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## Electronic and structural characterisation of polycrystalline platinum disulfide thin films†

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We employ a combination of scanning tunnelling microscopy (STM) and scanning tunnelling spectroscopy (STS) to investigate the properties of layered  $\text{PtS}_2$ , synthesised via thermally assisted conversion (TAC) of a metallic Pt thin film. STM measurements reveal the 1T crystal structure of  $\text{PtS}_2$ , and the lattice constant is determined to be  $3.58 \pm 0.03 \text{ \AA}$ . STS allowed the electronic structure of individual  $\text{PtS}_2$  crystallites to be directly probed and a bandgap of  $\sim 1.03 \text{ eV}$  was determined for a  $3.8 \text{ nm}$  thick flake at liquid nitrogen temperature. These findings substantially expand understanding of the atomic and electronic structure of  $\text{PtS}_2$  and indicate that STM is a powerful tool capable of locally probing non-uniform polycrystalline films, such as those produced by TAC. Prior to STM/STS measurements the quality of synthesised TAC  $\text{PtS}_2$  was analysed by X-ray photoelectron spectroscopy (XPS) and Raman spectroscopy. These results are of relevance to applications-focussed studies centred on  $\text{PtS}_2$  and may inform future efforts to optimise the synthesis conditions for thin film  $\text{PtS}_2$ .

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## 1. Introduction

The investigation of new two-dimensional (2D) materials has been of increasing interest since the experimental discovery of graphene<sup>1</sup> with many reports focusing on their interesting and unique properties which are of great interest for both fundamental research<sup>2–4</sup> and applications.<sup>5–8</sup> Due to the absence of a natural bandgap in graphene, research attention has turned towards 2D materials with a bandgap, particularly those which are stable in ambient conditions. Transition Metal Dichalcogenides (TMDs)<sup>9–11</sup> are a family of 2D materials that can be of semiconducting character with a layer-dependent bandgap.<sup>12,13</sup> The chemical formula for TMDs is generally expressed as  $\text{MX}_2$ , where M is a transition metal and X is a chalcogen (S, Se or Te). Each discrete layer of a TMD typically consists of a hexagonal layer of the M atoms which are held between two hexagonal layers of the X atoms. The layered structure of the crystals stems from van der Waals interactions between adjacent TMD layers. TMDs can have a variety of electronic properties, such as metallic, semiconducting, and superconducting, depending on their composition, crystal structure and thickness. The properties of 2D TMDs

are often significantly different from their bulk forms and can be fine-tuned for numerous applications.<sup>14</sup> Due to this versatility, TMDs are poised to play important roles in the fields of electronics,<sup>15</sup> optoelectronics,<sup>16</sup> electrochemical sensors<sup>17</sup> and catalysis,<sup>18</sup> to name but a few. One of the less-explored branches of TMD materials is noble-transition-metal dichalcogenides (NTMDs),<sup>19</sup> ( $\text{MX}_2$ ; M = Pt and Pd, and X = S, Se and Te) which are predicted to exhibit different structural and electronic properties in comparison with well-studied  $\text{MoS}_2$ .<sup>20,21</sup> One member of the NTMD group is platinum disulfide ( $\text{PtS}_2$ ), of which ultrathin films have recently been synthesised.<sup>22–24</sup> Theoretical and experimental studies have revealed that the size of the bandgap in  $\text{PtS}_2$  decreases with increasing film thickness,<sup>23–26</sup> going from  $\sim 1.7 \text{ eV}$  for monolayer  $\text{PtS}_2$  to  $\sim 0.25 \text{ eV}$  for the bulk state. Indirect bandgaps were observed for all layer thicknesses as opposed to  $\text{MoS}_2$  and  $\text{WS}_2$  where the monolayer forms possess a direct bandgap.<sup>25</sup> Relative positioning of the valence band maximum (VBM) and conduction band minimum (CBM) with respect to the Fermi level indicates that  $\text{PtS}_2$  is a p-type semiconductor.<sup>23–25</sup> Initial investigations of its properties suggest that  $\text{PtS}_2$  holds promise for use in diverse applications including photodetectors,<sup>27,28</sup> electrocatalysis,<sup>29</sup> photocatalysis,<sup>30</sup> thermoelectric devices,<sup>31</sup> nonlinear optics and laser photonics.<sup>32</sup> It should be noted that while  $\text{PtS}_2$  is a long-established material<sup>33,34</sup> its synthesis is not trivial. Under typical reaction conditions for the sulfurisation of Pt, the non-layered material  $\text{PtS}$ , rather than  $\text{PtS}_2$ , has been reported as the dominant product.<sup>35</sup> When considering the sulfurisation of Pt, the reaction parameters must be considered and carefully tuned to obtain an appreciable yield of  $\text{PtS}_2$ .<sup>34</sup>

