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Spectroscopy and photochemistry of copper nitrate clusters†

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The investigation of copper nitrate cluster anions $\text{Cu}(\text{II})_n(\text{NO}_3)_{2n+1}^-$, $n \leq 4$, in the gas phase using ultraviolet/visible/near-infrared (UV/vis/NIR) spectroscopy provides detailed insight into the electronic structure of the copper salt and its intriguing photochemistry. In the experimentally studied region up to 5.5 eV, the spectra of copper(II) nitrate exhibit a 3d–3d band in the vis/NIR and well-separated bands in the UV. The latter bands originate from Ligand-to-Metal Charge Transfer (LMCT) as well as $n-\pi^*$ transitions in the nitrate ligands. The clusters predominantly decompose by loss of neutral copper nitrate in the electronic ground state after internal conversion or *via* the photochemical loss of a neutral NO_3 ligand after a LMCT. These two decomposition channels are in direct competition on the ground state potential energy surface for the smallest copper nitrate cluster, $\text{Cu}(\text{II})(\text{NO}_3)_3^-$. Here, copper nitrate evaporation is thermochemically less favorable. Population of π^* orbitals in the nitrate ligands may lead to N–O bond photolysis. This is observed in the UV region with a small quantum efficiency, with photochemical loss of either nitrogen dioxide or an oxygen atom.

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Introduction

Copper and copper oxide nanoparticles are used in industry for many oxidation^{1–3} and reduction processes^{4,5} due to their favorable catalytic properties.^{6,7} As a coinage metal, copper exhibits a very distinct chemistry acting as an electron donor.^{8,9} It is particularly useful in methanol synthesis,^{10,11} carbon dioxide activation^{12,13} and hydrogen storage applications.^{14–16} Copper salts are typically used in the production of heterogeneous catalysts *via* a calcination process.^{17–20} Copper nitrate is the most widely used copper salt due to its stability, high availability and comparatively low cost.²¹

From a fundamental point of view, the photochemistry of nitrate is of high interest in its own right since it plays a key role in our atmosphere as trace compound.²² In the photochemistry of snow and ice, nitrate photolysis leads to very reactive products.^{23–25} The photochemistry of systems can be challenging to describe by quantum chemical methods since excited state calculations are still much more demanding than calculations on the electronic ground state.^{26,27} The photochemistry of nitrate has been investigated experimentally in aqueous solution, leading to very reactive OH radicals, among other products.²⁸ The photolysis of the isolated nitrate anion in the gas phase has only been

investigated theoretically, predicting N–O bond photolysis with the negative charge remaining on either fragment,²⁹ while experimental data on isolated nitrate or on nitrate ligands coordinated to metal centers is still missing. Well-defined gas-phase experiments can provide valuable insight into reaction mechanisms on a molecular level,^{30–32} while spectroscopy based experiments can yield detailed insight into the molecular and electronic structure.^{33,34}

In the gas phase, many molecular mechanisms involving the activation of methane,^{35,36} carbon dioxide^{37–40} and elementary steps of hydrogen storage applications^{41–43} have successfully been studied. Ultraviolet/visible/near-infrared (UV/vis/NIR) photodissociation spectroscopy is a very powerful tool to characterize the complex electronic structure of transition metal complexes and clusters in the gas phase.^{44–54} We recently investigated photochemical hydrogen evolution from hydrated magnesium,^{55,56} the effect of salt environments on the reactions and photolysis of organic substances,^{57–59} as well as the evolution of the hydration environment of a single electron⁶⁰ or a carbon dioxide radical anion⁶¹ using action spectroscopy.

Only recently, CuO^+ ligated with acetonitrile has been investigated spectroscopically, showing a very complex interaction with the additional ligand changing the biradical-oxygen character of bare $\text{Cu}(\text{I})^+\text{O}$ partially to $\text{Cu}(\text{II})^{2+}\text{O}^-$.⁶² Our investigation of copper formate by infrared multiple photon dissociation (IRMPD) revealed a complex vibrational signature involving Fermi interactions.^{63–65} The formation of formic acid upon IRMPD depends strongly on the available energy and the oxidation

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state of the copper centers. Furthermore, UV/vis photodissociation of copper(II) formate provided detailed insights into its photochemistry.⁶⁶ Internal conversion dominates after exciting 3d–3d excitations in the visible region and Ligand-to-Metal Charge Transfer (LMCT) excitations in the UV, which additionally allowed the photochemical loss of the neutral ligand.⁶⁶ In contrast, in a previous UV/vis study of $\text{Cu(II)(NO}_3)_3^-$ the 3d–3d excitations expected for the $[\text{Ar}]3d^9$ configuration of copper(II) have not been reported.⁶⁷ Two bands in the UV were observed in this study, along with a single decomposition fragment, $\text{Cu(I)(NO}_3)_2^-$.⁶⁷ However, decomposition of $\text{Cu(II)(NO}_3)_3^-$ in the electronic ground state following collisional activation revealed nitrate anion formation as a competing second channel, which is even preferred at lower energies.⁶⁸ Further decomposition of $\text{Cu(I)(NO}_3)_2^-$ results in $\text{CuO(NO}_3)^- + \text{NO}_2$ or $\text{Cu(I)NO}_3 + \text{NO}_3^-$.⁶⁸

Here, we investigate $\text{Cu(II)}_n(\text{NO}_3)_{2n+1}^-$, $n = 1-4$, using UV/vis/NIR action spectroscopy and quantum chemical calculations. Detailed insight into the electronic configuration, the photochemistry of copper nitrate as well as its decomposition on the ground state potential energy surface is obtained. In addition, the results provide valuable insight into the photochemistry of nitrate anions.

Experimental & theoretical methods

Mass selected copper nitrate anions $\text{Cu(II)(NO}_3)_3^-$ are introduced into the gas phase *via* electrospray ionization (ESI). For this purpose, copper nitrate (Sigma Aldrich) is dissolved at a concentration of about 5 mmol l^{-1} in a solution consisting of 50% water and 50% methanol (HPLC-Grade, Carl Roth). For larger clusters, isotopic enrichment is advantageous to increase the signal-to-noise ratio. To this end, the sample is prepared by reacting isotopically enriched 63-copper oxide (98.6% enrichment, Isotope JSC) at a concentration of 200 mmol l^{-1} with nitric acid at 60°C in a solution of 80% water and 20% nitric acid (70%, Carl Roth) to isotopically enriched $^{63}\text{Cu(II)(NO}_3)_2$. For ESI, this solution is then diluted in a 1:19 ratio with acetonitrile (HPLC grade, Merck), adding 3% in total volume of nitric acid to support the clustering to larger copper nitrate clusters. This allows to investigate up to $^{63}\text{Cu}_4(\text{NO}_3)_9^-$ in our experiment.

The ions formed by ESI are transferred to the cell of a commercial Bruker APEX Qe 9.4 Tesla Fourier-Transform Ion Cyclotron Resonance (FT-ICR) mass spectrometer described in more detail elsewhere.^{57,58,63} Here, the trapped clusters are irradiated from 225 to 1650 nm for up to 20 s with pulsed laser light at a repetition rate of 20 Hz provided by an EKSPALA NT342B optical parametric oscillator. The photodissociation spectra are corrected for fragmentation caused by black body infrared radiative dissociation (BIRD)^{69–74} and collisions with the background gas, while the photodissociation cross sections are calculated as described in detail before.^{66,75} Hereby, the beam profile is estimated as Gaussian beam and a correction factor is applied to the measured laser power after each mass spectrum to account for losses in beam steering optics and windows. An additional correction factor is applied to compensate for the

beam walk-off upon switching between the OPO regions, signal to idler at 710 nm, sum frequency generation to signal at 410 nm, and second harmonic generation to sum frequency generation at 296 nm, as described before.⁶⁶

For the ground-state calculations of copper nitrate clusters, density functional theory (DFT) within the B3LYP/def2TZVP level is used for geometry optimization, based on extensive benchmarking with the previously studied copper salt, copper formate.^{63–66} The good performance of the B3LYP method was independently assessed in the decomposition of copper nitrate.⁶⁸ Initially, several different structures are investigated with intact nitrate ligands and the nitrate photolysis products O^- and NO_2^- in different binding motifs towards the copper centers. All the calculated reaction energies are corrected for zero-point effects, and the stability of the ground state wave function was tested for all optimized structures.

