Reaction **Chemistry & Engineering**



PAPER

View Article Online



Cite this: React. Chem. Eng., 2020,

High redox performance of $Y_{0.5}Ba_{0.5}CoO_{3-\delta}$ for thermochemical oxygen production and separation†

M. Ezbiri, A. Reinhart, B. Huber, A. K. M. Allen, A. Steinfeld, B. Bulfin ** and R. Michalsky ** **

The efficient production and separation of oxygen is essential for numerous energy-intensive industrial applications in the fuel and mineral processing sectors. A thermochemical redox cycle is considered for separating oxygen from atmospheric air and other gas mixtures using solar or waste process heat. Based on electronic structure (DFT) computations $Y_{0.5}Ba_{0.5}CoO_{3-\delta}$ is selected as a redox material, which surpasses the redox performance of state-of-the-art Cu₂O. The thermochemical oxygen production is experimentally demonstrated by applying a temperature/pressure swing between 573 K at 0.2 bar O2 and 873 K at 1 bar O2. An energy balance shows the feasibility of using process waste heat from the solar thermochemical CO₂/H₂O splitting cycle and the potential to compete vis-a-vis with cryogenic distillation. Exploratory runs with a packed-bed reactor indicate the potential of both thermochemical oxygen production and separation for scale-up and industrial implementation.

Received 7th November 2019, Accepted 13th December 2019

DOI: 10.1039/c9re00430k

rsc.li/reaction-engineering

1. Introduction

Oxygen generation/removal is required for fuel and mineral processing and is mainly carried out by cryogenic distillation.¹⁻⁴ Other oxygen separation techniques include pressure swing adsorption using zeolites and carbon molecular sieves, and electrochemical separation using mixed ionic-electronic conducting membranes. 1-5 These separation technologies require an electrical input between 100 and 350 kWh per metric ton of O₂.²⁻⁶

An interesting alternative route to oxygen separation can be achieved via a thermochemical process, using a metal oxide redox cycle.7-9 Heat can be used to drive the reduction of a metal oxide releasing oxygen. The reduced metal oxide can then absorb oxygen from a gas stream such as air. BOC Linde developed such a process, with a temperature and pressure swing thermochemical redox cycle suppling oxygen to an oxycombustion process.10 This idea has been further developed in recent years, in particular in the search for new materials. For example, Cu₂O/CuO,^{7,11} Mn₃O₄/Mn₂O₃,¹¹ CoO/Co₃O₄ (ref. 11) and ABO3-type perovskites, 8,12,13 have all been explored as redox pairs for oxygen separation and production processes.

Perovskites in particular appear to be very promising for this application, as they can exhibit rapid kinetics at relatively low temperatures,9 and the thermodynamics can be tuned to optimize the energetics of the process.¹³ The perovskite SrFeO₃, and doped variants, can be considered the state of the art. 9,10 However, cobalt based perovskites such as SrCoO₃ appear to be more promising for lower temperature application, and they have the lowest energetic requirement in terms of the enthalpy change of the reduction reaction.8,13-15

In this paper, we consider both thermochemical oxygen separation (TOS) and thermochemical oxygen production (TOP) via redox cycles with a focus on perovskites with cobalt as the active redox ion. We apply electronic structure calculations on candidate redox materials and select $La_{0.5}Ba_{0.5}CoO_{3-\delta}$, $Pr_{0.5}Ba_{0.5}CoO_{3-\delta}$ and $Y_{0.5}Ba_{0.5}CoO_{3-\delta}$ for further experimental investigation of their O2 exchange capacity. We performed a comprehensive energy balance to show that the perovskite $Y_{0.5}Ba_{0.5}CoO_{3-\delta}$ can be applied in TOS driven by solar or waste heat and compare the results to the state-of-the-art cryogenic distillation. We also consider the combination of this process with the solar-driven splitting of CO₂ and H₂O, which commonly uses an inert gas for the removal of oxygen. 16-19 TOS can be used to recycle the inert gas, and waste heat from the solar fuel production can be integrated into the TOS processes. Furthermore, we conduct exploratory runs with a packed-bed reactor to elucidate the potential for scale-up and industrial implementation.

^a Department of Mechanical and Process Engineering, ETH Zürich, 8092 Zürich, Switzerland. E-mail: bulfinb@ethz.ch, michalskyr@ethz.ch

^b Solar Technology Laboratory, Paul Scherrer Institute, 5232 Villigen, Switzerland † Electronic supplementary information (ESI) available. See DOI: 10.1039/ c9re00430k

The TOS/TOP approach via 2-step redox reactions can be described by eqn (1a) and (1b) for generic perovskite with ABO₃ stoichiometry, where A and B are metal cations in twelve- and six-coordinated interstices, δ is the oxygen nonstoichiometry and $\Delta\delta$ is the oxygen exchange capacity ($\Delta\delta$ = $\delta_{\rm red}$ – $\delta_{\rm ox}$):

Oxidation:

$$ABO_{3-\delta_{red}} + \frac{\Delta \delta}{2}O_2(g) \rightarrow ABO_{3-\delta_{ox}} \tag{1a}$$

Reduction:

$$ABO_{3-\delta_{ox}} \rightarrow ABO_{3-\delta_{red}} + \frac{\Delta \delta}{2}O_2(g)$$
 (1b)

The redox cycle is performed by applying a temperature/pressure swing. During the oxidation step, oxygen is separated from a gas mixture through the oxidation of the perovskite at low pO_2 and low temperatures. During the reduction step, the perovskite is reduced using a gas mixture at high pO_2 and high temperatures. Redox conditions used in this study for both TOS and TOP are illustrated in Table 7.

2. Computational and experimental section

Electronic structure calculations

The grid-based projector-augmented wave (GPAW) code was employed for density functional theory computations. 19,20 Atomic configurations were handled in the atomic simulation environment (ASE), 21 with exchangecorrelation interactions treated by the revised Perdew-Burke-Ernzerhof (RPBE) functional.²² The optimizations of the atomic geometries were executed using the linesearch Broyden-Fletcher-Goldfarb-Shanno (BFGS) algorithm until the maximum force was less than 0.05 eV Å⁻¹. Convergence was achieved using a Fermi-Dirac smearing of 0.1 eV. The results were extrapolated to 0 K. The generalized gradient approximation (GGA) without a Hubbard U correction was used for all DFT calculations, in accordance to previous publications^{23–25} AA'BO₃-type perovskites were modeled with A₂B₂O₆ models containing a cubic ABO₃ unit cell, with one metal (Ba, Pr, La and Y) at the twelve-coordinated A-site interstices and one metal (Co) at the six-coordinated B-site interstices. All atoms of the bulk models were allowed to optimize their positions (relax). The Brillouin zone of these bulk models having periodic boundary conditions in all

Table 1 Structural properties calculated from Rietveld analysis of XRD patterns of LBCO, PBCO and YBCO

	Space group	wt% perovskite (-)	Volume (ų)
LBCO	$Pm\bar{3}m^{32}$	100	58.9
PBCO	$P4/mmm^{33}$	99.1	115.8
YBCO	$P4/mmm^{34}$	98.0	112.3

Table 2 Performance indicators of LBCO, PBCO and YBCO based on the TGA experiments shown with Fig. 2

	$T_{ m red,start}$ (K)		Reox. slope (μmol _O K ⁻¹)	$\delta_{ m end,reox} (m mol_O \ mol_{ABO_3}^{-1})$	T _{net,eq} (K)
LBCO	692	396	409	0.034	811
PBCO	530	313	373	0.055	725
YBCO	582	541	565	0.035	633

directions was sampled with $4 \times 4 \times 4$ k-points. All compositions were treated using spin-polarized calculations. The computed lattice constants are within ±1.6-2.6% of experimental values. The perovskite surfaces were modeled using the (010) facet with AO-termination, which is geometrically symmetric with the (001) and (100) facets dependent on composition, i.e. the thermodynamically most stable surface.26 The surface models contained the same number of atoms as the bulk models, i.e. stoichiometry A₂B₂O₆, of which the upper ABO₃ layer parallel to the surface was allowed to relax, while the lower layer was constrained to the bulk geometry. All surface structures were periodically repeated parallel to the surface, used 10 Å of vacuum perpendicular to the surface, and employed a k-point sampling of $4 \times 4 \times 1$. To model reduced perovskite surfaces, one third of the stoichiometric lattice oxygen was removed from the upper surface layer, resulting in A2B2O5 models corresponding to δ = 0.5.

