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## Large transition state stabilization from a weak hydrogen bond†

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A series of molecular rotors was designed to study and measure the rate accelerating effects of an intramolecular hydrogen bond. The rotors form a weak neutral O–H...O=C hydrogen bond in the planar transition state (TS) of the bond rotation process. The rotational barrier of the hydrogen bonding rotors was dramatically lower (9.9 kcal mol<sup>−1</sup>) than control rotors which could not form hydrogen bonds. The magnitude of the stabilization was significantly larger than predicted based on the independently measured strength of a similar O–H...O=C hydrogen bond (1.5 kcal mol<sup>−1</sup>). The origins of the large transition state stabilization were studied *via* experimental substituent effect and computational perturbation analyses. Energy decomposition analysis of the hydrogen bonding interaction revealed a significant reduction in the repulsive component of the hydrogen bonding interaction. The rigid framework of the molecular rotors positions and preorganizes the interacting groups in the transition state. This study demonstrates that with proper design a single hydrogen bond can lead to a TS stabilization that is greater than the intrinsic interaction energy, which has applications in catalyst design and in the study of enzyme mechanisms.

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## Introduction

Hydrogen bonds are key contributors to the large rate accelerations observed in enzyme and synthetic organocatalyst systems.<sup>1–8</sup> However, studying and measuring the kinetic effects of a hydrogen bond is challenging due to the instability and fleeting nature of transition states (TS). To address this problem, a series of molecular rotors 1–3 were synthesized, which measure the TS stabilizing effects of a weak intramolecular hydrogen bond on the rates of rotation (Fig. 1A). The rotors have an *N*-phenyl unit attached *via* a C–N single bond to a 5-membered imide ring. During bond rotation of the C–N single bond, R-groups at the *ortho*-position on the *N*-phenyl rotor are forced into close proximity to the imide carbonyl oxygens in the planar TS (Fig. 1B). The rate of rotation depends on the destabilizing steric and stabilizing non-covalent interactions between the R-groups and the imide carbonyls. Thus, the study of rotors 1–3 provide a simple and potentially accurate method of measuring the TS stabilizing intramolecular hydrogen bond.

Molecular rotors have been used as molecular machines to measure steric effects.<sup>9–11</sup> For example, Sternhell, Roussel, and Mazzanti developed molecular rotors 4, 5, and 6 to develop and compare new empirical steric parameters (Fig. 2).<sup>12–15</sup> The rotational barriers of the rotors were primarily determined by the steric size of the R-groups adjacent to the atropisomeric bond.<sup>9,12–22</sup> Deviations from the steric trends were attributed to the presence of stabilizing TS interactions such as hydrogen

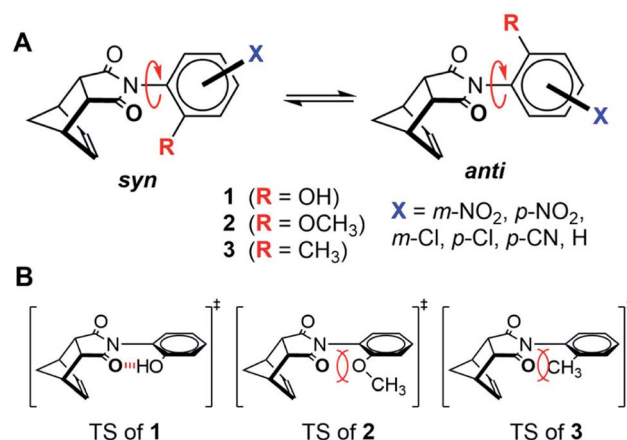


Fig. 1 (A) The *syn*–*anti* conformational equilibria of molecular rotors 1, 2, and 3 designed to isolate and measure the stabilization energy of the intramolecular TS hydrogen bond in rotor 1 *via* the rate of rotation. (B) Representations of the planar TS geometries for rotors 1, 2, and 3.

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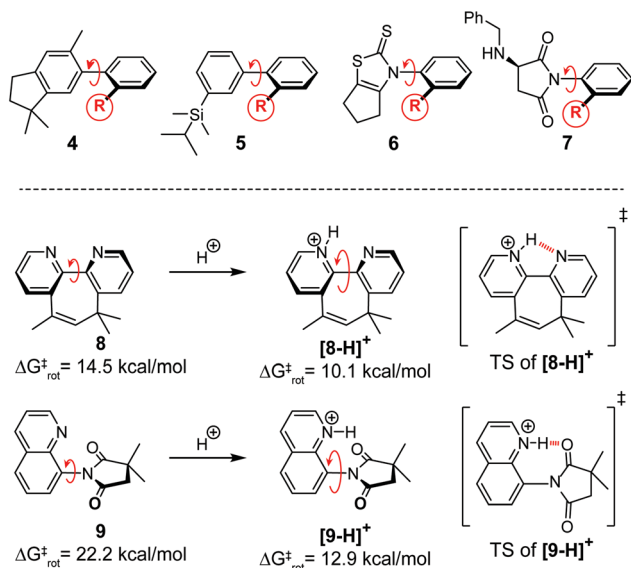