The excited states are investigated using Equation of Motion – Coupled Cluster Singles and Doubles (EOM-CCSD) calculations for $\text{Cu(NO}_3)_3^-$.^{76,77} The time dependent density functional theory (TD-DFT) with the BMK functional, TD-BMK, provided results close to the ones of EOM-CCSD, with small shifts towards lower energies. The aug-cc-pVDZ basis set seems sufficient to describe the electronic configuration as it yields results close to the triple-zeta basis set, see Table S1 and Fig. S5 in the ESI.† All excited states up to an excitation energy of at least 6 eV are considered. Due to computational efficiency, TD-BMK is applied in the investigation of $\text{Cu}_2(\text{NO}_3)_5^-$. The character of the electronic excitations is investigated for $\text{Cu(II)(NO}_3)_3^-$ using the Natural Transition Orbital (NTO) scheme.⁷⁸ The modelled excitation spectrum was obtained *via* a molecular dynamics run on the B3LYP/def2TZVP level of theory at a temperature of 300 K over the course of 101 ps at a step size of 100 a.u. (about 2.4 fs; note that no hydrogen atoms are present in the studied systems and thus a larger step size might be used). The first 8000 steps are discarded for thermalization. The electronic excitations are calculated every 50 steps using TD-BMK/aug-cc-pVDZ with 50 electronic states. The calculated peaks are convoluted with Gaussian functions at a width of 0.3 eV to produce the spectra.

All calculations were carried out with the Gaussian 16 program⁷⁹ in combination with the ABIN code⁸⁰ for the molecular dynamics.

Results and discussion

Photodissociation spectra

We start our investigation of the copper nitrate system with UV/vis/NIR action spectroscopy experiments on $\text{Cu(II)(NO}_3)_3^-$, $\text{Cu(II)}_2(\text{NO}_3)_5^-$, $\text{Cu(II)}_3(\text{NO}_3)_7^-$ and $\text{Cu(II)}_4(\text{NO}_3)_9^-$. The resulting total photodissociation cross sections σ are illustrated in Fig. 1. For the two smallest clusters, the calculated oscillator strength f of the electronic transitions is shown for the energetically lowest lying isomer along with the predicted structure.

For each cluster size, one can find two well-separated bands along with the flank of a third band in the deep UV above 5 eV. With the exception of $\text{Cu(II)(NO}_3)_3^-$, the low-energy band peaks



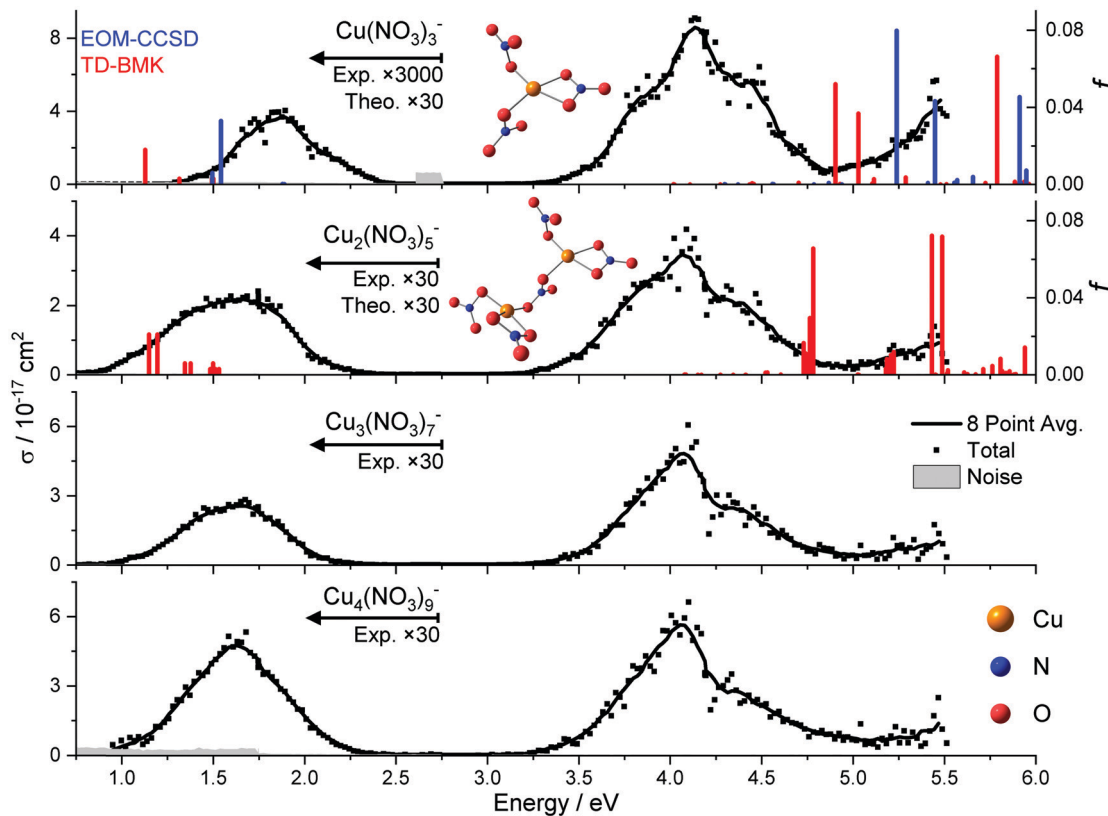


Fig. 1 The photodissociation cross section σ of $\text{Cu}(\text{II})(\text{NO}_3)_3^-$, $\text{Cu}(\text{II})_2(\text{NO}_3)_5^-$, $\text{Cu}(\text{II})_3(\text{NO}_3)_7^-$ and $\text{Cu}(\text{II})_4(\text{HCO}_2)_9^-$ along with the calculated electronic transitions for the energetically most favorable isomer of the two smallest clusters. The high level of theory EOM-CCSD/aug-cc-pVDZ//B3LYP/def2TZVP was used for $\text{Cu}(\text{II})(\text{NO}_3)_3^-$. TD-BMK/aug-cc-pVDZ//B3LYP/def2TZVP calculations were performed for $\text{Cu}(\text{II})(\text{NO}_3)_3^-$ and $\text{Cu}(\text{II})_2(\text{NO}_3)_5^-$. The calculated oscillator strength f is displayed by vertical lines. The section below 2.75 eV has been scaled for better visibility.

at about 1.6 eV in the NIR. It almost does not shift with cluster size, and its width changes only slightly. The band in the UV peaks at about 4.1 eV, featuring a roughly 30–60 times higher transition intensity than that of the vis/NIR band for $n \geq 2$. Also, the characteristic shape of the UV band hardly changes with cluster size. For $\text{Cu}(\text{II})(\text{NO}_3)_3^-$, the low-energy band peaks at about 1.9 eV, and is roughly 6000 times less intense than the UV band. In addition, its width is significantly reduced compared to the larger clusters. Both properties indicate that this band originates from a multi-photon process.