The Gibbs free energy change of oxygen vacancy formation $(\Delta G_V[O])$ was computed by:²⁶

$$\Delta G_{\rm V}[{\rm O}] = G_{\rm V} - (G_{\rm S} - G_{\rm O}^{\rm r}) \tag{2}$$

where $G_{\rm V}$ and $G_{\rm S}$ being the Gibbs free energies of formation of the perovskite surfaces with and without oxygen vacancies, respectively, while $G_{\rm O}^{\rm r}$ is the reference free energy of the liberated lattice oxygen taken as the energy difference of stable $\rm H_2O$ and $\rm H_2$ in the gas phase. Negative free energies correspond to exergonic reactions. In this work, the standard partial molar enthalpy change of the metal oxide bulk reduction at 298 K and 1 bar total pressure, $\Delta \overline{h}_{\rm O_3}^{\rm o*}$ is the

Table 3 Energy requirement per mol of O₂ removed from a 0.5%, 1% and 5% O₂-Ar mixture coming from a solar reactor at 1273 K. In case of no heat recuperation, $\Delta h_{\rm cycle}$ is calculated via $\Delta h_{\rm ha}$ + $\Delta h_{\rm red}$ + $\Delta h_{\rm hs}$, while $\Delta h_{\rm cycl} = \Delta h_{\rm ha} + \Delta h_{\rm red} + \Delta h_{\rm hs} - \left(\frac{\Delta h_{\rm ca} + \Delta h_{\rm cg} + \Delta h_{\rm ox}}{1.25} + \frac{\Delta h_{\rm cs}}{2}\right) \quad \text{is used in case of 50% solid and 80% gas heat recuperation}$

$p\mathrm{O}_2$	%-Fraction of heat recovered (–)		$\Delta h_{ m cycle}$	$\Delta h_{ m avail}$
(bar)	From solids	From gas	(kJ mol _{O2} ⁻¹)	(kJ mol _{O2} ⁻¹)
0.005	0	0	828	1480
0.005	50	80	-107	1480
0.01	0	0	1280	1483
0.01	50	80	91	1483
0.05	0	0	4217	1507
0.05	50	80	1371	1507

Table 4 Oxygen storage properties of YBCO compared to other state of the art materials in the literature. All data is from TGA analysis

Material	T_{low} – $T_{\mathrm{high}}\left[\mathrm{K}\right]$	ΔT [bar]	p_{low} – p_{high} [bar]	Δδ [-]
YBCO	563-773	210	0.01-0.2	0.064
SrFeO ₃ (ref. 9)	623-823	200	0.01 - 0.2	0.035
SrCoO ₃ ESI	600-900	300	0.01 - 0.2	0.029
Cu_2O (ref. 7)	1120-1450	330	0.01-0.01	0.4

Table 5 Extracted $\Delta\delta$ from PBR experiments at 523 K, 563 K, 603 K and 643 K and 0.01 bar pO2 during oxidation and at 773 K and 0.2 bar pO2 during reduction

Oxidation temperature (K)	Experimental oxidation period (min)	$\Delta\delta$ (mol _O mol _{ABO₃} ⁻¹)
523	225	0.083
563	175	0.064
603	130	0.043
643	100	0.020

descriptor for the standard partial molar Gibbs free energy change of the metal oxide bulk reduction and oxidation, $\Delta \overline{g}_{0}^{\circ *}$ and $-\Delta \overline{g}_{\Omega_0}^{\circ *}$, respectively. Tabulated thermochemical data, ²⁷ available for a limited number of mostly stoichiometric metal oxides, was used to compute $\Delta \overline{h}_{0}^{\circ *}$, based on previously reported scaling relations between $\Delta \overline{h}_{O_2}^{\circ*}$ and $\Delta G_V[O]$.²⁶ The linear scaling relation utilized in this work is:

$$\Delta G_{\rm V}[{\rm O}] = 4.49069 \times 10^{-3} \frac{\rm eV}{\rm kJ \; mol_{\rm O_2}^{-1}} \Delta \overline{h}_{\rm O_2}^{\rm o*} - 8.16401 \times 10^{-1} \; \rm eV \endalign{3}$$

In this study, all thermodynamic properties are defined per mole of diatomic oxygen. Details on how the DFTcomputed electronic energies were converted into Gibbs free energies at different temperatures and pO2, the reference energies, and details on the scaling relation²⁶ are given in ESI.†

Perovskite synthesis - perovskites were synthesized with a modified Pechini method, employing stoichiometric amounts of La(NO₃)₃·H₂O (Sigma Aldrich, 99.999%), Ba(NO₃)₂ (Alfa Aesar, 99%), Co(NO₃)₂·6 H₂O (Sigma-Aldrich, \geq 98%), Y(NO₃)₃ ·6 H₂O (Aldrich Chemistry, 99.8%), Pr(NO₃)₃·6 H₂O (Aldrich

Table 7 Redox conditions of oxygen separation (TOS) and production (TOP) in comparison to state-of-the-art cryogenic distillation

	TOS	TOP
pO _{2,red,in} (bar)	0.2	1
$pO_{2,red,out}$ (bar)	0.4	1
$pO_{2,ox,in}$ (bar)	0.01	0.2
$pO_{2,ox,out}$ (bar)	0.001	0.00035
$T_{\rm red}$ (K)	773	873
$T_{\text{ox}}\left(\mathbf{K}\right)$	563	573
$\Delta\delta$ (mol _O mol _{ABO} , ⁻¹)	0.064	0.114
$\Delta h_{\rm req}$ (kJ mol _{O2} ⁻¹)	0	377
State-of-the-art	Cryogenics	Cryogenics
$E_{ m th}$	50 kJ $\text{mol}_{N_2}^{-1}$ (ref. 6)	78 kJ $\text{mol}_{\text{O}_2}^{-1}$ (ref. 6)

Chemistry, 99.9%), C₂H₆O (Alcosuisse, >96.1% vol) and $C_6H_8O_7$ (Fluka Chemika, $\geq 99.5\%$). For the pellet production of $La_{0.5}Ba_{0.5}CoO_{3-\delta}$ (LBCO), $Pr_{0.5}Ba_{0.5}CoO_{3-\delta}$ (PBCO) and Y_{0.5}Ba_{0.5}CoO_{3-δ} (YBCO), the solid products were ground using mortar and pestle, calcined at 1273 K in air for 5 hours using 5 K min⁻¹ for heating and cooling and uniaxially pressed into pellets (5 metric tons, 8 mm in diameter). Granules of YBCO were manufactured by mixing the calcined powder with deflocculant agent (Dolapix CE 64, Zimmer & Schwartz) and deionized water. Granules (around 3 mm in diameter) of the resultant slurry were then pressed out of a syringe onto a platinum sheet and dried. LBCO and PBCO pellets were sintered in air at 1473 K for 5 h, while YBCO pellets, granules and calcined powder were sintered at 1323 K for 24 h. For all materials, a heating and cooling rate of 1 K min⁻¹ was employed.