Fig. 2 (Top) Literature examples of molecular rotors 4–7 used to measure steric effects of the *ortho*-substituent (R-group). (Bottom) Molecular machines 8 and 9, which form stabilizing hydrogen bonds in the planar transition state that reduce the rotational or isomerization barriers.

bonding and  $n \rightarrow \pi^*$  interactions.<sup>15,19–22</sup> For example, Rebek and co-workers designed molecular device 8 (Fig. 2), where the rate of isomerization is controlled by stabilizing TS interactions.<sup>17,18</sup> The rate of isomerization was greatly accelerated in the presence of a proton or metal ions which binds to the planar TS of the 2,2'-bipyridine unit. More recently, we developed molecular rotor 9 which has a greatly accelerated rate of rotation upon protonation due to the formation of a stabilizing TS hydrogen bond.<sup>20</sup>

Based on the above examples, the TS hydrogen bond in phenol rotor 1 was predicted to increase the rate of rotation.<sup>19,21</sup> Our expectations were that the increase in rate due to the hydrogen bond would be similar to the thermodynamic strength of the hydrogen bond. However, the magnitude of the TS stabilization (9.9 kcal mol<sup>−1</sup>) was 3 to 6.6 times larger than the measured strength of a hydrogen bond between the phenolic hydroxy group and a carbonyl (1.5 to 3.4 kcal mol<sup>−1</sup> in chlorinated organic solvents).<sup>23</sup> Therefore, the goal of this study was to verify and examine the origins of the unexpectedly large TS stabilization in rotor 1. We were particularly interested in whether hydrogen bonds could have ‘amplified’ effects on transition states that could be used in the design of new hydrogen bonding catalysts and could provide insight into the catalytic mechanism of enzymatic systems.

The use of molecular rotors to study TS interactions has a number of advantages. First, bond rotation is a simple and easy to measure kinetic process. The rate equation is unimolecular, which reduces the number of experimental variables and simplifies the kinetic analysis. In addition, the rate of bond rotation is easily and accurately measured using dynamic NMR methods such as lineshape analysis, coalescence temperature and exchange spectroscopy experiments.<sup>24,25</sup> Second, the bond

rotation transition states can be accurately modelled due to their relatively simple structure, rigid geometrical constraints, and minimal degrees of freedom. Finally, the TS hydrogen bond strengths can be systematically modulated using electron withdrawing substituents on the phenyl rotors (X = H, *p*-Cl, *m*-Cl, *p*-CN, *m*-NO<sub>2</sub>, *p*-NO<sub>2</sub>) that increase the acidity and hydrogen bond donating ability of the phenolic proton.<sup>26</sup>

As in previous systems, the rotational barriers of the molecular rotors were primarily determined by the steric interactions of the *ortho*-substituents (R-groups).<sup>9,15–21</sup> Therefore, the key to the analysis was separating the hydrogen bonding contributions from the steric contributions to the rotational barriers. Our first approach was to compare the rotational barriers of rotors that had similar steric TS interactions but varying hydrogen bonding abilities. Thus, the barriers for hydrogen bonding rotor 1 was compared with non-hydrogen bonding rotors 2 and 3. Rotor 1 has an *ortho*-OH group that can form an intramolecular hydrogen bond with the imide carbonyl oxygen. Control rotors 2 and 3 have *ortho*-OCH<sub>3</sub> and *ortho*-CH<sub>3</sub> groups that lack acidic hydrogens and, thus, cannot form TS hydrogen bonds. The similar steric sizes of the *ortho*-groups in the three rotors was established by comparison of their *B*-values, which is an empirical steric parameter developed by Mazzanti.<sup>13,14</sup> The *B*-values for the OH, OCH<sub>3</sub>, and CH<sub>3</sub> groups were similar at 5.4, 5.6, and 7.4 kcal mol<sup>−1</sup>, respectively.<sup>13,14</sup> Mazzanti’s steric parameter is particularly well-suited to our analysis as *B*-values are based on the rotational barriers for a series of very similar biphenyl rotors 5.

Rotors 1, 2, and 3 were synthesized *via* a one-step thermal condensation of the appropriately substituted aniline with *cis*-5-norbornene-*endo*-2,3-dicarboxylic anhydride.<sup>27,28</sup> The rotational barriers were measured by monitoring the rate of *syn*-*anti* interconversion using variable temperature <sup>1</sup>H NMR methods. The large difference in rotational barriers of hydrogen bonding 1 *versus* non-hydrogen bonding controls 2 was immediately evident by their rate of exchange (Fig. 3). At room temperature (25 °C), the protons for 1 were in fast exchange in the <sup>1</sup>H NMR spectra, as the peaks for the *syn*- and *anti*-rotamers were coalesced. In contrast, protons for rotors 2 were in slow exchange at 25 °C, as the *syn*- and *anti*-rotamers displayed separate sets of peaks.

