and (2) do not play a role, such that the signal of Fig. 1 can be attributed to the slower of reactions (1) and (2), in this case, obviously to reaction (1).

Absorption cross sections of SiF_2

Systematically inspecting final absorption levels of signals like Fig. 1 and 2, high-temperature absorption cross sections of SiF_2 were derived. Varying the temperature between 1500 and 1900 K (where the reaction was complete within the available observation time of about 1 ms), a major influence of temperature on the final absorption level could not be detected. On the one hand, this proved that the dissociation of $\text{Si}(\text{CH}_3)_2\text{F}_2$ was complete and no back-reaction took place. The final absorption level was found to be proportional to the reactant concentration which confirmed this conclusion. On the other hand, a temperature dependence of the absorption cross section of SiF_2 over this temperature range was only small. Slightly varying the wavelength (between 220 and 225 nm) indicated that the room-temperature band structure¹² was absent at the present elevated temperatures and that the SiF_2 spectrum now indeed was quasi-continuous. Assuming that each decomposing $\text{Si}(\text{CH}_3)_2\text{F}_2$ produces one SiF_2 , a maximum absorption cross section $\sigma(\text{SiF}_2, 222 \text{ nm}) = (2.45 \pm 0.5) \times 10^{-17} \text{ cm}^2$ (base e) was derived for temperatures near 1600 K. Fig. 3



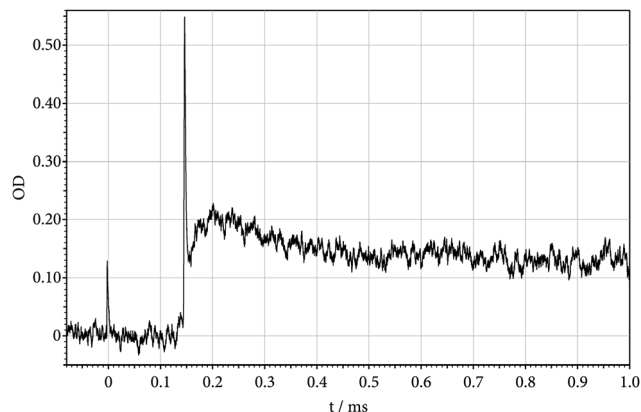


Fig. 2 As Fig. 1, but at 200 nm; superimposed absorptions of SiF_2 and CH_3 ($T = 1720$ K, $[\text{Ar}] = 8.2 \times 10^{-5} \text{ mol cm}^{-3}$, 210 ppm of $\text{Si}(\text{CH}_3)_2\text{F}_2$ in Ar).

shows the results for the wavelength dependence of $\sigma(\text{SiF}_2, \lambda)$ near 1600 and 3000 K (Table S1 of the ESI† shows experimental values of $\sigma(\text{SiF}_2, \lambda)$). Analogous to the observations for CF_2 from ref. 9, $\sigma(\text{SiF}_2, \lambda)$ has a Gaussian shape. The wavelength-integrated absorption cross section $I(\lambda) = \int \sigma(\text{SiF}_2, \lambda) d\lambda$ (between 195 and 255 nm) is equal to $4.3 (\pm 1) \times 10^{-23} \text{ cm}^3$, being close to the value $5.2 \times 10^{-23} \text{ cm}^3$ derived in ref. 12 from the experimental radiative lifetimes of SiF_2 from ref. 19.

