


 Cite this: *RSC Adv.*, 2020, **10**, 23662

High-performance near-infrared Schottky-photodetector based graphene/In₂S₃ van der Waals heterostructures†

 Long Chen, Zhenghan Li and Chaoyi Yan *

Two-dimensional (2D) β -In₂S₃ is a natural defective n-type semiconductor attracting considerable interest for its excellent photoelectronic performance. However, β -In₂S₃ based photodetectors exhibited a weak near-infrared photoresponse compared to visible wavelength in past reports. In this work, high-quality 2D β -In₂S₃ nanosheets were prepared by a space-confined chemical vapor deposition (CVD) method. Graphene/In₂S₃ van der Waals heterostructures were constructed to realize an enhanced near-infrared photodetection performance by a series of transfer processes. The photodetectors based on graphene/In₂S₃ van der Waals heterostructures through junction carrier separation exhibited a better infrared performance of high photoresponsivity (R_{light}) of 0.49 mA W⁻¹, external quantum efficiency (EQE) of 0.07%, and detectivity (D^*) of 3.05×10^7 jones using an 808 nm laser.

 Received 3rd March 2020
 Accepted 19th May 2020

DOI: 10.1039/d0ra02033h

rsc.li/rsc-advances

Introduction

Since the discovery of graphene, two-dimensional (2D) materials have demonstrated extremely impressive mechanical, electronic and optoelectronic properties of mechanical flexibility, adjustable bandgap and strong light–matter interactions.^{1–13} With the booming development of the 2D metal sulfides, multiphase In₂S₃ has attracted extensive attention and been used for high-performance applications in photodetectors, phase change memory devices and catalysis.^{14–17} There are three main crystalline structures of α -In₂S₃, β -In₂S₃, and γ -In₂S₃ which have been reported.¹⁴ In particular, β -In₂S₃ shows a distinguished performance in different devices like photodetectors, FETs, phase-change memory devices and P–N diodes.^{18–20} Their excellent properties of high photosensitivity and absorption coefficient make them most promising candidates for electronic and optoelectronic applications. Although β -In₂S₃ have a visible bandgap from 1.9 eV to 2.3 eV, CVD-grown β -In₂S₃ show detection capabilities from visible light to near-infrared light due to natural defect in past reports.¹⁴ However, the near-infrared light response performance of β -In₂S₃ is very weak compared with the visible band. Therefore, graphene/In₂S₃ van der Waals heterostructures^{21–23} were constructed by a series of transfer process to enhance the near-infrared photoelectric detection performance. The photodetectors based on graphene/In₂S₃ van der Waals heterostructures

showed excellent performance with a high R_{light} of 0.49 mA W⁻¹, EQE of 0.07%, and D^* of 3.05×10^7 jones.

Results and discussion

High-quality ultrathin β -In₂S₃ nanosheets obtained through a space-confined chemical vapor deposition (CVD) (detail in the “Experimental”) method. Fig. 1a–c show the bright field and dark field optical images of the as-grown ultrathin triangle 2D β -In₂S₃ nanosheets growing on mica substrate with different confined distance (d) of 35, 20 and 10 μm , respectively. When d is 20 μm , the edge size of triangular β -In₂S₃ nanosheets synthesized can be up to 40 μm with an atomically flat surface and a thin thickness of about 1.5 nm shown in the corresponding atomic force microscope (AFM) image (the inset of Fig. 1b). With the decreasing of d , the thickness of β -In₂S₃ can decrease to 1.1 nm with a regular triangular shape as shown in Fig. 1c. The confined distance dependence of In₂S₃ sheets thickness and size is shown in Fig. S1.†

In order to further research more characteristics of β -In₂S₃, Raman spectrum was used to study the vibrational mode of β -In₂S₃. Fig. 1d show the Raman spectrum of as-grown β -In₂S₃ nanosheets with different thickness. Five characteristic peaks located at 246 cm⁻¹, 270 cm⁻¹, 310 cm⁻¹, 328 cm⁻¹ and 369 cm⁻¹ of β -In₂S₃ nanosheets were excited by 532 nm laser, which is consistent with previous reports of β -In₂S₃.¹⁴ Among five peaks, the peak of 246 cm⁻¹, 310 cm⁻¹ and 369 cm⁻¹ belong to the A_{1g} vibration mode of β -In₂S₃, namely the out-of-plane vibration mode. The peak of 270 cm⁻¹ belongs to the E_g vibration mode, namely the in-plane vibration mode, and the peak of 328 cm⁻¹ belongs to the F_g vibration mode. As the thickness of 2D β -In₂S₃ nanosheet decreases, the strength of all peaks