The large difference in barriers between 1 and 2 were quantitatively measured using three separate dynamic <sup>1</sup>H NMR methods. Coalescence temperature, lineshape analysis, and EXSY all yielded similar barriers (Table 1). The barrier for 1 was 10.8–11.1 kcal mol<sup>−1</sup> and the barrier for 2 was 20.1–20.8 kcal mol<sup>−1</sup>. Thus, each method consistently measured a large difference in barrier ( $\Delta\Delta G_{(2-1)}^{\ddagger} = 9.1 - 9.9$  kcal mol<sup>−1</sup>) between rotors 2 and 1. The minor variations in the individual barriers can be attributed to the  $T\Delta S^{\ddagger}$  term of  $\Delta G^{\ddagger}$  which are characteristically small for rotational barriers.<sup>29,30</sup>

Several alternative hypotheses were explored for the large difference in barriers ( $\Delta\Delta G_{(2-1)}^{\ddagger}$ ). First, the higher barrier of 2 could be due to repulsion between the lone pairs on the oxygens of the carbonyl and *ortho*-OCH<sub>3</sub> groups (Fig. 1B).<sup>31–33</sup> In contrast, these destabilizing lone pair–lone pair (lp–lp) interactions are not formed in rotor 1 due to the formation of the TS hydrogen



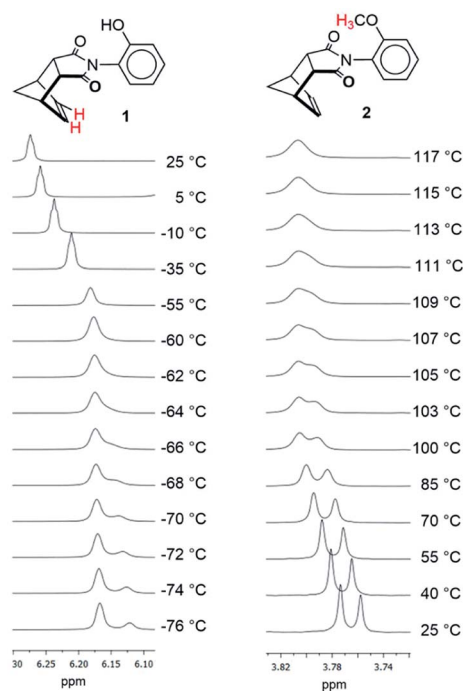


Fig. 3 Variable temperature  $^1\text{H}$  NMR of rotors **1** and **2** measured in  $\text{CD}_2\text{Cl}_2$  and  $\text{TCE-d}_2$ . The protons corresponding to the peaks in each set of spectra are highlighted in red in the structures.

bond. To assess the importance of the lp–lp interactions, a second control rotor **3** was examined which had an *ortho*- $\text{CH}_3$  group that could not form lp–lp interactions. The rotational barrier of control **3** (21.0 kcal mol $^{-1}$  by EXSY) was very similar to the barrier of non-hydrogen bonding control **2** (20.2 kcal mol $^{-1}$  by EXSY), which were consistent with the projected barriers based on the size of the *ortho*-groups from Mazzanti's steric parameters.<sup>22</sup> More importantly, the barriers of control **3** was also significantly higher than hydrogen bonding rotor **1** (11.1 kcal mol $^{-1}$ , by EXSY). Thus, the dramatically higher barrier for control rotor **2** does not appear to be due to lp–lp interactions.

The second hypothesis for the large rotational barrier difference for **1** and **2** was the different solvents ( $\text{CD}_2\text{Cl}_2$  and  $\text{TCE-d}_2$ ) used in the NMR barrier studies. Solvents with very different freezing and boiling points were required because of the large differences in the coalescence temperatures of rotors **1**

and **2**. The first argument against solvent effects is the similarity in solvent polarity and hydrogen bonding ability of  $\text{CD}_2\text{Cl}_2$  and  $\text{TCE-d}_2$  as both are chlorinated organic solvents.<sup>19</sup> The second argument against the solvent effect hypothesis is the relative insensitivity of rotational barriers to solvent environment. For example, Kishikawa examined the solvent effect on the rotational barrier for a series of very similar *N*-phenylimide rotors **7** over a much broader range of solvent polarities (Fig. 2).<sup>34</sup> The differences in barrier measured between a very non-polar (toluene) and polar solvent (DMSO) was only 1.2 kcal mol $^{-1}$ , which was significantly smaller than the observed difference in barrier for **1** and **2**.<sup>34</sup> The possibility that the higher barrier of control rotor **2** was due to hydrogen bonding in the ground state to residual water in the samples was examined. We examined the solvent effects of the rotational barrier of similar *N*-arylimide rotor with an *ortho*-benzyl ether group in a previous study.<sup>19</sup> The barrier of the *ortho*-benzyl ether rotor remained constant when measured in different solvent systems, even those that form strong hydrogen bonds such as acetic acid and triethylamine. This suggests that the high barrier of control rotor **2** is not due to solvent effects. Finally, computational estimates of the barriers made in the absence of solvent were able to accurately reproduce the large barrier difference between **1** and **2** (*vide infra*), providing support that the difference was not due to the different solvent environments.