A representation of the complete wavelength and temperature dependence of σ in Sulzer–Wieland form²⁰

$$\sigma(\nu, T) \approx \sigma_{\max} [\tanh(\theta_0/2T)]^{1/2} \exp\{-\tanh(\theta_0/2T) [(\nu - \nu_0)/\Delta\nu_0]^2\} \quad (3)$$

(with $\nu = 1/\lambda$) requires information on the four parameters σ_{\max} , θ_0 , ν_0 , and $\Delta\nu_0$. Fitting the data of Fig. 3 to eqn (3) leads to the parameters $\sigma_{\max} \approx 2.87 \times 10^{-17} \text{ cm}^2$, $\nu_0 \approx 45045 \text{ cm}^{-1}$, and $\Delta\nu_0 \approx 1785 \text{ cm}^{-1}$. The determination of the parameter θ_0 requires experiments over larger temperature ranges. In the present case, experiments were extended up to temperatures of 3200 K where $\text{Si}(\text{CH}_3)_2\text{F}_2$ decomposes in less than a μs . Fig. 4 shows an example for 3060 K. Instead of staying constant as in

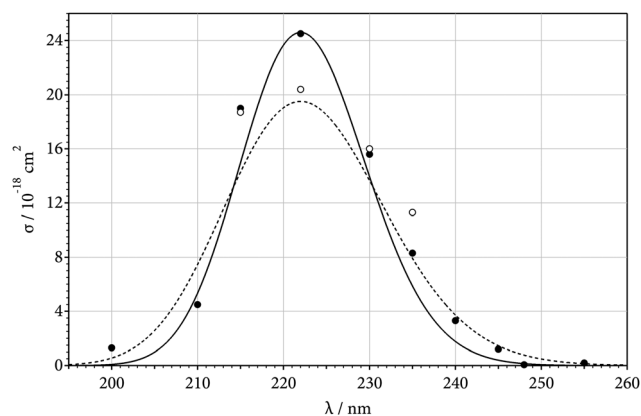


Fig. 3 Absorption cross sections σ of SiF_2 (experimental points from the present work, T near 1600 K, \bullet , and 3000 K, \circ ; representation of $\sigma(T)$ by eqn (3) for $T = 1600$ K, solid line, and 3000 K, dashed line).

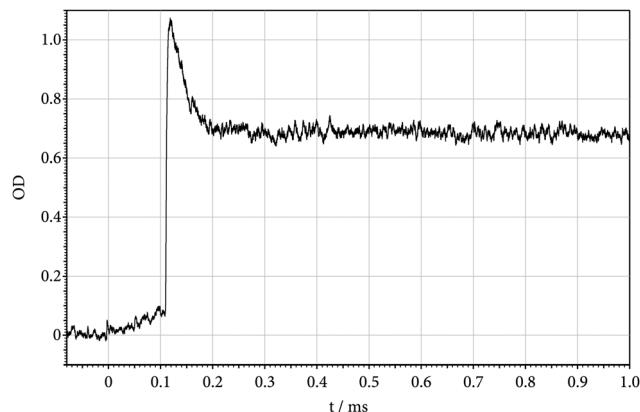


Fig. 4 As Fig. 1, but at higher temperatures (210 ppm of $\text{Si}(\text{CH}_3)_2\text{F}_2$ in Ar; $T = 1470$ K and $[\text{Ar}] = 1.6 \times 10^{-5} \text{ mol cm}^{-3}$ behind incident shock wave; $T = 3060$ K and $[\text{Ar}] = 8.1 \times 10^{-5} \text{ mol cm}^{-3}$ behind reflected shock wave; time scale behind incident shock wave compressed by a factor of 3.3).

Fig. 1, the signal here decreases by secondary reactions which will be analyzed later. Nevertheless, the absorption cross section of SiF_2 can be determined before this decay becomes a problem (e.g., after 10 μs in Fig. 4). Selected results are included in Fig. 3 and compared with a Sulzer–Wieland plot using $\theta_0 \approx 3000$ K. This value of θ_0 is relatively uncertain. It is mainly based on measurements at 222 nm and it is definitely larger than the corresponding value⁹ for the spectrum of CF_2 (it should be mentioned that the present results are consistent with values derived from the thermal dissociation of SiF_4 which will be reported separately, thus supporting the described analysis). It should also be mentioned that absorption-time signals at wavelengths larger than 240 nm increasingly deviate from Fig. 1. Apparently, absorptions from other species here are superimposed on the absorptions from SiF_2 , such as analyzed below.