State Key Laboratory of Electronic Thin Films and Integrated Devices, University of Electronic Science and Technology of China, Chengdu 610054, China. E-mail: cyan@uestc.edu.cn

† Electronic supplementary information (ESI) available. See DOI: 10.1039/d0ra02033h



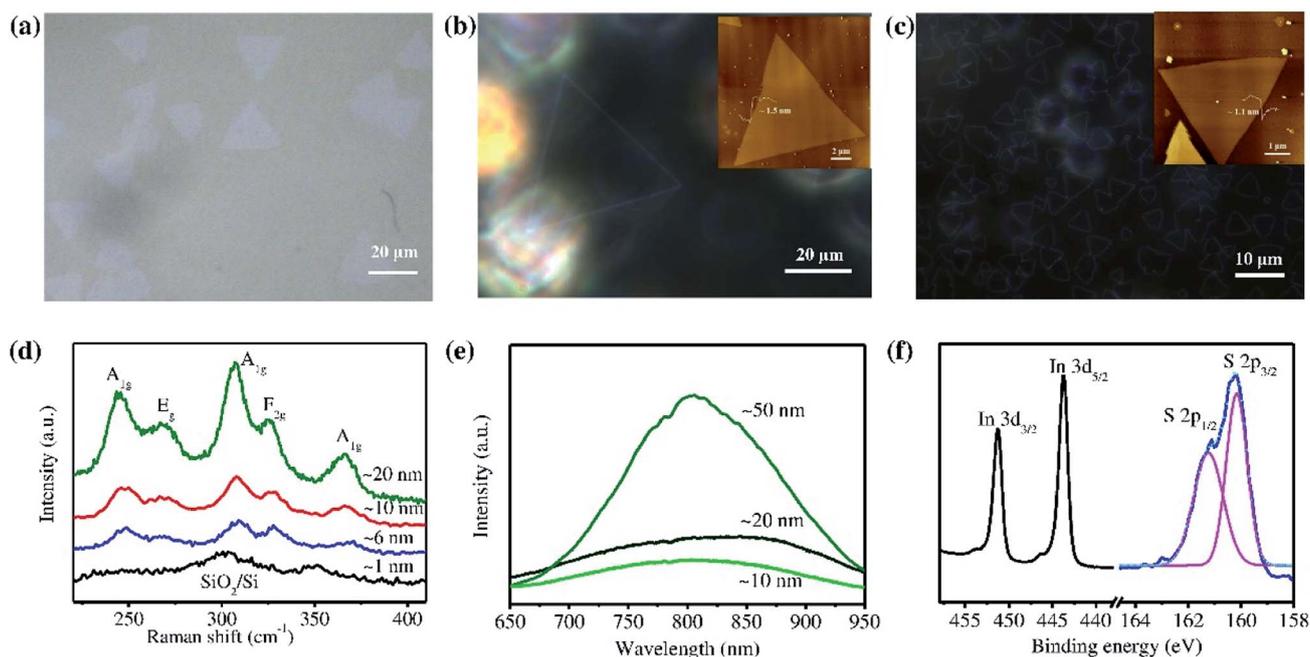


Fig. 1 Optical images of typical as-synthesized β - In_2S_3 nanosheets under different confined space (d): (a) $d = 35 \mu\text{m}$. (b) $d = 20 \mu\text{m}$, the inset is the corresponding AFM image. (c) $d = 10 \mu\text{m}$, the inset is the corresponding AFM image. (d) Raman spectrum of β - In_2S_3 sheets with different thickness of 1, 6, 10, and 20 nm. (e) The PL spectrum of β - In_2S_3 sheets with different thickness of 10, 20, and 50 nm. (f) XPS spectrum of as-synthesized β - In_2S_3 sheets.

decreases, but the peak position of each characteristic peak does not change. When the thickness of the sample was reduced to 1 nm, there are almost no characteristic peaks of β - In_2S_3 due to the greatly reduced scattering center, which is similar to other 2D III–VI group materials such as InSe ,²⁴ GaSe ²⁵ and GaTe .²⁶ Fig. 1d shows the PL spectrum of β - In_2S_3 nanosheets with different thickness obtained under the irradiation of incident light with a wavelength of 532 nm. In the wavelength range of 700–900 nm, β - In_2S_3 nanosheets have a wide emission spectrum, which is derived from the defect transition (E_d) generated by the recombination of the sample's intrinsic sulphur (S) vacancy and In vacancy. The luminescence intensity generated by defects increases with the increase of sample thickness, because the concentration of intrinsic defects in samples increases with the increase of thickness. However, due to quenching effect, PL spectrum cannot show information of intrinsic band gap (E_g) of β - In_2S_3 . Fig. 1e shows the X-ray photoelectron spectroscopy (XPS) spectra of In_{3d} and S_{2p} . It is found that the peak of In_{3d} are located in the 443.7 eV ($\text{In}_{3d_{3/2}}$) and 451.3 eV ($\text{In}_{3d_{5/2}}$) and the peak of the S_{2p} in 160.3 eV ($\text{S}_{2p_{3/2}}$) and 161.5 eV ($\text{S}_{2p_{1/2}}$). The results of the XPS were consistent with the previous β - In_2S_3 XPS results.¹⁴ Full XPS spectrum as shown in Fig. S2† demonstrated the as-grown β - In_2S_3 nanosheets obtained by the space-confined CVD method have high chemical purity.