The inability of the lp–lp and solvent hypotheses to explain the lower barrier of rotor **1** left the intramolecular hydrogen bond as the most likely explanation. Therefore, the next set of experiments examined the correlation between the strength of the hydrogen bond and the TS stabilization. The hydrogen bond strength of the phenolic OH was systematically modulated by attaching electron withdrawing substituents of varying strengths ( $\text{X} = \text{H}, p\text{-Cl}, m\text{-Cl}, p\text{-CN}, m\text{-NO}_2, p\text{-NO}_2$ ) to the *N*-phenyl rotor (Fig. 1A). The substituents increase the acidity and hydrogen bonding ability of the OH proton.<sup>26</sup> An analogous series of similarly substituted control rotors **2** were prepared, and their barriers were measured to assess the substituent effects in the absence of the intramolecular hydrogen bond.

The ability of the substituents to modulate the hydrogen bond strength of the phenolic OH was separately assessed using a series of phenol·NMP (*N*-methylpyrrolidinone) complexes (Fig. 4, top). These bimolecular complexes form the same  $\text{OH}\cdots\text{O}=\text{C}$  hydrogen bond as the TS hydrogen bond in rotor **1**. The association energies ( $\Delta G_{\text{Ka}}$ ) of the substituted complexes with

Table 1 Experimentally measured and calculated rotational barriers<sup>a</sup>

Rotor	$\Delta G^\ddagger$ lineshape <sup>b</sup>	$\Delta G^\ddagger$ coalescence <sup>b</sup>	$\Delta G^\ddagger$ EXSY <sup>b</sup>	$\Delta G^\ddagger_{\text{calc}}$ <sup>c</sup>	$\Delta E^\ddagger_{\text{calc}}$ <sup>c</sup>
<b>1</b>	10.8	10.9 <sup>d</sup>	11.1	10.1	10.1
<b>2</b>	20.1	20.8 <sup>e</sup>	20.2	20.1	18.6
<b>3</b>	19.1	21.4 <sup>f</sup>	21.0	20.3	19.2
<b>1*</b>	NA	NA	NA	NA	19.7
$\Delta\Delta G^\ddagger_{(2-1)}$	9.3	9.9	9.1	10.0	9.6

<sup>a</sup> All energy in kcal mol $^{-1}$ . <sup>b</sup> Measured by VT  $^1\text{H}$  NMR. <sup>c</sup> Geometry optimized at (B3LYP-D3(BJ)/def2-TZVP) level. Energies calculated at (B2GP-PLYP-D3(BJ)/def2-TZVP) level. <sup>d</sup> Coalescence occurred at  $-55^\circ\text{C}$ . <sup>e</sup> Coalescence occurred at  $113^\circ\text{C}$ . <sup>f</sup> Coalescence occurred at  $137.5^\circ\text{C}$ .



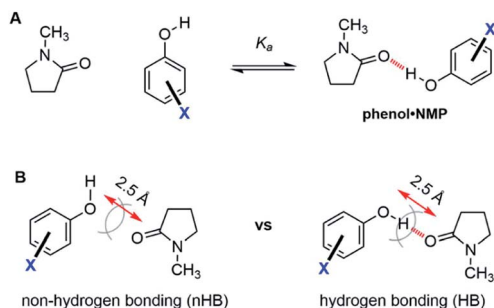


Fig. 4 (A) Bimolecular phenol·NMP equilibrium used to measure the ability of the substituents (X) to modulate the hydrogen bond strength. (B) Non-hydrogen bonding (nHB) and hydrogen bonding (HB) bimolecular phenol·NMP geometries used in the SAPT analysis of the non-covalent interactions.

substituted phenols (X = H, *p*-Cl, *m*-Cl, *p*-CN, *m*-NO<sub>2</sub>, *p*-NO<sub>2</sub>) were measured by <sup>1</sup>H NMR titration in CD<sub>2</sub>Cl<sub>2</sub>.

The association energy ( $\Delta G_{Ka}$ ) of the unsubstituted phenol·NMP complex (X = H) confirmed that the kinetic effects of the TS hydrogen bonds in the molecular rotors were significantly larger than the strength of the hydrogen bonding interaction. The hydrogen bond strengths in the phenol·NMP complex were  $-1.5$  to  $-3.4$  kcal mol<sup>-1</sup>, which were consistent with previous measures of the phenol hydrogen bond strengths in chlorinated organic solvents.<sup>23</sup> More importantly these values were considerably lower than the hydrogen bonding effects on the rotational barrier of rotor 1 ( $\Delta\Delta G_{(2-1)}^{\ddagger} = 9.1 - 9.9$  kcal mol<sup>-1</sup>), even when taking into account the difference in strength of an inter- versus intramolecular interaction (which have been estimated to be 1.4 kcal mol<sup>-1</sup>).<sup>35</sup>

The association energies of the substituted phenol·NMP complexes confirmed the ability of the substituents systematically modulate the strength of the phenol hydrogen bond in a predictable manner. A strong correlation was observed between the electron withdrawing abilities of the substituents and the hydrogen bonding interaction energies as seen by the linear Hammett plot with a negative slope (Fig. 5).