Solid-state analysis

Room-temperature X-ray diffraction (XRD) on YBCO, PBCO and LBCO and high-temperature XRD (HT-XRD) on YBCO were performed in the Bragg Brentano geometry using Cu Kα radiation (20-80° 2θ , 0.01° min⁻¹ scan rate, 45 kV/20 mA output, PANalytical/X'Pert MPD/DY636, Philips). HT-XRD scans were executed between 308 and 923 K in 0.2 bar O2 in N₂ during heating up and 0.01 bar O₂ in Ar during cooling down using a platinum sheet as sample holder. Basic structural data of the room temperature measurement was obtained by multiphase Rietveld analysis (Jana2006). The morphology of pellets and granules was analyzed via

Table 6 Energy balance of TOP in comparison to state-of-the-art cryogenic distillation. $E_{\rm el}$ and $E_{\rm th}$ denote the electrical and thermal energies, respectively, while Δh_{reg} and Δh_{recup} denote the required energy and the energy recuperated by solid and gas heat recuperation, respectively

TOP				
%-Fraction of heat r	ecovered (-)			
From solids 50	From gas 80	Δh_{req} (kJ $\mathrm{mol_{O_2}}^{-1}$) 844	$\begin{array}{c} \Delta h_{\rm recup} \ ({\rm kJ \ mol_{O_2}}^{-1}) \\ 466 \end{array}$	$\Delta h_{\text{cycle}} \left(\text{kJ mol}_{\text{O}_2}^{-1} \right)$ 377
Cryogenic distillatio	n			
Purity 99.5%	$E_{\mathrm{el}} \left(\mathrm{kJ} \; \mathrm{mol}_{\mathrm{O_{2}}} \right)^{-}$	1)	$E_{\text{th}} \text{ (kJ mol}_{\text{O}_2}^{-1}\text{)}$	

scanning electron microscopy (SEM; HITACHI Tabletop Microscope TM-1000). Brunauer-Emmett Teller Micrometrics TriStar 3000 N₂ absorption measurements were performed to acquire information about the specific surface area of the YBCO granules. The BET results, XRD patterns and SEM images are presented in ESI.†

Screening

PBCO, **LBCO** and YBCO were screened using thermogravimetric analysis (TGA) between 303 and 1273 K using +2 and -2 K min⁻¹ in 0.2 bar O₂ in N₂ during heating and in 0.005 bar O2 in Ar during cooling. The slow heating and cooling rates were chosen to ensure thermodynamic equilibrium at all times. To correct for buoyancy, blank runs were performed using the same measurement conditions employed for the experimental runs.

Thermodynamic analysis

The thermodynamic analysis was performed to extract the reduction enthalpy, $\Delta h_{\rm red}$, using the approach of Scheffe et al.:28 Oxygen release and uptake of a perovskite can generally be described as:

$$ABO_3 \hookrightarrow ABO_{3-\delta} + \frac{\delta}{2}O_2(g)$$
 (4)

From eqn (4) we can derive the oxygen partial pressure dependence (pO_2) on the equilibrium constant K:

$$K = (pO_2)^{0.5} \tag{5}$$

and its relation to the standard Gibbs free energy change $(\Delta \overline{g}_{\mathcal{O}_2}^{\circ})$:

$$\Delta \overline{g}_{O_2}^{\circ}(\delta, T) = -2RT \ln K = -RT \ln(pO_2/p^{\circ})$$
 (6)

where p° is the standard pressure at 1 bar. $\Delta g_{\Omega_{2}}^{\circ}$ can further be derived from the standard partial molar enthalpy change of reduction, $\Delta \overline{h}_{O_2}^{\circ}$, and the standard partial molar entropy change of reduction, $\Delta \bar{s}_{O_2}^{\circ}$:

$$\Delta \overline{g}_{\Omega_{2}}^{\circ}(\delta, T) = \Delta \overline{h}_{\Omega_{2}}^{\circ}(\delta) - T \Delta \overline{s}_{\Omega_{2}}^{\circ}(\delta) \tag{7}$$

Combining eqn (6) and (7), we obtain $\Delta \overline{h}_{0_2}^{\circ}$ and $\Delta \overline{s}_{0_2}^{\circ}$ as a function of δ by determining the slope and intercept of $-\ln(pO_2/p^\circ)$ as a function of T^{-1} for a constant δ (cf. eqn (8)):

$$-\ln\left(\frac{pO_2}{p^{\circ}}\right) = \frac{\Delta \overline{h}_{O_2}^{\circ}(\delta)}{RT} - \frac{\Delta \overline{s}_{O_2}^{\circ}(\delta)}{R}\bigg|_{\delta = \text{const}}$$
(8)

Finally, integrating $\Delta \overline{h}_{O_2}^{\circ}$ over the initial and final δ , we obtain the reduction enthalpy, $\Delta h_{\rm red}$ (cf. eqn (9)). $\Delta h_{\rm red}$ is directly related to the solar input required to drive the redox reaction.

$$\Delta h_{\rm red} = \frac{\int_{\delta_{\rm i}}^{\delta_{\rm f}} \Delta \overline{h}_{\rm O_2}^{\circ}(\delta)}{\delta_{\rm f} - \delta_{\rm i}} d\delta \tag{9}$$

where δ_i and δ_f are the oxygen nonstoichiometries before and after reduction, respectively.

To determine $\Delta h_{\rm red}$, TGA measurements were conducted by heating up to 1073 K in 0.2 bar O₂ in N₂ and cooling down in 0.005, 0.01, 0.05 and 0.2 bar O2 in Ar, using slow rates of +2 K min⁻¹ and -2 K min⁻¹ to ensure thermodynamic equilibrium at all times (cf. Fig. 2).

Molar heat capacity measurements

The molar heat capacity (c_p) of YBCO was measured with a heat flux differential scanning calorimetry (DSC, SENSYS evo DSC, Setaram). The sample, in the form of crushed pellet pieces with a total weight of 178.5 mg, was subjected to an Ar flow (99.998%, Carbagas) of 5 mL min⁻¹ between 373 and 1073 K using a heating rate of +10 K min⁻¹. To evaluate the measurement accuracy, we further acquired the c_p of CeO₂ and compared the result to previously published data²⁹ (cf. Fig. S11†), revealing a discrepancy on average of only 1.94%. All experimental curves were fitted using the least-squares fitting method.

Oxygen nonstoichiometry cycling measurements

Oxygen nonstoichiometry (δ) was measured using a thermogravimetric analyzer (TGA, STA 409/C/3, Netzsch). YBCO granules (112 mg) were placed on a platinum crucible supported with an Al₂O₃ rod on the microbalance of the TGA (±0.1 μg). The granules were thereafter exposed to a gas flow (constant flow rate of 200 mL min⁻¹) with specified pO₂ that was adjusted via mixing O₂ (99.5%, Messer) and N₂ (99.999%, Carbagas) using three electronic mass flow controllers (MFC, MFC400, Netzsch; accuracy $\pm 1\%$, precision ± 1 mL min⁻¹). For the TOS process, the mass change of the granules was recorded during 38 consecutive redox cycles at 773 K and 0.2

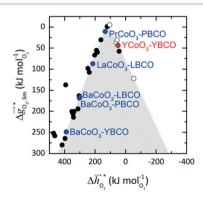


Fig. 1 The limiting Gibbs free energy change of the reduction reaction $\left(\Delta \overline{g}_{O_2,\, lim}^{\, o*}\right)$ as a function of the standard partial molar enthalpy change of the oxide reduction at 298 K $(\Delta \overline{h}_{O_2}^{o*})$. All data is at 1 bar. The coloured compositions represent perovskites, with blue marked materials limited by reduction, and red marked materials limited by oxidation. The ABO_3 -type perovskite written before the $AA'BO_3$ -type perovskite is the surface layer of the AA'BO₃-type perovskite. Black and white round symbols represent solid metal oxides and metal/metal oxide pairs limited by their oxide reduction and oxidation, respectively.

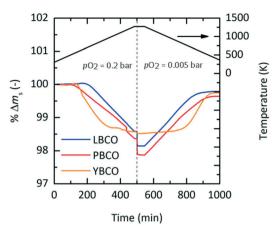


Fig. 2 Percent weight change vs. time from 298 to 1273 K, changing from 0.2 bar O_2 to 0.005 bar O_2 at the dashed line for LBCO, PBCO and YBCO. Heating and cooling rates were set to +2 K min⁻¹ and -2 K min⁻¹, respectively.

bar O_2 in N_2 for reduction and 563 K and 0.01 bar O_2 in Ar for oxidation using +20 K min⁻¹ and -20 K min⁻¹, respectively. For the TOP process, the mass change of the samples was recorded during 50 consecutive redox cycles at 873 K and 1 bar O_2 for reduction and 573 K and pO_2 and 0.2 bar O_2 in N_2 for oxidation using +20 K min⁻¹ and -20 K min⁻¹, respectively. To correct for buoyancy, blank runs were performed using the same measurement conditions employed for the experimental runs.