In order to identify the crystal structure of 2D β - In_2S_3 , transmission electron microscope (TEM) was used to analyse the transferred β - In_2S_3 on the Cu grid. Fig. 2a show the low-magnification TEM image of a typical triangular β - In_2S_3 sheet.

It can be seen that the surface of the sample is clean and free of impurities. In the HRTEM image of β - In_2S_3 nanosheet, clear lattice diffraction spots show good crystallinity of β - In_2S_3 (shown in Fig. 2b). The spacing of 0.38 nm along the parallel line can be corresponds to the (220) crystal plane of β - In_2S_3 . And the corresponding the fast Fourier transform (FFT) was shown in the inset of Fig. 2b. X-ray energy dispersive spectrometer (EDS) was used to analyse the distribution of In and S elements of In_2S_3 nanosheets (Fig. 2c and d). In elements and S elements are evenly distributed on the entire triangle β - In_2S_3 sample. It has been proved that the high crystallization quality and clean surface β - In_2S_3 nanocrystals have been prepared by the CVD method.

Although β - In_2S_3 has a direct bandgap from 1.9 eV to 2.3 eV belong to visible wavelength, β - In_2S_3 -based photodetectors have a broad photoresponse range from visible to near-infrared wavelength due to the unique defective structure. However, β - In_2S_3 based photodetectors have a weaker light detection performance under the infrared radiation compared to these under the visible wavelength radiation in the past report.¹⁴ Graphene/ In_2S_3 -based Schottky-photodetectors constructed to enhance infrared photoelectric detection performance. Fig. 3a shows the current–voltage curve (I – V curve) of β - In_2S_3 -based photodetectors under the light wavelength of 808 nm and dark. Typical optical image of β - In_2S_3 based on photodetectors is shown in the inset of Fig. 3a. The dark current of β - In_2S_3 based on photodetectors is about 1.76 pA and the on/off ratio is only of 1.54. Fig. 3b shows the current–time (I – T) curve under the incident light irradiation with different light power density