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Transmission Electron Microscopy (TEM) is a powerful and widely-used technique for the characterisation of 2D materials. It is capable of atomic resolution and can be used to investigate the crystal structure, as well as the presence and quantity of defects.<sup>24,27,36</sup> However, direct measurements of the electronic band structure are not straightforward *via* TEM. Another limit associated with TEM is the requirement of additional sample preparation, namely that the material to be investigated must first be transferred onto a TEM grid, a potentially destructive process which can introduce contaminants. Moreover, recent reports have indicated that damage by the TEM electron beam can be a problematic issue for TMDs.<sup>37-39</sup> On the other hand, scanning tunnelling microscopy and spectroscopy (STM/STS) are non-destructive techniques, which can be used to map morphology and probe the electronic structure of samples<sup>40-43</sup> and, provided a conductive pathway can be established through the sample, do not require a transfer process.

In this letter we study the electronic properties and atomic structure of  $\text{PtS}_2$  films synthesised *via* thermally assisted conversion (TAC) of pre-deposited Pt films. The  $\text{PtS}_2$  films were also characterised by X-ray photoelectron spectroscopy (XPS) and Raman spectroscopy. We carried out STM measurements to investigate the crystal structure of  $\text{PtS}_2$  and determined the in-plane lattice constant to be  $3.58 \pm 0.03 \text{ \AA}$ . Scanning Tunneling Spectroscopy (STS) was employed in order to probe the electronic properties of the  $\text{PtS}_2$ . From STS measurements the electronic bandgap of a 3.8 nm thick  $\text{PtS}_2$  crystal, measured at 77 K, was determined to be around 1.03 eV.

## 2. Experimental

$\text{PtS}_2$  samples were synthesised through the TAC of pre-deposited Pt films. A 2 nm thick metal Pt film was deposited by argon-ion sputtering (Gatan Precision Etching and Coating System (PECS)) onto a conductive pyrolytic carbon (PyC) film supported on a  $\text{SiO}_2/\text{Si}$  substrate.<sup>44</sup>

A low-temperature STM from Createc was utilised in this study. The STM and STS measurements were conducted in ultra-high vacuum (UHV) at liquid-nitrogen temperature (77 K). Single-crystalline W tips<sup>45</sup> with a (001) orientation, prepared by electrochemical etching in NaOH solution, were employed. These tips were sharpened in UHV by argon-ion sputtering to remove surface oxides. All STM images were recorded in constant-current mode. While little thermal noise is expected at low temperature, it should be noted that no drift corrections have been applied to any STM images presented here. STS spectra ( $dI/dV$ ) were acquired by differentiation of the current-voltage  $I(V)$  spectra. Prior to and post grid STS measurements, the surface positions were verified by STM imaging to confirm the location had remained fixed.

XPS spectra were recorded using a PHI VersaProbe III instrument. Spectra were recorded at a pass energy of 26 eV with a monochromatic Al  $\text{K}\alpha$  (1486.6 eV) source and dual-beam charge neutralisation with charge correction to the C 1s core-level at 284.8 eV using the CasaXPS software. XRD was acquired with a Bruker D8 Discover, equipped with a monochromated Cu K-alpha source.

