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# BiMnPO<sub>5</sub> with ferromagnetic Mn<sup>2+</sup>–(μ-O)<sub>2</sub>–Mn<sup>2+</sup> units: a model for magnetic exchange in edge-linked Mn<sup>2+</sup>O<sub>6</sub> octahedra†

Rukang Li,<sup>†</sup> Ronald I. Smith<sup>‡</sup> and Colin Greaves<sup>‡</sup>

**Magnetic interactions within Mn–(μ-O)<sub>2</sub>–Mn pairs are crucial to the function of some essential enzymes and catalysts, but their nature is unclear. Neutron diffraction reveals that similar units in BiMnPO<sub>5</sub> show ferromagnetic coupling which has been rationalized by density functional theory modelling and calculations of magnetic exchange energies. The results are important to many solid state and biological systems.**

Magnetic Mn–O–Mn interactions play an important role in solid state, organometallic and biological chemistry. Perovskite related LaMnO<sub>3</sub> derivatives show a giant magneto-resistance (GMR) effect and AMnO<sub>3</sub> (A = Bi, Y) show multiferroic properties, which are fundamentally important for the new generation of integrated electronics.<sup>1</sup> The interaction is also crucial to the production of oxygen by the manganese containing enzyme photosystem II (PSII), hence the beginning of life on earth.<sup>2</sup> PSII has therefore initiated much research on itself<sup>3</sup> and the synthesis of organometallic biomimetic catalysts with multinuclear Mn–O–Mn cores as the active sites for splitting H<sub>2</sub>O.<sup>4</sup>

The properties of manganese oxides and complexes, especially magnetic, electric or catalytic, are related to local structures, oxidation states and spin states of the ions. For example, the functional Mn<sub>4</sub>CaO<sub>5</sub> core for photosynthesis in PSII involves spin-dependent states and it is postulated that separate high- and low-spin states control the formation of the active high-spin state.<sup>5</sup>

Such spin states are controlled by direct Mn–Mn interactions or superexchange *via* bridging oxygen ligands. The nature of exchange interactions of 3d metal ions can often be explained by the Goodenough–Kanamori–Anderson (GKA) rules;<sup>6</sup> *e.g.* for Mn(II) (d<sup>5</sup>), the direct exchange involves overlap of half-filled d orbitals and is therefore always antiferromagnetic (AFM), whereas the exchange through the bridging O is strongly dependent on the Mn–O–Mn bridging angle  $\theta$ . The simplest superexchange is ferromagnetic (FM) and occurs from the accidental orthogonality of 2p orbitals on the O bridges which is a maximum at  $\theta = 90^\circ$ . As a result, the overall exchange depends on local structure, and predictions are hindered by the existence of more complex exchange mechanisms.

Some enzyme research suffers significantly from the complexity of Mn–O–Mn magnetic interactions. Simple Mn–(μ-O)<sub>2</sub>–Mn pairs are found in enzymes (*e.g.* arginase, catalase, and Ribonucleotide Reductase (RNR)), which are found in most living organisms including humans.<sup>7</sup> Various studies on these and their biomimetic counterparts have greatly enhanced our understanding of their structure and function.<sup>8</sup> Magnetically, binuclear Mn(II, II) pairs generally show AFM coupling to give an  $S = 0$  unit, but some synthetic analogues with various ligands show FM behaviour.<sup>9</sup> Complementary experimental and theoretical studies have probed the exchange mechanisms within the pairs but no reliable roadmap exists to allow reliable predictions.<sup>9h</sup>

Our studies of materials containing MnO chains with FM interactions drew our attention to the current lack of appropriate signposts to the resultant magnetic order. Although Mn<sup>2+</sup> is not widespread in solid oxides, we became aware that BiMnPO<sub>5</sub> contains Mn<sub>2</sub><sup>2+</sup>O<sub>10</sub> pairs formed by edge-shared MnO<sub>6</sub> octahedra,<sup>10</sup> which therefore closely resemble the units in the biological systems. Given that the Mn–O–Mn angle (100.9°) is sufficiently greater than 90° to weaken the FM superexchange, time-of-flight (TOF) neutron powder diffraction (NPD) surprisingly revealed that the Mn ions within each Mn<sub>2</sub>O<sub>10</sub> pair have FM coupling. Here, the NPD results are reported alongside density functional theory (DFT) calculations (with a hybrid functional) for a model molecule Mn<sub>2</sub>O<sub>2</sub>(H<sub>2</sub>O)<sub>8</sub>; importantly,