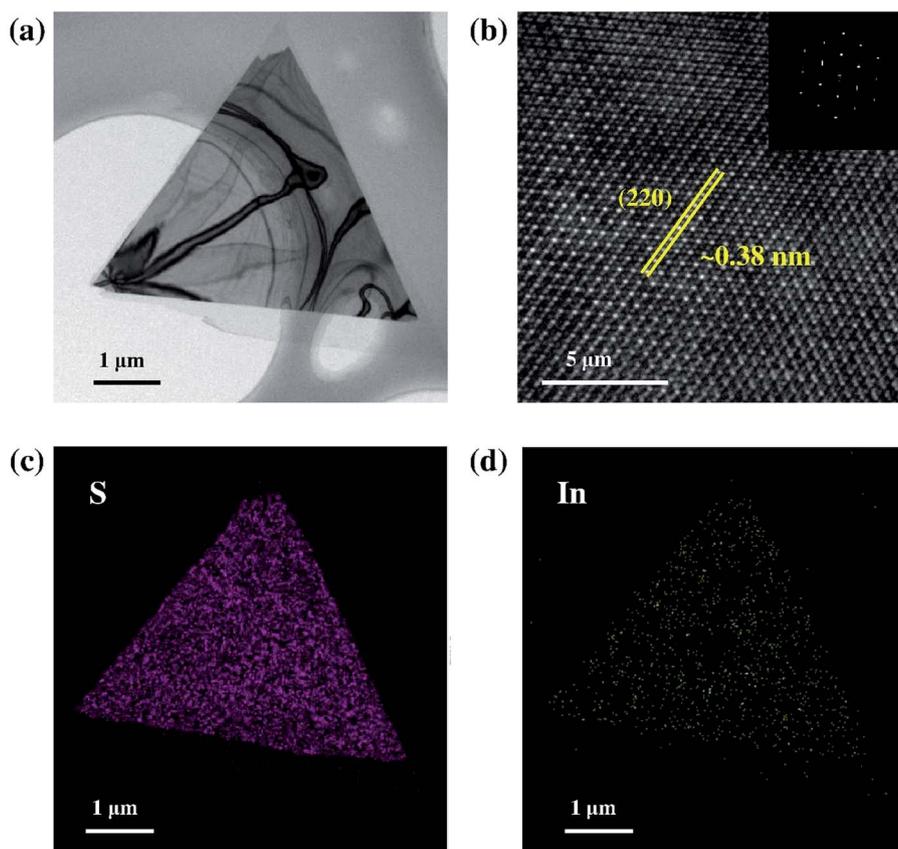


Fig. 2 TEM characterizations of the β - In_2S_3 nanosheet. (a) The low-magnification TEM image of a typical triangular β - In_2S_3 sheet. (b) HRTEM image of β - In_2S_3 sheet, the inset shows the corresponding the fast Fourier transform. (c) In and (d) S elemental mapping of the β - In_2S_3 nanosheet in (a).

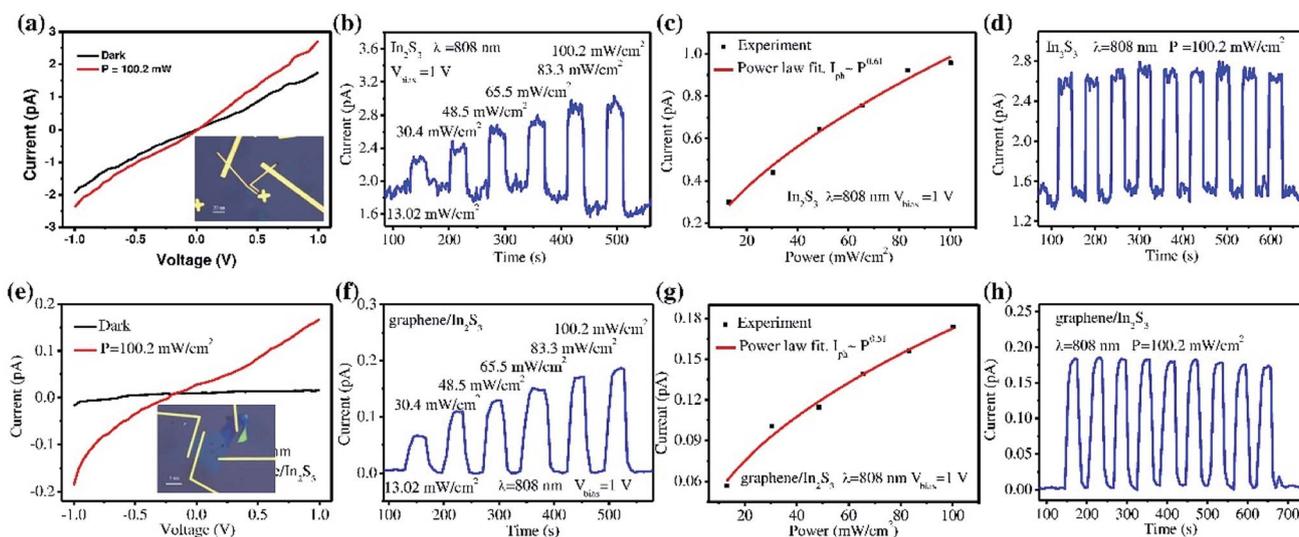


Fig. 3 Optoelectronic properties of β - In_2S_3 -based and graphene/ In_2S_3 -based photodetectors: the current–voltage curve (I – V curve) under a laser radiation of 808 nm and dark (a) β - In_2S_3 -based photodetector and (e) graphene/ In_2S_3 -based photodetector. The inset is typical optical image of β - In_2S_3 -based and graphene/ In_2S_3 -based photodetectors, respectively. The current–time (I – T) curve under the incident light irradiation with different light power density, $V_{\text{bias}} = 1$ V: (b) β - In_2S_3 -based photodetector and (f) graphene/ In_2S_3 -based photodetector. Photocurrent as a function of illumination power: (c) β - In_2S_3 -based photodetector and (g) graphene/ In_2S_3 -based photodetector. Time-resolved photo-response of photodetectors at a bias voltage of 1 V and illumination power of 100.2 mW cm^{-2} : (d) β - In_2S_3 -based photodetector and (h) graphene/ In_2S_3 -based photodetector.



when $V_{\text{bias}} = 1$ V. The photocurrent can be expressed by a power law of $I_{\text{ph}} \sim P^{\theta}$. The θ value of $\beta\text{-In}_2\text{S}_3$ based on photodetectors is about 0.61 (shown in Fig. 3c), which indicates that the near-infrared light response of $\beta\text{-In}_2\text{S}_3$ nanosheets to 808 nm wavelength is mainly derived from the defect state. The defect state in the material is easier to capture the photocarrier with the increasing of optical power density, resulting in the saturation of devices' photocurrent and the nonlinearly curve shown in Fig. 3c.^{14,27–32} Fig. 3d shows the stability of the $\beta\text{-In}_2\text{S}_3$ based photodetectors with periodic switching tests. Obviously, the photocurrent and dark current of $\beta\text{-In}_2\text{S}_3$ based photodetectors have a large noise.