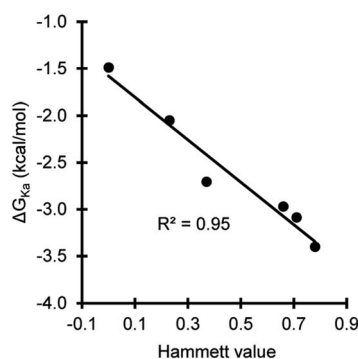


Fig. 5 Hammett plot of the association energies ( $\Delta G_{Ka}$ ) for the phenol·NMP hydrogen bonding complex versus the electrostatic Hammett parameter values for the phenolic substituents.

The substituent effects measured for the phenol·NMP complexes ( $\Delta G_{Ka}$ ) were used to assess the influence of the hydrogen bond strength on the TS stabilizing effects in rotor 1. First, the rotational barriers for similarly substituted rotors 1 and 2 were measured (X = H, *p*-Cl, *m*-Cl, *p*-CN, *m*-NO<sub>2</sub>, *p*-NO<sub>2</sub>). As expected, the barriers for rotor 1 decreased with increasing strength of the intramolecular hydrogen bond from 10.9 kcal mol<sup>-1</sup> (X = H) to 9.7 kcal mol<sup>-1</sup> (X = *p*-NO<sub>2</sub>). This trend was confirmed by the linear correlation between the rotational barriers and the biomolecular association energies ( $\Delta G_{Ka}$ ) for similarly substituted (Fig. 6A). However, the magnitude of the decrease was much smaller than expected. The slope of the trendline for rotor 1 was 0.63, which means that a change in intermolecular hydrogen bond strength of 1.0 kcal mol<sup>-1</sup> lead to only a decrease in rotational barrier of 0.63 kcal mol<sup>-1</sup>. In addition, the TS stabilization attributable to the hydrogen bond was even smaller as the control rotor 2 which cannot form

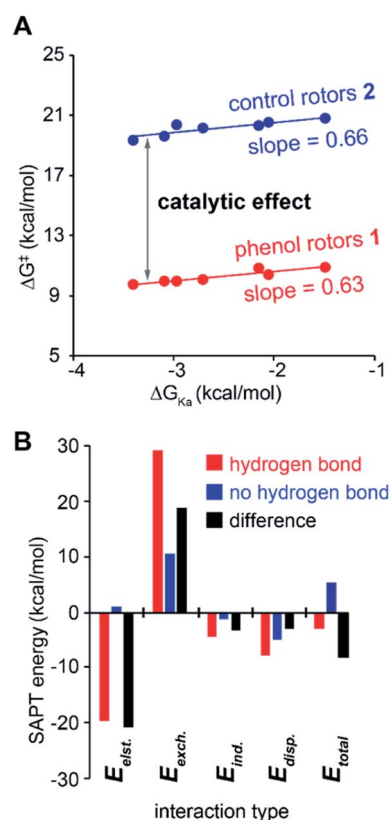


Fig. 6 (A) Correlation plot between the <sup>1</sup>H NMR lineshape analysis measured rotational barriers ( $\Delta G^{\ddagger}$ ) of substituted rotors 1 and 2 (left to right, X = *p*-NO<sub>2</sub>, *m*-NO<sub>2</sub>, *p*-CN, *m*-Cl, *p*-*t*-Amyl, *p*-Cl, and H) versus the measured association energies ( $\Delta G_{Ka}$ ) of similarly substituted phenol·NMP complexes. The difference in the rotational barriers between 1 and 2 corresponds to the TS stabilization of the hydrogen bond. (B) Calculated component ( $E_{elst}$  = electrostatic,  $E_{exch}$  = exchange,  $E_{ind}$  = induction, and  $E_{disp}$  = dispersion) and total SAPT (CCSD/cc-pVQZ) interactions energies ( $E_{total}$  = sum of all component energies) of the phenol·NMP complex (X = H) fixed in hydrogen bonding (HB) and non-hydrogen bonded (nHB) geometries. The difference energies ( $\Delta E$ ) between the HB and nHB geometries provides a measure of the hydrogen bond stabilizing effects.





TS hydrogen bonds displayed nearly identical substituent effects (Fig. 6A, blue circles). Surprisingly, the hydrogen bond induced TS stabilization defined as  $\Delta\Delta G_{(2-1)}^\ddagger$  remained constant ( $9.9 \pm 0.3 \text{ kcal mol}^{-1}$ ) across the series of substituted rotors.

A possible explanation for the similar substituent effects in rotors 1 and 2 is the ability of the substituent on the *N*-phenyl ring to modulate the rotational barrier by through-bond conjugation effects. In the planar TS, the more strongly electron withdrawing substituents of the *N*-phenyl ring form stabilizing resonance structures with the nitrogen of the imide. Support for the through-bond hypothesis was provided by the observation by Kishikawa and co-workers for a series of substituted *N*-phenylsuccinimide rotor 7.<sup>34</sup> The rotational barriers of 7 had similar magnitude substituent effects even with a much less polarizable *ortho*-methyl group. Thus, through-bond substituent effects were observed but were relatively small. More importantly, similar magnitude effects were observed for rotors 1 and 2. Therefore, the large difference in barrier between rotors 1 and 2 cannot be attributed to these through-bond effects.

