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Understanding the role of conjugation length on the self-assembly behaviour of oligophenyleneethynylenes†

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Oligophenyleneethynylenes (OPEs) are prominent building blocks with exciting optical and supramolecular properties. However, their generally small spectroscopic changes upon aggregation make the analysis of their self-assembly challenging, especially in the absence of additional hydrogen bonds. Herein, by investigating a series of OPEs of increasing size, we have unravelled the role of the conjugation length on the self-assembly properties of OPEs.

Self-assembled structures of π -conjugated systems have become an important class of functional materials with manifold potential applications ranging from optoelectronics to life sciences. Particularly, linear oligomeric π -systems have attracted widespread attention by virtue of their exceptional optical and electronic properties.2 Examples thereof include oligophenylenes (OPs),3 oligophenylenevinylenes (OPVs),4 oligophenyleneethynylenes (OPEs),⁵ oligoanilines (OAs)⁶ and oligothiophenes.7 Among them, OPEs have been particularly employed because of their excellent photophysical and electronic properties, which have been exploited for many interesting applications, such as light emitting diodes (LEDs),8 field-effect transistors (FETs),9 organic photovoltaic devices,10 sensing,11 imaging¹² and thermoelectric materials.¹³ As for many other classes of π -conjugated systems, the optical properties of OPEs can be finely tuned by chemical functionalization.¹⁴

In addition to their well-known functional properties, efforts were also taken to modulate the self-assembly behaviour of OPEs in solution, owing to the importance of supramolecular structures of these systems in organic electronics.¹⁵ However, an inherent disadvantage of the oligomeric OPE backbone, at least compared to other types of dyes, is its comparatively

smaller spectroscopic changes during self-assembly, 16 which makes detailed supramolecular investigations challenging. This can be explained by the smaller, less polar and more conformationally flexible carbon-only aromatic surface of the OPEs, which decreases the association constant compared to typical dyes with fused aromatic rings, heteroatoms or polar substituents. 17 For this reason, additional functional groups for hydrogen bonding are typically attached to the OPEs core in order to enhance their aggregation propensity and, in turn, facilitate their supramolecular analysis by spectroscopy. 16,18

However, despite numerous reports on OPE-based hydrogen-bonded assemblies, it is surprising that no systematic studies are available for OPEs where aromatic interactions are the chief driving force for aggregation. To bridge this knowledge gap, we herein unravel the effect of π -conjugation length on the optical and self-assembly behaviour of OPEs that are primarily governed by aromatic interactions. Our studies allow to unveil the relative contribution of the OPE core to supramolecular polymerization by establishing a structure-selfassembly relationship. Accordingly, we designed and synthesized linear OPEs with different π -length (OPE3-7) to compare the photophysical and self-assembly behaviour in solution (Scheme 1).

The synthesis of OPE3-7 was accomplished using multi-step synthetic protocols based on previous reports as described in the electronic supplementary information (ESI†). All molecules have been characterized using spectroscopic techniques (¹H and ¹³C NMR, FTIR, HRMS (ESI†)).

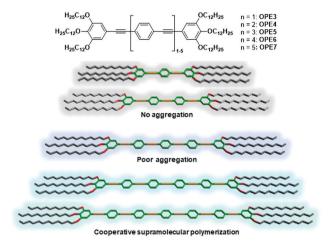
Solvent-dependent UV-Vis and fluorescence studies allowed us to identify chloroform (CHCl₃) and methylcyclohexane (MCH) as the good and poor solvents, respectively (Fig. S1-S5, ESI†). Based on that, the effect of conjugation length on the photophysical properties of the OPEs was monitored by UV-Vis and emission spectroscopy dissolving **OPE3-7** in CHCl₃ ($c = 1 \times 1$ 10⁻⁵ M). As expected, increasing the conjugation length from **OPE3** to **OPE6** induces a progressive red shift in the absorption maximum (from 337 to 363 nm) and a concomitant increase in

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Chemical structure of OPE3-7 (top) and classification of their selfassembly behaviour in solution depending on the conjugation length (bottom).