Correspondingly, Fig. 3e is the I - V curve of heterojunction optoelectronic device irradiated by incident light at 808 nm wavelength and dark. The apparent difference is that the heterojunction device shows a large switching ratio of 11.16 and a smaller dark current of 0.015 pA at 1 V. The Raman spectrum with two kinds of material characteristic peak demonstrates that the vertical structure was constructed between the low layer graphene and $\beta\text{-In}_2\text{S}_3$ (Fig. S3†). Fig. 3f shows the light response curve under the irradiation of incident light with different light power density and the corresponding fitting value θ was 0.51 (Fig. 3g). The sub-linear relationship between the photocurrent and the optical power density curve may be caused by the intrinsic defects of $\beta\text{-In}_2\text{S}_3$ and the adsorption state at the interface with graphene and substrate. For stability testing, heterojunction devices exhibit more stable switching and image noise is much lower than that of individual devices (Fig. 3h).

In order to show the performance of heterogeneous junction device is improved comparing with that of the $\beta\text{-In}_2\text{S}_3$ based device, the key parameters of photodetector including R_{light} , EQE, and D^* were calculated and defined as follows:

$$R_{\text{light}} = I_{\text{light}} - I_{\text{dark}}/P_{\text{light}}S$$

$$\text{EQE} = hcR_{\text{light}}/e\lambda$$

$$D^* = R_{\text{light}}S^{1/2}/(2eI_{\text{dark}})^{1/2}$$

where I_{light} is photocurrent, I_{dark} is dark current, P_{light} is optical power density, S is effective device area, h is Planck's constant, c is the speed of light, e is electronic charge, λ is light wavelength. For $\beta\text{-In}_2\text{S}_3$ based on photodetectors, the R_{light} , EQE, D^* can be calculated to be 0.42 mA W⁻¹, 0.06% and 1.73×10^6 jones with the P_{light} of 13.02 mA cm⁻², $S = 6 \mu\text{m}^2$, $V_{\text{bias}} = 1$ V, $I_{\text{ph}} = 0.31$ pA. For the photodetector based on graphene/ In_2S_3 heterojunction, the R_{light} , EQE, D^* were calculated to be 0.49 mA W⁻¹, 0.07% and 3.05×10^7 jones at the bias of 5 V with $I_{\text{light}} = 0.057$ pA, $P_{\text{light}} = 13.02$ mW cm⁻², and $S = 10 \mu\text{m}^2$, respectively. Compared with $\beta\text{-In}_2\text{S}_3$ based photodetector, the D^* of heterojunction get a significant promotion with an order of magnitude indicating that the heterojunction device is more sensitive to near-infrared light.

The PL spectrum of graphene/ In_2S_3 heterojunction is tested to study the mechanism of improving the infrared detection performance (Fig. 4a). It can be seen that the PL spectrum of graphene/ In_2S_3 heterojunction is similar to that of $\beta\text{-In}_2\text{S}_3$, but

the strength of PL peak is significantly weakened and the quenching effect appears. The PL mapping as shown in Fig. 4b further illustrates this point. This phenomenon can be attributed that the carrier recombination is avoided due to the photogenerated carriers were effectively separated at the interface of $\beta\text{-In}_2\text{S}_3$ and graphene and flows into graphene and $\beta\text{-In}_2\text{S}_3$ at the junction area, respectively, which is shown in Fig. 4c. For photodetectors, due to the high light transmittance of graphene, the light absorption of heterojunction mainly comes from $\beta\text{-In}_2\text{S}_3$. When photons are absorbed by the heterojunction, the electrons separate from the holes, which are captured by $\beta\text{-In}_2\text{S}_3$, and the electrons flow to graphene. The captured holes in $\beta\text{-In}_2\text{S}_3$ can act as local grids in the graphene channel to regulate its charge transfer. The difference of electron-hole transport time and the increase of the lifetime of the photocarrier is the reason why the heterojunction has a great response to the near-infrared light.

