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Synthesis and structural studies on (*E*)-3-(2,6-difluorophenyl)-1-(4-fluorophenyl)prop-2-en-1-one: a promising nonlinear optical material†

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A new fluorinated chalcone (*E*)-3-(2,6-difluorophenyl)-1-(4-fluorophenyl)prop-2-en-1-one was synthesized in 90% yield and crystallized by a slow evaporation technique. Its full structural characterization and purity were determined by scanning electron microscopy, infrared spectroscopy, gas chromatography-mass spectrometry, ¹H, ¹³C and ¹⁹F nuclear magnetic resonance, thermal gravimetric analysis (TGA), differential scanning calorimetry (DSC), Raman microspectroscopy, UV-Vis absorption spectroscopy, single crystal X-ray diffraction (XRD) and Hirshfeld surface (HS) analysis. The fluorinated chalcone crystallized in centrosymmetric space group *P*2₁/*c* stabilized by the C–H⋯O and C–H⋯F interactions and the $\pi\cdots\pi$ contact. The crystalline environment was simulated through the supermolecule approach where a bulk with 378 000 atoms was built. The electric parameters were calculated at the DFT/CAM-B3LYP/6-311++G(d,p) level as function of the electric field frequency. The macroscopic parameters such as linear refractive index and third-order nonlinear susceptibility ($\chi^{(3)}$) were calculated, and the results were compared with experimental data obtained from the literature. The $\chi^{(3)}$ -value for the chalcone crystal is $369.294 \times 10^{-22} \text{ m}^2 \text{ V}^{-2}$, higher than those obtained from a few similar types of molecule, showing that the chalcone crystal can be considered as a nonlinear optical material. Also, molecular theoretical calculations such as infrared spectrum assignments, frontier molecular orbital analysis and MEP were implemented, revealing that the most positive region is around the hydrogen atoms of the aromatic rings, and electrophilic attack occurs on the carbonyl group.

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1. Introduction

Nonlinear optical (NLO) materials are useful, playing an essential role in nonlinear optics, and in particular they have a great impact on information technology and industrial applications such as second harmonic generation (SHG) optical modulation, electro-optic switches, optical logic, and others.¹

Their properties and applications depend on the degree of charge separation (polarization) induced by laser light.² In this regard, organic molecules have been constantly reported for their relevant NLO properties^{3–8} due to the presence of a delocalized π system and by different donor or acceptor substituents on the molecular structure.

Among the organic nonlinear materials (NLO) studied so far,^{9–12} chalcones have been widely reported, making use of different donor and acceptor substituents.^{13–19} Chalcones are simple chemical scaffolds of many naturally occurring compounds that are studied due to their structural versatility, consisting of two substituted aromatic rings joined by an α,β -unsaturated ketone.^{20,21} Thus, this family has attracted much interest not only from the nonlinear materials (NLO) perspective but also due to its broad and interesting biological activities.²¹ Chalcones have become even more sought after because they are easily obtained by extraction from natural sources or chemical synthesis through Claisen–Schmidt condensation reaction under basic or acid conditions.^{22,23}

This study focuses on the synthesis, crystal growth and extensive characterization of a fluorinated chalcone (*E*)-3-(2,6-difluorophenyl)-1-(4-fluorophenyl)prop-2-en-1-one (FCH) by

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scanning electron microscopy (SEM), infrared spectroscopy (IR), gas chromatography-mass spectrometry (GC-MS), ^1H -, ^{13}C - and ^{19}F -nuclear magnetic resonance (NMR), thermal gravimetric analysis (TGA), differential scanning calorimetry (DSC), solid state UV-Vis absorption spectroscopy, Raman microspectroscopy, and single crystal X-ray diffraction (XRD). Hirshfeld surface (HS) studies using a model with fluorine atoms as substituents with intense electronegativity characteristics, and structural analysis, were performed aiming to understand structural factors such as torsions, bonding distances, and intermolecular interactions.

Moreover, the linear and nonlinear optical parameters of the FCH crystal were calculated using the supermolecule (SM)^{24–30} approach at DFT/CAM-B3LYP/6-311++G(d,p) level as a function of the electric field frequency. The third-order nonlinear susceptibility ($\chi^{(3)}$) was calculated, and the obtained result for FCH ($369.294 \times 10^{-22} \text{ m}^2 \text{ V}^{-2}$) was compared with one for another organic crystal. The high value for the third-order nonlinear susceptibility qualifies the FCH crystal as a nonlinear optical material. In order to gain insight into the molecular properties, we calculated the molecular orbital frontiers (FMO), molecular electrostatic potential (MEP) and vibrational assignment at the M062X/6-311++G(d,p) level of theory.

2. Experimental and computational procedures

Synthesis and crystallization

Chemicals and solvents required for FCH synthesis were obtained from commercial sources and used without further purification. In a volumetric flask, 4'-fluoroacetophenone (1 mmol; 138 mg) and 2,6-difluorobenzaldehyde (1 mmol; 142 mg) were added in ethanol (1 mL). Then, pulverized KOH (1 mmol; 56.1 mg) was added and the reaction mixture was kept under manual shaking for 3 min, at 25 °C. The reaction's progress was monitored by TLC (Silica gel 60 UV₂₅₄) plate using a mixture of solvent (ethyl acetate/hexane 3 : 7) as an eluent. At the end, a white precipitate was obtained and collected by vacuum filtration followed by a crystallization process using ethanol as solvent. (*E*)-3-(2,6-difluorophenyl)-1-(4-fluorophenyl) prop-2-en-1-one (FCH) was obtained in 90% yield (0.9 mmol; 236 mg). The crystal growth of FCH was carried out in dichloromethane in a conical flask with a known volume of solvent and chalcone. The mixture was kept at 25–30 °C for slow evaporation for 72 h until the crystals formed. Infrared (IR) spectra (Fig. S1†) were recorded using a Perkin Elmer-8400S FT-IR (400–4000 cm^{-1}) with KBr pellet technique. IR (KBr): 3048, 1650, 1580, 1390, 860, 770 cm^{-1} .

Scanning electron microscopy

Microscopic structures of FCH were observed by a scanning electron microscopes (SEM, Hitachi model TM-3030 Plus). The sample was fixed on the carbon adhesive tape in an aluminum stub and, subsequently, observed in the microscope. The SEM image is presented in Fig. 2.

Gas chromatography-mass spectrometry

GC chromatogram (Fig. S2†) and mass spectrum (GC-MS) (Fig. S3†) were recorded in a Shimadzu QP2010 Ultra equipped with capillary column CBP-5 ($30 \times 0.25 \times 0.25$), with injection volume of 1.0 μL , in split mode and helium as drag gas with 1.0 mL min^{-1} flow. For the analysis the injector temperature was 280 °C and the detector 310 °C. The initial oven temperature was 100 °C (held for 2 min), followed by a heating ramp of 30 °C min^{-1} till 300 °C (held for 10 min). MS m/z : 262, 243, 167, 139, 123 and 95.

Nuclear magnetic resonance

A Bruker Avance 500 MHz NMR spectrometer was used to obtain ^1H -, ^{13}C -, and ^{19}F -NMR spectra (Fig. S4–S6†) with TMS as internal standard and CDCl_3 as solvent. The corresponding shifts reveal the purity of the compound. ^1H -, ^{13}C -, and ^{19}F -NMR spectra were used to determine the sample purity and structural conformation of FCH : $\text{C}_{15}\text{H}_9\text{F}_3\text{O}$ (262.2 g mol^{-1}); white solid, 90% yield and mp 108.3 °C; ^1H NMR (500 MHz, CDCl_3) δ 7.86 (d, $J = 16.2$ Hz, 1H), 7.63 (d, $J = 16.2$ Hz, 1H), 7.58 (d, $J = 1.0$ Hz, 1H), 7.28–7.19 (m, 2H), 6.87 (t, $J = 8.6$ Hz, 2H), 6.51 (dd, $J = 3.5$, 1.7 Hz, 1H); ^{13}C NMR (126 MHz, CDCl_3) δ 177.98 (s), 162.06 (dd, $J = 256.0$, 6.8 Hz), 153.59 (s), 146.87 (s), 145.95 (s), 131.30 (t, $J = 11.1$ Hz), 129.79 (t, $J = 2.2$ Hz), 126.82 (t, $J = 8.5$ Hz), 118.03 (s), 112.56 (s), 111.91 (dd, $J = 21.4$, 4.7 Hz); ^{19}F NMR (376 MHz, CDCl_3) δ –105.13, –109.69.