Next, the origins of the large TS stabilizing effect in rotor 1 were examined computationally using functionals and basis sets previously identified as providing accurate barriers for substituted biaryl rotors.<sup>29</sup> The GS and TS geometries of the series of the substituted rotors 1 and 2 were optimized at the B3LYP-D3(BJ)/def2-TZVP level of theory and verified by vibrational analysis. Single point energies were then calculated on the B3LYP-D3(BJ) optimized structures at the B2GP-PLYP-D3(BJ)/def2-TZVP level of theory. From the entropies and enthalpies, the Gibbs free energies ( $\Delta G_{\text{calc}}^\ddagger$ ) were calculated at the individual coalescence temperature for each rotor. The calculated barriers accurately reproduced the experimental barriers for the substituted rotors including the large difference between rotors 1 and 2 (Table 1). The calculated barriers only slightly underestimated the experimental barrier by an average of  $0.63 \text{ kcal mol}^{-1}$ . Taking this systematic error into account, the calculated barriers reproduced the experimental barriers with an accuracy of  $\pm 0.19 \text{ kcal mol}^{-1}$ . For example, the calculated barriers for the unsubstituted ( $X = \text{H}$ ) rotors 1 and 2 were  $10.1$  and  $20.1 \text{ kcal mol}^{-1}$  versus the NMR line shape analysis barriers by of  $10.8$  and  $20.1 \text{ kcal mol}^{-1}$ .

The accuracy of the calculated barriers suggests that theoretical model was accurately reproducing the TS and GS geometries (Fig. 7). In the TS of rotor 1, a well-defined intramolecular hydrogen bond was formed with O–H $\cdots$ O bond angles of  $165.76^\circ$  and a very short O $\cdots$ O distance of  $2.50 \text{ \AA}$  (Fig. 3). For comparison, the equilibrium hydrogen bonding distance for an O–H $\cdots$ O hydrogen bond is typically  $2.8$  to  $2.9 \text{ \AA}$ .<sup>36</sup> The short O $\cdots$ O distance of the TS hydrogen bond did not vary across the substituent series ( $2.49 \text{ \AA}$ ,  $\pm 0.005$ ) even with the strongest electron withdrawing substituents such as *p*-NO<sub>2</sub> or *m*-NO<sub>2</sub>. Interestingly, the TS geometry of the non-hydrogen bonding 2 was nearly identical to the TS of the hydrogen bonding 1. For example, the average O $\cdots$ O distances of 2 was  $2.48 \pm 0.003 \text{ \AA}$ , again with very little variance across the substituted series. The similarity in the TS distances and geometries of 1 and 2 demonstrates the ability of the rigid rotor

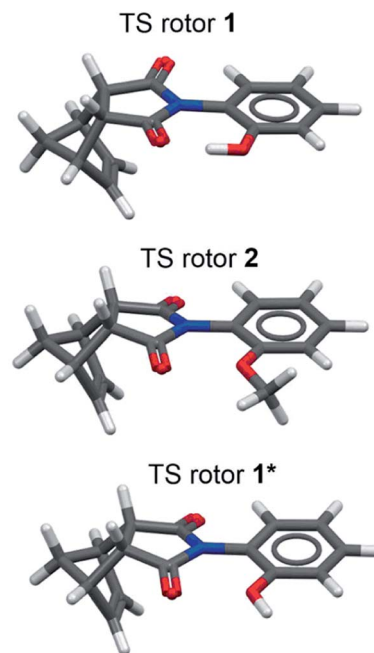


Fig. 7 The optimized TS structures (B3LYP-D3(BJ)/def2-TZVP) and energies (B2GP-PLYP-D3(BJ)/def2-TZVP) for rotors 1 and 2 and rotor 1\* constrained with the phenol hydrogen fixed in a non-hydrogen bonding geometry.

framework to precisely position the interacting groups in the TS.

Two important conclusions were drawn from the computational studies. First, the non-hydrogen bonding rotor 2 was confirmed as a good control to isolate the hydrogen bonding effects in rotor 1. The large difference in barrier between rotors 1 and 2 were due the formation of the intramolecular hydrogen bond in 1 and not due to steric differences between the OH and OCH<sub>3</sub> groups. To confirm the similarity in steric size of the OH and OCH<sub>3</sub> groups, a geometrically constrained phenol rotor 1\* was calculated containing an OH group by constrained in a non-hydrogen bonding geometry with the phenolic proton pointing away from the imide carbonyl (Fig. 7, bottom structure). The barrier for 1\* was calculated and compared with the previously calculated barriers for 1 and 2. Due to the constraint, the  $\Delta G_{\text{calc}}^\ddagger$  values for 1\* were not readily estimated. Therefore, the  $\Delta E_{\text{calc}}^\ddagger$  for the barriers of the rotors were compared (Table 1, column 6). The barrier of 1\* ( $\Delta E_{\text{calc}}^\ddagger = 19.7 \text{ kcal mol}^{-1}$ ) was very similar to the non-hydrogen bonding control 2 ( $\Delta E_{\text{calc}}^\ddagger = 18.6 \text{ kcal mol}^{-1}$ ). The similarity in barrier confirms that, in the absence of the hydrogen bond, the OH and OCH<sub>3</sub> groups have similar steric effects. In addition, 1\* had a much higher barrier ( $19.7 \text{ kcal mol}^{-1}$ ) than unconstrained 1 ( $10.1 \text{ kcal mol}^{-1}$ ), which provided further support for the dominant role of the hydrogen bond in lowering the barrier for rotor 1.

