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1. Introduction

As a prototypical electron-deficient element next to carbon in the periodical table, boron has a rich chemistry characterized with delocalized multicenter-two-electron (mc-2e) bonds in both bulk allotropes and polyhedral molecules.¹ Boron-based materials have found wide applications in field emissions, supercapacitors, optical absorptions, photodetectors, and etc.²⁻⁶ A bottom-up approach has received considerable attention in the past two decades to investigate the structural transitions from small boron nanoclusters to boron nanomaterials. Persistent joint photoelectron spectroscopy (PES) and firstprinciples theory investigations in the past two decades by Wang and coworkers have unveiled an unexpectedly rich landscape for size-selected boron clusters from planar or quasiplanar $B_n^{-/0}$ (n = 3-38, 41 and 42) to cage-like borospherenes $D_{2d} B_{40}^{-/0}$ and $C_3/C_2 B_{39}^{-}$. Multiple low-lying isomers appear to compete and coexist in most gas-phase boron $B_n^{-/0}$ clusters, starting from $n = 7.7^{-10}$ Ion-mobility measurements in combination with density functional theory (DFT) calculations, on the other hand, have shown that boron cluster monocations (B_n^+) possess double-ring tubular structures in the size range

Perfect cubic La-doped boron clusters $La_6 B_{24}$ [La@B₂₄]^{+/0} as the embryos of low-dimensional lanthanide boride nanomaterials[†]

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La-doped boron nanoclusters have received considerable attention due to their unique structures and bonding. Inspired by recent experimental observations of the inverse sandwich D_{8h} La₂B₈ (1) and tripledecker C_{2v} La₃B₁₄⁻ (2) and based on extensive global searches and first-principles theory investigations, we present herein the possibility of the perfect cubic La-doped boron clusters O_h La₆ \mathcal{B} [La@B₂₄]⁺ (3, ¹A_{1g}) and O_h La₆ \mathcal{B} [La@B₂₄] (4, ²A_{2g}) which appear to be the embryos of the metallic one-dimensional La₁₀B₃₂ (5) nanowire, two-dimensional La₃B₁₀ (6) nanosheet, and three-dimensional LaB₆ (7) nanocrystal, facilitating a bottom-up approach to build cubic lanthanide boride nanostructures from gas-phase clusters. Detailed molecular orbital and bonding analyses indicate that effective $(d-p)\sigma$, $(d-p)\pi$ and (d-p) δ covalent coordination interactions exist in La₆ \mathcal{B} [La@B₂₄]^{+/0} (3/4) clusters, while the 1D La₁₀B₃₂ (5), 2D La₃B₁₀ (6), and 3D LaB₆ (7) crystals exhibit mainly electrostatic interactions between the trivalent La centers and cubic B₂₄ frameworks, with weak but discernible coordination contributions from La (5d) \leftarrow B (2p) back-donations. The IR and Raman spectra of La₆ \mathcal{B} [La@B₂₄]^{+/0} (3/4) and band structures of La₁₀B₃₂ (5) and La₃B₁₀ (6) are computationally simulated to facilitate their future characterizations.

> between n = 16-25¹¹ unveiling another structural domain for boron nanoclusters. Transition-metal-doping has proven to induce dramatic structural and bonding pattern changes to boron clusters. Typical examples include the transition-metalcentered monocyclic boron wheel clusters $M@B_n$ (Co@B₈⁻, Ru(B_9^{-} , and Ta(B_{10}^{-}) with the maximum planar coordination number of CN = 10 and transition-metal-centered boron double-ring tubular clusters M@B_n⁻ (Co@B₁₆⁻, Rh@B₁₈⁻, and $Ta(B_{20})$ with the record tubular coordination number of CN = 20 known in experiments.¹²⁻¹⁸ A family of inverse sandwich dilanthanide boron complexes D_{7h} La₂B₇⁻, D_{8h} La₂B₈ (1), and D_{9h} La₂B₉⁻ with a perfect B_n ring (n = 7, 8, 9) sandwiched between two La atoms were observed in 2018 in PES measurements.19,20 The first inverse triple-decker di-lanthanide boron cluster C_{2v} La₃B₁₄⁻ (2) was discovered in 2019 in a joint experimental and theoretical investigation²¹ in which two conjoined $La_2B_8(1)$ inverse sandwiches share a B-B edge and a La vertex in the titled La-B₈-La-B₈-La motif. Boron clusters as typical Wankel motor molecules with fluxional bonds also exhibit structural fluxionalities where metal dopants play an important role.²²⁻²⁵ Our group very recently predicted the possibility of the inverse sandwich double-ring tubular molecular rotor C_{2h} La₂& $[B_2(a)B_{18}]$ which possesses the smallest core-shell structure reported in boron clusters.26