Thermogravimetric analysis

TGA was performed in a Perkin Elmer Pyris equipment, 1 TGA model, installed at Universidade Estadual de Goiás. The evaluation was carried out with 1.506 mg sample mass in an alumina crucible. The program was an initial temperature of 25 °C reaching till 600 °C with heating rate of 10 °C min^{-1} under nitrogen purge gases (20 mL min^{-1} flow).

Differential scanning calorimetry

Differential scanning calorimetry (DSC, 214 Polyma, Netzsch) measurements were carried out in platinum crucibles sealed without drilling the lid with 2.16 mg of sample. The temperature range was 20–200 °C applied at a heating rate of 10 °C min^{-1} , under a dynamic 5.0 N_2 atmosphere of 40 mL min^{-1} . For these analyses, correction was performed to eliminate background. Fig. S7† displays the DSC traces of synthesized chalcone with a sharp exothermic peak at 108.3 °C for heating and 42.4 °C for cooling. The representative DSC thermogram of chalcone (Fig. S7 and S8†) indicates direct melting of crystal phase to isotropic liquid phase since the thermogram showed single thermal cycles. The enthalpy variations (ΔH) during heating and cooling cycles were 28.4 kJ mol^{-1} and –22.04 kJ mol^{-1} , respectively.

UV-Vis and Raman spectroscopy

A UV-Vis spectrum was obtained using a Cary 5000 UV-Vis-NIR spectrometer. Raman measurements were performed on the WITec Alpha 300 RAS microscope (WITec, Ulm, Germany). The excitation wavelength was 785 nm, detection range was 100–

4000 cm⁻¹. The spectra were collected with 20× magnification objective (Zeiss, Jena, Germany). Spectra were recorded with the integration time of 30 s and 2 accumulations. Obtained spectra were processed using WITec Project FOUR and Origin 2016 software.

Crystallographic characterization

The powder diffraction measurements were performed in a Bruker D8 Discover diffractometer. Monochromatic radiation from a copper anode tube coupled to a Johansson monochromator for K α 1 operating at 40 kV and 40 mA was used, Bragg Brentano configuration θ -2 θ , one-dimensional detector Lynxeye®, a range of 2 θ of 5° to 80°, with a pitch of 0.01°. The samples were kept in rotation at 15 rpm during the measurement. Powder diffraction patterns are shown in Fig. S9.† A single chalcone crystal was selected and the single diffraction data were collected on a Bruker APEX-II CCD diffractometer with a monochromated MoK α (λ = 0.71073 Å) at 296.15 K, of the 2490 reflections collected [2087 reflections with $I > 2\sigma(I)$ and 403 independent refined parameters].

The structural solution was carried out with the ShelXS³¹ and ShelXL³² program suite on the Olex² platform.³³ The direct method was applied to solve the structure and least square minimization to refine the molecular model. The ellipsoid diagram was generated from ORTEP³⁴ and the packing diagram and supramolecular representations were generated using Mercury³⁵ and Crystal Explorer 17.5.³⁶ Geometric parameters were obtained using PLATON.³⁷ Non-hydrogen atoms were anisotropically refined and hydrogen atoms were refined using a riding model with a distance of aromatic C–H = 0.93 and $U_{iso}(H) = 1.2U_{eq}$ (hydrogen atoms involved in intermolecular interactions were released). Refinement parameters are shown in Table 1 and crystallographic information files were deposited in the Cambridge Structural Database under code 1997399.

Hirshfeld surface analysis

The Hirshfeld surface can be understood as a spatial map, developed with the intention of defining the space occupied by

a molecule in the crystal, so that this graphical visualization tool makes it possible to describe the intermolecular interactions and contacts through color mapping that identifies regions and describes specific surface properties. In this case, blue is for longer contacts, and red indicates shorter contacts, facilitating visual analysis from molecular interactivity patterns that provide important information about the interactions.^{38,39} Dividing the space in the crystal into specific regions, where the distribution of electrons from one atom to the molecule (ρ_{molecule}) dominates the analogous sum over the crystal (ρ_{crystal}), which is located in one region of the crystal-line space around the molecule that is defined by $W(r) = 0.5$.^{38,39} Since $W(r)$ is a weight function defined for a molecule in a crystal, derived from the atomic electron density $\rho_i(r)$, it is defined in eqn (1):

$$W(r) = \frac{\sum_{i \in \text{molecule}} \rho_i(r)}{\sum_{i \in \text{crystal}} \rho_i(r)} \quad (1)$$

On the Hirshfeld surface a function of d_{norm} distance (normalized contact distance) is defined in eqn (2), which is related in terms of d_i (distance from a point on the surface to the nearest nucleus inside the surface), where the molecule acts as stronger intermolecular donor contacts and d_e (distance from a point on the surface to the nearest nucleus outside the surface), which is where the molecules act as stronger intermolecular contact receptors, as a function of Van der Waals radius (r_i^{vdw} and r_e^{vdw}), was used for intermolecular interactions and to fingerprints plots. It is an informative and quantitative summary of each type of intermolecular contact in the crystal, through the relative surface area corresponding to each type of interaction present in the molecule.^{38–40} Also, the shape index on the Hirshfeld surface auxiliary allowed us to identify characteristic packing modes.^{17,24,38–40}

$$d_{\text{norm}} = \frac{d_i - r_i^{\text{vdw}}}{r_i^{\text{vdw}}} + \frac{d_e - r_e^{\text{vdw}}}{r_e^{\text{vdw}}} \quad (2)$$

Computational procedures

The calculations of the electronic structure were carried out using the Gaussian09 program package⁴¹ for chalcone. The chalcone geometry was optimized (Table S1†) using M062X exchange-correlation functional and basis set 6-311++G(d,p). M062X is a hybrid meta exchange-correlation functional with double the nonlocal exchange (2×). This functional is parametrized for nonmetal and for long noncovalent range interactions.⁴² The absence of imaginary frequencies shows that the optimized geometry is found in a local minimum on the potential energy hypersurface. The highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO), which account for the ability to donate and accept electrons, were carried out by Kohn–Sham⁴³ analysis. MEP is used to assign the nucleophilic and electrophilic site within the molecule. The main assignments of infrared frequencies and

Table 1 Crystallographic data and structure refinement for FCH

Crystal data	Chalcone
Chemical formula	C ₁₅ H ₉ F ₃ O
Molecular weight	262.22
Space group	<i>P</i> 2 ₁ / <i>c</i>
<i>a</i> , <i>b</i> , <i>c</i> (Å)	11.2133(5), 9.1206(4), 12.1021(5)
α , β , γ (°)	90, 100.356(2), 90
<i>V</i> (Å ³)	1217.54(9)
<i>Z</i>	4
Radiation type	Mo K α
μ (mm ⁻¹)	0.120
$R[F^2 > 2\sigma(F^2)]$	0.0460
$wR(F^2)$	0.1356
<i>S</i>	1.122
No. of reflections	2490
No. of parameters	187



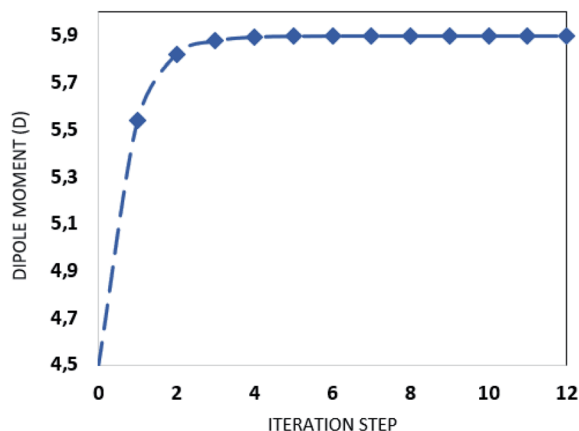


Fig. 1 Dipole moment as function of the iteration steps.

Raman spectra were performed in conjunction with quantum chemical calculations.