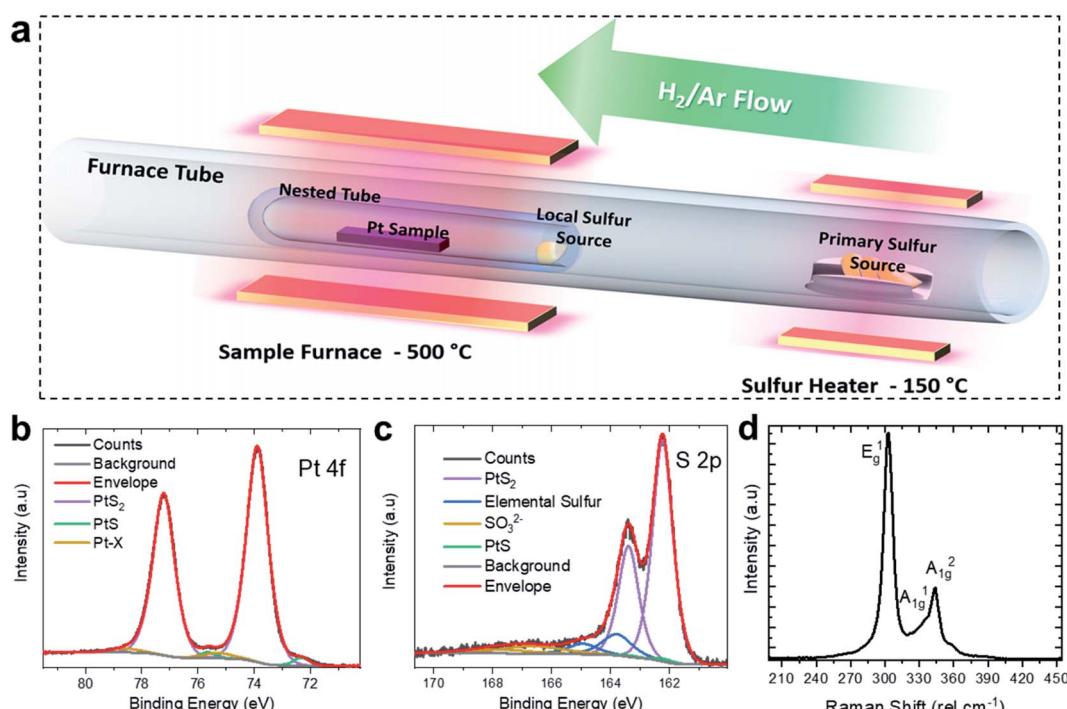


Fig. 1 Characterization of the  $\text{PtS}_2$  film. (a) Schematic of the sulfurization process with the nested tube configuration. (b) & (c) Fitted XPS of  $\text{PtS}_2$  film demonstrating the Pt 4f and S 2p core levels respectively. (d) Raman spectrum of a  $\text{PtS}_2$  film presenting the three Raman-active vibrational modes. These two methods help to confirm the successful synthesis of  $\text{PtS}_2$ .



Raman spectroscopy was performed with a WITec Alpha 300 R confocal Raman microscope and a 532 nm excitation energy. Spectra were gathered by averaging across  $>10$  discrete point spectra with a laser power  $< 200 \mu\text{W}$ .

### 3. Results and discussion

#### 3.1. Synthesis of 1T-PtS<sub>2</sub>

The synthesis of the PtS<sub>2</sub> films was performed through TAC of pre-deposited Pt films. It is notable that the synthesis temperatures of Pt-based TMDs by TAC are much lower than those associated with W or Mo films,<sup>13,20,46,47</sup> which lowers the thermal budget and broadens the potential application space by enabling growth on diverse substrates. A 2 nm thick metal Pt film was deposited by argon-ion sputtering onto a PyC film on a SiO<sub>2</sub>/Si substrate.<sup>44</sup> The Pt film was annealed for 1 hour at 500 °C in a quartz-tube furnace with a localised high saturation of S vapour and a 150 sccm 10% H<sub>2</sub>/Ar flow at 300 mbar (see Fig. 1(a)). This arrangement was based on the nested configuration described by Xu *et al.*<sup>36</sup> An additional local sulfur charge is placed at the mouth of the inner tube, at the edge of the primary hot-zone, to increase the S vapour concentration within. Previous research has shown that at lower pressure, or lower sulfur content, films with PtS stoichiometry are obtained.<sup>35</sup> This simple method avoids the need for exotic,

expensive or hazardous precursors and is essentially a scalable process suitable for producing large-area polycrystalline films.