<sup>a</sup> Beijing Centre for Crystal Research and Development, Key Laboratory of Functional Crystals and Laser Technology, Technical Institute of Physics and Chemistry, Chinese Academy of Sciences, Beijing 100190, China

<sup>b</sup> Center of Materials Science and Optoelectronics Engineering, University of Chinese Academy of Sciences, Beijing 100049, China

<sup>c</sup> ISIS Pulsed Neutron and Muon Source, Rutherford Appleton Laboratory, Harwell Campus, Didcot OX11 0QX, UK

<sup>d</sup> School of Chemistry, University of Birmingham, Birmingham B15 2TT, UK. E-mail: c.greaves@bham.ac.uk

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**Table 1** Structural parameters of BiMnPO<sub>5</sub> from NPD data at ambient temperature

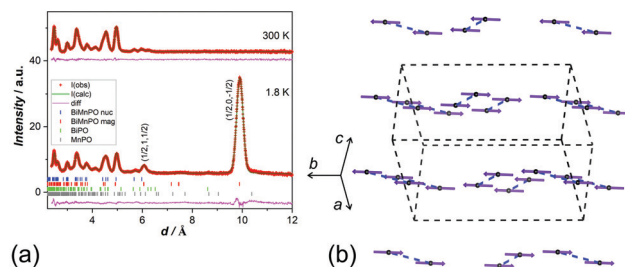
Atoms	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>	<i>U</i> <sub>iso</sub> /Å <sup>2</sup>
Bi	0.18688(6)	0.09468(4)	0.11839(8)	0.00686(9)
Mn	0.809(2)	0.09219(9)	0.3685(2)	0.0049(2)
P	0.0168(1)	0.35033(6)	0.2125(1)	0.0058(1)
O1	0.16341(9)	0.42087(4)	0.4332(1)	0.0093(1)
O2	0.9869(1)	0.22847(6)	0.3163(1)	0.0131(2)
O3	0.83079(9)	0.41908(7)	0.1083(1)	0.0119(1)
O4	0.9855(1)	0.02294(4)	0.7530(1)	0.0072(1)
O5	0.0947(1)	0.32697(6)	0.9830(2)	0.0130(2)

Space group: *P*<sub>2</sub><sub>1</sub>/*n*, *a* = 7.43229 (8), *b* = 11.3559 (1), *c* = 5.35785 (5) Å, β = 109.189(1)° with χ<sup>2</sup> = 0.93.

we propose a phenomenological mechanism to explain this observation and suggest that it could enhance understanding of Mn–Mn interactions in other oxides and also biological systems.

NPD data from a sample of BiMnPO<sub>5</sub> (details in ESI†) at 300 K were used for Rietveld refinement<sup>11</sup> and indicated impurities of ~18 wt% of Bi<sub>2.76</sub>O<sub>2</sub>[PO<sub>3.64</sub>][PO<sub>3.5</sub>]<sup>12</sup> and ~8 wt% of β'-Mn<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub>.<sup>13</sup> Because of the magnetic properties of the latter, magnetic characterization of the primary phase required the unique phase-specific abilities of neutron diffraction. All three phases were therefore included in the refinements, which included data from three discrete detector banks with average 2θ positions of 145°, 35° and 14°. Table 1 lists the refined parameters and Fig. 1(a) shows the structure containing Mn<sub>2</sub>O<sub>10</sub> pairs linked by PO<sub>4</sub> groups; each pair is formed from two identical MnO<sub>6</sub> octahedra, edge-linked by O4 (Fig. 1(b)). The mid-point between the two Mn<sup>2+</sup> ions in each pair is an inversion centre so the ions are linked by symmetric short/long bonds of *a* = 2.192 Å, *b* = 2.263 Å, with angles of θ<sub>1</sub> = θ<sub>2</sub> = 100.9°; the Mn–Mn separation is 3.435 Å. The symmetry and structural parameters are important for the computational considerations (*vide infra*).