Unimolecular dissociation of $\text{Si}(\text{CH}_3)_2\text{F}_2$

In the second part of our experiments, the kinetics of $\text{Si}(\text{CH}_3)_2\text{F}_2$ dissociation was explored. The time dependence of the absorption-time profile of Fig. 1 corresponds to a first-order process, i.e. $[\text{SiF}_2](t) = [\text{Si}(\text{CH}_3)_2\text{F}_2](t = 0) \{1 - \exp(-kt)\}$ with $k = 8.3 \times 10^3 \text{ s}^{-1}$. In cases where $[\text{SiF}_2](t)$ did not approach the final level $[\text{SiF}_2](t = \infty)$ sufficiently well within the available observation time, i.e. at temperatures below about 1500 K, the initial rate of absorption increase could also be evaluated to derive k . This required the use of absorption cross sections σ from eqn (3). Then, $\text{dOD}(t)/\text{d}t = \sigma l [\text{SiF}_2](t = \infty)k$ (with the optical density $\text{OD} = \sigma l [\text{SiF}_2]$ and the optical path length $l = 9.4 \text{ cm}$ of our arrangement) also led to rate constants k . This evaluation was necessary in particular for measurements behind incident shock waves. Fig. 4 shows an example. As our modelling (see the ESI†) predicted an only weak dependence of k on $[\text{Ar}]$, only the comparison of measurements behind incident and reflected shock waves provided a sufficiently large variation of bath gas concentrations to draw meaningful conclusions on the shape of falloff curves $k([\text{Ar}])$ (the accessible



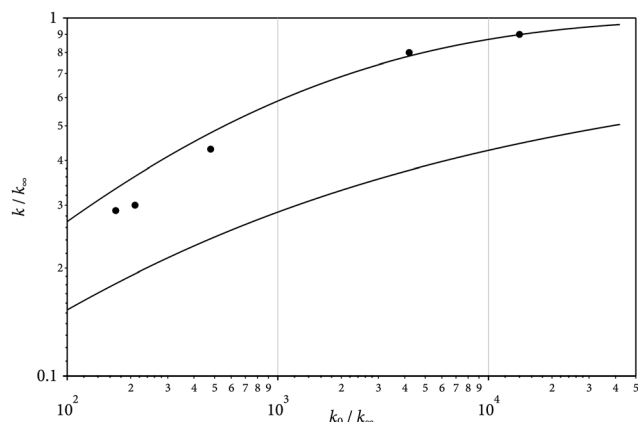


Fig. 5 Doubly-reduced representation of falloff curves $k([Ar])$ for the unimolecular dissociation of $Si(CH_3)_2F_2$ (representation of $k([Ar])/k_\infty$ as a function of k_0/k_∞ with k_0 from eqn (5) and k_∞ from eqn (4); upper solid line: modelling with eqn (8), lower solid line: modelling with eqn (7); experimental points from left to right: $T/K = 1400, 1500, 1600, 1700, 1800$, respectively).

range was from $[Ar] \approx 10^{-5} \text{ mol cm}^{-3}$ in incident waves to $10^{-4} \text{ mol cm}^{-3}$ in reflected waves). Fig. 5 compares two alternative representations of $k([Ar])$ with experimental results from incident and reflected shock waves (a doubly-reduced representation of k/k_∞ vs. k_0/k_∞ was chosen in order to include results from different temperatures; the used k_0 and k_∞ are from the modelled expressions given below). A more complete representation of experimental data is provided by the Arrhenius plots of Fig. 6 (for constant $[Ar]$; small mismatches of the experimental $[Ar]$ from the given values were accounted for by the $[Ar]$ -dependences of Fig. 5; the scatter of about $\pm 20\%$ of the points in Fig. 5 and 6 is larger than the systematic uncertainty of the measurements). Within the scatter, measured and modelled rate constants agree. As the reaction was studied not far from the high-pressure limit, the agreement mostly confirms the quality of the quantum-chemistry based