Nonlinear optical properties

In order to calculate the nonlinear optical (NLO) properties of the FCH crystal, the supermolecule (SM) approach was employed at the DFT/CAM-B3LYP/6-311++G(d,p) level to build a bulk with $15 \times 15 \times 15$ unit cell, with 4 asymmetric units in each unit cell totalizing 13 500 molecules, each with 28 atoms; thus the bulk was composed of 378 000 atoms. The SM is an iterative process that considers the atoms of the molecules that surround the isolated molecule as specific charges. First, in the SM approach using ChelpG software, the electrostatic potential of the isolated molecule in vacuum was fitted, and the partial atomic charges were calculated. Then an iterative process began by the substitution of the each atom in the unit cells generated by the atom with the partial atomic charge calculated in the previous step. In each step of the approach the electrical properties of the embedded molecules are calculated, and the process finishes with the dipole moment ($\mu = \sqrt{\mu_x^2 + \mu_y^2 + \mu_z^2}$) convergence. The SM method has been successfully employed in several recent works to simulate the crystalline environment polarization effects on the electrical properties of organic molecules.^{44–48} Fig. 1 shows the convergence of the dipole moment as a function of the iterative steps, and it can be seen that the crystalline environment polarization effect is significant, increasing the μ -value from 4.5 D for isolated molecules (step 0) for 5.9 D for embedded molecules, a perceptual increase of $\sim 31\%$.

In our calculations the average linear polarizability ($\langle\alpha\rangle$) and the static average second hyperpolarizability were calculated using the following expression in eqn (3) and (4).

$$\langle\alpha\rangle = \frac{\alpha_{xx} + \alpha_{yy} + \alpha_{zz}}{3} \quad (3)$$

$$\langle\gamma(0;0,0,0)\rangle = \frac{1}{5} [\gamma_{xxxx} + \gamma_{yyyy} + \gamma_{zzzz} + 2(\gamma_{xxyy} + \gamma_{xxzz} + \gamma_{yyzz})] \quad (4)$$

The Kleimann symmetry was employed since the medium optical dispersion was not taken into account. The third-order nonlinear susceptibility can be written in the following form (eqn (5)).

$$\chi^{(3)}(-\omega; \omega, \omega, -\omega) = \frac{(n(\omega)^2 + 2)^4}{3} \frac{N \langle\gamma(-\omega; \omega, \omega, -\omega)\rangle}{\epsilon_0 V} \quad (5)$$

The intensity-dependent refractive index (IDRI) second hyperpolarizability is given by eqn (6).

$$\langle\gamma(-\omega; \omega, \omega, -\omega)\rangle \cong 2\langle\gamma(-\omega; \omega, 0, 0)\rangle - \langle\gamma(0; 0, 0, 0)\rangle \quad (6)$$

Also using the Clausius–Mossotti relationship the linear refractive index (n) can be calculated through the eqn (7).

$$\frac{n^2 - 1}{n^2 + 2} = \frac{4\pi N}{3} \langle\alpha\rangle, \quad (7)$$

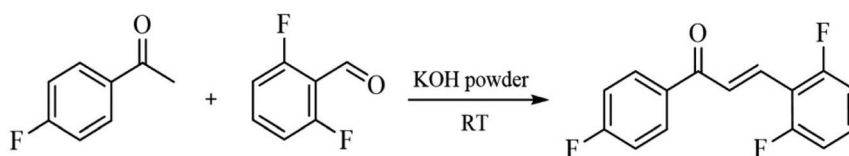
where N is the volumetric density of a molecule in the unit cell. All calculations were performed using Gaussian-09.⁴⁹

3. Results and discussions

Synthesis and crystal growth

FCH molecule was synthesized by Claisen–Schmidt condensation procedure. For this reaction, 4'-fluoroacetophenone, 2,6-difluorobenzaldehyde, and KOH were reacted by manual shaking using ethanol as a green solvent (Scheme 1). FCH was obtained as white solid (90% yield and mp 108.3 °C).

A solubility study shows that FCH is not soluble in hexane, moderately soluble in ethanol, and highly soluble in dichloromethane, acetone and chloroform. High quality crystals of FCH (Fig. 2a) were obtained by solvent evaporation method using dichloromethane (CH_2Cl_2), at 25–30 °C. Surface morphology of FCH was investigated by scanning electron microscope (Fig. 2b). From the analysis of the SEM micrograph (30 \times), it can be seen that the crystal has non-uniform distribution.



Scheme 1 Synthesis of FCH.



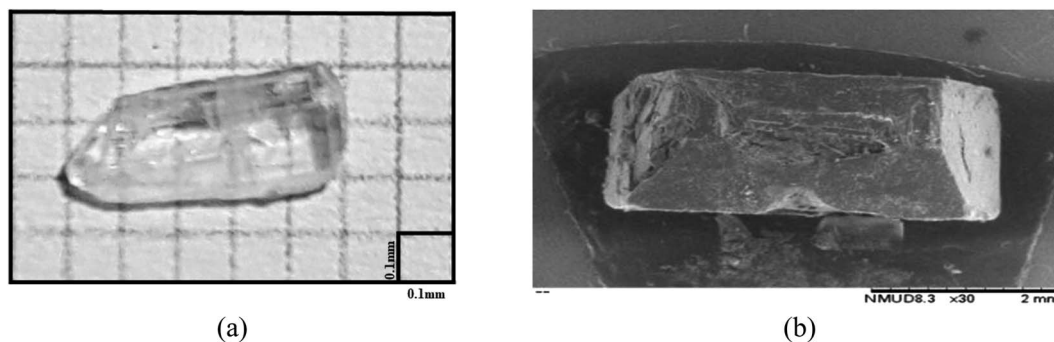


Fig. 2 (a) Single crystal photograph of FCH; (b) SEM micrograph of FCH.

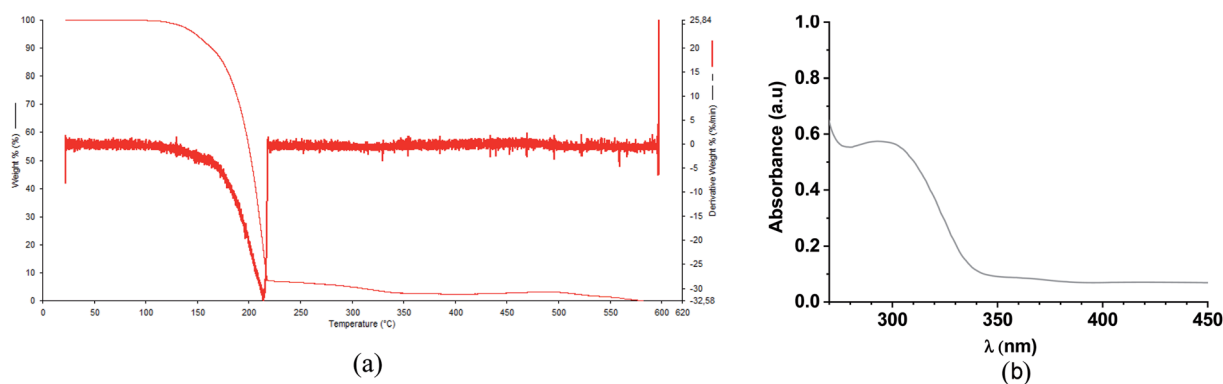


Fig. 3 (a) TG/TGA and (b) UV-Vis absorption spectrum of FCH.

Structural characterization

In order to confirm the molecular structure and to determine the purity of the fluorinated chalcone (FCH), GC-MS analysis (Fig. S2 and S3†) and ^1H -, ^{13}C -, and ^{19}F -NMR (Fig. S4–S6†) were performed. From the analysis of the GC chromatogram (Fig. S2†), a single peak is observed at 6.69 min, whose mass spectrum (Fig. S3†) is characteristic for *E*-(2,6-difluorophenyl)-1-(4-fluorophenyl)prop-2-en-1-one (FCH) fragments (m/z 262, 243, 167, 139, 123 and 95). Moreover, the high purity and molecular structure of FCH crystals were also determined by ^1H -

^{13}C -, and ^{19}F -NMR and described below: ^1H -NMR (500 MHz, CDCl_3) δ 7.86 (d, $J = 16.2$ Hz, 1H), 7.63 (d, $J = 16.2$ Hz, 1H), 7.58 (d, $J = 1.0$ Hz, 1H), 7.28–7.19 (m, 2H), 6.87 (t, $J = 8.6$ Hz, 2H), 6.51 (dd, $J = 3.5, 1.7$ Hz, 1H); ^{13}C -NMR (126 MHz, CDCl_3) δ 177.98 (s), 162.06 (dd, $J = 256.0, 6.8$ Hz), 153.59 (s), 146.87 (s), 145.95 (s), 131.30 (t, $J = 11.1$ Hz), 129.79 (t, $J = 2.2$ Hz), 126.82 (t, $J = 8.5$ Hz), 118.03 (s), 112.56 (s), 111.91 (dd, $J = 21.4, 4.7$ Hz); ^{19}F -NMR (376 MHz, CDCl_3) δ -105.13, -109.69.