Packed bed reactor experiments

YBCO was tested in a packed bed reactor (PBR) setup, in accordance to previous experimental campaigns investigating gas-solid reactions.7 PBR experiments for the TOS process were performed with a vertically inserted quartz tube with a diameter of 1 cm in an electric furnace (Carbolite, VS 17/250). Gas flows were introduced from the bottom at 20 mL min⁻¹ and controlled using two electronic MFC (EL-Flow Select, Bronkhorst; accuracy ±0.5%). The oxygen concentration in the outlet gas flow was monitored by gas chromatography (GC, Agilent Micro GC, CP-4900). Two K-type thermocouples were inserted in the packed bed. From the bottom to the top, the reactor bed consisted of approximately 2 cm of compressed Al₂O₃ wool, 0.5 cm quartz sand with a diameter of 200-300 µm, 15 cm YBCO powder (8 g) mixed with quartz sand (16 g) and 1 cm pure quartz sand on top. A total of 15 consecutive cycles were executed. Reduction was always performed at 773 K and 0.2 bar O2 in N2 and held isothermally for 45 min. 30 min prior to cooling down, 0.2 bar O2 in N2 was exchanged to 0.01 bar O2 to ensure full gas replacement in the PBR. Oxidation was performed in 0.01 bar O2 at 563 K (cycles 1-4, 240 min isothermal), 523 K (cycles 5-7, 300 min isothermal), 603 K (cycles 8-12, 180 min isothermal) and 643 K (cycles 13-15, 180 min isothermal). 15 min prior to heating up to 773 K, 0.01 bar O2 in Ar was exchanged to 0.2 bar O₂ in N₂ to ensure full gas replacement.

For the TOP process an Inconel tube with a diameter of 1 cm in an electric furnace (Carbolite, VS 17/250) was employed. Gas flows were introduced from the bottom. Oxidation was performed at 5 mL min⁻¹, 0.2 bar O₂ in N₂ and 573 K during 240 min, while reduction was done at 10 mL min⁻¹, 873 K and 1 bar O₂ during 160 min. 60 and 30 min prior to cooling down from 873 to 573 K and prior to heating up from 573 to 873 K, respectively, respective gases were exchanged to ensure full gas replacement within the PBR. The oxygen concentration and gas flows were monitored by GC (Agilent Micro GC, CP-4900) and two electronic MFCs (EL-Flow Select, Bronkhorst; accuracy ±0.5%), respectively. An additional MFC (F-201CV, Bronkhorst; accuracy ±0.5%) was used to record absolute gas flows at the outlet of the PBR. Two K-type thermocouples were inserted in the packed bed. From the bottom to the top, the reactor bed consisted of approximately 2 cm of compressed Al₂O₃ wool, 0.5 cm quartz sand with a diameter of 200-300 µm, 13 cm YBCO granules (9 g) and 1 cm pure quartz sand on top. A total of 9 consecutive cycles were executed.

3. Results and discussion

Materials design

It was previously shown that metal oxide redox materials can be designed by plotting the limiting partial molar Gibbs free energy, $\Delta \overline{g}_{{\rm O}_{2,{\rm lim}}}^{\,\circ*}$ as a function of the descriptor $\Delta \overline{h}_{{\rm O}_{2}}^{\,\circ*}$, *i.e.* the partial molar enthalpy at room temperature.8,26 The limiting Gibbs free energy $\Delta \overline{\overline{g}}_{\mathrm{O}_{2,\mathrm{lim}}}^{\,\mathrm{o}*}$ is defined as the larger Gibbs free energy change of either the reduction or oxidation reaction at the conditions to be used in the process. In other words, it is Gibbs free energy change of the reaction step, either oxidation or reduction, which is limiting the overall redox process. Here, $\Delta \overline{g}_{O_{2, lim}}^{\, \circ *}$ and $\Delta \overline{h}_{O_{2}}^{\, \circ *}$ were computed for 32 solid metal oxides and six metal/metal oxide pairs at 563 K and 0.01 bar O_2 for oxidation and 773 K and 0.2 bar O_2 for reduction. Fig. 1 shows $\Delta \overline{g}_{O_2, lim}^{~o*}$ as a function of $\Delta \overline{h}_{O_2}^{o*}$ near the intersection of $\Delta \overline{g}_{0, \text{lim}}^{\circ *}$. Compounds with balanced, and thus favourable, redox behaviour are located near the top of the resulting volcano-shaped curve, i.e. they bind oxygen strongly enough at lower temperatures, but weakly enough to be reduced at moderately higher temperatures. Hence, the separation of oxygen from gas mixtures with a lower pO2 than the pO_2 during the metal oxide reduction can be characterized by an intermediately strong bonding of the lattice oxygen.⁸ Ideally, $\Delta \overline{g}_{O_2, lim}^{\circ *}$ is exergonic, which is only reached by RhO/Rh2O and PbO/Pb3O4, rare and toxic materials, respectively. Various perovskite compositions, including LBCO, PBCO and YBCO, were therefore computed to tailor inexpensive and non-toxic redox materials (cf. Fig. 1), by modelling stoichiometric ABO₃(010) and oxygen-deficient ABO_{2.5}(010) with the grid-based projector-augmented wave (GPAW) and atomic simulation environment (ASE) electronicstructure code. $^{19-21}$ $\Delta \overline{g}_{\mathrm{O}_{2,\,\mathrm{lim}}}^{\,\,\mathrm{o}*}$ and $\Delta \overline{h}_{\mathrm{O}_{2}}^{\,\,\mathrm{o}*}$ were then calculated

from the scaling of $\Delta G_{V}[O]$ and the redox energetics of the bulk oxides (cf. ESI†). 8,26,30 The surface layer of each AA'BO₃type perovskite is ideally composed of either ACoO3 or A'CoO₃. Fig. 1 displays the results of both possibilities, i.e. BaCoO₃-YBCO denotes that for YBCO, BaCoO₃ is on the surface. Interestingly, YCoO₃-YBCO is the only composition for which oxidation is the limiting reaction.

To verify the DFT-predicted trends and find a suitable material for thermochemical redox processes, LBCO, PBCO and YBCO were tested experimentally at different sets of redox conditions, presented in the next subsections.

Structural characterization

Fig. S1† displays the XRD spectra of sintered LBCO, PBCO and YBCO which were all identified to be in the perovskitetype structure. While LBCO was determined to be cubic, PBCO and YBCO revealed a tetragonal structure (cf. Table 1). From the XRD scans we see that decreasing ionic radii of the A-site dopants ($La^{3+} > Pr^{3+} > Y^{3+}$), result in increasing structural lattice distortions along with a shift towards higher 2θ angles (cf. Fig. S1†).31 Detected impurities originate from BaO in PBCO and from Y2O3 and YBaCo2O5 in YBCO.

Furthermore, the high-temperature XRD (HT-XRD) scans YBCO shown with Fig. S3† illustrate that this compound stays stable throughout the whole temperature range between 308 and 923 K in 0.2 and 0.01 bar O2 during heating and cooling, respectively.

Experimental screening

To investigate the established redox trends by DFT, TGA experiments were conducted with LBCO, PBCO and YBCO between 303 and 1273 K (cf. Fig. 2). A pO2 of 0.2 bar was chosen during heating, while 0.005 bar O2 was taken during cooling. Heating and cooling rates were both set to a low value of +2 K min⁻¹ and -2 K min⁻¹ such that thermodynamic equilibrium at all times can be assumed. We see that the three compositions all reduce and oxidize, with PBCO reaching the largest reduction extent.

To evaluate which of the three materials performs best, several performance indicators based on the experiments shown with Fig. 2 were applied (cf. Table 2). $T_{\text{red.start}}$ corresponds to the temperature at which the perovskites start to reduce, $\delta_{\rm end,reox}$ corresponds to the final oxygen nonstoichiometry of the reoxidized state and $T_{\rm net,eq}$ corresponds to the temperature that has to be exceeded during reduction to achieve a net $\Delta\delta$ (cf. Fig. S6†). The relative oxygen nonstoichiometries can be calculated from the data shown with Fig. 2 with:

$$\Delta \delta = \Delta m_{\rm s} \frac{M_{\rm s}}{M_{\rm O}} \tag{10}$$

where Δm_s is the relative weight loss at equilibrium, and M_s and $M_{\rm O}$ are the molar mass of the sample and of monoatomic oxygen, respectively. M_s was taken as the molar mass of the fully oxidized perovskite, ABO₃.