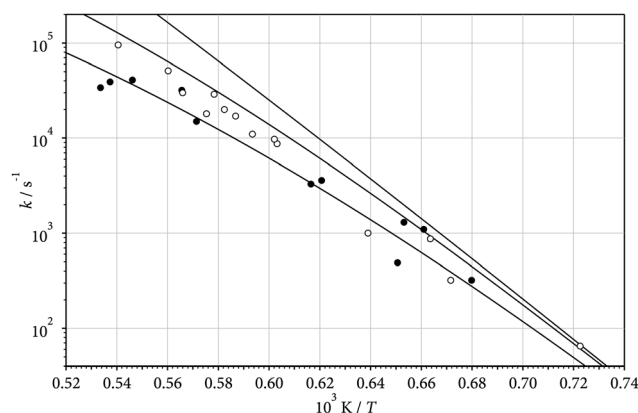


Fig. 6 Arrhenius plots of rate constants $k(T)$ for the unimolecular dissociation of $Si(CH_3)_2F_2$ at $[Ar] \approx 10^{-5}$ (\circ) and 10^{-4} (\bullet) mol cm^{-3} , between $T = 1370$ and 1890 K (modelled lines: representation of falloff curves by eqn (8), from bottom to top for $[Ar] = 10^{-5} \text{ mol cm}^{-3}$, $[Ar] = 10^{-4} \text{ mol cm}^{-3}$, and k_∞ , see the ESI†).

calculation of k_∞ (see the ESI†). Because of the uncertainty of the used collisional energy transfer parameters (see the ESI†), the modelling of the low-pressure rate constants k_0 is less certain. Its influence on the derived high-pressure constants k_∞ , however, is only weak. In conclusion, the experimental data are remarkably consistent with the high-pressure rate constants such as modelled in the ESI†. These can be expressed by

$$k_\infty = 1.24 \times 10^{19} (T/2000 \text{ K})^{-6.63} \exp(-58400 \text{ K}/T) \text{ s}^{-1} \quad (4)$$

Low-pressure rate constants were modelled as

$$k_0 \approx [Ar] 2.94 \times 10^{25} (T/2000 \text{ K})^{-25.04} \exp(-61980 \text{ K}/T) \text{ cm}^3 \text{ mol}^{-1} \text{ s}^{-1} \quad (5)$$

The representation of the falloff curves of Fig. 5 has employed expressions of the form

$$k([Ar])/k_\infty = [x/(1+x)]F(x) \quad (6)$$

with $x = k_0/k_\infty$ and “broadening factors” $F(x)$. Either “symmetric broadening factors” $F(x)$ (i.e., $F(x) = F(1/x)$) of the form proposed in ref. 21,

$$\log F(x) \approx \log F_{\text{cent}}/[1 + (\log x/N)^2] \quad (7)$$

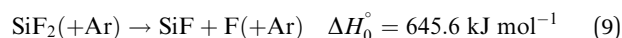
(with $N \approx 0.75\text{--}1.27 \log F_{\text{cent}}$ and system-specific “center broadening factors” F_{cent}^{22}), or “asymmetric broadening factors” $F(x)$ (i.e. $F(x) \neq F(1/x)$) were used, the latter being of the form proposed in ref. 14 and 15

$$F(x) \approx (1 + x/x_0)/[1 + (x/x_0)^n]^{1/n} \quad (8)$$

(with $n = [\ln 2/\ln(2/F_{\text{cent}})] [1 - b + b(x/x_0)^q]$, $q = (F_{\text{cent}} - 1)/\ln(F_{\text{cent}}/10)$, and the parameters x_0 and b close to $x_0 = 1(\pm 0.1)$ and $b = 0.2(\pm 0.05)$). The comparison of the two alternative expressions for falloff curves with the experimental data in Fig. 5 suggests that the asymmetric form of $F(x)$, i.e. Eqn (8), near to the high-pressure limit performs much better than the symmetric form of $F(x)$, i.e. Eqn (7).