the molar absorption coefficient (ε) (Fig. 1a). **OPE7** also follows this trend in terms of red shifted absorption compared to OPE6 (Fig. 1a). However, unexpectedly, the absorption intensity of OPE7 is decreased compared to that of OPE6, and a minor shoulder band at longer wavelengths is also noticeable. These results can be ascribed to a weak aggregation of OPE7 because of its larger π -surface compared to OPE6, leading to the observed weakened absorption. Emission studies also showed a red shift in the emission maximum (410 to 428 nm) upon increasing conjugation length from OPE3 to OPE5 (Fig. 1b). This is accompanied by a gradual decrease in the emission intensity (Fig. 1b), possibly due to the higher rotational degrees of freedom upon increasing the number of phenylethynyl units.¹⁹ Further increase in the π -length results in negligible shifts in the emission maximum; however, an additional shoulder band becomes apparent for OPE6-7 at ca. 410 nm (Fig. 1c), which might be the result of short oligomer formation. This solvent- and conjugation length-dependent absorption and

emission properties are also visible to the naked eye, as shown in the photographs of the corresponding solutions (Fig. 1d).

Subsequently, we examined the potential aggregation behaviour of OPE3-7 in MCH. In analogy to the UV-Vis spectra in CHCl₃, a simultaneous red shift and increase in absorption was observed when moving from OPE3 to OPE5 (Fig. 1e), indicating no to poor aggregation in MCH. On the other hand, OPE6 and **OPE7** showed more significant changes, particularly **OPE7**. For this derivative, a broad shoulder band at ca. 410 nm that spreads up to ca. 450 nm appears, which points to aggregation in MCH. Therefore, as expected, increasing the conjugation length enhances the aggregation tendency via aromatic interactions. Further, emission studies for OPE3-7 in MCH were performed under identical experimental conditions. For OPE3-4, a strong and structured emission was observed, indicating a molecularly dissolved state (Fig. 1f). Notably, moving from OPE4 to OPE5 leads to a significant drop in emission intensity, suggesting the initial stages of an aggregation event for OPE5, which does not occur for OPE3,4 under the same experimental conditions. Unsurprisingly, further increasing the π -length to **OPE6** and **OPE7** causes a further emission quenching due to stronger aggregation (Fig. 1f). This phenomenon is further characterized by the appearance of an additional shoulder for OPE6-7 at ca. 540 nm in emission experiments (Fig. 1g). The absence of such shoulder band for OPE5 might be due to only partial or incomplete aggregation in MCH under these conditions. The proposed aggregation for OPE6 and OPE7 in MCH under the selected experimental conditions is also visible by the naked eye, as only the MCH solutions of these two derivatives show a light-yellow color whereas the solutions of OPE3-5 are colorless (Fig. 1h). In accordance with the appreciable luminescence observed in Fig. 1h, OPE3-7 exhibited fluorescence quantum yields $(\Phi_{\rm F})$ of 40-50% in MCH and CHCl₃ (Table S1, ESI†).

Subsequently, we conducted concentration-dependent UV-Vis and fluorescence studies to further examine the selfassembly. As expected, OPE3-5 display no significant changes

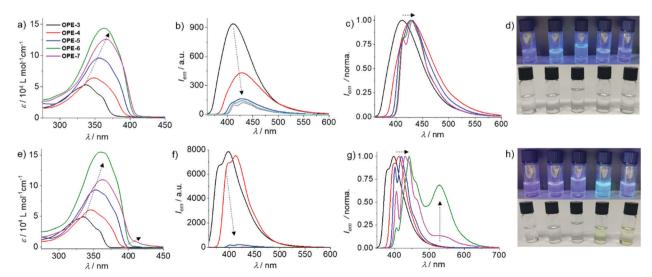


Fig. 1 (a and e) UV-Vis, (b and f) fluorescence (λ_{exc} = 335–365 nm), (c and g) normalized emission spectra and (d and h) photographs under UV and daylight of **OPE3-7** in CHCl₃ (a-d) and MCH (e-h) ($c = 1 \times 10^{-5}$ M, 298 K).