In summary, an ideal material for TOS has a low $T_{\text{net.eq}}$, resulting from a low $T_{\text{red,start}}$ and steep reduction and reoxidation slopes, i.e. fast reaction rates. Thus, the results suggest that YBCO is the most promising material from the three tested perovskites, having a T_{net,eq} of only 633 K (cf. Table 2).

The following two sections focus on experiments performed for thermochemical oxygen separation and production using YBCO at different sets of redox conditions, summarized in Table 7.

Thermodynamic analysis

In this section, we report the acquisition of $\Delta \overline{h}_{O_2}^{\circ}$ and $\Delta h_{\rm red}$ (cf. eqn (8) and (9)). $\Delta h_{\rm red}$ is directly related to the energy input required to drive the redox reaction,28 and is thus crucial to establish an energy balance. Fig. 3a shows the $\Delta\delta$ of YBCO at 0.005 bar O2, 0.01 bar O2, 0.05 bar O2 and 0.2 bar O₂ between 400 and 1073 K. As expected, we see that a higher pO2 shifts the thermodynamic equilibrium towards higher oxygen nonstoichiometries, i.e. lower reduction extents. The relative non-stoichiometry values $\Delta\delta$ can also be seen to converge towards a maximum reduction extent of around $\Delta\delta$ = 0.2 mol_{O} $\text{mol}_{\text{ABO}_3}^{-1}$, for all oxygen partial pressures.

To explain this maximum $\Delta \delta$, we need to consider the possible reduction pathways of the perovskite. Yttrium will always remain in a 3+ oxidation state, and barium in a 2+ state. The fully oxidised perovskite would then have oxidation states $Y_{0.5}^{3+}$ Ba $_{0.5}^{2+}$ Co $_{0.5}^{3+}$ Co $_{0.5}^{4+}$ O $_{3}^{2-}$, which could be considered as a solid solution of the two perovskites YCoO₃ and BaCoO3. The formation of oxygen vacancies could then take place by the reduction of Co⁴⁺ to Co³⁺, or Co³⁺ to Co²⁺. DFT calculations of YCoO3 reduction performed by the authors suggest an enthalpy of reduction of $\Delta \overline{h}_{O_2}^{\circ}$. This is much larger than the experimental values measured here (Fig. 3b) and also means that reduction of the perovskite via Co³⁺ to Co²⁺, would require higher temperatures. Indeed, experimental results obtained by Mei et al. for the reduction of $LaCoO_{3-\delta}$ via Co^{3+} to Co^{2+} transitions show only very small non-stoichiometry of δ < 0.001, for the most reducing

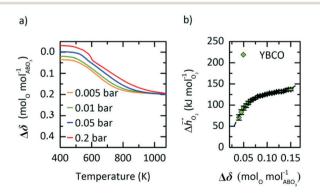


Fig. 3 a) Oxygen nonstoichiometry, δ , vs. temperature during reoxidation of YBCO at 0.005, 0.01, 0.05 and 0.2 bar O_2 , and b) standard partial molar enthalpy change of reduction, $\Delta \overline{h}_{\mathrm{O}_2}^{\circ}$, as a function of δ for YBCO. Error bars are $\pm 2\sigma$ confidence intervals of the slope of $-\ln(pO_2/p^\circ)$ versus T^{-1} (cf. Fig. S7†).

conditions seen in Fig. 3a) of T = 1073 K and $pO_2 = 0.005$ bar. On the other hand, reduction via Co⁴⁺ to Co³⁺ occurs readily at low temperatures in SrCoO₃,⁸ and other perovskites containing Co⁴⁺. ¹⁴ Additionally, these Co⁴⁺ perovskites exhibit non-stoicheometry even at ambient conditions when cooled under air, which explains the positive values of $\Delta\delta$ seen in Fig. 3a. It is therefore likely that the majority of the reduction takes place via Co⁴⁺ to Co³⁺ transitions. The value of $\Delta\delta \approx 0.2$ that the nonstoicheometry converges towards could then correspond to complete reduction of Co⁴⁺ to Co³⁺ giving a reduced perovskite with a stoichiometry of Y_{0.5}Ba_{0.5}CoO_{2.75}. This would mean that the zero point in Fig. 3a is Y_{0.5}Ba_{0.5}CoO_{2.95}, but as their may be some reduction of Co³⁺, we cannot with certainty state the exact stoichiometry.

From the data shown with Fig. 3a, S7† and eqn (8), $\Delta \overline{h}_{0}^{\circ}$ was deduced (cf. Fig. 3b). The good linear fit of the data shown with Fig. S7† $(R^2 > 0.94)$ confirms that $\Delta \overline{h}_{0}^{\circ}$, does not have strong dependence on temperature at the investigated conditions. Fig. 3b shows increasing $\Delta \overline{h}_{O_2}^{\circ}$ with increasing δ values, as observed by Babiniec et al.³⁵ for $\text{La}_x\text{Sr}_{1-x}\text{Co}_y\text{M}_{1-y}$ $O_{3-\delta}$ (M = Mn, Fe) and Takacs et al. ³⁶ for $La_{0.6}M_{0.4}MnO_{3-\delta}$ (M = Sr, Ca). The data in Fig. 3a may be kinetically limited at lower temperatures (T < 600 K), so that the Van't Hoff method may not be accurate in this region. The steep drop in Fig. 3b at lower $\Delta\delta$ < 0.07, may be a result of this. Symbols correspond to the measured δ values shown with Fig. 3, while the dashed line corresponds to a third-degree fit $(R^2 > 0.99)$. Error bars are $\pm 2\sigma$ confidence intervals of the slope of $-\ln(pO_2/p^\circ)$ versus T^{-1} (cf. Fig. S7†).

Applying eqn (9) to the result shown with Fig. 3b, we obtain a reduction enthalpy change of $\Delta h_{\rm red}$ = 121 kJ mol_{O₂}⁻¹ for YBCO at a $\Delta\delta$ of 0.064 mol_O mol_{ABO}.

Thermochemical oxygen separation

In the following sections we present results for a TOS processes using YBCO, including an energy balance, cycle stability runs in the TGA, and a packed bed reactor demonstration.

Energy balance

Using $\Delta h_{\rm red}$ and measured $c_{\rm p}$ (cf. Fig. S10†) of YBCO, an energy flow diagram for a redox cycle separating oxygen from a 1% O2-Ar stream is depicted in Fig. 4. The energy requirement per mole of O_2 , Δh_{cycle} , is calculated from the thermodynamic data of YBCO. c_p values of gas mixtures are approximated using the pure component properties obtained from NIST-JANAF tables.37 Assumed temperature/pressureswing operating conditions are 773 K and 0.2 bar O2 for reduction and 563 K and 0.01 bar O2 for oxidation. The choice of these conditions was made based on the knowledge acquired by the thermodynamic assessment (cf. Fig. 3). All involved equations used for the calculations of the different components are shown in ESI.† According to the net energy balance (cf. Fig. 4), the required energy to separate O2 from a 1% O_2 -Ar gas stream amounts to 1.82 MJ_{th} mol_{O2}⁻¹, while 1.83 MJ mol₀, ⁻¹ is available in form of sensible heat from gas and solid heat exchange. In the present study, TOS is considered for the cleaning of the outlet stream from a solar reactor, which can be assumed to be at around 1273 K at the currently employed solar reactor conditions with the state-ofthe-art material CeO2.38 This heat can be directly extracted in the TOS redox cycle using a printed-circuit heat exchanger, achieving an effectiveness of up to 97%.39 The extractable sensible heat $(\Delta h_{\text{avail}})$ is considered as the sensible heat of the 1% O2-Ar gas mixture between 1273 and 563 K and it amounts to 1.48 MJ mol₀, ⁻¹. In this scenario, no heating of the incoming 1% O_2 -Ar gas mixture is necessary ($\Delta h_{\rm hg} = 0$) and thus, Δh_{cycle} is 1.28 MJ mol₀, ⁻¹. Table 3 summarizes the energy requirements, assuming 50% solid and 80% gas heat recuperations and no heat recuperation at all, for O2-Ar gas mixtures of 0.005, 0.01 and 0.05 bar O2 coming from a solar reactor at 1273 K. We see from Table 3 that with 0.005 and

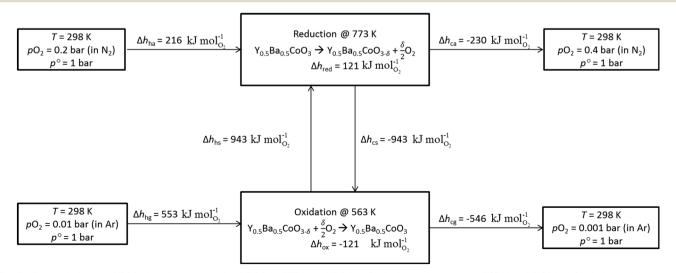


Fig. 4 Energy balance of TOS via redox cycle using YBCO at 773 K and 0.2 bar O2 during reduction and 563 K and 0.01 bar O2 during oxidation.