Fig. 2(a) shows a very strong (1/2, 0, −1/2) magnetic peak at 1.8 K, which critically shows that the high-spin Mn<sup>2+</sup> ions have FM coupling within each pair; AFM coupling between the pairs gives an enlarged unit cell. The magnetic moments, 4.46(2) μ<sub>B</sub> at 1.8 K, are aligned with the unique *b*-axis; the order is shown in Fig. 2(b) with further information provided in Table 2 and



**Fig. 2** (a) Fitted NPD patterns from BiMnPO<sub>5</sub> collected at 300 K and 1.8 K in the 2θ = 14° detector bank. The 300 K data are offset by ~40 units for clarity; (b) magnetic order with reference to the nuclear unit cell (*P*<sub>2</sub><sub>1</sub>/*n*) showing moments aligned along ±[010] with FM coupling within the Mn–Mn pairs; FM layers perpendicular to [10–1] are AFM to adjacent layers and enlarge the unit cell.

**Table 2** Structural parameters of BiMnPO<sub>5</sub> from NPD data at 1.8 K

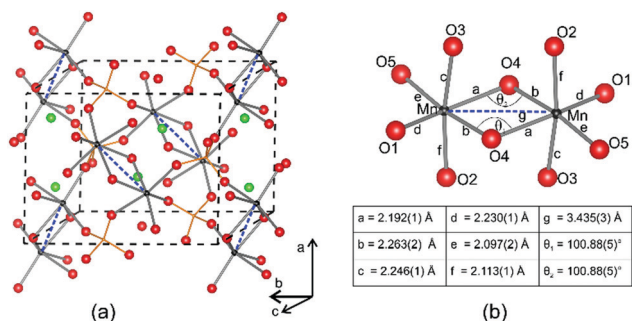
Atoms	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>	<i>U</i> <sub>iso</sub> /Å <sup>2</sup>
Bi	0.18567(1)	0.09554(7)	0.1197(1)	0.0021(2)
Mn	0.8102(2)	0.0912(2)	0.3688(3)	0.00032(3)
P	0.0189(2)	0.3496(1)	0.2103(2)	0.0030(2)
O1	0.1633(1)	0.4220(1)	0.4330(2)	0.0034(2)
O2	0.9887(2)	0.22688(8)	0.3131(2)	0.0053(2)
O3	0.8305(1)	0.4179(1)	0.1075(2)	0.0036(2)
O4	0.9850(2)	0.02305(9)	0.7524(2)	0.0031(2)
O5	0.0938(2)	0.32842(9)	0.9778(2)	0.0054(2)

Space group: *P*<sub>2</sub><sub>1</sub>/*n*, *a* = 7.4456(1), *b* = 11.3408(2), *c* = 5.3532(1) Å, β = 109.250(1)° with χ<sup>2</sup> = 1.19; wR<sub>p</sub> = 0.0160 (35° detector), 0.0062 (145° detector), 0.0206 (14° detector). Magnetic model: magnetic space

group *P*<sub>a</sub>2<sub>1</sub>/*c* with transformation matrix  $\begin{pmatrix} 1 & 0 & -1 \\ 0 & 1 & 0 \\ 1 & 0 & 1 \end{pmatrix}$  to give *a*' = 10.5060(2), *b*' = 11.3408(2), *c*' = 7.6034(1) Å, β = 109.588(1)°. Moment along [010] with μ = 4.46(2) μ<sub>B</sub> on Mn at (0.9707, 0.9088, 0.1605) and μ = −4.46(2) μ<sub>B</sub> on Mn at (0.4707, 0.9088, 0.6605).