The TG/TGA evaluation of FCH under nitrogen purge gases is shown in ESI†. The thermogram of chalcone shows that the thermal decomposition occurs between 100 °C and 250 °C

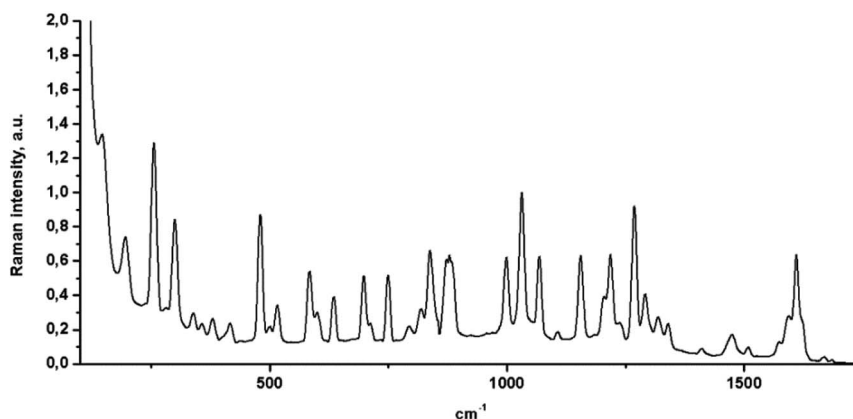


Fig. 4 Raman spectrum of FCH. Raman peak at 1032 cm^{-1} was used to normalize the spectrum.



(Fig. 3a). The decomposition occurs in a well-defined event represented by a single DTG peak reaching 0% of mass loss. Fig. 3b shows the UV-Vis absorption of FCH. Due to the presence of $n-\pi^*$ transition on the structure of the fluorinated chalcone FCH, an absorption band in the UV-region (280–340 nm) was observed.⁵⁰

Raman microspectroscopy was also used to characterize the FCH, and its spectrum was normalized at 1032 cm^{-1} (Fig. 4). From the analysis of Fig. 4, C–H stretching vibrations were observed at $1029, 1070, 1157, 1217$ and 1269 cm^{-1} with a strong Raman intensity. The C–H out-of-plane deformations were observed at $747, 838, 876$, and 996 cm^{-1} for FCH crystal. The C–F Raman band can be observed between 560 and 610 cm^{-1} . Additionally, C–C stretching vibrations also were observed at $1293, 1320, 1337, 1411, 1477$, and 1509 cm^{-1} . A sharp band at 1608 cm^{-1} was also observed and assigned to ethylenic bridge vibrations. The C=O vibrations were also observed as a weak Raman band at 1670 cm^{-1} . This weak Raman band may be

explained by intermolecular hydrogen bonding between the C=O group and a phenyl ring from the FCH.

The ORTEP diagram with the asymmetric unit is presented in Fig. 5. The molecule has in its structure three fluorine atoms, in the *ortho* and *para* positions. The carbonyl group C7=O1 appears in an (*E*)-configuration with respect to the stereochemistry C8=C9. The molecular coplanarity is indicated by C5–C6–C7–O1, C6–C7–C8–C9 and C8–C9–C10–C15 dihedral angles respectively are 178.5° , -178.89° and -178.0° . Planarity of chalcones is correlated with π conjugation in molecule, being evidenced in the literature, which contributes to their use as potential optical devices.⁵¹

The chalcone compound is crystallized from the $P2_1/c$ centrosymmetric monoclinic space group with $Z = 4$, which represents 4 molecules in the unit cell. The final structural refinement data and the main crystallographic parameters are presented in Table 1. In addition, other details such the principal bond distance, angles and dihedral angles are shown in Table 2. Table 2 shows that the experimental and theoretical

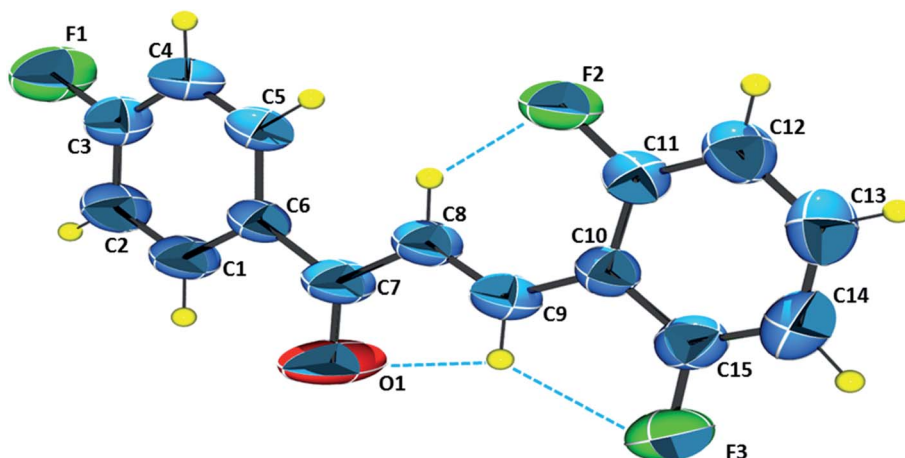


Fig. 5 ORTEP diagram for FCH. The ellipsoids are represented at 50% of probability level with the atomic numbering scheme. Hydrogen atoms are in arbitrary radii. Intramolecular interactions C8–H8...F2, C9–H9...F3 and C9–H9...O1 are represented.

Table 2 Relevant experimental and theoretical bond length (Å), bond angles ($^\circ$) and dihedral angles ($^\circ$) for FCH

	Experimental	Theoretical		Experimental	Theoretical
O1–C7	1.213(2)	1.212	F2–C11–C12	117.13(16)	117.5
F1–C3	1.360(3)	1.340	F3–C15–C10	117.4(17)	117.9
F2–C11	1.354(2)	1.341	O1–C7–C8	120.4(2)	121.5
F3–C15	1.349(3)	1.339	C8–C9–C10–C11	2.3(3)	0.7
C5–C6	1.392(2)	1.396	C8–C9–C10–C15	$-178.0(2)$	-179.3
C6–C7	1.488(3)	1.498	C5–C6–C7–O1	178.5(2)	162.7
C8–C9	1.313(3)	1.339	C6–C7–C8–C9	$-178.8(18)$	171.1
C10–C11	1.392(2)	1.398	C1–C6–C7–C8	178.7(17)	163.3
C8–C9–C10	122.0(13)	128.0	C5–C6–C7–C8	0.8(3)	-18.0
C9–C10–C15	120.6(16)	119.6	F1–C3–C4–C5	$-178.8(19)$	-179.3
O1–C7–C6	120.1(19)	117.8	F2–C11–C12–C13	$-179.2(17)$	179.8
C6–C7–C8	119.4(15)	118.1	C9–C10–C15–F3	0.3(3)	-0.05
C1–C6–C7	118.4(15)	117.8	O1–C7–C8–C9	1.8(3)	9.7
C5–C6–C7	123.0(16)	122.8	O1–C7–C8–H8	178.0(1)	168.8
F1–C3–C4	118.7(19)	118.6	C6–C7–C8–H8	$-3.0(1)$	-10.3



Table 3 Hydrogen-bond geometry (Å, °) obtained from structural analysis for FCH

D–H...A	D–H	H...A	D...A	D–H...A	Symmetry code
C5–H5...O1	0.97	2.50	3.2974(1)	139.0(16)	$x, 1/2 - y, 1/2 + z$
C8–H8...F2	0.94	2.26	2.8457(1)	119.6(16)	Intramolecular
C9–H9...F3	0.91	2.37	2.7590(1)	105.7(15)	Intramolecular
C9–H9...O1	0.91	2.44	2.7753(1)	101.6(14)	Intramolecular
C14–H14...F1	0.95	2.54	3.4769(2)	173.0(2)	$1 + x, 1 + y, z$
C4–H4...F2	0.89	2.616	3.464	158.56(2)	$1 - x, -1/2 + y, 1.5 - z$
C1–H1...O1	0.96	2.42	2.767	101.3	Intramolecular