0.01 bar O₂, $\Delta h_{\rm avail} > \Delta h_{\rm cycle}$. With 0.05 bar O₂, $\Delta h_{\rm avail} <$ Δh_{cycle} when no gas and solid heat recuperation is applied. This is the only tested scenario where an additional energy input would be necessary to run the process. Oxidation at 0.05 bar O2 is more energy intensive than in case of lower pO₂ values, because of the larger amount of reactive material to be heated. At high pO2 values, a higher quantity of oxygen has to be separated per redox cycle. This has a significant impact on the energy needed to heat the solid, $\Delta h_{\rm hs}$, which increases from 334 and 668 to 3342 kJ mol₀. ⁻¹ for 0.005 and 0.01 to 0.05 bar O_2 , respectively. The larger $\Delta\delta$ at higher pO_2 values, i.e. here it increases from 0.053 and 0.064 to 0.090 mol_O mol_{ABO₂}⁻¹ upon changing from 0.005 and 0.01 to 0.05 bar O2, respectively, cannot compensate for the increased $\Delta h_{\rm hs}$. Furthermore, the energy needed to heat the air, $\Delta h_{\rm ha}$, increases from 77 and 153 to 767 kJ molo, 150 of 0.005 and 0.01 to 0.05 bar O2, respectively, since the required molar amount of air is roughly half the molar amount of perovskite, according to Hänchen et al. 7 Δh_{avail} at 0.05 bar O_2 increases slightly compared to 0.01 and 0.005 bar O_2 , because the c_p of O2 (29.4 J mol K-1) is larger than theone of Ar (20.8 J mol K⁻¹). Summarizing, the energy balance is favored at rather low pO_2 and heat recuperation is essential at high pO_2 . However, at low pO_2 oxidation rates are expected to decrease.

Thermochemical cycling experiments for TOS

To verify if YBCO is suitable in an industrial implementation, porous granules were produced according to the procedure described in the experimental section. Production of the granules was much more efficient than pellet production, while still having a diameter large enough to avoid significant channelling effects and pressure drops if employed in a PBR. Fig. 5 shows the $\Delta\delta$ of YBCO granules cycled 38 times in a TGA between 563 K and 0.01 bar O2 during oxidation and 773 K and 0.2 bar O2 during reduction. We see that $\Delta \delta$ stays stable throughout all executed cycles, reaching values of around 0.061 mol_O mol_{ABO}. It is assumed that the slightly lower value of 0.061 vs. 0.064 mol_O mol_{ABO}, observed with the pellets (cf. Fig. 3), might originate from the two additional small XRD peaks around

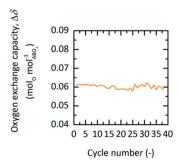


Fig. 5 Oxygen exchange capacity, $\Delta \delta$, vs. cycle number of YBCO granules at 0.01 bar O2 and 563 K during oxidation and 0.2 bar O2 at 773 K during reduction for TOS.

28° (indexed by *, cf. Fig. S2†), which might be due to an ordering of excess oxygen forming a superstructure. 40

The oxygen storage capacity $\Delta \delta$, of YBCO, looks very good when compared to other materials studied for TOS and TOP processes, as can be seen in Table 4. YBCO has approximately double the oxygen storage capacity of both SrFeO3 and SrCoO₃ over similar temperature and pressure ranges. Cu₂O (ref. 7) on the other hand, has a large oxygen storage due to the stoichiometric nature of the reaction, but requires much higher temperatures. It is clear that the low temperature redox conditions employed with YBCO would not work for Cu₂O, since Ezbiri et al.8 previously showed that Cu₂O irreversibly oxidizes to CuO, when cycled between 900 and 600 K at 0.2 and 0.035 bar O2, respectively. We therefore concluded for applications in the lower temperature range investigated here that YBCO has superior performance compared to the state of the art.

Packed bed reactor experiments for TOS

PBR experiments were executed to test the potential of YBCO for an industrial setting. To minimize the effects of channelling and pressure drops in the PBR, the previously discussed YBCO granules were initially used. However, the available setup did not allow for YBCO granules to reach thermodynamic equilibrium, possibly due to mass transfer limitations in the gas phase and the lower specific area of the granules (cf. Fig. S13†). Hence, the results executed with the YBCO powder are shown. Fig. 6 displays the measured pO_2 at the outlet of the PBR along with the temperature and the pO2 at the inlet of the PBR as a function of time. Fig. 6 shows oxidation at 563 K; Fig. S12† shows oxidation at 523, 603 and 643 K. For all these oxidation temperatures,

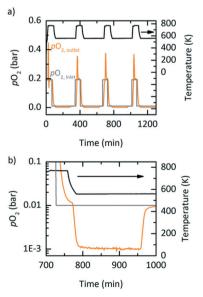


Fig. 6 a) A number of cycles in the PBR showing temperature, pO2 at the outlet (orange) and pO2 at the inlet of the PBR (grey) using YBCO at 563 K and 0.01 bar pO2 for oxidation and 773 K and 0.2 bar pO2 for reduction. b) A single cycle from a).

thermodynamic equilibrium was reached during oxidation. We see from Fig. 6b that during oxidation, the $p{\rm O}_2$ drops from 0.01 bar ${\rm O}_2$ to 0.001 bar ${\rm O}_2$ over a steady period of 175 minutes, which can clearly be attributed to the reoxidation reaction of YBCO. 0.001 bar ${\rm O}_2$ is the maximum allowable ${\rm O}_2$ concentration in recycled inert gas. Reaching 955 minutes in Fig. 6b, the $p{\rm O}_2$ quickly increases back to 0.01 bar ${\rm O}_2$, which means that the whole reactor bed has been oxidized. An amount of 3.13 litres of ${\rm O}_2$ is separated during the oxidation at 563 K, in each of the cycles, corresponding to 0.064 molo ${\rm mol}_{{\rm ABO}_3}^{-1}$ as observed in the TGA (cf. Table 5). Table 5 further shows the oxidation time and $\Delta\delta$ of the PBR experiments at oxidation temperatures of 523, 603 and 643 K. As expected, we see that the time until oxidation and the $\Delta\delta$ values increase with decreasing oxidation temperature.

Thermochemical oxygen production

In the following sections we present results for a TOP processes using YBCO, again including an energy balance, cycle stability runs in the TGA, and a packed bed reactor demonstration of the process.