Our experimental spectrum of $\text{Cu}(\text{II})(\text{NO}_3)_3^-$ is comparable to the earlier measurement of Kaufman and Weber, covering the photon energy range of 3.0–5.6 eV.⁶⁷ They observed a very broad band around 4.5 eV, with a hint of a shoulder at 3.8 eV. This band lies in the onset of a very intense transition deeper in the UV. However, the reflectron-time-of-flight (RETOF) setup used by Kaufman and Weber lacks a thermalization region, and they report metastable decay of some ions. This indicates that the ions in their experiment had significantly higher internal energy than in the present work, where the ions are thermalized to room temperature in the hexapole collision cell and by exchange of infrared radiation in the almost collision-free environment of the ICR cell. This explains the significantly broader, less structured band in the RETOF work. Due to the experimental setup, only one laser pulse can be applied for

photodissociation in the RETOF work, which probably renders the very weak absorption in the vis/NIR region inaccessible. To detect this band, we irradiated the ions for up to 20 s at a pulse repetition rate of 20 Hz, *i.e.*, with up to 400 laser pulses.

Excited state calculations

Theoretical calculations on the B3LYP/def2TZVP level of theory predict one energetically low-lying isomer for $\text{Cu}(\text{II})(\text{NO}_3)_3^-$ with C_2 symmetry, shown in Fig. 1. In this structure, one nitrate acts as a bidentate ligand, while the other two nitrate ligands form one short Cu–O bond with the metal center. Minimizing Coulomb repulsion between these two nitrate ligands and maximizing the Coulomb attraction with the Cu(II) center determines their alignment relative to one another. Similar structural motifs are present in the lowest energy structure found for $\text{Cu}(\text{II})_2(\text{NO}_3)_5^-$. Both Cu(II) centers have one bidentate ligand. The Cu(II) centers are connected *via* a bridging bidentate nitrate ligand, while the remaining nitrate ligands form one Cu–O bond. Overall, the calculated structures reflect the preferred square-planar coordination of the $[\text{Ar}]3d^9$ configuration of Cu(II).

The highest level of theory that we managed to apply to the $\text{Cu}(\text{II})(\text{NO}_3)_3^-$ ion was EOM-CCSD/aug-cc-pVDZ. While the calculations predict well-separated bands in the vis/NIR and within the UV, the energetic position deviates from experiment.



Even on the EOM-CCSD/aug-cc-pVDZ level of theory, the discrepancy of the first band is about 0.3 eV, and the shift is even higher with at least 1.1 eV in the UV if we consider only spin-allowed transitions between doublet states. While zero-point effects will shift the calculated energies of dissociative excited states to somewhat lower values due to the kinetic energy correction,⁸¹ one can expect these effects to be in the range of about 0.2–0.4 eV. TD-DFT calculations on the TD-BMK/aug-cc-pVDZ//B3LYP/def2TZVP level of theory reproduce the trends of EOM-CCSD/aug-cc-pVDZ relatively well for a transition metal complex, with small shifts towards lower energies.

For the larger cluster, $\text{Cu(II)}_2(\text{NO}_3)_5^-$, the theoretical calculations at the minimum-energy geometry of the ground state, *i.e.* in the Franck–Condon (FC) point, predict a very similar situation compared to $\text{Cu(II)}(\text{NO}_3)_3^-$, with minor energetic shifts. This can be expected since a nitrate ligand inserts between the two copper ions in the energetically most favorable isomer, see Fig. 1. For this cluster size, the experimental and theoretical intensities agree relatively well between the vis/NIR and the UV band, which is additional evidence for the involvement of a multiple-photon dissociation in the vis/NIR band of $\text{Cu(II)}(\text{NO}_3)_3^-$. The energetic position of the vis/NIR band in the TD-DFT calculations matches well with our experiment after considering the shift towards higher energies in the higher level method. The absence of a copper–copper bond further explains why the spectra remain largely unchanged for larger copper nitrate clusters. The same behavior was observed previously with copper formate clusters.⁶⁶ However, the computationally predicted position of the UV band here still deviates significantly from experiment.

Molecular dynamics of $\text{Cu(II)}(\text{NO}_3)_3^-$ performed at 300 K using TD-BMK/aug-cc-pVDZ//B3LYP/def2TZVP suggest that the first three intense UV transitions merge to one broad UV band, as displayed in Fig. 2. The UV band is shifting to slightly lower excitation energies compared to the excitations in the FC point. Another high-energy band arises in these calculations from significantly more intense transitions even deeper in the UV. Interestingly, the structure is very flexible in the dynamics.

In particular, a Cu–O bond to the bidentate ligand may break, and one of the monodentate ligands may turn bidentate by forming a new Cu–O bond. This mechanism affords complete randomization of the ligands. Overall, the modelled absorption spectra exhibit a comparable shape to the experimental photodissociation action spectrum of clusters with $n \geq 2$, with transition intensities in the same order of magnitude, while the experimental cross section in the vis/NIR of $\text{Cu(II)}(\text{NO}_3)_3^-$ is significantly smaller. In this energy range, the cross section is lower because the photodissociation of $\text{Cu(II)}(\text{NO}_3)_3^-$ requires the sequential absorption of two photons in order to gain sufficient energy to decompose, which is discussed in detail below.

A possible explanation for the energetic discrepancy of the transitions in copper nitrate anions in the UV are multi-reference effects, which a single-reference based method like EOM-CCSD and TD-BMK cannot describe properly. Copper as transition metal is well-known to be difficult to calculate, and potential charge-transfer excitations involving the nitrate anion with many degenerate excited states pose additional challenges to theory. Multi-reference calculations are required to properly describe these excitations, with many other states being close in this energy range.^{27,82} However, single-reference based methods worked relatively well for a different copper salt, copper formate.⁶⁶

An alternative explanation for the discrepancy is inaccurate oscillator strengths due to the involvement of quartet states of copper nitrate clusters, which lie in the energy range of the experimental UV band. Upon considering spin–orbit coupling, they might mix with the intense, spin-allowed excitations deeper in the UV and borrow a part of their oscillator strength.^{83,84} However, the theoretical description of the multi-reference character and spin–orbit coupling remains challenging, and our attempts in this direction were thus far not successful. For a meaningful multi-reference calculation, a relatively large active space including copper's active 3d, 4s and 4p orbitals along with all participating nitrate orbitals is required. Such calculations with a reliable method, for example the Multi-Reference Configuration Interaction (MRCI) level of theory, are computationally not feasible for a system of this size. In addition,

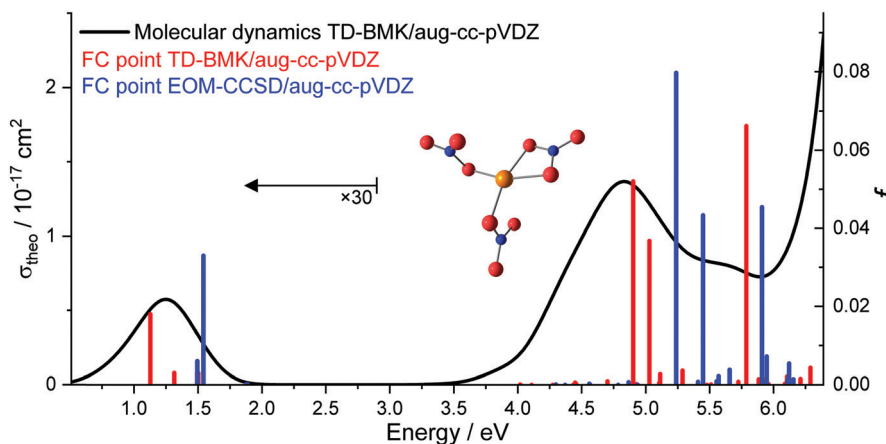


Fig. 2 The modelled photodissociation cross section σ_{theo} of $\text{Cu(II)}(\text{NO}_3)_3^-$ via a molecular dynamics run at 300 K over 101 ps at the TD-BMK/aug-cc-pVDZ//B3LYP/def2TZVP level of theory along with the oscillator strength of the electron excitations in the FC point for comparison.



calculations including spin-orbit coupling in open-shell systems are currently not implemented across the quantum chemical program packages accessible to us in a way that would allow for a treatment of even the smallest copper nitrate cluster anion. Therefore, a definitive assignment of the bands remains challenging. The interpretation of the experiments here relies on comparison to similar cases, while the theoretical results using single-reference based methods without spin-orbit coupling can only yield semi-quantitative support.