ESI†, especially Fig. S1 and S2 (ESI†). The magnetic unit cell (symmetry *P*<sub>a</sub>2<sub>1</sub>/*c*) is derived from the standard setting (*P*<sub>2</sub><sub>1</sub>/*c*) of *P*<sub>2</sub><sub>1</sub>/*n* that is reported for BiMnPO<sub>5</sub>.<sup>10</sup> As previously indicated, most Mn–(μ-O)<sub>2</sub>–Mn interactions are AFM, but NPD definitively shows that in BiMnPO<sub>5</sub> they are FM. There are no reports of clear FM coupling in similar pairs in the solid state, although MnV<sub>2</sub>O<sub>6</sub> has infinite FM chains of edge-linked MnO<sub>6</sub> octahedra.<sup>14</sup> In order to rationalize the magnetic order, we have modelled the Mn<sub>2</sub>O<sub>10</sub> pair in BiMnPO<sub>5</sub> with the hypothetical molecule Mn<sub>2</sub>O<sub>2</sub>(H<sub>2</sub>O)<sub>8</sub>, shown in Fig. 3(a), using CRYSTAL 17 code with a hybrid functional (details in ESI†).<sup>15</sup> Although a similar approach has been employed previously,<sup>16</sup> in this case the geometry was constrained: all Mn–O bonds were held equal at approximately the average length in BiMnPO<sub>5</sub>, and for this ideal model only the Mn–Mn separation was varied (and hence the correlated Mn–O4–Mn angle in Fig. 1(b)) to examine the effect of direct magnetic exchange on the final ordered state.

The calculations revealed that the Mn–Mn exchange changes from AFM at smaller Mn–Mn separations to FM with a maximum FM exchange *J* at 3.4 Å (102°), and a trend to return to AFM at Mn–Mn > 3.7 Å (115°), Fig. 3(b). More importantly, it predicts that at the experimental Mn–Mn separation of 3.435 Å in



**Fig. 1** (a) The structure of BiMnPO<sub>5</sub> and (b) structural details for the Mn<sub>2</sub>O<sub>10</sub> units. Spheres represent Bi (green), Mn (black), P (orange), O (red); Mn–Mn interactions are blue dotted lines.



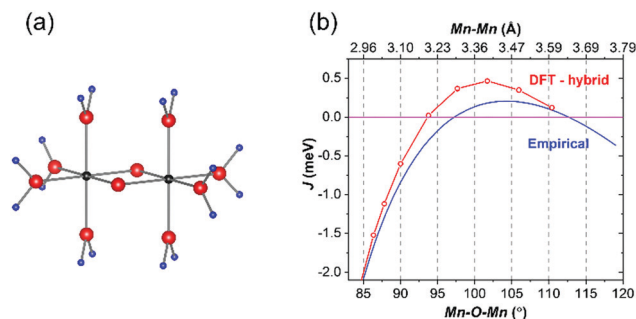


Fig. 3 (a) The  $\text{Mn}_2\text{O}_2(\text{H}_2\text{O})_8$  molecule (Mn black, O red, H blue) used to model the exchange energies for  $\text{Mn}_2\text{O}_{10}$  pairs in  $\text{BiMnPO}_5$ ; (b) the calculated overall exchange,  $J$ , against Mn-O-Mn angle and Mn-Mn distance.

$\text{BiMnPO}_5$ , the  $\text{Mn}_2\text{O}_{10}$  pair is almost optimal for FM coupling. In addition, the crossover from FM to AFM exchange at low separation occurs at 3.19 Å ( $93^\circ$ ). This result is highly significant since GKA rules cannot indicate the dominant exchange in this region because the conflicting AFM (direct exchange) and FM (superexchange) interactions are both optimized. Our calculations are supported by data for organic complexes as well as solid state materials, which suggest that more than 75% of Mn-( $\mu\text{-O}$ )<sub>2</sub>-Mn units are AFM. Indeed, our results are consistent with  $\text{MnV}_2\text{O}_6$ <sup>14</sup> and  $\text{MnSb}_2\text{O}_4$ <sup>17</sup> which both contain chains of edge-linked  $\text{MnO}_6$  octahedra where the intra-chain magnetic exchange is dominant. Whereas  $\text{MnV}_2\text{O}_6$  (Mn-Mn = 3.53 Å, Mn-O-Mn =  $105^\circ$ ) is FM,  $\text{MnSb}_2\text{O}_4$  (Mn-Mn = 3.00 Å, Mn-O-Mn =  $89^\circ$ ) is AFM. It is also interesting to note that the FM exchange in the binuclear  $\text{Mn}^{2+}$ -( $\mu\text{-N}$ )<sub>2</sub>- $\text{Mn}^{2+}$  units in  $[\text{Mn}(\text{terpy})(\text{N}_3)_2]$ <sup>18</sup> with Mn-Mn 3.53 Å and Mn-N-Mn =  $104.6^\circ$  is also consistent with the results for the O-bridged units.