High-resolution XPS, Raman spectroscopy and X-ray diffraction (XRD) were utilised to confirm the synthesis of PtS<sub>2</sub>. The characteristic XPS core-levels for PtS<sub>2</sub> are the Pt 4f and the S 2p (see Fig. 1(b) and (c)), both regions show dominant doublet pairs corresponding to PtS<sub>2</sub>. The Pt 4f<sub>7/2</sub> peak is centred at 73.9 eV while the S 2p<sub>3/2</sub> is at 162.2 eV, consistent with the expected values.<sup>48</sup> From the fitting, the small shoulder at  $\sim 72$  eV is attributed to a low level of Pt(II) contribution. Comparing the relative peak areas of the PtS<sub>2</sub> components yields a stoichiometry of PtS<sub>1.89</sub>.<sup>36</sup> As noted earlier, a lack of sulfur during synthesis will result in the synthesis of PtS, but despite evidence of an excess of elemental sulfur, there are still indications of trace PtS present in the films. The formation of polycrystalline PtS<sub>2</sub> is also supported by XRD of the thin-film, as shown in Fig. S1 of the ESI.†

Raman spectroscopy using 532 nm excitation yielded the spectrum shown in Fig. 1(d). PtS<sub>2</sub> has three main characteristic Raman modes, all of which are clearly distinguishable in this spectrum. These are the out-of-plane E<sub>g</sub><sup>1</sup> mode at 303 cm<sup>-1</sup> and the in-plane modes A<sub>1g</sub><sup>2</sup> and A<sub>1g</sub><sup>1</sup> at 343 and  $\sim 334$  cm<sup>-1</sup>, respectively. The small shoulder  $\sim 360$  cm<sup>-1</sup> is likely due to the polycrystalline nature of the film or the presence of PtS, as this feature is not observed for in the Raman spectra of exfoliated PtS<sub>2</sub>.<sup>24,31</sup> Otherwise this spectrum is consistent