Excitation character

The character of the excitations is analyzed for $\text{Cu(II)(NO}_3)_3^-$ using natural transition orbitals on the TD-BMK/aug-cc-pVDZ//B3LYP/def2TZVP level of theory, see Fig. 3. The first band in the experiment corresponds to the first four calculated transitions. These are 3d–3d excitations from one of the four doubly occupied 3d orbitals into the singly occupied $3d_{xy}$ orbital of Cu(II), resulting in the excited states D_1 – D_4 . Interestingly, the lowest excitation energies are calculated for excitation from the $3d_{z^2}$ and $3d_{x^2-y^2}$ orbital, leading to excited states D_1 and D_2 , respectively. In this case, the electron density in the xy plane, *i.e.*, the plane of coordination, is not significantly increased. Significantly more energy is required for excitation from the $3d_{xz}$ and $3d_{yz}$ orbital, which feature nodes in the xy plane, corresponding to the excited states D_3 , D_4 . In this case, the electron density at the copper center increases in the excited state in the xy plane, which weakens the coordinate bonds with the nitrate ligands.

The lowest lying excited states in the UV, D_5 and D_6 , exhibit a biradical character and share the configuration with the quartet states $Q_{1,2}$. They correspond to $n-\pi^*$ excitations of the nitrate ligands, which energetically fit well to the first UV band. While they exhibit a relatively low oscillator strength, they play an important role in the photochemistry of the system, see below. Deeper in the UV beyond 6 eV, more intense transitions within

the nitrate ligands can be found, which could contribute to the flank at the upper energy limit of the experiment.

The three lowest-lying strong transitions with oscillator strengths $f > 0.01$ correspond to a UV band in the experiment, involving excited states D_8 , D_9 and D_{12} . They have LMCT character from the doubly occupied n orbitals of the nitrate ligands towards the singly occupied 3d orbital in copper(II). The interpretation of the excitation character involving LMCT excitations along with excitations within the nitrate ligand deeper in the UV agrees with the previous assignment of the $\text{Cu(II)(NO}_3)_3^-$ spectrum by Kaufman and Weber.⁶⁷ However, with spin-orbit coupling also the low-lying quartet states might play a role in the first UV band. Comparison to copper(II) formate clusters with similar absorption bands in the vis and early UV shows very analogous excitation characters, with energetically low-lying 3d–3d excitations along with LMCT excitations in the UV while the high-energy band in the UV observed here is missing.⁶⁶ This indicates that the latter excitation arises from excitations within the nitrate ligand, since the electronic structure of the formate ligand differs considerably, and excitations are much higher in energy.

As the copper centers do not interact directly with each other in larger clusters and the experimental spectra are very similar, the character of the excitations is not expected to change with cluster size.

Photochemistry

To investigate the photochemistry of copper nitrate clusters, selected partial photodissociation cross sections of $\text{Cu(II)(NO}_3)_3^-$ and $\text{Cu(II)}_2(\text{NO}_3)_5^-$ are displayed in Fig. 4a and b, respectively. In the low-energy band, $\text{Cu(II)(NO}_3)_3^-$ only decomposes into the energetically most favorable products, namely NO_3^- and $\text{Cu(I)(NO}_3)_2^-$ *via* reaction (1) and (2), respectively, see Table 1. However, two photons are required since the photon energy lies far below the energy threshold of the reactions. In the UV band, $\text{Cu(I)(NO}_3)_2^-$ is clearly dominant. $\text{Cu(I)(NO}_3)_2\text{O}^-$ is most likely

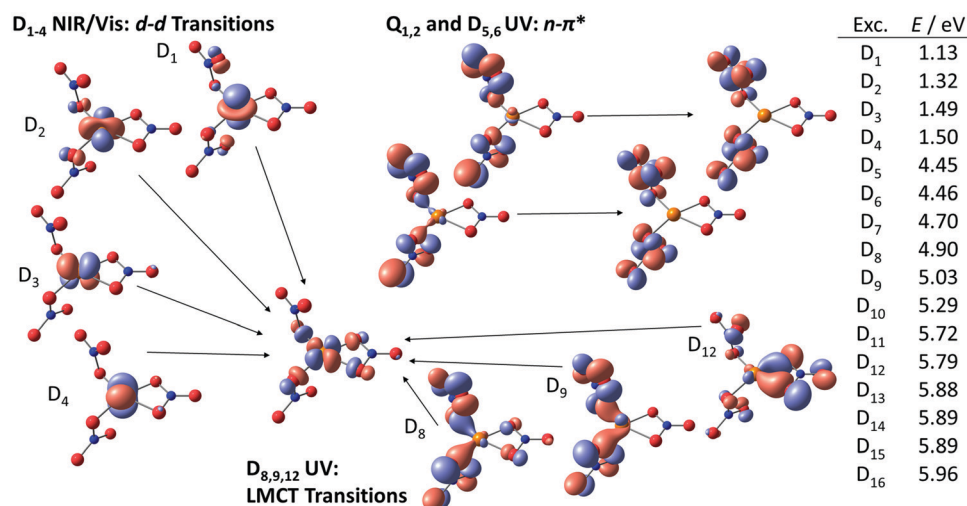


Fig. 3 The original and final electron orbitals are shown for the first six electronic excitations along with the most intense UV transitions below 6.0 eV in $\text{Cu(II)(NO}_3)_3^-$. Natural transition orbitals are calculated at the TD-BMK/aug-cc-pVDZ//B3LYP/def2TZVP level of theory. For nearly degenerate transitions, orbitals of different symmetry mix; then, only the most important component is shown.