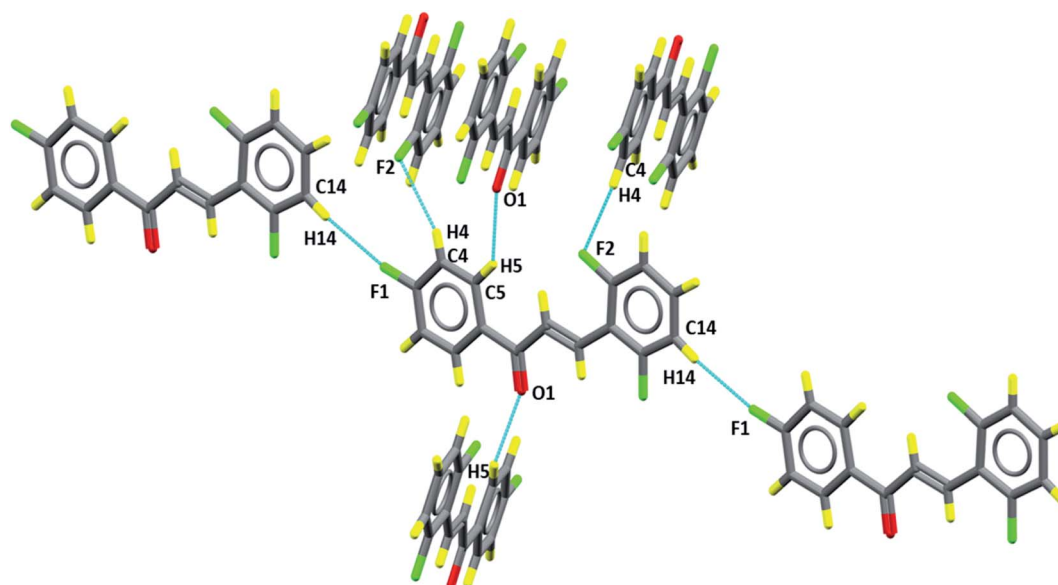


Fig. 6 Representation of intermolecular interactions C5–H5...O1, C14–H14...F1 and C4–H4...F2.

selected geometric parameters are in good agreement. The C5–C6–C7–O1, O1–C7–C8–C9 and O1–C7–C8–H8 theoretical dihedral angles are 162.7°, 9.7° and 168.8°, respectively, while experimental dihedral angles are 178.5°, 1.8° and 178.0°, respectively, as shown in Table 2. These differences can be explained by taking into account that the theoretical results were obtained in the gas phase, not taking into account the interactions that occur in the crystalline environment. The intermolecular interactions presented in all crystal structures are listed in Table 3.

The packing diagram, in Fig. 6, shows the motifs in the crystalline lattice. It was observed that crystal structure is stabilized by intermolecular interactions (Table 3), since C–H groups are hydrogen bonding donors, C5–H5...O1 which can be described as $C_1^1(5)$, involving a carbonyl group (Fig. 6) which results in the formation of linear chains running along the c axis, while for C4–H4...F2 and C14–H14...F1 the graph set is given as $C_1^1(10)$ and $C_1^1(12)$, involving, respectively, fluorine F2 and F1, which are located in Cg2 (C10–C15) and Cg1 (C1–C6) aromatic rings, respectively. These intermolecular interactions describe supramolecular architectures.

The Hirshfeld surface contributes to the understanding of packaging, promotes the analysis of intermolecular interactions and enables the recognition of predominant bonds. The HS mapped over d_{norm} (ranging from -0.0635 to 3392 Å) is shown in Fig. 7. In it, the compound is stabilized by hydrogen bonds, on which the red spots display the d_{norm} identifying the region in which interactions C5–H5...O1/O1...H5–C5 (Fig. 7a) of the FCH are depicted. In addition, interactions C14–H14...F1/F1...H14–C14 (Fig. 7b) and C4–H4...F2/F2...H4–C4 (Fig. 7c), despite being flagged as weak, are relevant for packaging stabilization. Thus, planarity is correlated with conjugation throughout the molecule.

Shape index is a graphical tool that reveals hydrophobic interactions and aids to interpret the interactions that occur in system π . The interaction $\pi \cdots \pi$ (Fig. 8a) is displayed by red and blue triangles, where occur the overlapping of two aromatic rings occurs; for this interaction, the aromatic ring centroids were calculated (Fig. 8b), with the distance between two rings of 3.856 Å, providing greater stability to the structure.

The quantitative summary is generated from the graph of (d_e vs. d_i) resulting from the contributions of each interaction that is exposed in the supramolecular arrangement. The interactions



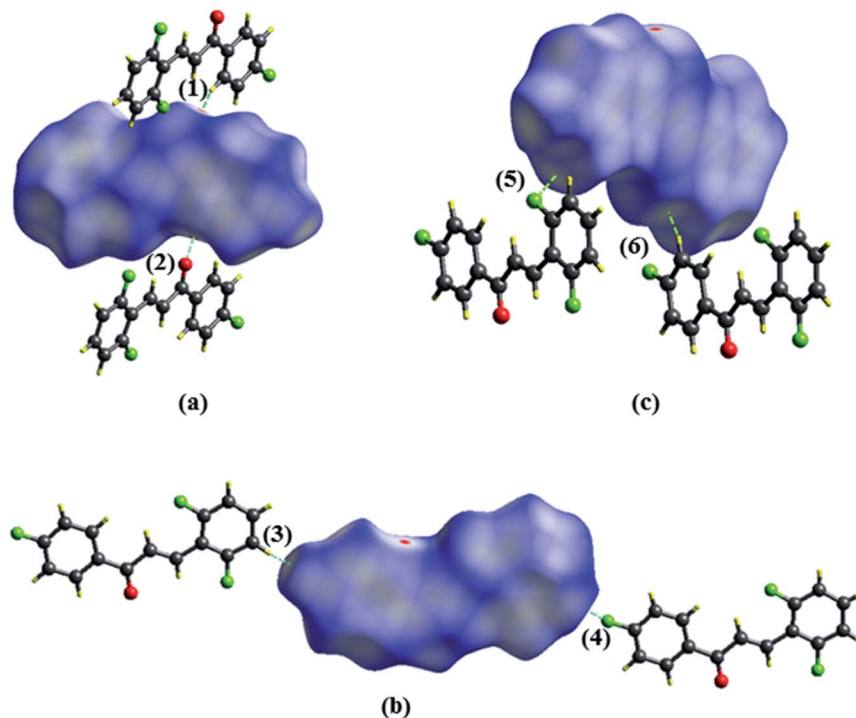


Fig. 7 Hirshfeld surfaces plotted for FCH (a) C5–H5...O1 (1)/O1...H5–C5–F1 (2), (b) F1...H14–C14 (3)/C14–H14...F (4) and (c) C4–H4...F2 (5)/F2...H4–C4 (6). Dotted lines were used to represent hydrogen bonds.

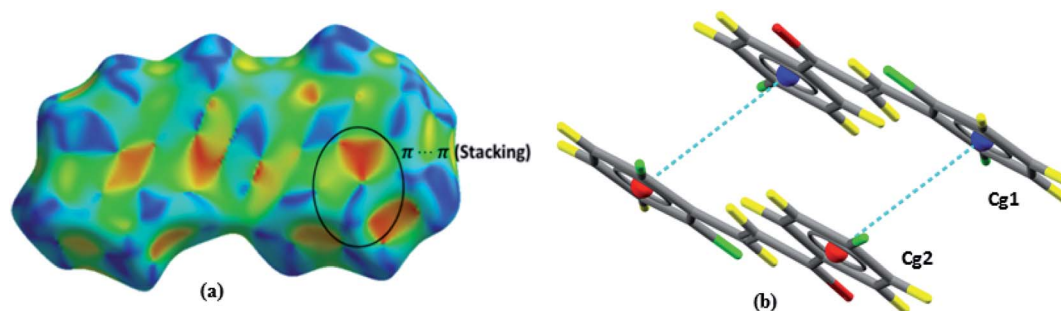


Fig. 8 Shape index surfaces of (a) evidencing $\pi \cdots \pi$ interactions. Representation of $\pi \cdots \pi$ interactions (b).

are highlighted in fingerprint plots (Fig. 9), in which most contacts are related to the interaction $\text{H} \cdots \text{F}$ (31.3%), which indicates the $\text{C}-\text{H} \cdots \text{F}$ contacts. As this is an organic compound, the $\text{H} \cdots \text{H}$ interactions had a high value (21.9%), followed by the $\text{C} \cdots \text{H}$ connection (16.4%). Another characteristic of chalcone is the presence of $\text{C} \cdots \text{C}$ interactions (10.13%) present in the center, which helps in understanding $\pi \cdots \pi$ bonding. $\text{O} \cdots \text{H}$ interactions (9.9%) are characterized by peaks at the bottom of the fingerprint plots and the contacts $\text{C} \cdots \text{F}$ (6.6%). It is observed that these interactions become important for the stabilization of the supramolecular arrangement. Additionally, solid-state UV-Vis absorption spectrum of FCH clearly shows a maximum absorption peak at ~ 409 nm, similar to the chalcone-type molecules (Fig. 3).