Kinetics of SiF_2 reactions

The experiments described so far, which characterize the formation of SiF_2 (and CH_3) in the unimolecular dissociation of $Si(CH_3)_2F_2$, finally were extended to higher temperatures where the primary dissociation is so rapidly complete that it cannot be resolved any longer. Fig. 4 gives an example for the reflected shock wave. If only reactions (1) and (2) would take place, the absorption signal behind the reflected shock wave then would remain constant. Instead, one observes a decrease to a new steady level. The decrease of the signal is much faster than expected for the thermal dissociation of SiF_2 , i.e.



(for this reaction, a rate constant of $k_{9,0} = [Ar] 2 \times 10^{16} (T/1000 \text{ K})^{-1.34} \exp(-72910 \text{ K}/T) \text{ cm}^3 \text{ mol}^{-1} \text{ s}^{-1}$ has been modelled analogous to the calculations described in the present ESI† such that SiF_2 would have a half-life of about 7 s; likewise, the final absorption level of Fig. 4 cannot correspond to a dissociation equilibrium $SiF_2 \leftrightarrow SiF + F$). An unambiguous interpretation of



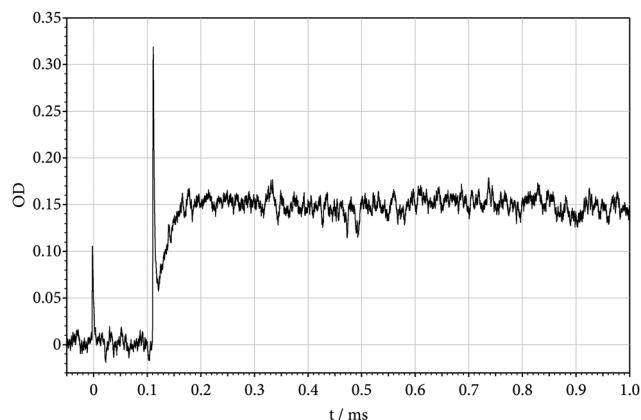
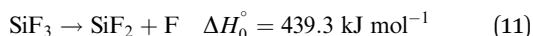


Fig. 7 As Fig. 1, but at 265 nm, showing the formation of SiF in the consumption of SiF₂ by reaction (12) (reflected shock wave with $T = 3080$ K, $[\text{Ar}] = 3.8 \times 10^{-5} \text{ mol cm}^{-3}$, 210 ppm of Si(CH₃)₂F₂ in Ar).

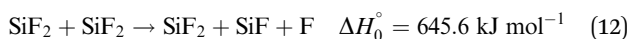
the signal, instead, is provided by measurements at wavelengths where an absorption from SiF₂ can be neglected. Fig. 3 indicates that, even at the temperature of Fig. 4, an absorption signal from SiF₂ should be negligible at wavelengths larger than about 250 nm. Fig. 7 gives an example for a wavelength of 265 nm and nearly the same temperature as in Fig. 4. The initial decay of the SiF₂ signal from Fig. 4 now is mirrored by an absorption increase in Fig. 7. The rate of the initial decay of SiF₂ in Fig. 4 and the formation of a reaction product in Fig. 7 both were found to increase proportional to $[\text{SiF}_2](t = 0)$, *i.e.* the observation corresponds to a bimolecular process. In addition, the rate constant for this process was found to have an only small positive temperature coefficient. These observations suggest that SiF₂ is consumed by a reaction



A modelling of the rate constant for unimolecular dissociation of SiF₃ analogous to that described in the present ESI† on the other hand, indicates that SiF₃ under the conditions of Fig. 4 and 7 should rapidly dissociate by