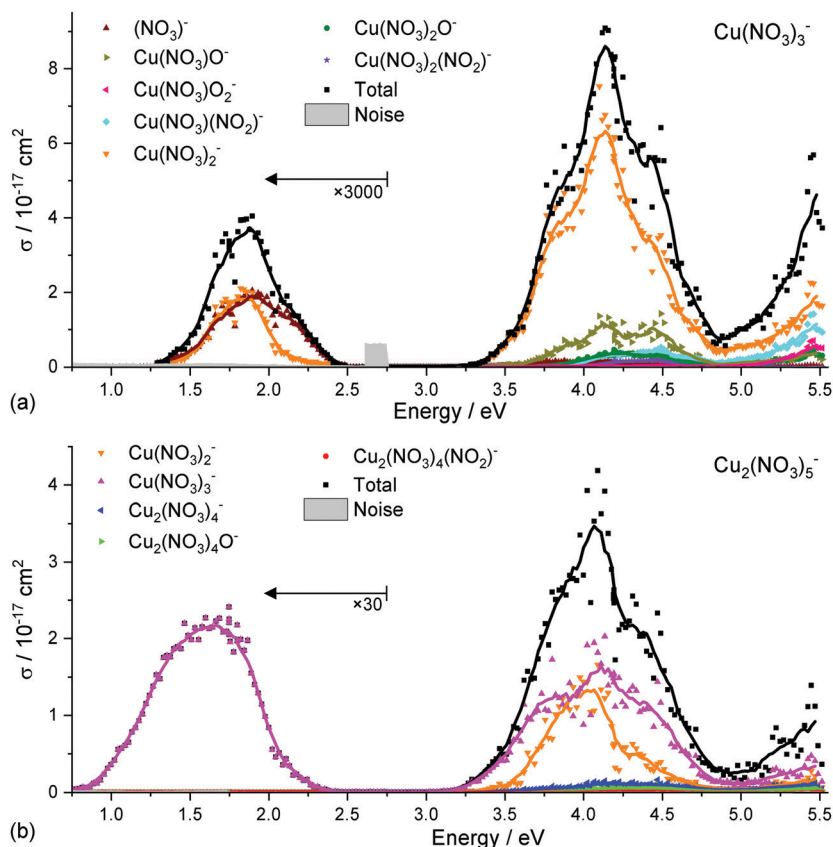


Fig. 4 The experimental spectra of $\text{Cu(II)(NO}_3)_3^-$ in (a) and $\text{Cu(II)}_2(\text{NO}_3)_5^-$ in (b) as total dissociation cross section plus selected partial cross sections. A running average of eight data points is used in addition to a multiplication factor of 3000 and 30, respectively, below 2.75 eV for visibility.

Table 1 Appearance energy E_A in eV of the decomposition channels shown in Fig. 3 when the running average of the partial cross section surpasses the noise level along with the calculated theoretical reaction energy E_{theo} . Calculated at the B3LYP/def2TZVP level of theory

Reaction	Reactant	Products	E_A/eV	$E_{\text{theo}}/\text{eV}$
(1)	$\text{Cu(NO}_3)_3^-$	$\text{NO}_3^- + \text{Cu(NO}_3)_2$	1.3	2.00
(2)		$\text{Cu(NO}_3)_2^- + \text{NO}_3$	1.4	2.05
(3a)		$\text{Cu(NO}_3)_2\text{O}^- + \text{NO}_3 + \text{NO}_2$	3.3	4.78
(3b)		$\text{Cu(NO}_3)_2\text{O}^- + \text{N}_2\text{O}_5$	3.3	4.09
(4)		$\text{Cu(NO}_3)_2\text{O}^- + \text{NO}_2$	3.5	2.57
(5)		$\text{Cu(NO}_3)_2(\text{NO}_2)^- + \text{O}(^3\text{P}_2)$	3.9	3.39
(6)		$\text{Cu(NO}_3)_2\text{O}_2^- + 2(\text{NO}_2)$	3.4	3.82
(7)		$\text{Cu(NO}_3)(\text{NO}_2)^- + \text{NO}_3 + \text{O}(^3\text{P}_2)$	3.7	5.52
(8)	$\text{Cu}_2(\text{NO}_3)_5^-$	$\text{Cu(NO}_3)_3^- + \text{Cu(NO}_3)_2$	0.7	0.83
(9a)		$\text{Cu(NO}_3)_2^- + \text{Cu(NO}_3)_2 + \text{NO}_3$	3.2	2.87
(9b)		$\text{Cu(NO}_3)_2^- + \text{Cu(NO}_3)_3$	3.2	2.58
(10)		$\text{Cu}_2(\text{NO}_3)_4^- + \text{NO}_3$	3.3	1.85
(11)		$\text{Cu}_2(\text{NO}_3)_4\text{O}^- + \text{NO}_2$	3.8	1.85
(12)		$\text{Cu}_2(\text{NO}_3)_4(\text{NO}_2)^- + \text{O}(^3\text{P}_2)$	4.0	3.33

formed sequentially *via* reaction (3a) after absorption of a second photon, as the photon energy is not sufficient for reaction (3b) at the start of the band.

$\text{Cu(II)}_2(\text{NO}_3)_5^-$ decomposes in the 3d–3d band exclusively to $\text{Cu(II)(NO}_3)_3^-$ *via* the energetically most favorable reaction (8) by loss of neutral $\text{Cu(II)(NO}_3)_2$, requiring only 0.83 eV. In the UV band, however, this product exhibits significant sequential fragmentation into $\text{Cu(I)(NO}_3)_2^-$ *via* reaction (9a), afforded by

the excess photon energy and potentially also secondary photo-fragmentation of the dominant $\text{Cu(II)(NO}_3)_3^-$ product. Direct loss of $\text{Cu(III)(NO}_3)_3$ is also possible with the available photon energy *via* reaction (9b).

It is noteworthy that the copper(I) nitrate products can sequentially decompose predominantly *via* NO_2 loss leading towards the copper oxide CuO_2^- or even Cu_2O_3^- while copper(II) nitrate predominantly loses a neutral nitrate radical or copper nitrate, see Fig. S1–S4 (ESI[†]) for a full fragment list across all investigated clusters. For clusters with a single copper center, these decomposition pathways compete with the direct loss of the negatively charged nitrate ligand. The predominantly observed decomposition pathways of $\text{Cu(II)(NO}_3)_3^-$ are matching the fragmentation observed in CID experiments, with $\text{Cu(I)(NO}_3)_2\text{O}^-$ followed by Cu(II)O_2^- as sequential products of initial $\text{Cu(I)(NO}_3)_2^-$ formation, which compete with NO_3^- loss.⁶⁸ Furthermore, the predominant generation of $\text{Cu(I)(NO}_3)_2^-$ in the UV is in agreement with the result from Kaufman and Weber,⁶⁷ who only observed the $\text{Cu(I)(NO}_3)_2^-$ ion with single-pulse irradiation. The predominant fragmentation of $\text{Cu(II)(NO}_3)_3^-$ and $\text{Cu(II)}_2(\text{NO}_3)_5^-$ *via* the energetically most favorable channels points towards fast internal conversion upon photoexcitation, followed by statistical decomposition in the ground state. This decomposition pathway along with the cluster size and oxidation state dependence of the decomposition mirrors the behavior of the previously studied copper formate clusters.^{63,64,66}



The branching ratio of reactions (1) and (2) in the low-energy band of $\text{Cu(II)(NO}_3)_3^-$ in Fig. 4a is 50 : 50 up to ~ 1.7 eV. Towards higher energies, the NO_3^- fragment becomes strongly preferred. Reaction (1) is the energetically most favorable decomposition channel and requires 2.00 eV while reaction (2) is slightly higher at 2.05 eV. The decomposition requires more energy than provided by a photon at the center of the band, pointing towards a multiple-photon process. This explains why the intensity of the low-energy band is significantly lower than in the spectra of larger clusters, when a single photon is sufficient to evaporate a neutral copper nitrate unit. With the excess energy available after two photon excitation, statistical decomposition into the two almost isoenergetic products explains the 50 : 50 branching ratio. The high-energy flank of the vis/NIR band exhibits a clear shoulder, indicating that specific transitions are responsible for the preferential formation of NO_3^- in this range. As discussed above, the D_3 , D_4 excitations result in an increased electron density in the plane of coordination, which leads to a repulsion foremost of the bidentate nitrate ligand. This would tentatively rationalize the preferential formation of NO_3^- as a consequence of the ultrafast dynamics initiated by the excitation.

In the UV bands, however, NO_3^- formation does not play a significant role, the predominant decomposition channel is by far reaction (2) leading to $\text{Cu(I)(NO}_3)_2^-$. This is explained in a straightforward manner by the LMCT character of the excitations in this energy range or by an accessible conical intersection towards these states. The charge transfer states photochemically prepare a $(\text{NO}_3)[\text{Cu(I)(NO}_3)_2]^-$ structure which can then easily decompose *via* reaction (2). Direct nitrate loss can also be observed in low intensities in the UV bands of $\text{Cu(II)}_2(\text{NO}_3)_5^-$,

reaction (10), where it is energetically very unfavorable compared to copper nitrate loss *via* reaction (8).