The effects of the hydrogen bond on the TS 1 were next examined using symmetry-adapted perturbation theory (SAPT) analysis, which separates the interaction energies into electrostatic, exchange, induction, and dispersion components. SAPT is designed for intermolecular interactions. Thus, the method



was applied to study the O–H⋯O hydrogen bonding interaction in the phenol·NMP complex (Fig. 4B). The interacting atoms of the phenol and NMP amide were constrained in a planar geometry, and the O-to-O distance was fixed at 2.5 Å to mimic the TS distances in the rotors. Again, the hydrogen bonding effects were isolated by comparing hydrogen bonding (HB) and non-hydrogen bonding (nHB) geometries of the biomolecular complex. In the HB complex, the phenol proton was unconstrained and formed an intermolecular hydrogen bond with the NMP carbonyl oxygen. In the nHB complex, the phenol proton was fixed in a non-hydrogen bonding position pointing away from the NMP carbonyl. The total SAPT interaction energies ( $E_{\text{total}}$ ) and component energies (electrostatic ( $E_{\text{elst}}$ ), exchange ( $E_{\text{exch}}$ ), induction ( $E_{\text{ind}}$ ), and dispersion ( $E_{\text{disp}}$ )) were calculated using CCSD(T)/cc-VQZ for the HB complex, nHB complex, and the difference energy (nHB – HB) (Fig. 6).

The ability of the SAPT analyses of the phenol·NMP complexes to provide insight into the hydrogen bonding effects in the molecular rotors was confirmed by the similarity with previous trends. First, the SAPT component analysis of the HB complexes (Fig. 6B, red bars) matched previous component analyses of hydrogen bonding interactions.<sup>37</sup> Specifically, the hydrogen bonding interaction in the phenol·NMP complex (Fig. 6A, red bars) was made up of large opposing attractive ( $-32.1 \text{ kcal mol}^{-1}$ ) and repulsive ( $+29.2 \text{ kcal mol}^{-1}$ ) terms, which largely cancel to yield a weakly stabilizing interaction ( $E_{\text{total}} = -2.9 \text{ kcal mol}^{-1}$ ). The attractive term is dominated by the electrostatic component with smaller contributions from the induction, and dispersion components. The repulsive term is made up entirely of the exchange component. Second, the bimolecular complexes showed the same discrepancy between the apparent strength of the hydrogen bonding interaction and its effect on the stability of the complexes. Specifically, the hydrogen bonding interaction energy, as measured by the difference in interaction energies ( $E_{\text{SAPT } \Delta}$ , Fig. 6A black bars) for the nHB and HB complexes, was significantly larger than the hydrogen bonding interaction energy ( $E_{\text{SAPT HB}}$ , Fig. 6A red bars).

Analysis of the component energies revealed that the origin of the discrepancy in the  $E_{\text{SAPT } \Delta}$  and  $E_{\text{SAPT HB}}$  energies was due to the repulsive exchange component of the hydrogen bonding interaction. The variation in the repulsive term is evident from a comparison of the SAPT component energies. Due to the magnitude of the repulsive component, even a small difference of 33% in the  $E_{\text{exch}}$  component of the HB complex (Fig. 6, red bars) and the difference energy (Fig. 6, black bars) has a large impact. By comparison, the attractive terms ( $E_{\text{elst}}$ ,  $E_{\text{ind}}$ , and  $E_{\text{disp}}$ ) in the HB complex and the difference energies do not differ significantly.

The disparity in the exchange components is due to the constraints imposed on the biomolecular complexes, which mimic the rigid framework of the molecular rotors. These constraints 'prepay' the repulsive steric interactions by holding the heavy atom oxygens in close proximity in the hydrogen bonding and non-hydrogen bonding complexes. The repulsive interactions of the oxygen atoms in the OH⋯O=C interaction of the HB complex (2.50 Å) make up approximately one-third of

the overall repulsive component of the hydrogen bonding interaction. Therefore, the difference energy between the HB and nHB complexes contains all of the attractive terms of the hydrogen bonding interaction but only two-thirds of the repulsive term. This helps explain how the effect of the hydrogen bond, which is simulated by the different energy between the nHB and HB complex, is significantly larger than the hydrogen bonding energy which is measured by the interaction energy in the HB complex. The rigid aromatic framework of the molecular rotors imposes similar position and distance constraints of the interacting groups in the planar TS. Thus, the trends and interaction energies observed in the bimolecular complexes should be similar to those in the molecular rotors.