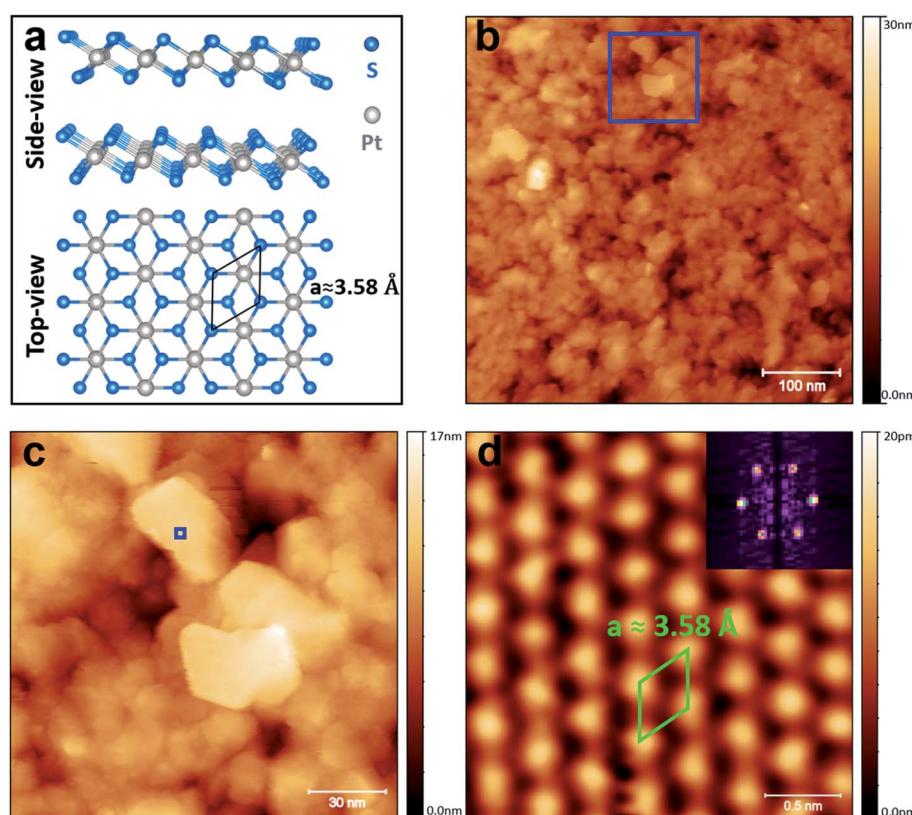


Fig. 2 STM images of PtS<sub>2</sub> film. (a) Structure model of 1T-PtS<sub>2</sub> with side- and top-views. (b) Large-scale STM image of the PtS<sub>2</sub> showing the polycrystalline film ( $500 \times 500 \text{ nm}^2$ ,  $V = 1.5 \text{ V}$  and  $I = 120 \text{ pA}$ ). (c) STM image of the area which is depicted in (b) (blue box), demonstrating flakes of PtS<sub>2</sub> with partially atomically-clean surface ( $150 \times 150 \text{ nm}^2$ ,  $V = 1.5 \text{ V}$  and  $I = 120 \text{ pA}$ ). (d) Atomically-resolved STM image, with FFT inset, of the 1T-PtS<sub>2</sub> surface indicates an interatomic distance of  $3.58 \pm 0.03 \text{ \AA}$  ( $2.5 \times 2.5 \text{ nm}^2$ ,  $V = 1.4 \text{ V}$  and  $I = 400 \text{ pA}$ ).



with high-purity  $\text{PtS}_2$  with relatively small grain sizes yielding broader Raman peaks when compared to mechanically-exfoliated  $\text{PtS}_2$ .<sup>24,31,36,49</sup>

### 3.2. STM and STS characterisation of 1T- $\text{PtS}_2$

STM was used to investigate the atomic structure of the TAC-grown  $\text{PtS}_2$  film as shown in Fig. 2. The 1T atomic-structure model of  $\text{PtS}_2$  is schematically shown in Fig. 2(a) (side- and top-views). The large-area scan in Fig. 2(b) reveals that the films are polycrystalline consisting of crystalline grains (flakes) of  $\text{PtS}_2$  with approximate lateral dimensions up to 40 nm. The root-mean square (RMS) roughness of the image 2(b) is around 2.6 nm. Overall, the surface morphology is similar to that of TAC-grown  $\text{PtSe}_2$ .<sup>50,51</sup> It can be seen from Fig. 2(c) (area depicted in Fig. 2(b) by the blue square) that there are some atomically-clean areas on the surface of the flakes, as can be judged from the sharp straight edges of the flakes. This indicates the stability of the  $\text{PtS}_2$  during transfer, under ambient conditions, from the growth furnace to the STM chamber, which is in contrast with the behaviour of some 2D materials such as black phosphorous.<sup>52</sup> However, some contamination is present on the surface, most likely due to the exposure to ambient conditions. The crystal angles observed advocate hexagonal symmetry, which implies the presence of  $\text{PtS}_2$ , rather than cubic symmetry, which would suggest a significant presence of  $\text{PtS}$ .<sup>35,53</sup> Fig. 2(d)

displays an atomic-resolution image of the defect-free  $\text{PtS}_2$  surface (area indicated by the blue square in Fig. 2(c)), verifying the 1T crystal structure of the  $\text{PtS}_2$  grains.<sup>24,36</sup> The in-plane interatomic distance and a monolayer step height for the  $\text{PtS}_2$  of  $3.58 \pm 0.03 \text{ \AA}$  and  $5.5 \pm 0.06 \text{ \AA}$ , respectively were obtained from this atomically-resolved STM images which is in good agreement with previous theoretical and experimental reports.<sup>24,27,30,32,54</sup> The Fast Fourier Transform (FFT) of the atomically resolved STM image is presented in inset of Fig. 2(d): the surface exhibits clearly ordered structures with the  $(1 \times 1)$  symmetry of  $\text{PtS}_2$ . The spatial periodicity of  $a = 3.58 \text{ \AA}$  in the  $\text{PtS}_2$  crystal was measured from FFT. It should be noted that no surface preparation under UHV (annealing or ion sputtering) was carried out prior to STM/STS investigation.