Experimental

CVD growth of $\beta\text{-In}_2\text{S}_3$

Space-confined CVD is used for synthesis of 2D ultrathin $\beta\text{-In}_2\text{S}_3$ nanosheets. The quartz tube with the dual-temperature zone tubular furnace was used to support the stable growing environment as shown in Fig. S4.† Before growing process, 15 mg InI and 0.1 g S powder as precursors were placed in two Al_2O_3 porcelain boats, which were placed at the first heating zone of the upstream of tube furnace with an apart distant of 9 cm. Fresh fluorophlogopite mica sheets were selected as the growth substrate due to highly passivated and flat surface with extremely low migration energy of adsorbed atoms. "Face to face" superimposed mica sheets were placed on the center of heating zone 2. Before the reaction starts, the inside of the quartz tube is pumped into a vacuum. And then, the high-purity argon is used to expel the air in the quartz tube. This process was called "washing" process. After "washing" process, 60 sccm of high-purity argon gas was injected at a standard atmospheric pressure as the reaction carrier gas. The heating zone 1 heated up to 680 °C in 44 min and was kept for 10 min, while the heating zone 2 heated up to 420 °C in 42 min and was kept for 14 min. After heating zone 1 heats up to 680 °C and maintains for two minutes, push the upstream InI and S powder into the insulation zone of heating zone 1. After the reaction, the system will be cooled naturally to room temperature.

Device fabrication

As-grown $\beta\text{-In}_2\text{S}_3$ were transferred from mica to Si substrates by standard wet transfer process. PMMA is spun onto the mica substrate to protect $\beta\text{-In}_2\text{S}_3$ samples at 3500 rad s⁻¹ in 60 s. Then, the surface tension of water is used to peel the organic film off the mica substrate. After that, remove the organic film from the water to a silicon substrate. The silicon substrates with organic film were heated at 135 °C for 15 minutes to steam the water and stick to the film. Then, organic film was removed by acetone solution. The construction of heterogeneous structures requires a mechanical transfer process with polyvinyl alcohol



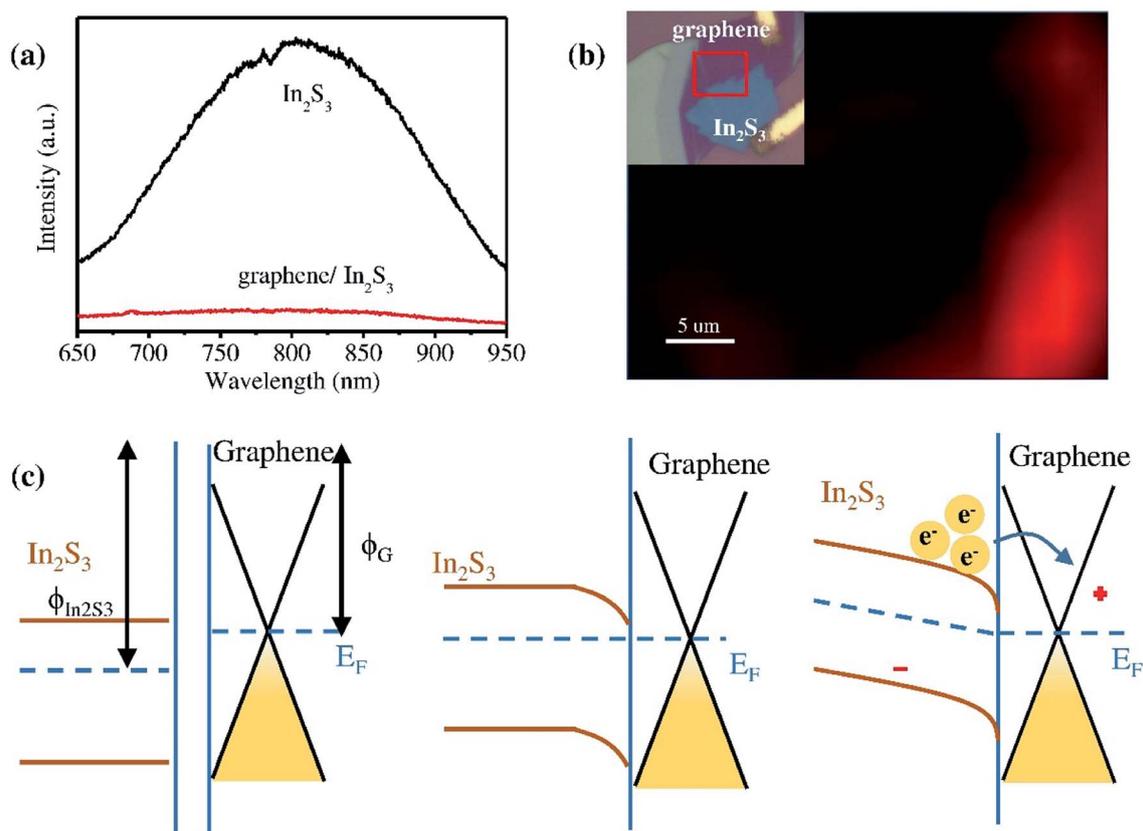


Fig. 4 (a) PL spectrum of the β - In_2S_3 sheet and graphene/ In_2S_3 heterojunction. (b) PL mapping of graphene/ In_2S_3 heterojunction, the inset shows the optical image of graphene/ In_2S_3 heterojunction. (c) Schematic diagram of electronics band structure of graphene/ In_2S_3 heterojunction.