In order to rationalize these calculations, we have quantified the significant exchange pathways including the empirical model proposed by Kahn for the case of strict orthogonality.<sup>19</sup> These interactions are given in Fig. S3 (ESI†). In the special case of Mn-( $\mu\text{-O}$ )<sub>2</sub>-Mn with  $90^\circ$  Mn-O-Mn links, strict orthogonality can be fulfilled by appropriate orbital ordering (Fig. S3, ESI†, case III), as observed in  $\text{O}_3\text{Cu}-(\mu\text{-O})_2\text{-VO}_3$ ,<sup>20</sup> inducing FM coupling. Following the proposal by Khomskii *et al.*,<sup>21</sup> we considered the main contributions for the cases shown in Fig. S3 (ESI†) using typical parameters (all/eV):

$$U_{\text{dd}} = 4.5, \quad U_{\text{pp}} = 2, \quad \Delta = 3, \quad t_{\text{pd}\sigma} = 0.9, \quad t_{\text{pd}\pi} = 0.48 \text{ and} \\ t_{\text{dd}} = 0.33, \quad J_{\text{H}}^{\text{d}} = 0.9, \quad J_{\text{H}}^{\text{p}} = 1.2 \quad (1)$$

The three major contributions are from the direct exchange ( $d_{xy}\text{-}d_{xy}$ , case I), Mn-O-Mn  $180^\circ$  superexchange ( $d_{x^2-y^2}\text{-O}2p\text{-}d_{x^2-y^2}$ , case II) and Kahn's strict orthogonality ( $d_{x^2-y^2}\text{-O}2p\text{-}d_{xy}$ , case III) which have components  $J_{\text{dd}} = -24.2$  meV,  $J_{\text{S}}^{\text{O}} = -25.3$  meV and  $J_{\text{K}}^{\pi} = 7.2$  meV, respectively. Cases IV, V, and VI have only minor contributions which tend to cancel each other out. If we further assume rigid Mn-O bonds and consider the relations  $t_{\text{dd}} \sim \frac{r_{\text{d}}^3}{D^5}$  and  $J_{\text{S}}^{\text{O}} = J_{\text{S}}^{\text{O}} \sin^2 \theta + J_{\text{S}}^{\text{O}} \cos^2 \theta$  ( $D$  is the

Mn-Mn pair distance) previously proposed,<sup>22</sup> the total magnetic exchange, as a function of bond angle  $\theta$  is:

$$J^{\theta} = J_{\text{dd}} \left( \frac{\sin 45^\circ}{\sin \theta/2} \right)^{10} + 2 J_{\text{K}}^{\pi} \sin^2 \theta + J_{\text{S}}^{\text{O}} \cos^2 \theta \quad (2)$$

For these parameters, the overall empirical exchange interaction as a function of Mn-Mn distance (and the dependent Mn-O-Mn angle) is shown in Fig. 3(b); the general trend agrees not only with the DFT calculations but also, and more importantly, with the experimental observations.