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a)1.00 b) 1.00 0.75 0.75 0.50 0.50 / SQF 0.25 SQF 0.25 0.00 300 350 400 450 500 350 400 450 c) 15 9 E/ 104 L mol-1cm-1 ε/ 10⁴ L mol⁻¹cm⁻¹ 0 6 0 300 350 400 450 500 550 300 350 400 450 500 550 λ / nm f) e) 1 0.8 0.8 0.6 0.6 α_{agg} 0.4 0.2 0.2

Fig. 2 Concentration-dependent UV-Vis spectra of (a) OPE6 and (b) OPE7 in MCH at 298 K. Variable temperature (VT)-UV-Vis spectra of (c) **OPE6** ($c = 1 \times 10^{-3}$ M) and (d) **OPE7** ($c = 3 \times 10^{-4}$ M) in MCH. Global fitting of the cooling curves derived from the corresponding VT experiments of (e) **OPE6** ($c = 1 \times 10^{-3} - 4 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) **OPE7** ($c = 3 \times 10^{-3}$ M, $\lambda = 405$ Nm) and (f) $c = 10^{-3}$ M, $\lambda = 405$ Nm) and (f) $c = 10^{-3}$ M, $\lambda = 10^{-3}$ M, $10^{-4} - 8 \times 10^{-4} \text{ M}, \lambda = 411 \text{ nm}$

290 300

270 280 0

280

290

300

in the absorption spectra upon increasing concentration from 5×10^{-5} M to 1×10^{-3} M (Fig. S6-S8, ESI†). These results support the previous hypothesis that these derivatives exhibit only poor aggregation in MCH, presumably due to the shorter OPE core. In contrast, OPE6 showed clear signs of aggregation in MCH upon increasing the concentration from 2.5×10^{-6} M to 100×10^{-6} M, as evident from the slight red shift in the absorption maximum and the emergence of a shoulder band at around 410 nm (Fig. 2a and Fig. S9a, ESI†). Similarly, a remarkable red shift in emission (440 nm to 530 nm) was observed for OPE6 in MCH upon increasing the concentration from 7.5×10^{-6} M to 1×10^{-3} M (Fig. S9b and c, ESI†). The changes in the absorption and emission studies upon increasing concentration suggest both a slipped and twisted arrangement of the monomers during aggregation.²⁰ Finally, in analogy to OPE6, OPE7 displays a red shifted shoulder when the concentration is increased from 5×10^{-6} M to 500×10^{-6} M, diagnostic of aggregate formation (Fig. 2b and Fig. S10a, ESI†). Notice that above 1×10^{-4} M the aggregate band becomes sharp, which can be due to the precise organization of OPE7 with largest π -system. Additionally, a red shift in the emission was also observed for OPE7 with increasing concentration, indicating a similar aggregation mode as that observed for OPE6 (Fig. S10b, ESI†). However, in contrast to OPE6, OPE7 precipitates over time above 1×10^{-4} M, which is concomitant with a drop in fluorescence, as evident in the photographs shown in Fig. S10c (ESI†).

Further, we performed variable temperature (VT)-UV-Vis experiments for OPE3-7 in MCH to elucidate the thermodynamics of their supramolecular polymerization. As expected, no changes in the UV-Vis spectrum were observed for OPE3 and **OPE4** ($c = 1 \times 10^{-3}$ M) upon varying temperature from 293 K to 263 K (Fig. S11, ESI†). On the other hand, a noticeable broadening of the UV-Vis spectra was observed for OPE5 in MCH under similar experimental conditions, however with no shoulder band formation (Fig. S12a, ESI†). These findings suggest the initial stages of an aggregation event, which may also provide a plausible explanation for the reduced emission intensity of OPE5 compared to OPE3,4 (vide supra). Interestingly, raising the concentration of **OPE5** to 2×10^{-3} M and further to 5×10^{-3} M (where spectral saturation is rather significant) allowed us to identify the red shifted shoulder band at ca. 410 nm that was also characteristic for the aggregates of OPE6 and OPE7 (Fig. S12c and d, ESI†). However, due to the high concentrations required to assemble OPE5 (above ca. 4 mM), a quantitative thermodynamic analysis of this derivative is not possible. VT-1H-NMR experiments were also performed for **OPE3–5** in MCH- d_{14} (c = 5 mM, Fig. S13, ESI†), which, unlike OPE6,7, are sufficiently soluble for these studies. In agreement with previous observations, the π -stacking tendency increases in the order OPE3 \rightarrow OPE4 \rightarrow OPE5, as evident from the more significant broadening of the aromatic signals at low temperatures (263 K-253 K) upon increasing the π -length.