(PVA) and polydimethylsiloxane (PDMS). Firstly, the PVA/PDMS/glass slide was pasted onto the target low-layer graphene sample and heated at 55°C for 3 minutes to peel the graphene from the silicon substrate. After that, the graphene/PVA/PDMS/glass slide was pasted onto the target β - In_2S_3 sample and heated at 70°C until the PDMS can be separated from PVA. After cooling to room temperature, the sample was immersed in deionized water to remove PVA.

The photodetectors based on transferred β - In_2S_3 and graphene/ In_2S_3 junction were fabricated by standard electron beam lithography (EBL) process. Patterning PMMA films were fabricated by EBL and the Cr/Au electrodes (10 nm/50 nm) were evaporated by thermal evaporation methods. Acetone was eventually used to remove excess photoresist and gold attached to photoresist.

Characterizations

The β - In_2S_3 nanosheets and graphene/ In_2S_3 junctions were characterized by the optical microscope (BX51, OLMPUS). The thickness and morphology of In_2S_3 and junctions are measured by the atomic force microscope (AFM, Demension icon, Bruker). A confocal Raman/PL system (Alpha 300RS+, WITec) is used to characterize the Raman and PL spectra of In_2S_3 nanosheets, graphene and junctions at room

temperature with a 532 nm laser. X-ray photoelectron spectroscopy (XPS, Axis Supra, Shimadzu) is adopted to analyze the electronic structure of In_2S_3 sheets. The transmission electron microscope (TEM, Tecnai G2 F20 S-Twin TMP, FEI) equipped with an X-ray energy dispersive spectrometer (EDS) is used to characterize the crystal structure of In_2S_3 . Photoelectric properties of devices based on In_2S_3 and graphene/ In_2S_3 junction are measured by a characterization system combined with semiconductor testing instrument (4200, Keithley), a probe station (CRX-6.5K, Lake Shore) and a laser source with adjustable wavelength of 532, 808 and 1550 nm.

Conclusions

In summary, we conducted a space-confined CVD method for the controlled synthesis of ultrathin 2D β - In_2S_3 flakes. Furthermore, the graphene/ In_2S_3 heterogeneous structures are implemented through a series of transfer methods. High infrared detection was realized with an excellent R_{light} of 0.49 mA W^{-1} , EQE of 0.07%, and D^* of 3.05×10^7 by junction carrier separation. Such a good detection performance based on 2D graphene/ In_2S_3 opens up opportunities for the emerging metal sulfide in future application of electronics and optoelectronics.



Conflicts of interest

There are no conflicts to declare.