Another way to explain the larger than expected effects of the hydrogen bond is to compare the repulsive components in the two systems shown in Fig. 4. The top set of structures (Fig. 4A) is an equilibrium between hydrogen bonding phenol and NMP molecules. On the right-hand side of the equilibrium, the molecules are close together forming a hydrogen bonding interaction. On the left-hand side, the two molecules are far apart (left) and cannot form any interactions. Thus, the equilibrium energy ( $\Delta G_{\text{Ka}}$ ), which is the difference in energy between the right and left side of the equilibrium arrow, measures all of the attractive and repulsive components of the hydrogen bonding interaction. By comparison, the bottom set of structures (Fig. 4B) is representative of the kinetic measurements in our molecular rotor systems where the interacting groups are rigidly constrained by the *N*-arylimide framework. The structure on the right is the HB complex forms similar hydrogen bonding interactions as the above equilibrium system. The control structure on the left does not form a hydrogen bonding interaction. However, the interacting groups are still positioned in close proximity due to the geometric constraints and therefore still contains significant repulsive interactions. The oxygens of the phenol and carbonyl are closer (2.50 Å) than the sum of their VDW radii (3.04 Å). Therefore, the difference energy ( $\Delta E_{\text{HB}} - \Delta E_{\text{nHB}}$ ) contains all of the attractive components of the hydrogen bond but only a fraction of the repulsive components. Therefore, the difference energy for the bottom set of constrained structure can be considerably larger than  $\Delta G_{\text{Ka}}$  for the top set of equilibrium structures.

These molecular rotors demonstrate an alternative strategy for enhancing the kinetic effects of a hydrogen bonds *via* modulation of the repulsive component. More typically hydrogen bonding interaction are attenuated by modulation the attractive component. For example, the attractive electrostatic component of the hydrogen bond in the phenol·NMP complex was systematically strengthened using electronegative substituents.

A survey of the literature identified other examples (Fig. 8) of hydrogen bonds being modulated *via* the repulsive component. In these systems, like the molecular rotors, the interacting groups are fixed in close proximity in the hydrogen bonding and non-hydrogen bonding control structures. A common feature was the surprising strength or influence of the hydrogen bonds. The first example is proton sponge, 1,8-diaminonaphthalene



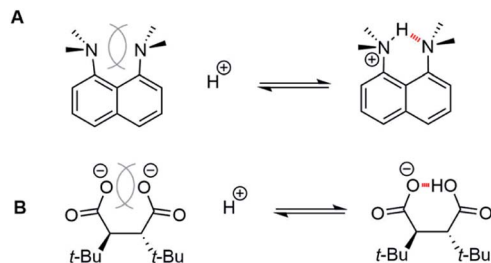


Fig. 8 Protonation equilibria for (A) proton sponge and (B) (±)- $\alpha,\alpha'$ -di-*tert*-butylsuccinate which form strong intramolecular hydrogen bonds due to the strong destabilizing repulsive interactions in the unprotonated structures.<sup>40</sup>

(Fig. 8A). The unusually high proton affinity of proton sponge is not due to the strength of the hydrogen bonding interacting in the protonated structure. Instead, computational studies have attributed the high proton affinity to the relief of the extreme strain in the unprotonated diamine.<sup>38</sup> A more recent example was provided by Perrin *et al.* (Fig. 8B). (±)- $\alpha,\alpha'$ -di-*tert*-butylsuccinate also shows a very high proton affinity, which was initially attributed to the strong intramolecular hydrogen bond. However, more careful analysis found that the high proton affinity was due to the reduction of the large repulsive interactions in the non-hydrogen bonding dianion structure.<sup>39</sup>

## Conclusions

The study of the kinetic effects of a hydrogen bond using molecular rotor **1** demonstrates that a single neutral hydrogen bond can have a TS stabilization that appears to be many times larger than the thermodynamic strength of the hydrogen bond. The origins of the enhanced TS stabilization were examined, which provided new design strategies for hydrogen bonding catalysts. The traditional approach has been to optimize the attractive component of hydrogen bonding interactions. However, the molecular rotors demonstrate that large rate accelerations can be affected by reducing the large repulsive energy term of the hydrogen bond. A key question is whether the trends observed for the intramolecular hydrogen bonds in the molecular rotors are also relevant to bimolecular catalytic systems, which generally have longer atom–atom distances and greater flexibility. The extensive analysis of enzyme catalysis by Warshel suggest that similar strategies of reducing the repulsive component of non-covalent interactions are operative in biological systems.<sup>41</sup> Specifically, Warshel has hypothesized that a significant portion of the large catalytic rate enhancements can be attributed to the ability of the enzyme framework to preorganize the interacting and catalytic groups in the transition state. Thus the protein framework positions polar and charged groups in close proximity overcoming the repulsive forces and creating a high energy ground state that is closer to the TS energy. Small molecule catalysts could also be designed to prepay repulsive energy during binding of the catalyst and ligand prior to the reaction preceding. Like in the enzyme, by forming a complex with key reacting groups positioned

appropriately so the energy penalty is paid during the initial complex formation, a catalyst can more effectively drive a reaction forward.

## Conflicts of interest

There are no conflicts to declare.

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