Energy balance

This section presents the energy balance established for TOP, similar to the energy balance illustrated in Fig. 4. Here, a different set of redox conditions is employed (cf. Table 7) with 673 K and 0.2 bar $\rm O_2$ during oxidation and 873 K and 1 bar $\rm O_2$ during reduction. Since this process is not coupled to another high-temperature technology, $\Delta h_{\rm hg}$ cannot be taken as zero, i.e.

$$\Delta h_{\mathrm{cycl}} = \Delta h_{\mathrm{ha}} + \Delta h_{\mathrm{red}} + \Delta h_{\mathrm{hs}} + \Delta h_{\mathrm{hg}} - \left(\frac{\Delta h_{\mathrm{ca}} + \Delta h_{\mathrm{cg}} + \Delta h_{\mathrm{osi}}}{1.25} + \frac{\Delta h_{\mathrm{cs}}}{2}\right)$$

assuming 80 and 50% gas and solid heat recuperation, respectively. With YBCO and the employed redox conditions, $\Delta h_{\rm cycle}$ amounts to 377 kJ ${\rm mol_{O_2}}^{-1}$, while state-of-the-art cryogenic distillation requires a thermal energy of 78 kJ ${\rm mol_{O_2}}^{-1}$, assuming a conversion factor of 0.4 (for state-of-the-art coal-fired power plants⁴¹) between electric ($E_{\rm el}$) and thermal energy ($E_{\rm th}$) (cf. Table 6). We thus conclude that TOP using YBCO at the used redox conditions might only compete with cryogenic distillation if coupled to another high-temperature technology, such as oxy-fuel combustion, 2,10 where the cryogenic air separation unit is a major economic drawback. Other possible technologies, where the wider use of pure oxygen could reduce ${\rm CO_2}$ emissions, include integrated gasification combined cycles and the partial oxidation of methane. 5

Thermochemical cycling experiments for TOP

Using TGA, 50 redox cycles were performed with YBCO granules at 573 K and 0.2 bar O_2 during oxidation and at 873 K and 1 bar O_2 during reduction. We see that $\Delta\delta$ stays stable at around 0.1 mol_O mol_{ABO₃}⁻¹ throughout all executed redox cycles (*cf.* Fig. 7).

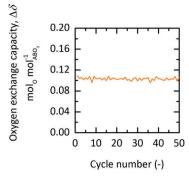


Fig. 7 Oxygen exchange capacity, $\Delta\delta$, vs. cycle number of YBCO at 0.2 bar O₂ and 573 K during oxidation and 1 bar O₂ at 873 K during reduction for TOP.

Packed bed reactor experiments for TOP

Exploratory experimental runs with a packed bed reactor were executed using YBCO granules, where in this case the granules were indeed suitable as they did reach thermodynamic equilibrium. This was likely due to the higher temperature operation when compared to TOS. Fig. 8 shows the absolute gas flow and pO_2 measured at the outlet of the reactor vs. time for nine consecutive redox cycles. 5 and 10 mL min⁻¹ gas flows were used during oxidation and reduction, respectively. The changes in absolute gas flow during oxidation and reduction are a clear indication of the reactivity of YBCO. We further note that the pO_2 does not change during reduction since it is at a 100% at the inlet already. During oxidation, however, a decrease in pO_2 from the set point of 0.2 bar O_2 is observed.

Table 7 gives an overview of the used redox conditions, the resulting $\Delta\delta$ and the required energies of TOS and TOP in comparison to state-of-the-art cryogenic distillation.

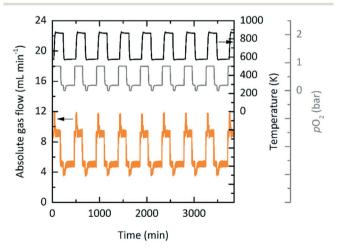


Fig. 8 Consecutive redox cycles for TOP with a packed-bed reactor. Absolute gas flow (orange) and pO_2 (grey) measured at the outlet of the packed bed reactor vs. time using YBCO at 573 K and 0.2 bar O_2 during oxidation and 873 K and 1 bar O_2 during reduction. The absolute gas flows at the inlet were set to 5 and 10 mL min⁻¹ during oxidation and reduction, respectively.

Paper

Using a descriptor-based design principle, we predicted promising redox thermodynamics of the three AA'BO₃-type perovskites La_{0.5}Ba_{0.5}CoO_{3- δ}, Pr_{0.5}Ba_{0.5}CoO_{3- δ} and Y_{0.5}Ba_{0.5}-CoO_{3- δ}. These compounds were further evaluated using dynamic oxygen exchange measurements, with Y_{0.5}Ba_{0.5}CoO_{3- δ} being identified as most ideal for thermochemical oxygen separation and production. This compound shows the fastest reaction rates, lowest reaction temperatures and highest oxygen exchange capacities compared to those of the tested perovskites in this work, and the state-of-the-art materials SrCoO₃ and SrFeO₃. The TOS processes was experimentally demonstrated using Y_{0.5}Ba_{0.5}CoO_{3- δ} and an energy balance indicates that the it could in principle be driven by process waste heat from the solar splitting of CO₂/H₂O. Finally, consecutive redox cycles executed with a packed-bed reactor indicated the potential of

both thermochemical oxygen production and separation for

Nomenclature

LBCO

L min⁻¹

 $pO_{2,ox,in}$

 $pO_{2,ox,out}$

 $pO_{2,red,in}$

 $pO_{2,red,out}$

 $M_{\rm O}$

 $M_{\rm s}$ $m_{\rm s}$

 pO_2

scale-up and industrial implementation.

ca	Cooling air
cg	Cooling gas
cs	Cooling solid
DFT	Density functional theory
$E_{ m el}$	Electrical energy (kJ mol _{O2} ⁻¹)
E_{th}	Thermal energy (kJ mol _{O2} ⁻¹)
$\Delta \overline{g}_{\mathrm{O}_2}^{\circ}$	Standard partial molar Gibbs free energy change
_	of metal oxide reduction (kJ mol _{O2} ⁻¹)
$\Delta \overline{g}_{\mathrm{O}_{2}}^{\circ *}$	Computed standard partial molar Gibbs free
	energy change of metal oxide reduction (kJ $mol_{O_2}^{-1}$)
$\Delta \overline{g}_{\mathrm{O}_{2},\mathrm{lim}}^{\circ}$	Limiting standard partial molar Gibbs free energy
	change of metal oxide reduction (kJ $mol_{O_2}^{-1}$)
$\Delta G_{ m V}[{ m O}]$	Gibbs free energy change of oxygen vacancy
	formation (eV)
Δh	Enthalpy (kJ $\operatorname{mol}_{\operatorname{O}_2}^{-1}$)
$\Delta \overline{h}_{\mathrm{O}_2}^{\circ}$	Standard partial molar enthalpy change of metal
	oxide reduction (kJ mol _{O2} ⁻¹)
ha	Heating air
hg	Heating gas
hs	Heating solid

$p_{\rm tot}$	System pressure (bar)
p°	Standard pressure (bar)
PBCO	$Pr_{0.5}Ba_{0.5}CoO_{3-\delta}$
PBR	Packed bed reactor
R	Universal gas constant (J mol _O ⁻¹ K ⁻¹)
SEM	Scanning electron microscopy
$T_{\rm ox}$	Oxidation temperature (K)
$T_{\rm net,eq}$	Net equilibrium temperature (K)
$T_{\rm red}$	Reduction temperature (K)
TOP	Thermochemical oxygen production
TOS	Thermochemical oxygen separation
TGA	Thermogravimetric analyzer
ΔT	Difference between oxidation and reduction
	temperatures (K)
XRD	X-ray diffraction
YBCO	$Y_{0.5}Ba_{0.5}CoO_{3-\delta}$
δ	Degree of oxygen nonstoichiometry at
	thermodynamic equilibrium (-)
$\delta_{ m ox}$	Degree of oxygen nonstoichiometry after oxidation
	at thermodynamic equilibrium (-)
$\delta_{ m red}$	Degree of oxygen nonstoichiometry after
	reduction at thermodynamic equilibrium (-)

Conflicts of interest

 $\Lambda\delta$

There are no conflicts of interest to declare.

Acknowledgements

We gratefully acknowledge the financial support by the Swiss Competence Center Energy & Mobility, the Helmholtz-Gemeinschaft Deutscher Forschungszentren (Virtuelles Institut SolarSyngas), the European Research Council under the European Union's ERC Advanced Grant (SUNFUELS-No. 320541), and the European Union's Horizon 2020 research and innovation programme (Project SFERA-III – Grant Nr. 823802). The DFT calculations were performed at the High-Performance Computation cluster of ETH Zürich.