In the UV bands of $\text{Cu(II)(NO}_3)_3^-$ and $\text{Cu(II)}_2(\text{NO}_3)_5^-$, fragments which correspond to decomposition of the nitrate ligand *via* loss of NO_2 or an oxygen atom are observed in smaller amounts. Namely $\text{Cu(NO}_3)_2\text{O}^-$, $\text{Cu(NO}_3)_2(\text{NO}_2)^-$, $\text{Cu(NO}_3)_2\text{O}_2^-$ and $\text{Cu(NO}_3)(\text{NO}_2)^-$ are observed as products of $\text{Cu(II)(NO}_3)_3^-$, reactions (4)–(7). These product channels were not observed in previous experiments employing single-pulse laser irradiation.⁶⁷ For $\text{Cu(II)}_2(\text{NO}_3)_5^-$, $\text{Cu}_2(\text{NO}_3)_4\text{O}^-$ and $\text{Cu}_2(\text{NO}_3)_4(\text{NO}_2)^-$ are found, reaction (11) and (12), among many more fragments in the larger clusters, see Fig. S1–S4 (ESI[†]) for a comprehensive list of fragments and partial photodissociation spectra.

To investigate these decomposition channels computationally, we followed the two N–O dissociation coordinates for a nitrate ligand of $\text{Cu(II)(NO}_3)_3^-$ in the lowest-lying quartet state Q_1 , $d(\text{Cu(NO}_3)_2\text{O}^- - \text{NO}_2)$ in Fig. 5a and $d(\text{Cu(NO}_3)_2(\text{NO}_2)^- - \text{O})$ in Fig. 5b. The lowest-lying excited states in the UV region are the $D_{5,6}$ and $Q_{1,2}$ states which exhibit a biradical character. This allows a relaxed PES scan in the lowest lying state of the quartet manifold to directly investigate the photochemistry of the nitrate ligand in the excited state, without employing computationally expensive excited state methods like EOM-CCSD. In the quartet state, one nitrate ligand has two unpaired electrons and relaxes to a non-planar structure, which corresponds to a local minimum on the Q_1 surface as depicted in Fig. 5. Due to the biradical character of the NO_3^- ligand, the N–O bonds are slightly lengthened by up to about 0.1 Å. This local minimum lies substantially below the excitation energy in the FC point, red triangle down in Fig. 5. The relaxation from the FC point to

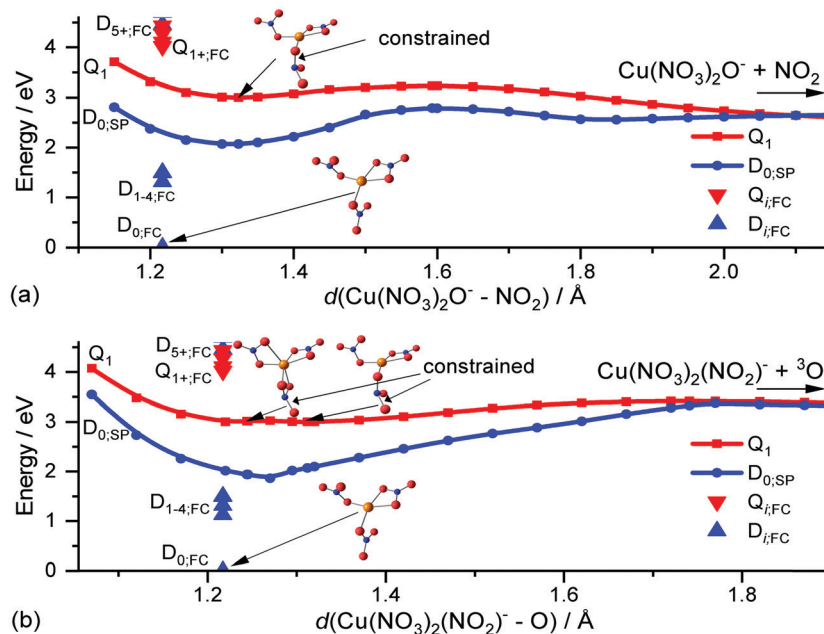


Fig. 5 Relaxed potential energy surface scan of $\text{Cu(II)(NO}_3)_3^-$ in the Q_1 state and the corresponding single-point energies of D_0 , $D_{0,SP}$, on the B3LYP/def2TZVP level of theory b-splined along the decomposition coordinate of a nitrate ligand; (a) $d(\text{Cu(NO}_3)_2\text{O}^- - \text{NO}_2)$ and (b) $d(\text{Cu(NO}_3)_2(\text{NO}_2)^- - \text{O})$. Additionally, the energy of the ground state D_0 in the FC point and the corresponding vertical excitation energies $Q_{i,FC}$ and $D_{i,FC}$ on the TD-BMK/aug-cc-pVDZ//B3LYP/def2TZVP level of theory is shown along with structures in the Q_1 minima and the FC point.



the local minimum thus releases significant excess energy, which is more than sufficient to overcome the minor energetic barriers of about 0.2 eV and 0.4 eV along the dissociation coordinates in Fig. 5a and b, leading to NO₂ and O loss, respectively.

Therefore, after an excitation in the UV, a nitrate ligand can decompose *via* reaction (4) and (5), respectively, if the excited electron ends up in its π^* orbital, either directly *via* an $n\text{-}\pi^*$ transition or *via* internal conversion through one or more conical intersections following an LMCT transition. Alternatively, the photochemical loss of the nitrate ligand as a neutral radical is possible if the electron is instead transferred to the copper center. As a third option, the ion can relax through a conical intersection into the manifold of the D₀₋₄ states, where all nitrate ligands are in their ground state configuration and the excitation is localized on the Cu(II) center. This leads eventually to statistical decomposition on the ground state potential energy surface. The latter two decomposition pathways are very similar to the Cu(II) formate system, where fast internal conversion into the ground state has been observed in addition to direct loss of the formate ligand as formoxyl radical after a LMCT excitation.⁶⁶ Furthermore, the observed decomposition of the nitrate ligand along an N–O bond is in good agreement with the theoretically predicted decomposition channels of NO₃[−] into reactive NO₂[−] or O[−] anions.²⁹ Thus, excitation of copper nitrate in the UV is also an indirect probe of the photochemistry of the nitrate anion. Similar N–O bond photolysis may occur in atmospherically more relevant species containing nitrate, like nitric acid trihydrate aerosols^{85,86} or sea-salt aerosols with traces of nitrate.⁸⁷

Conclusion

The observed vis/NIR band of Cu(II)_n(NO₃)_{2n+1}[−], $n \leq 4$, up to 2.5 eV corresponds to a 3d–3d band, consistent with a [Ar]3d⁹ electronic configuration of the copper centers without copper–copper bonding interactions. For $n \geq 2$, two Cu(II) centers are bridged by a bidentate nitrate ligand. Two bands in the UV, starting around 3.25 eV, arise from a complex interaction of the excited states involving a LMCT character and $n\text{-}\pi^*$ transitions of the nitrate ligands. Since the transitions are strongly localized, the spectra of all clusters are very similar, especially for $n \geq 2$. The only pronounced exception is the weakness of the d–d band for $n = 1$, which occurs because two photons are required for dissociation in this case.

The photochemistry is dominated by fast internal conversion followed by statistical decomposition in the ground state. Photochemical loss of a nitrate radical occurs through LMCT excitations in the UV or internal conversion from higher lying states into the LMCT states. The fragmentation of copper nitrate in the ground state is highly size and oxidation state dependent. While large clusters predominantly lose neutral Cu(II) nitrate, nitrate radical evaporation is competitive in small clusters, *e.g.* Cu(II)(NO₃)₃[−] to Cu(I)(NO₃)₂[−]. For Cu(I) nitrate clusters, decomposition of the nitrate ligands results in copper nitrate–oxide mixtures, which can eventually lead to pure copper oxides like Cu_mO_{m+1}[−] ($m = 1, 2$) in sequential photolysis steps.