Molecular modeling description

The root of the mean squared (RMS) value, predicted by Mercury software, which accounts for the difference between theoretical and experimental geometries, was 0.00934. The overlap geometries show a good agreement of the geometric parameters (Fig. 10). The aromatic ring (Cg1) displays a theoretical dihedral twist C1–C6–C7–C8 and C5–C6–C7–C8 of 163.3° and -18.0° that differ by experimental dihedral angles, whose value is 178.7° and 0.8° , respectively (Table 2). This contribution to solid state planarity is indicated by $\pi \cdots \pi$ (Cg1...Cg2) inter-molecular interaction.

Fig. 11 shows the frontier molecular orbital (FMO) obtained at the M062X/6-311++G(d,p) level of theory. The FMO analysis of bond-antibonding interactions is taken into consideration by filled (donor) Lewis base and empty (acceptor) Lewis acid. The



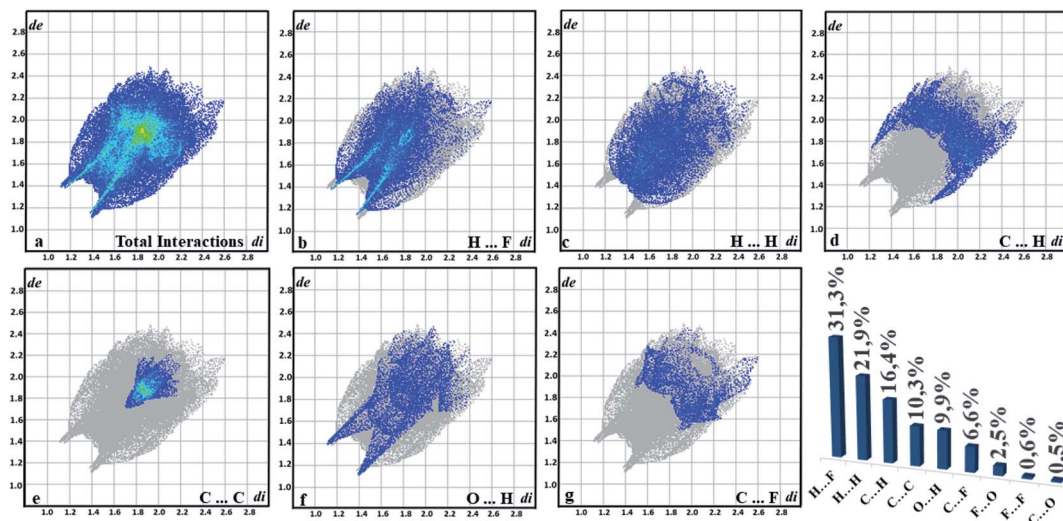


Fig. 9 Fingerprint plots for FCH.

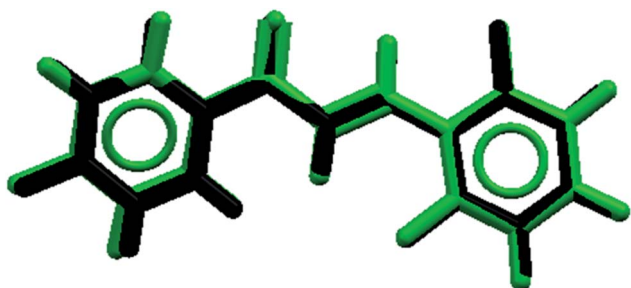


Fig. 10 Overlapping of X-ray geometric parameters (black) and theoretical calculation (green) structures for FCH.

HOMO as the electron donor is localized on the π bond, and the olefin region appears as a π bonding orbital, which is characteristic of the nucleophilic region. The HOMO energy is $-192.06 \text{ kcal mol}^{-1}$. The LUMO orbital is a π antibonding orbital localized on the $\text{C7}=\text{O1}$ carbonyl group. The LUMO energy is $-41.39 \text{ kcal mol}^{-1}$, indicating that the FCH is an electrophilic species.

The MEP topology is a physicochemical tool that identifies the chemical reactivity of the molecular systems.⁵² The electrostatic potential can be determined using eqn (8):

$$V(r) = \sum_{\alpha} \frac{Z_{\alpha}}{|r - R_{\alpha}|} - \int \frac{\rho(r')}{|r - r'|} dr' \quad (8)$$

Where $V(r)$ is the potential energy by a positive unit charge at r ; Z_{α} is the nuclear charge α located at R_{α} , and $\rho(r')$ is the electron density at point r' . To comprehend the MEP done for FCH we adopted the tridimensional MEP representation (Fig. 12), where red specifies the most negative region, which is concentrated on the oxygen atom (O1) of the carbonyl group ($\text{C7}=\text{O1}$), with value about $-31.06 \text{ kcal mol}^{-1}$. On the other hand, the blue region specified the most positive potential (strongest repulsion), and the isovalue potential energy of about $21.84 \text{ kcal mol}^{-1}$ is around the hydrogen atoms in Cg2 ($\text{C10}-\text{C15}$). These results identify the electrophilic and nucleophilic site attacks and provide an understanding of where FCH interacts intermolecularly.

Table 4 displays some selected theoretical and experimental IR absorption bands. Fig. 13 compares the experimental and theoretical IR spectra graphically. The theoretical vibrational

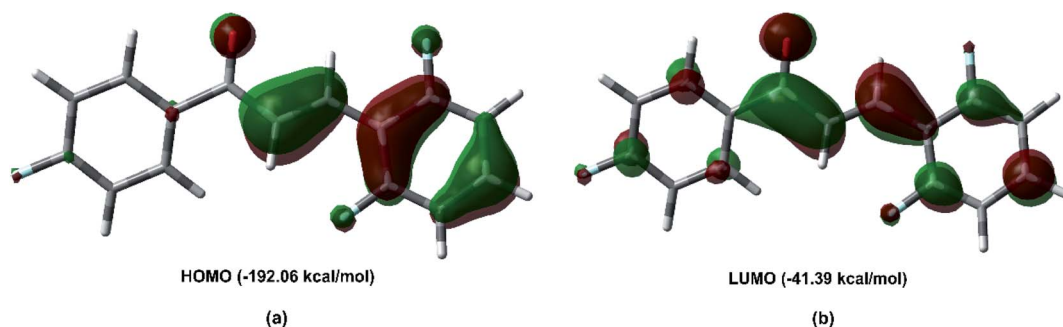


Fig. 11 Frontier molecular orbitals for FCH (a) the HOMO " π " bonding orbital and (b) the LUMO " π " antibonding orbital derived from Kohn–Sham analysis at M062X/6-311++G(d,p) level of theory with the isovalue of 0.04 atomic units.