The sequence of reactions (10) and (11), *i.e.*



then corresponds to a process which is catalyzed by the reactant SiF₂ and which is much faster than the slow thermal dissociation of SiF₂ by reaction (9). We found no evidence against the assumption that reaction (12) proceeds until SiF₂ is completely consumed and converted to SiF + F. In this case, the final absorption levels of Fig. 4 and 7 can be attributed exclusively to SiF and high-temperature absorption cross sections of SiF can also be derived. Values of $\sigma/10^{-17} \text{ cm}^2 = 1.0, 2.2, 1.3, 0.9, 0.5, 0.3, 0.4, 0.5, 0.8, 0.6$, and 0.2 were determined near 3000 K for wavelengths of 200, 210, 220, 230, 240, 250, 260, 270, 280, 290, and 300 nm, respectively. It is known that SiF has numerous band systems from the vacuum-UV to the red (see a summary in ref. 23). At high temperatures, hot bands from these systems overlap into a broad quasi-continuum,

extending beyond that from SiF₂, but intense enough to be observed. The oscillator strengths of the band systems of SiF and SiF₂ in the ESI† were modelled to be of similar magnitude, which appears consistent with the present observations.

Evaluating SiF₂ consumption and SiF formation from experiments like Fig. 7, led to rate constants k_{10} in the range 10^{11} – $10^{12} \text{ cm}^3 \text{ mol}^{-1} \text{ s}^{-1}$ between $T = 1900$ and 3500 K, respectively. Experiments with varying reactant concentrations led to similar values which supported the proposed interpretation. One may finally ask for the fate of the F atoms from the net reaction (12). This question could not be answered here. It may be that leftover C₂H₆ and CH₃ from the precursor act as a sink for these atoms. Evidence for an interference with the described mechanism of reactions (10) and (11) was not found. It should finally be mentioned that absorption signals like Fig. 7 at higher temperatures and higher reactant concentration show a decrease with time, before another increase sets in. These observations are similar as those found in the fluorocarbon system.¹¹ An analogous interpretation by secondary reactions like $\text{SiF} + \text{SiF} \rightarrow \text{Si}_2\text{F} + \text{F}$, followed by $\text{Si}_2\text{F} \rightarrow \text{Si}_2 + \text{F}$, would appear possible, but cannot be confirmed at this stage. More details of the suggested autocatalytic reaction sequence of reactions (10) and (11) clearly have to be explored.

Conclusions

The present work illustrated that the thermal dissociation of Si(CH₃)₂F₂ is a suitable source for generating SiF₂ under high-temperature conditions such as studied in shock waves. On the one hand, this allowed to record and calibrate the temperature- and wavelength-dependence of UV absorption cross sections of SiF₂. The wavelength-integrated absorption cross section here was found to be consistent with the value derived from the radiative lifetime of the species at room temperature.¹² In future work on high-temperature reactions of SiF₂, the absorption cross sections from eqn (3) will serve for quantitative determinations of SiF₂ concentrations.

In addition to the study of the UV spectrum of SiF₂, the thermal dissociation of Si(CH₃)₂F₂ could be studied under conditions where the reaction is unimolecular. The reaction was found to be not far from its high-pressure limit. A quantum-chemistry based modelling of the rate constant gave results in close agreement with the experiments, which confirmed the reliability of the modelling approach. The falloff curves of the unimolecular reaction could best be represented with asymmetric broadening factors in the form suggested in ref. 14 and 15.

It was finally suggested that the consumption of SiF₂ under the applied conditions did not proceed by thermal unimolecular dissociation, but by an autocatalytic process, *i.e.* via a sequence of the steps $\text{SiF}_2 + \text{SiF}_2 \rightarrow \text{SiF} + \text{SiF}_3$, followed by $\text{SiF}_3 \rightarrow \text{SiF}_2 + \text{F}$. An analogy to the reaction sequence $\text{CF}_2 + \text{CF}_2 \rightarrow \text{CF} + \text{CF}_3$, followed by $\text{CF}_3 \rightarrow \text{CF} + \text{F}$, as observed in high-temperature fluorocarbon chemistry appears obvious.



Conflicts of interest

There are no conflicts of interest to report.

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