A very intriguing feature in the photochemistry of copper nitrate clusters is that N–O bond photolysis occurs after population of a π^* orbital in the UV bands. These channels lead to the decomposition of the corresponding NO₃[−] ligand and the formation of reactive species.

Author contributions

Tobias F. Pascher: data curation; formal analysis; investigation; project administration; visualization; writing – original draft; Milan Ončák: conceptualization; methodology; software; supervision; validation; writing – review & editing; Christian van der Linde: data curation; methodology; project administration; visualization; writing – review & editing; Martin K. Beyer: conceptualization; funding acquisition; software; supervision; writing – review & editing.

Data availability

The data that supports the findings of this study are available within the article and its ESI.†

Conflicts of interest

There are no conflicts of interest to declare.

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References

- 1 D. Delimaris and T. Ioannides, *Appl. Catal., B*, 2009, **89**, 295.
- 2 S. Royer and D. Duprez, *ChemCatChem*, 2011, **3**, 24.
- 3 M. S. Kamal, S. A. Razzak and M. M. Hossain, *Atmos. Environ.*, 2016, **140**, 117.
- 4 G. Xie, Z. Liu, Z. Zhu, Q. Liu, J. Ge and Z. Huang, *J. Catal.*, 2004, **224**, 42.
- 5 P. Granger and V. I. Parvulescu, *Chem. Rev.*, 2011, **111**, 3155.
- 6 C. E. Elwell, N. L. Gagnon, B. D. Neisen, D. Dhar, A. D. Spaeth, G. M. Yee and W. B. Tolman, *Chem. Rev.*, 2017, **117**, 2059.
- 7 M. B. Gawande, A. Goswami, F.-X. Felpin, T. Asefa, X. Huang, R. Silva, X. Zou, R. Zboril and R. S. Varma, *Chem. Rev.*, 2016, **116**, 3722.
- 8 A. J. Jordan, G. Lalic and J. P. Sadighi, *Chem. Rev.*, 2016, **116**, 8318.
- 9 H. Häkkinen, M. Moseler and U. Landman, *Phys. Rev. Lett.*, 2002, **89**, 33401.
- 10 A. Álvarez, A. Bansode, A. Urakawa, A. V. Bavykina, T. A. Wezendonk, M. Makkee, J. Gascon and F. Kapteijn, *Chem. Rev.*, 2017, **117**, 9804.

- 11 K. C. Waugh, *Catal. Lett.*, 2012, **142**, 1153.
- 12 J. Artz, T. E. Müller, K. Thenert, J. Kleinekorte, R. Meys, A. Sternberg, A. Bardow and W. Leitner, *Chem. Rev.*, 2018, **118**, 434.
- 13 A. Goepfert, M. Czaun, J.-P. Jones, G. K. Surya Prakash and G. A. Olah, *Chem. Soc. Rev.*, 2014, **43**, 7995.
- 14 M. Grasemann and G. Laurenczy, *Energy Environ. Sci.*, 2012, **5**, 8171.
- 15 D. Mellmann, P. Sponholz, H. Junge and M. Beller, *Chem. Soc. Rev.*, 2016, **45**, 3954.
- 16 R. S. Dhayal, W. E. van Zyl and C. W. Liu, *Dalton Trans.*, 2019, **48**, 3531.
- 17 S. Fujita, S. Moribe, Y. Kanamori, M. Kakudate and N. Takezawa, *Appl. Catal., A*, 2001, **207**, 121.
- 18 R. Prasad and P. Singh, *Catal. Sci. Technol.*, 2013, **3**, 3326.
- 19 J. M. Dumas, C. Geron, A. Kribii and J. Barbier, *Appl. Catal.*, 1989, **47**, L9–L15.
- 20 R. Prasad and P. Singh, *Bull. Chem. React. Eng. Catal.*, 2011, **6**, 63.
- 21 M. Gao, R. Ye, W. Shen and B. Xu, *Org. Biomol. Chem.*, 2018, **16**, 2602.
- 22 J. Schuttlefield, G. Rubasinghege, M. El-Maazawi, J. Bone and V. H. Grassian, *J. Am. Chem. Soc.*, 2008, **130**, 12210.
- 23 A. M. Grannas, A. E. Jones, J. Dibb, M. Ammann, C. Anastasio, H. J. Beine, M. Bergin, J. Bottenheim, C. S. Boxe, G. Carver, G. Chen, J. H. Crawford, F. Dominé, M. M. Frey, M. I. Guzmán, D. E. Heard, D. Helmig, M. R. Hoffmann, R. E. Honrath, L. G. Huey, M. Hutterli, H. W. Jacobi, P. Klán, B. Lefer, J. McConnell, J. Plane, R. Sander, J. Savarino, P. B. Shepson, W. R. Simpson, J. R. Sodeau, R. von Glasow, R. Weller, E. W. Wolff and T. Zhu, *Atmos. Chem. Phys.*, 2007, **7**, 4329.
- 24 H.-W. Jacobi and B. Hilker, *J. Photochem. Photobiol., A*, 2007, **185**, 371.
- 25 R. E. Honrath, M. C. Peterson, S. Guo, J. E. Dibb, P. B. Shepson and B. Campbell, *Geophys. Res. Lett.*, 1999, **26**, 695.
- 26 S. Mai and L. González, *Angew. Chem., Int. Ed.*, 2020, **59**, 16832.
- 27 H. Lischka, D. Nachtigallová, A. J. A. Aquino, P. G. Szalay, F. Plasser, F. B. C. Machado and M. Barbatti, *Chem. Rev.*, 2018, **118**, 7293.
- 28 P. Warneck and C. Wurzinger, *J. Phys. Chem.*, 1988, **92**, 6278.
- 29 O. Svoboda and P. Slavíček, *J. Phys. Chem. Lett.*, 2014, **5**, 1958.
- 30 H. Schwarz, *Coord. Chem. Rev.*, 2017, **334**, 112.
- 31 H. Schwarz, *Catal. Sci. Technol.*, 2017, **7**, 4302.
- 32 H. Schwarz, S. Shaik and J. Li, *J. Am. Chem. Soc.*, 2017, **139**, 17201.
- 33 J. Roithová, *Chem. Soc. Rev.*, 2012, **41**, 547.
- 34 F. Gunzer, S. Krüger and J. Grotemeyer, *Mass Spectrom. Rev.*, 2019, **38**, 202.
- 35 H. Schwarz, P. González-Navarrete, J. Li, M. Schlangen, X. Sun, T. Weiske and S. Zhou, *Organometallics*, 2017, **36**, 8.
- 36 H. Schwarz, *Angew. Chem., Int. Ed.*, 2011, **50**, 10096.
- 37 J. M. Weber, *Int. Rev. Phys. Chem.*, 2014, **33**, 489.
- 38 E. Barwa, M. Ončák, T. F. Pascher, A. Herburger, C. van der Linde and M. K. Beyer, *Chem. – Eur. J.*, 2020, **26**, 1074.
- 39 E. Barwa, T. F. Pascher, M. Ončák, C. van der Linde and M. K. Beyer, *Angew. Chem., Int. Ed.*, 2020, **59**, 7467–7471.
- 40 T. F. Pascher, E. Barwa, C. van der Linde, M. K. Beyer and M. Ončák, *Theor. Chem. Acc.*, 2020, **139**, 127.
- 41 A. Zavras, M. Krstić, P. Dugourd, V. Bonačić-Koutecký and R. A. J. O'Hair, *ChemCatChem*, 2017, **9**, 1298.
- 42 R. A. J. O'Hair, A. Mravak, M. Krstić and V. Bonačić-Koutecký, *ChemCatChem*, 2019, **11**, 2443.
- 43 M. Krstić, Q. Jin, G. N. Khairallah, R. A. J. O'Hair and V. Bonačić-Koutecký, *ChemCatChem*, 2018, **10**, 1173.
- 44 S. Xu, J. E. T. Smith, S. Gozem, A. I. Krylov and J. M. Weber, *Inorg. Chem.*, 2017, **56**, 7029.
- 45 S. Xu, J. E. T. Smith and J. M. Weber, *J. Phys. Chem. A*, 2016, **120**, 7650.
- 46 S. Xu, J. E. T. Smith and J. M. Weber, *J. Chem. Phys.*, 2016, **145**, 24304.
- 47 C. S. Byskov, J. M. Weber and S. B. Nielsen, *Phys. Chem. Chem. Phys.*, 2015, **17**, 5561.
- 48 M. Förstel, K. M. Follow, K. Saroukh, E. A. Najib, R. Mitric and O. Dopfer, *Angew. Chem., Int. Ed.*, 2020, **59**, 21403.
- 49 M. Förstel, W. Schewe and O. Dopfer, *Angew. Chem., Int. Ed.*, 2019, **58**, 3356.
- 50 B. K. A. Jaeger, M. Savoca, O. Dopfer and N. X. Truong, *Int. J. Mass Spectrom.*, 2016, **402**, 49.
- 51 M. A. Ashraf, J. Kozubal and R. B. Metz, *J. Chem. Phys.*, 2018, **149**, 174301.
- 52 M. D. Johnston, M. R. Gentry and R. B. Metz, *J. Phys. Chem. A*, 2018, **122**, 8047.
- 53 W. L. Pearson, C. Copeland, A. Kocak, Z. Sallese and R. B. Metz, *J. Chem. Phys.*, 2014, **141**, 204305.
- 54 J. S. Daluz, A. Kocak and R. B. Metz, *J. Phys. Chem. A*, 2012, **116**, 1344.
- 55 M. Ončák, T. Taxer, E. Barwa, C. van der Linde and M. K. Beyer, *J. Chem. Phys.*, 2018, **149**, 44309.
- 56 T. Taxer, M. Ončák, E. Barwa, C. van der Linde and M. K. Beyer, *Faraday Discuss.*, 2019, **217**, 584.
- 57 N. K. Bersenkovitsch, M. Ončák, C. van der Linde, A. Herburger and M. K. Beyer, *Phys. Chem. Chem. Phys.*, 2018, **20**, 8143.
- 58 N. K. Bersenkovitsch, M. Ončák, J. Heller, C. van der Linde and M. K. Beyer, *Chem. – Eur. J.*, 2018, **24**, 12433.
- 59 N. K. Bersenkovitsch, M. Ončák, J. Heller, T. F. Pascher, C. van der Linde and M. K. Beyer, *Phys. Chem. Chem. Phys.*, 2020, **22**, 12028–12038.
- 60 A. Herburger, E. Barwa, M. Ončák, J. Heller, C. van der Linde, D. M. Neumark and M. K. Beyer, *J. Am. Chem. Soc.*, 2019, **141**, 18000.
- 61 A. Herburger, M. Ončák, C.-K. Siu, E. G. Demissie, J. Heller, W. K. Tang and M. K. Beyer, *Chem. – Eur. J.*, 2019, **25**, 10165.
- 62 M. Srnc, R. Navrátil, E. Andris, J. Jašík and J. Roithová, *Angew. Chem., Int. Ed.*, 2018, **57**, 17053.
- 63 T. F. Pascher, M. Ončák, C. van der Linde and M. K. Beyer, *ChemPhysChem*, 2019, **20**, 1420.
- 64 T. F. Pascher, M. Ončák, C. van der Linde and M. K. Beyer, *ChemistryOpen*, 2019, **8**, 1453–1459.
- 65 T. F. Pascher, M. Ončák, C. van der Linde and M. K. Beyer, *J. Chem. Phys.*, 2020, **153**, 184301.