We next examined the self-assembly behaviour of OPE6 by VT-UV-Vis. Notably, upon cooling a hot solution of OPE6 $(c = 1 \times 10^{-3} \text{ M})$ in MCH from 323 K to 263 K, the monomeric band at 360 nm decreases at the expense of the previously described sharp aggregate band at 410 nm (Fig. 2c). Similarly, VT-UV-Vis cooling experiments of OPE7 in MCH also lead to aggregation (Fig. 2d); however, lower concentrations are required ($c = 3 \times 10^{-4}$ M) due to its higher aggregation propensity resulting from the larger π -surface. Distinct UV-Vis spectral changes corresponding to aggregate formation were also observed for OPE6 and OPE7 in VT-UV-Vis at different concentrations (Fig. S14 and S15, ESI,† respectively). The plots of α_{agg} vs. temperature extracted from these experiments were globally fitted to the cooperative (nucleation-elongation) model (Fig. 2e and f).21 Both derivatives show a comparable, yet relatively low degree of cooperativity, which can be explained by the absence of additional non-covalent interactions such as hydrogen bonding.²² This effect also explains the low values of elongation constant (K_e) (in the range of 10^2 – 10^3 M⁻¹) as well as low negative ΔG° values (-8.62 and -16.56 kJ mol⁻¹ for OPE6 and OPE7, respectively; Tables S2 and S3, ESI†). The additional stability of the OPE7 aggregates compared to OPE6 can be ascribed to the increase in conjugation length.

Ultimately, atomic force microscopy (AFM) for both aggregates **OPE6** and **OPE7** at 2.5×10^{-4} M on mica allowed us to visualize entangled networks of fibers with 3-5 nm height and 200-300 nm width (Fig. 3). The fact that the heights are considerably smaller than the molecular length of OPE6,7 suggests a tilted and possibly Communication ChemComm

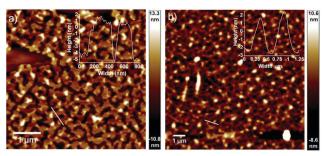


Fig. 3 Height AFM images of (a) **OPE6** and (b) **OPE7** ($c = 2.5 \times 10^{-4}$ M) on mica.

slipped stacking of the OPEs on the substrate. Scanning Electron Microscopy (SEM) on silicon wafer further confirmed the fiber formation of OPE6 and OPE7 (Fig. S18, ESI†).

In conclusion, we have unravelled the relationship between conjugation length and self-assembly behaviour in linear OPEs. Spectroscopic studies demonstrate a red shift both in absorption and emission upon increasing the conjugation length. The aggregation propensity in nonpolar media also follows this trend (OPE3 → OPE7). At mM concentrations, a minimum of five aromatic rings (OPE5) is required to drive aggregation via aromatic interactions. This is most pronounced for the largest derivatives of the series (OPE6,7), leading to emissive fiber-like aggregates. Although the self-assembly mainly occurs by aromatic and solvophobic interactions, both derivatives follow the cooperative mechanism. These results contrast with most selfassemblies that are mainly stabilized by aromatic interactions (without hydrogen bonding), indicating that possible planarization events of the conformationally unrestricted OPE core can turn on cooperative effects. Our studies may contribute to a better understanding of the structure/property relationship in π -conjugated systems.

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Conflicts of interest

There are no conflicts to declare.

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