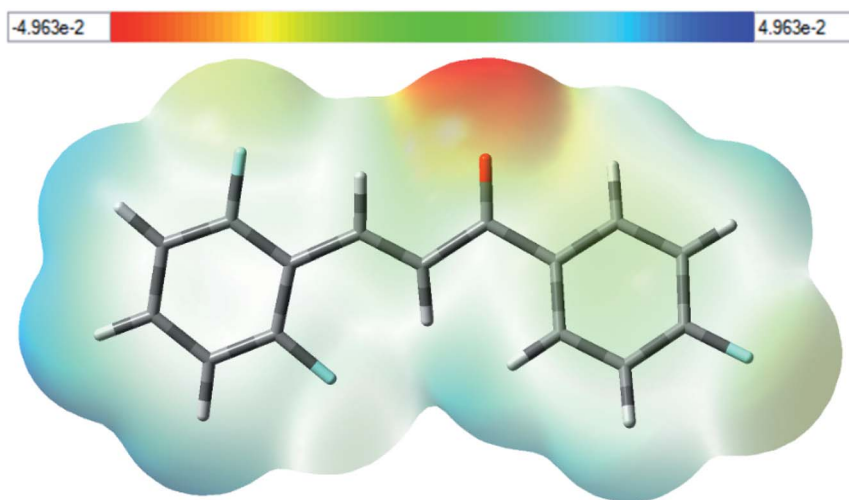


Fig. 12 The molecular electrostatic potential surface mapped for FCH. The red-colored region of the carbonyl group, and the blue-colored region around the hydrogen atoms. The density isovalue of 4.0×10^{-4} electrons per bohr³ was used to make the molecular electrostatic potential surfaces.

frequency values were scaled by 0.943 (ref. 53) for a better fit with the experimental results. In chalcones, the resonance effects of α,β -unsaturated ketone must be recognized; the stretching of $C_{sp^2}-H$ absorbs in the $3106-3077\text{ cm}^{-1}$, while the computed wavenumbers are in the range of $3078-3006\text{ cm}^{-1}$. Experimentally, the theoretical measurements of the carbonyl group, made supposing the molecule was in gas phase, absorb at 1695 cm^{-1} , while the experimental measurements, in which the decrease of the vibration frequency value occurs due to the molecular hydrogen interaction, absorb at 1666 cm^{-1} .⁵⁴ The FT-IR spectrum of FCH indicates a good correlation to the theoretical data obtained in the gas phase

NLO properties

The DFT/CAM-B3LYP/6-311++G(d,p) results for the average linear polarizability ($\langle\alpha(-\omega;\omega)\rangle$), the IDRI second hyperpolarizability ($\langle\gamma(-\omega;\omega,\omega,-\omega)\rangle$), the linear refractive index ($n(\omega)$) and the third-order nonlinear susceptibility

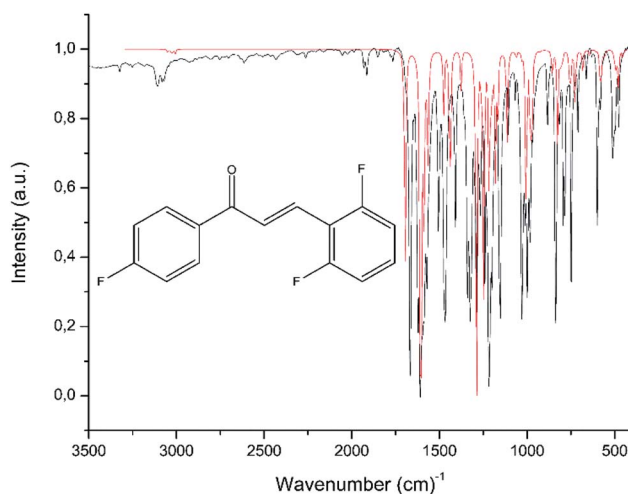


Fig. 13 The theoretical (red) and experimental (black) overlapped the FT-IR spectrum of FCH.

Table 4 Experimental and theoretical assignments for FCH. The calculation was carried out at the M062X/6-311++G(d,p) level of theory, in the gas phase^a

Vibrational mode	Exp. freq. ^b	Scaled freq. ^{b,c}
$\nu(C-H)_{Ar, Alk}^d$	3106-3077	3078-3006
$\nu(C=O)$ and $\nu(C=C)_{Alk}^d$	1666	1695
$\nu(C=C)_{Ar, Alk}^d$	1621-1572	1608-1563
$\nu(C=C)_{Ar}^d$; $\delta(C=C)_{Ar}^d$ and $\delta(C-H)_{Ar}^d$ in plane	1507-1410	1475-1436
$\nu(C=C)_{Ar}^d$; $\delta(C-H)_{Ar}^d$ in plane	1341	1376
$\nu(C=C)_{Ar}^d$ and $\delta(C-H)_{Alk}^d$ in plane	1290-1241	1267-1243
$\delta(C-H)_{Ar, Alk}^d$ in plane	1153-1000	1200-1004
$\delta(C=C)_{Ar}^d$ and $\delta(C-H)_{Ar, Alk}^d$ out plane	990-974	978-926
$\delta(C=C)_{Ar}^d$ and $\delta(C-H)_{Ar, Alk}^d$ out plane	893-765	864-762
$\delta(C-F)$	745-730	674
$\delta(C(O)-C=C)$	530-500	580

^a ν = stretching; δ = bending. ^b cm^{-1} . ^c Scale factor 0.943. ^d Ar = aromatic ring. ^e Alk = alkene.



Table 5 Numerical results for $\langle\alpha(-\omega;\omega)\rangle$ (10^{-24} esu), $\langle\gamma(-\omega;\omega,\omega,\omega)\rangle$ (10^{-36} esu), $n(\omega)$ and $\chi^{(3)}(-\omega;\omega,\omega,\omega)$ (10^{-22} m² V⁻²) for the FCH crystal

ω	$\langle\alpha(-\omega;\omega)\rangle$	$\langle\gamma(-\omega;\omega,\omega,\omega)\rangle$	$n(\omega)$	$\chi^{(3)}(-\omega;\omega,\omega,\omega)$
0.000	26.211	45.6651	1.6409	125.406
0.003	26.214	45.9547	1.6410	126.202
0.004	26.216	46.0031	1.6411	126.335
0.006	26.223	46.1417	1.6413	126.715
0.007	26.227	46.2321	1.6414	126.964
0.008	26.232	46.3365	1.6416	127.250
0.009	26.237	46.4551	1.6417	127.576
0.010	26.243	46.5879	1.6419	127.941
0.024	26.397	49.9849	1.6467	137.270
0.043	26.831	60.1717	1.6601	165.245
0.050	27.073	66.2793	1.6677	182.018
0.060	27.495	77.6933	1.6811	213.363
0.070	28.035	93.7787	1.6984	257.537
0.072	28.159	97.7381	1.7024	268.410
0.080	28.722	116.878	1.7207	320.972
0.086	29.191	134.474	1.7362	369.294
0.090	29.600	151.173	1.7499	415.154
0.095	30.140	174.773	1.7680	479.965
0.100	30.750	204.679	1.7890	562.093

$\langle\chi^{(3)}(-\omega;\omega,\omega,-\omega)\rangle$ as a function of the electric frequency are shown in Table 5. In the frequency range of 0.0 a.u. $\leq \omega \leq$ 0.10 a.u., the $\langle\alpha(-\omega;\omega)\rangle$ -values and $\langle\gamma(-\omega;\omega,\omega,-\omega)\rangle$ -values increased 17.31% and 348.22%, respectively, as compared with its static values. The linear refractive index values increased smoothly with the increasing frequency, from the static value (1.6409) up $\omega = 0.05$ a.u. (1.6677) the enhancing at the $n(\omega)$ -value was only of 1.63%, and from this point up to $\omega = 0.10$ a.u. the enhancement was of 7.3%. The third-order nonlinear susceptibility presents a significant increase due to the electric field frequency of 348.22%, going from the static value of 125.046×10^{-22} m² V⁻² to 562.093×10^{-22} m² V⁻².