- 66 T. F. Pascher, M. Ončák, C. van der Linde and M. K. Beyer, *Chem. – Eur. J.*, 2020, **26**, 8286–8295.
- 67 S. H. Kaufman and J. M. Weber, *J. Phys. Chem. A*, 2014, **118**, 9687.
- 68 C. J. M. Pruitt and D. J. Goebbert, *J. Phys. Chem. A*, 2015, **119**, 4755.
- 69 R. C. Dunbar, *Mass Spectrom. Rev.*, 2004, **23**, 127.
- 70 T. Schindler, C. Berg, G. Niedner-Schatteburg and V. E. Bondybey, *Chem. Phys. Lett.*, 1996, **250**, 301.
- 71 W. D. Price, P. D. Schnier, R. A. Jockusch, E. F. Strittmatter and E. R. Williams, *J. Am. Chem. Soc.*, 1996, **118**, 10640.
- 72 D. Thölmann, D. S. Tonner and T. B. McMahon, *J. Phys. Chem.*, 1994, **98**, 2002.
- 73 M. Sena and J. M. Riveros, *Rapid Commun. Mass Spectrom.*, 1994, **8**, 1031.
- 74 B. S. Fox, M. K. Beyer and V. E. Bondybey, *J. Phys. Chem. A*, 2001, **105**, 6386.
- 75 A. Herburger, C. van der Linde and M. K. Beyer, *Phys. Chem. Chem. Phys.*, 2017, **19**, 10786.
- 76 H. Koch and P. Jørgensen, *J. Chem. Phys.*, 1990, **93**, 3333.
- 77 J. F. Stanton and R. J. Bartlett, *J. Chem. Phys.*, 1993, **98**, 7029.
- 78 R. L. Martin, *J. Chem. Phys.*, 2003, **118**, 4775.
- 79 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery Jr, J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, *Gaussian 16, Revision A.03*, Gaussian Inc., Wallingford CT, 2016.
- 80 D. Hollas, O. Svoboda and M. Ončák, ABIN, Source code available at <https://github.com/PHOTOX/ABIN>.
- 81 S. Y. Lee, R. C. Brown and E. J. Heller, *J. Phys. Chem.*, 1983, **87**, 2045.
- 82 F. Neese, T. Petrenko, D. Ganyushin and G. Olbrich, *Coord. Chem. Rev.*, 2007, **251**, 288.
- 83 M. Gerloch and E. C. Constable, *Transition Metal Chemistry: The Valence Shell in d-Block Chemistry*, Wiley-VCH, Weinheim, New York, 1994.
- 84 T. Schwab, M. Niedermaier, G. A. Zickler, M. Ončák and O. Diwald, *Chem. – Eur. J.*, 2020, **26**, 16049.
- 85 F. Weiss, F. Kubel, Ó. Gálvez, M. Hoelzel, S. F. Parker, P. Baloh, R. Iannarelli, M. J. Rossi and H. Grothe, *Angew. Chem., Int. Ed.*, 2016, **55**, 3276.
- 86 A. K. Bertram and J. J. Sloan, *J. Geophys. Res.*, 1998, **103**, 13261.
- 87 P. Kasibhatla, T. Sherwen, M. J. Evans, L. J. Carpenter, C. Reed, B. Alexander, Q. Chen, M. P. Sulprizio, J. D. Lee, K. A. Read, W. Bloss, L. R. Crilley, W. C. Keene, A. A. P. Pszenny and A. Hodzic, *Atmos. Chem. Phys.*, 2018, **18**, 11185.