In order to understand the electronic nature of the TAC-grown  $\text{PtS}_2$ , STS was used to map the local density of states (LDOS) on the pristine surface of  $\text{PtS}_2$ . A large-scale STM image showing grains of  $\text{PtS}_2$  is shown in Fig. 3(a). The inset in Fig. 3(a) shows the line profile of the  $\text{PtS}_2$  flake (around 3.8 nm) where STS measurements were performed. It should be noted that thickness of the flake is an approximation since STM measures the density of states. Fig. 3(b) shows a high-resolution STM image ( $4 \times 4 \text{ nm}^2$ ) of the atomically clean  $\text{PtS}_2$  surface. Fig. 3(c) illustrates grid STS<sup>55</sup> measurements performed on the same area as in Fig. 3(b).  $40 \times 40$  individual  $I(V)$  measurements were

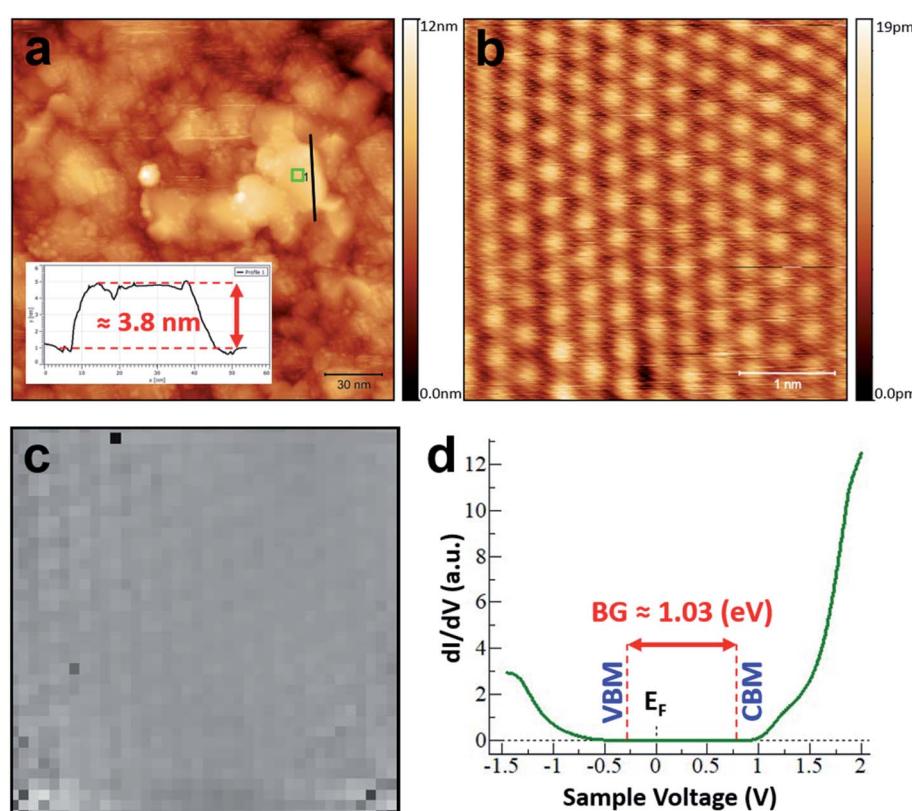


Fig. 3 STS of  $\text{PtS}_2$  film. (a) Large-area STM image with  $\sim 3.8 \text{ nm}$  thick flake of  $\text{PtS}_2$  where STS analysis was performed. (b) Atomically-resolved STM image of the  $\text{PtS}_2$  surface ( $4 \times 4 \text{ nm}^2$ ,  $V = 1.4 \text{ V}$  and  $I = 400 \text{ pA}$ ). A grid spectroscopy measurement was performed on the area depicted in (b). (c)  $40 \times 40$  point grid spectroscopy measurement demonstrates the tunneling current value (taken from the  $I(V)$ ) at each spectroscopy point on the grid at a bias voltage value of  $-1.45 \text{ V}$ . (d) An averaged  $dI/dV$  spectra of the whole grid, which indicates that the bandgap is around  $1.03 \text{ (eV)}$ .