Table 5 shows the results of the linear refractive index and the third-order nonlinear susceptibility for the crystal of FCH at electric field frequency of $\omega = 0.086$ a.u. ($\lambda = 532$ nm). In addition, several experimental results obtained by Z-scan are presented by comparison. As can be seen, the

Table 7 Partial charges in isolated and embedded molecules of the FCH

A1				
	Atom	Isolated	Embedded	$\Delta\%$
1	F	-0.2449	-0.2430	0.75
2	F	-0.2339	-0.2550	9.05
3	F	-0.2806	-0.3058	8.99
4	C	-0.1937	-0.1963	1.35
5	C	-0.1373	-0.1692	23.21
6	C	0.3701	0.3554	3.97
7	O	-0.5520	-0.5957	7.92
8	C	0.1584	0.1706	7.67
9	H	0.0974	0.0812	16.64
10	C	-0.0825	-0.0806	2.20
11	H	0.1242	0.1603	29.09
12	C	-0.4356	-0.4508	3.48
13	H	0.1894	0.2003	5.76
14	C	-0.0034	0.0232	777.33
15	H	0.1102	0.0965	12.47
16	C	0.3483	0.3485	0.05
17	C	-0.2560	-0.2532	1.10
18	H	0.1678	0.1868	11.33
19	C	0.4328	0.4231	2.24
20	C	0.6396	0.6760	5.68
21	C	-0.2877	-0.2808	2.39
22	H	0.1759	0.1855	5.43
23	C	-0.2702	-0.2717	0.56
24	H	0.1718	0.1770	3.05
25	C	-0.0235	-0.0078	66.96
26	H	0.1214	0.1304	7.38
27	C	-0.2861	-0.3029	5.86
28	H	0.1799	0.1981	10.13
	Ring A			
	Ring B			
	Ring C			

$\langle\chi^{(3)}(-\omega;\omega,\omega,-\omega)\rangle$ -value for FCH (369.294×10^{-22} m² V⁻²) is greater than the other experimental values presented in Table 6. This value is 1.33 times the $\chi^{(3)}$ -value measured for the crystal (2E)-3-(3-methylphenyl)-1-(4-nitrophenyl)prop-2-en-1-one (3MPNP)⁵⁵ and 185.58 times the value for (2E)-1-(3-bromophenyl)-3-[4-(methylsulfanyl) phenyl]prop-2-en-1-one.⁵⁶

Table 6 The linear refractive index and third-order susceptibility $\times 10^{-22}$ (m² V⁻²) for several organic crystals

	λ (nm)	$n(\omega)$	$\chi^{(3)}(-\omega;\omega,\omega,-\omega)$
FCH molecule (present work)	532	1.736	369.294
(2E)-3-(3-Methylphenyl)-1-(4-nitrophenyl)prop-2-en-1-one ⁵⁵	532	1.418	277.100
(2E)-3-(3-Methylphenyl)-1-(4-nitrophenyl)prop-2-en-1-one ⁴⁷	532	2.002	176.400
4,6-Dichloro-2-(methylsulfonyl)pyrimidine ⁴⁵	532	1.613	56.740
(E)-3-(2-Bromophenyl)-1-(2-((phenylsulfonyl)amine)-phenyl)prop-2-en-1-one ⁴⁴	1064	1.680	25.700
1-(5-Chlorothiophen-2-yl)-3-(2,3-dimethoxyphenyl)prop-2-en-1-one ^{55,57}	532	1.594	23.830
1-(5-Chlorothiophen-2-yl)-3-(2,3-dichlorophenyl)prop-2-en-1-one ⁵⁸	532	—	16.210
2-(4-Methylphenoxy)-N-[(1E)-(4-nitrophenyl)methylene]acetohydrazide ⁵⁹	532	—	10.240
1-(4-Aminophenyl)-3-(3,4,5-trimethoxyphenyl)prop-2-en-1-one ⁶⁰	532	—	8.700
(2E)-3-[4-(Methylsulfanyl)phenyl]-1-(4-nitrophenyl)prop-2-en-1-one ⁵⁶	800	1.363	2.370
(2E)-1-(4-Bromophenyl)-3-[4-(methylsulfanyl)phenyl]prop-2-en-1-one ⁵⁶	800	1.365	2.300
(2E)-1-(3-Bromophenyl)-3-[4-(methylsulfanyl)phenyl]prop-2-en-1-one ⁵⁶	800	1.360	1.990



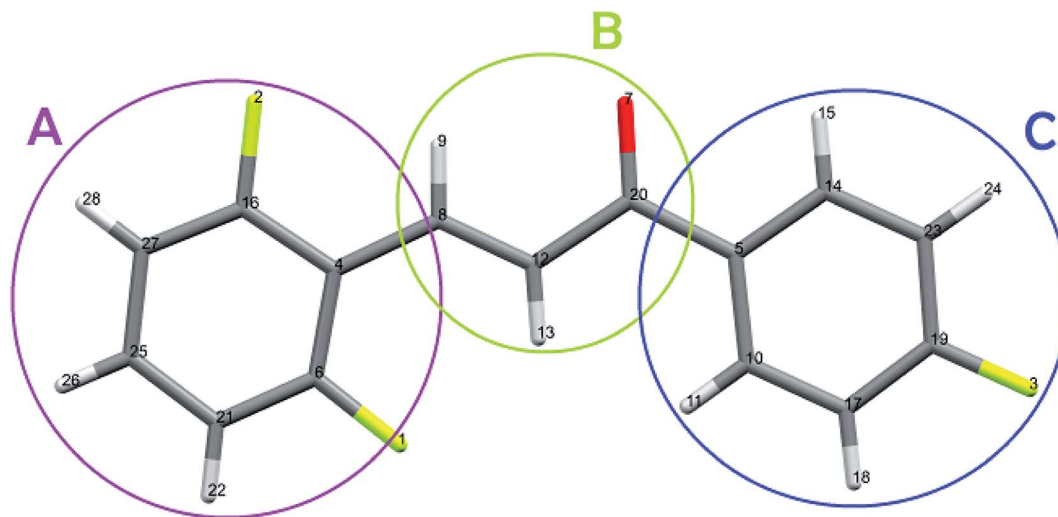


Fig. 14 Highlight of the rings A, B and C, with the atoms numbered as presented in Table 7.

The atomic charges of the FCH range between $-0.5520e$ (O7) and $0.6396e$ (C20) for isolated molecules, and between $-0.5957e$ (O7) and $0.6760e$ (C20) for embedded molecules, as can be seen in Table 7.

The carbon atoms bonded to the oxygen atom (C20) and to the fluorine (C6, C16 and C19) are positive due to the negativity of the neighboring atom. The crystalline polarization effect was more significant on the charges of the carbon atoms C14, C25, and C5, with increasing of 77%, 67% and 23% respectively (Fig. 14). All the hydrogen atoms have naturally positive charge and only atoms H15 and H9 had their charges decreased due to the embedded packing process.

4. Conclusion

In the present work, a new fluorinated chalcone derivative (FCH) was synthesized and crystals were grown by using slow evaporation method. Its structural characterization and determination of purity were carried out by scanning electron microscopy, infrared spectroscopy, gas chromatography-mass spectrometry, ^1H , ^{13}C and ^{19}F nuclear magnetic resonance, differential scanning calorimetry (DSC), Raman microspectroscopy, UV-Vis absorption spectroscopy, and single crystal X-ray diffraction (XRD). The Hirshfeld surface and fingerprints were performed to identify and quantify the intermolecular interactions that occur in the molecule, in order to provide information about crystal packing and considerable values in the quantification of interactions $\text{H}\cdots\text{F}$ (31.3%) by influence of fluorine substituents $\text{C}\cdots\text{C}$ (10.13%), which evidence the interaction $\pi\cdots\pi$ confirmed by the shape index. With this investigation, it is possible to report that crystal packing is stabilized by intermolecular factors, for example, interactions $\text{C}-\text{H}\cdots\text{O}$ and $\text{C}-\text{H}\cdots\text{F}$. The FMO indicates that FCH chalcone is an electrophilic species. MEP is consistent with the susceptible electrophilic attack around the oxygen atom of the carbonyl group. In general, the theoretical IR absorptions bands

are consistent with the experimental IR-FT bands and Raman spectra.

The crystalline environment of the FCH chalcone crystal was simulated through the supermolecular approach where a bulk with 378 000 atoms was built. The electric parameters such as the dipole moment, average linear polarizability and second hyperpolarizabilities corresponding to the Kerr effect and IDRI process were calculated at DFT/CAM-B3LYP/6-311++G(d,p) level as a function of the electric field frequency. The macroscopic parameters such as the linear refractive index and the third-order nonlinear susceptibility were also calculated, and the results compared with experimental data were obtained from literature. At $\omega = 0.086$ a.u. the $\chi^{(3)}$ -value for the Chalcone crystal is of $369.294 \times 10^{-22} \text{ m}^2 \text{ V}^{-2}$; this value is 1.33 times the experimental value for the 3MPNP⁵⁵ and hundreds of times higher than some values experimentally obtained for other organic compounds, as shown in Table 6. Therefore, the chalcone crystal must be considered as a potential material to use in the fabrication of photonic devices.

Conflicts of interest

There are no conflicts to declare.

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