The empirical model predicts that Mn-( $\mu\text{-O}$ )<sub>2</sub>-Mn pairs with bond angles close to  $90^\circ$  will be subject to strongly AFM exchange rather than FM, which can be attributed to rapidly increasing  $d^5\text{-}d^5$  direct AFM interactions caused by a shortening of the Mn-Mn distance with reduction in bond angle. At higher bond angles, the expansion in Mn-Mn separation reduces the direct orbital overlap ( $\propto D^{-5}$ ) and enhances the significance of the FM interactions *via* the strict orthogonality case III. At much higher bond angles, the  $180^\circ$  AFM superexchange (case II) eventually dominates and Fig. 3(b) shows that FM behaviour is predicted when:

$$3.287 \text{ \AA} (\theta = 97.3^\circ) < D < 3.645 \text{ \AA} (\theta = 112.7^\circ) \quad (3)$$

An incomplete survey of literature reporting magnetic characterization of Mn-( $\mu\text{-O}$ )<sub>2</sub>-Mn pairs shows that the Mn-Mn coupling constants ( $J^{\theta}$ ) generally comply with the above predictions. In recent compilations of dinuclear complexes,<sup>9h,23</sup> rare examples of FM pairs are reported, but no satisfactory explanations are provided. Since previous theoretical calculations for different models have always predicted AFM order, the present study provides an impetus to find new structures with FM pairs and hence develop our understanding of their functionality in biological materials. An extension of our theoretical investigation on the  $\text{Mn}_2\text{O}_2(\text{H}_2\text{O})_8$  model with asymmetric bonding geometries at the  $\text{Mn}_2\text{O}_2$  core shows that Kahn's case III mechanism for FM exchange can tolerate different Mn-O bond lengths for a given bridging O, as long as a center of symmetry exists between the Mn ions (as is the case for  $\text{BiMnPO}_5$ ). To a lesser extent, different bond angles ( $\theta_1$  and  $\theta_2$  in Fig. 1(b)), are also acceptable but this necessarily destroys the inversion center.  $\alpha\text{-MnMoO}_4$  contains unusual FM clusters of four edge-linked  $\text{Mn}^{2+}\text{O}_6$  octahedra, but the symmetry does not allow direct comparison with the results presented here and another mechanism may also be involved.<sup>24</sup> These structural modifications can explain the examples of AFM coupling which violate the general indications of eqn (3). An interesting test for our model is  $\text{KMnBP}_2\text{O}_7(\text{OH})_2$  which has centrosymmetric Mn-( $\mu\text{-O}$ )<sub>2</sub>-Mn units with Mn-Mn = 3.34 Å and Mn-O-Mn =  $98.9^\circ$ .<sup>25</sup> The low temperature magnetic order is AFM but becomes FM in a field of 7 T. It is suggested that the individual  $\text{Mn}_2$  pairs are internally AFM but our modelling suggests the alternative description in which FM  $\text{Mn}_2$  units experience weak external exchange which allows full FM behaviour at 7 T. Low temperature NPD data would be needed to confirm the magnetic structure.

To conclude, the most significant aspect of the structure of  $\text{BiMnPO}_5$  is the presence of Mn-( $\mu\text{-O}$ )<sub>2</sub>-Mn pairs, which also





occur in various enzymes and biomimetic catalysts. It is unusual that the magnetic interaction within the dimanganese(II, II) pair is ferromagnetic, whereas the pairs are antiferromagnetic coupled through Mn–O–P–O–Mn type super–super-exchange. DFT calculations for a model  $\text{Mn}_2\text{O}_2(\text{H}_2\text{O})_8$  molecule suggest that Mn–( $\mu\text{-O}$ )<sub>2</sub>–Mn pairs can be FM only within a certain range of bond lengths and angles. This treatment, however, applies only to highly symmetric units and AFM is dominant where distortions occur. Nevertheless, calculations of the energies of the various exchange pathways provides a rationalization of the observed magnetic behaviour and can be extended to biological and biomimetic systems.

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## Conflicts of interest

There are no conflicts to declare.

## Notes and references

§ Normalized NPD datasets are available from: Greaves, Colin; Li, Rukang (2020), “BiMnPO<sub>5</sub>”, Mendeley Data, VI, DOI: 10.17632/93m5634wv9.1. For raw data, use DOI: 10.5286/ISIS.E.RB15094-3 for 1.8 K; replace 3 with 6 for ambient data.

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