Oxygen exchange capacity (mol_O mol_{ABO₃}⁻¹)

References

- 1 R. M. Thorogood, Gas Sep. Purif., 1991, 5, 83-94.
- 2 K. Zhang, J. Sunarso, Z. P. Shao, W. Zhou, C. H. Sun, S. B. Wang and S. M. Liu, *RSC Adv.*, 2011, 1, 1661–1676.
- 3 T. Burdyny and H. Struchtrup, Energy, 2010, 35, 1884-1897.
- 4 A. R. Smith and J. Klosek, Fuel Process. Technol., 2001, 70, 115-134.
- 5 P. M. Geffroy, J. Fouletier, N. Richet and T. Chartier, *Chem. Eng. Sci.*, 2013, 87, 408–433.
- 6 F. G. Kerry, Industrial Gas Handbook: Gas Separation and Purification, CRC Press, 2010.
- 7 M. Hänchen, A. Stiel, Z. R. Jovanovic and A. Steinfeld, *Ind. Eng. Chem. Res.*, 2012, 51, 7013–7021.
- 8 M. Ezbiri, K. M. Allen, M. E. Gàlvez, R. Michalsky and A. Steinfeld, *ChemSusChem*, 2015, **8**, 1966–1971.

 $La_{0.5}Ba_{0.5}CoO_{3-\delta}$

Molar mass of O (g mol⁻¹)

Mass of reactive sample (mg)

Oxygen partial pressure (bar)

experiment during oxidation (bar)

experiment during oxidation (bar)

experiment during reduction (bar)

experiment during reduction (bar)

Molar mass of reactive sample (g mol⁻¹)

Standard liter per minute at 298 K and 1 bar

Oxygen partial pressure at the inlet of the

Oxygen partial pressure at the outlet of the

Oxygen partial pressure at the inlet of the

Oxygen partial pressure at the out let of the

- 9 B. Bulfin, J. Lapp, S. Richter, D. Gubàn, J. Vieten, S. Brendelberger, M. Roeb and C. Sattler, Chem. Eng. Sci., 2019, 203, 68-75.
- 10 K. Krishnamurthy, D. Acharya and F. Fitch, Pilot-Scale Demonstration of a Novel, Low-Cost Oxygen Supply Process and its Integration with Oxy-Fuel Coal-Fired Boilers, 2008, Tech. rep., Boc Group, Inc., The (A Delaware Corp).
- 11 B. Moghtaderi, Energy Fuels, 2010, 24, 190-198.
- 12 J. Vieten, B. Bulfin, F. Call, M. Lange, M. Schmücker, A. Francke, M. Roeb and C. Sattler, J. Mater. Chem. A, 2016, 4(35), 13652-13659.
- 13 J. Vieten, B. Bulfin, P. Huck, M. Horton, D. Guban, L. Zhu, Y. Lu, K. A. Persson, M. Roeb and C. Sattler, Energy Environ. Sci., 2019, 12(4), 1369-1384.
- 14 J. Dou, E. Krzystowczyk, A. Mishra, X. Liu and F. Li, ACS Sustainable Chem. Eng., 2018, 6(11), 15528-15540.
- 15 A. M. Narayanan, R. Parasuraman and A. M. Umarji, Ind. Eng. Chem. Res., 2018, 57(43), 14749-14757.
- 16 C. L. Muhich, B. W. Evanko, K. C. Weston, P. Lichty, X. H. Liang, J. Martinek, C. B. Musgrave and A. W. Weimer, Science, 2013, 341, 540-542.
- 17 L. J. Venstrom, R. M. De Smith, Y. Hao, S. M. Haile and J. H. Davidson, Energy Fuels, 2014, 28, 2732-2742.
- W. Villasmil, M. Brkic, D. Wuillemin, A. Meier and A. Steinfeld, J. Sol. Energy Eng., 2014, 136, 011016.
- 19 J. Enkovaara, C. Rostgaard, J. J. Mortensen, J. Chen, M. Dulak, L. Ferrighi, J. Gavnholt, C. Glinsvad, V. Haikola, H. A. Hansen, H. H. Kristoffersen, M. Kuisma, A. H. Larsen, L. Lehtovaara, M. Ljungberg, O. Lopez-Acevedo, P. G. Moses, J. Ojanen, T. Olsen, V. Petzold, N. A. Romero, J. Stausholm-Moller, M. Strange, G. A. Tritsaris, M. Vanin, M. Walter, B. Hammer, H. Hakkinen, G. K. H. Madsen, R. M. Nieminen, J. Nørskov, M. Puska, T. T. Rantala, J. Schiotz, K. S. Thygesen and K. W. Jacobsen, J. Phys.: Condens. Matter, 2010, 22, 253202.
- 20 J. J. Mortensen, L. B. Hansen and K. W. Jacobsen, Phys. Rev. B: Condens. Matter Mater. Phys., 2005, 71, 035109.
- 21 S. R. Bahn and K. W. Jacobsen, Comput. Sci. Eng., 2002, 4, 56-66.
- 22 B. Hammer, L. B. Hansen and J. K. Nørskov, Phys. Rev. B: Condens. Matter Mater. Phys., 1999, 59, 7413-7421.

- 23 A. J. Medford, J. Wellendorff, A. Vojvodic, F. Studt, F. Abild-Pedersen, K. W. Jacobsen, T. Bligaard and J. K. Norskov, Science, 2014, 345, 197-200.
- 24 M. García-Mota, A. Vojvodic, H. Metiu, I. C. Man, H. Y. Su, J. Rossmeisl and J. K. Nørskov, ChemCatChem, 2011, 3, 1607-1611.
- M. Ezbiri, M. Takacs, B. Stolz, J. Lungthok, A. Steinfeld and 25 R. Michalsky, J. Mater. Chem. A, 2017, 5(29), 15105-15115.
- R. Michalsky, V. Botu, C. M. Hargus, A. A. Peterson and A. Steinfeld, Adv. Energy Mater., 2014, 4, 1401082.
- 27 I. Barin, Thermochemical Data of Pure Substances, 1993.
- J. R. Scheffe and A. Steinfeld, *Energy Fuels*, 2012, **26**, 1928–1936.
- L. B. Pankratz, Thermodynamic Properties of Elements and Oxides, USA, 1982.
- 30 R. Michalsky, D. Neuhaus and A. Steinfeld, Energy Technol., 2015, 3, 784-789.
- 31 I. O. Troyanchuk, N. V. Kasper, D. D. Khalyavin, H. Szymczak, R. Szymczak and M. Baran, Phys. Rev. B: Condens. Matter Mater. Phys., 1998, 58, 2418-2421.
- 32 T. Nakajima, M. Ichihara and Y. Ueda, J. Phys. Soc. Jpn., 2005, 74, 1572-1577.
- 33 M. M. Seikh, V. Pralong, O. I. Lebedev, V. Caignaert and B. Raveau, J. Appl. Phys., 2013, 114, 013902.
- 34 T. Vogt, P. M. Woodward, P. Karen, B. A. Hunter, P. Henning and A. R. Moodenbaugh, Phys. Rev. Lett., 2000, 84, 2969-2972.
- 35 S. M. Babiniec, E. N. Coker, J. E. Miller and A. Ambrosini, Sol. Energy, 2015, 118, 451-459.
- 36 M. Takacs, M. Hoes, M. Caduff, T. Cooper, J. R. Scheffe and A. Steinfeld, Acta Mater., 2016, 700-710.
- 37 W. G. Mallard and P. J. Linstrom, webbook.nist.gov, retrieved February 2016), Gaithersburg MD, 1998.
- 38 D. Marxer, P. Furler, J. Scheffe, H. Geerlings, C. Falter, V. Batteiger, A. Sizmann and A. Steinfeld, Energy Fuels, 2015, 29, 3241-3250.
- 39 J. H. Jeong, L. S. Kim, J. K. Lee, M. Y. Ha, K. S. Kim and Y. C. Ahn, presented in part at the ASEM Turbo Expo 2007: Power for land, sea and air, 2007.
- 40 W. Zhou, C. T. Lin and W. Y. Liang, Adv. Mater., 1993, 5, 735-738.
- 41 J. Bugge, S. Kjaer and R. Blum, Energy, 2006, 31, 1437-1445.