conducted on the grid with dimensions of  $4 \times 4 \text{ nm}^2$ . At each point of the grid, the spectra were averaged over three individual  $I(V)$ -measurements. The tip was moved in constant-current mode with the scanning parameters  $V = 1.4 \text{ (V)}$  and  $I = 400 \text{ (pA)}$  between points. The small variation observed in the tunnelling current at different voltages is due to differences in the electronic structure. The voltage was swept between  $-1.45 \text{ V}$  and  $+2 \text{ V}$ . Fig. 3(c) represents the tunnelling current from the  $I(V)$  at each spectroscopy point at a bias voltage value of  $-1.45 \text{ V}$ . Slightly brighter points correspond to a higher tunnelling current value at  $-1.45 \text{ V}$ . A pixelated outline of the atomically-resolved features can be observed from the grid STS in Fig. 3(c).

It is known that  $dI/dV$  spectra (numerical derivative of  $I(V)$ ) are proportional to the LDOS. The averaged  $dI/dV$  spectrum, which was acquired on the pristine terrace of  $\text{PtS}_2$ , is depicted in Fig. 3(d). The spectrum shown is an average of a few thousand individual spectra on the atomically clean surface. The VBM is positioned at  $0.27 \pm 0.02 \text{ eV}$  below the Fermi level, and the CBM is positioned at  $0.76 \pm 0.02 \text{ eV}$  above the Fermi level. Thus, the obtained bandgap of  $\text{PtS}_2$  here is around  $1.03 \text{ (eV)}$ , which is slightly larger than the bandgap predicted for this thickness in theoretical studies.<sup>25,56</sup> The LDOS of the  $\text{PtS}_2$  flake is sharp in the conduction band, while in the valence band the LDOS increases gradually. The relative location of the Fermi level with respect to the CBM and VBM reveals p-type semiconducting behaviour for our  $\text{PtS}_2$  flake.<sup>23,25</sup> Such a bandgap would be suitable for many applications including thermoelectric devices<sup>31</sup> and infrared photodetectors.<sup>27,28</sup> The beauty of this material is that the bandgap can be tuned by controlling the thickness of the  $\text{PtS}_2$  film. Additionally, low-temperature synthesis of the  $\text{PtS}_2$  allows growth on diverse substrates, such as glass and polyimide,<sup>51,57</sup> and thus broadens the potential application area. Our results establish that STM and STS are powerful techniques that can be applied to the characterisation of TAC-grown  $\text{PtS}_2$ .

## 4. Conclusions

Our results verify that  $\text{PtS}_2$  can be preferentially grown by the TAC process and we have demonstrated that STM-based techniques can be applied to selectively investigate the atomic structure and electronic properties of polycrystalline  $\text{PtS}_2$  films synthesised in this way. XPS, XRD and Raman spectroscopy confirmed the successful growth of  $\text{PtS}_2$  with the correct stoichiometry. STM measurements allowed us to resolve the crystal structure of the TAC-grown  $\text{PtS}_2$  and determine a lattice constant of  $3.58 \pm 0.03 \text{ \AA}$ . The bandgap for a flake with thicknesses  $3.8 \text{ nm}$ , obtained via STS measurements, was found to be  $1.03 \text{ eV}$ . The relative position of the Fermi level indicated p-type semiconducting behaviour. These fundamental results can contribute to future studies focused on  $\text{PtS}_2$  and its application in devices.

## Conflicts of interest

There are no conflicts to declare.

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