Notes and references

- C. Gong, Y. Zhang, W. Chen, J. Chu, T. Lei, J. Pu, L. Dai, C. Wu, Y. Cheng, T. Zhai, L. Li and J. Xiong, *Adv. Sci.*, 2017, **4**, 1700231.
- A. Raja, L. Waldecker, J. Zipfel, Y. Cho, S. Brem, J. Ziegler, M. Kulig, T. Taniguchi, K. Watanabe, E. Malic, T. Heinz, T. Berkelbach and A. Chernikov, *Nat. Nanotechnol.*, 2019, **14**, 832–837.
- M. Wang, S. Cai, C. Pan, C. Wang, X. Lian, Y. Zhuo, K. Xu, T. Cao, X. Pan, B. Wang, S. Liang, J. Yang, P. Wang and F. Miao, *Nat. Electron.*, 2018, **1**, 130.
- L. Li, W. Han, L. Pi, P. Niu, J. Han, C. Wang, B. Su, H. Li, J. Xiong, Y. Bando and T. Zhai, *InfoMat*, 2019, **1**, 54.
- J. Kim, S. Baik, S. Ryu, Y. Sohn, S. Park, B. Park, J. Denlinger, Y. Yi, H. Choi and K. Kim, *Science*, 2015, **6249**, 723.
- A. Splendiani, L. Sun, Y. Zhang, T. Li, J. Kim, C. Chim, G. Galli and F. Wang, *Nano Lett.*, 2010, **10**, 1271.
- O. Lopez-Sanchez, D. Lembke, M. Kayci, A. Radenovic and A. Kis, *Nat. Nanotechnol.*, 2013, **8**, 497.
- H. Santos, A. Ayuela, L. Chico and E. Artacho, *Phys. Rev. B*, 2012, **85**, 245430.
- T. Hung, A. Rustagi, S. Zhang, P. Upadhyaya and Z. Chen, *InfoMat*, DOI: 10.1002/inf2.12095.
- T. Alonso-Lanza, A. Ayuela and F. Aguilera-Granja, *Phys. Chem. Chem. Phys.*, 2016, **18**, 21913.
- A. Ayuela Fernández, W. Jaskólski, H. Santos Expósito and L. Chico, *New J. Phys.*, 2014, **16**, 083018.
- C. W. Tan, S. P. Xu, Z. J. Tan, L. Z. Sun, J. X. Wu, T. R. Li and H. L. Peng, *InfoMat*, 2019, **1**, 54.
- C. Gong, Y. Zhang, W. Chen, J. Chu, T. Lei, J. Pu, L. Dai, C. Wu, Y. Cheng, T. Zhai, L. Li and J. Xiong, *Adv. Sci.*, 2017, **4**, 1700231.
- W. Huang, L. Gan, H. Yang, N. Zhou, R. Wang, W. Wu, H. Li, Y. Ma, H. Zeng and T. Zhai, *Adv. Funct. Mater.*, 2017, **27**, 1702448.
- X. Xie and G. Shen, *Nanoscale*, 2015, **7**, 5046.
- R. Sharma, Y. Chouryal, S. Chaudhari, J. Saravanakumar, S. Dey and P. Ghosh, *ACS Appl. Mater. Interfaces*, 2017, **9**, 11651.
- Y. Gao, S. Zhang, X. Bu and Y. Tian, *Catal. Today*, 2019, **1**, 271.
- R. H. Bube and W. H. McCarroll, *J. Phys. Chem. Solids*, 1959, **4**, 333.
- L. Bhira, H. Essaidi, S. Belgacem, G. Couturier, J. Salardenne, N. Barreaux and J. C. Bernede, *Phys. Status Solidi A*, 2000, **2**, 427.
- R. Sumi, A. Warriar and C. Vijayan, *J. Phys. D: Appl. Phys.*, 2014, **10**, 105103.
- Y. Liu, N. O. Weiss, X. Duan, H. Cheng, Y. Huang and X. Duan, *Nat. Rev. Mater.*, 2016, **1**, 16042.
- J. Lu, A. Wei, Y. Zhao, L. Tao, Y. Yang, Z. Zheng, H. Wang, D. Luo, J. Liu, L. Tao, H. Li, J. Li and J. Xu, *ACS Photonics*, 2018, **5**, 4912.
- G. Rao, X. Wang, Y. Wang, P. Wangyang, C. Yan, J. Chu, L. Xue, C. Gong, J. Huang, J. Xiong and Y. Li, *InfoMat*, 2019, **1**, 272.
- X. Yuan, L. Tang, S. Liu, P. Wang, Z. Chen, C. Zhang, Y. Liu, W. Wang, Y. Zou, C. Liu, N. Guo, J. Zou, P. Zhou, W. Hu and F. Xiu, *Nano Lett.*, 2015, **5**, 3571.
- O. Makkawi, Y. Qiu, W. Feng and A. Ping, *J. Nanosci. Nanotechnol.*, 2016, **9**, 9813.
- Z. Yang, W. Jie, C.-H. Mak, S. Lin, H. Lin, X. Yang, F. Yan, S. Lau and J. Hao, *ACS Nano*, 2017, **4**, 4225.
- P. Hu, Z. Wen, L. Wang, P. Tan and K. Xiao, *ACS Nano*, 2012, **6**, 5988.
- E. Liu, M. Long, J. Zeng, W. Luo, Y. Wang, Y. Pan, W. Zhou, B. Wang, W. Hu and Z. Ni, *Adv. Funct. Mater.*, 2016, **26**, 1938.
- B. Jin, P. Huang, Q. Zhang, X. Zhou, X. Zhang, L. Li, J. Su, H. Li and T. Zhai, *Adv. Funct. Mater.*, 2018, **28**, 1800181.
- J. Yao, Z. Zheng and G. Yang, *Adv. Funct. Mater.*, 2017, **27**, 1701823.
- F. Wang, T. Gao, Q. Zhang, Z. Y. Hu, B. Jin, L. Li, X. Zhou, H. Li, G. Van Tendeloo and T. Zhai, *Adv. Mater.*, 2019, **31**, 1806306.
- P. Hu, L. Wang, M. Yoon, J. Zhang, W. Feng, X. Wang, Z. Wen, J. C. Idrobo, Y. Miyamoto and D. B. Geohegan, *Nano Lett.*, 2013, **13**, 1649.