> Bulk lanthanide hexaboride LaB_6 (7) has been widely used as cathode materials due to its numerous excellent properties such as low work function, ultrahigh hardness, high chemical

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inertness, and high melting point.27-30 It possesses a typical cubic CaB₆-type lattice $(Pm\bar{3}m, O_h^{-1})$ characterized with a 3D network constituted of B₆ octahedrons, with the interstitial locations filled by trivalent La atoms.³¹ Interestingly, each unit cell in LaB₆ (7) contains a B₂₄ cubic framework which possesses six equivalent B₈ rings on the surface sharing twelve B-B dumbbells on the edges, with each B₈ ring sandwiched between two neighboring La atoms. The strong structural similarity between the gas-phase inverse sandwich La_2B_8 (1) complex and cubic three-dimensional (3D) LaB₆ (7) crystal has been previously noticed.³¹ LaB₆ nanowires in cubic LaB₆ (7) lattice were also recently confirmed to be active electrochemical materials for supercapacitors.³² However, a cubic embryo of the LaB_6 (7) lattice still remains unknown to date in a bottom-up approach starting from gas-phase clusters. Such a cubic embryo may also be used to form one-dimensional (1D) lanthanide boride nanowires and two-dimensional (2D) lanthanide boride nanosheets.

Keep the inspiration in mind and based on extensive global minimum (GM) searches and first-principles theory calculations, we predict herein the highly stable perfect cubic La-doped boron clusters $O_h \text{La}_6\&[\text{La}(B_{24}]^+ (3, {}^{1}A_{1g}) \text{ and } O_h \text{La}_6\&[\text{La}(B_{24}]^+ (3, {}^{2}A_{2g}) \text{ which consist of six equivalent La}_{2}B_8 (1) inverse sandwiches embedded in a cube sharing one La atom at center. More interestingly, La_6&[La(B_{24}]^{+/0} (3/4) clusters with a B_{24} framework turn out to be the cubic embryos of the metallic 1D La_{10}B_{32} (5) nanowire, 2D La_3B_{10} (6) nanosheet, and 3D LaB_6 (7) nanocrystal in a bottom-up approach. The newly obtained ferromagnetic 1D La_{10}B_{32} (5) and nonmagnetic 2D La_3B_{10} (6) are expected to serve as efficient electronical and optical nanomaterials.$

2. Theoretical procedure

Extensive GM searches were performed on $La_7B_{24}^+$ monocation using the TGmin2 code³³ at the DFT level, in combination with manual structural constructions based on the experimentally observed La_2B_8 (1) and $La_3B_{14}^-$ (2). Over 3000 trial structures were explored in both singlet and triplet states at PBE/TZVP. The low-lying isomers were then fully optimized at the PBE0 (ref. 34) and TPSSh³⁵ levels with the 6-311+G* basis set³⁶ for B and Stuttgart relativistic small-core pseudopotential for $La^{37,38}$ using the Gaussian 09 program suite.³⁹ Low-lying isomers of the open-shell neutral La_7B_{24} were obtained from the corresponding lowest-lying structures of $La_7B_{24}^+$.

The GM search on 1D La₁₀B₃₂ (5) was performed with the EOU program developed by our group.⁴⁰ A GM search based on the PSO technique implemented in the widely used Particle Swarm Optimization (CALYPSO)⁴¹ package was executed for 2D La₃B₁₀ (6). The calculations on 1D La₁₀B₃₂ (5), 2D La₃B₁₀ (6), and 3D LaB₆ (7) nanostructures were performed using the Vienna *ab initio* simulation package (VASP),^{42,43} within the framework of projector augmented wave (PAW) pseudopotential method^{44,45} and PBE generalized gradient approximation (GGA).^{46,47} The Coulomb-corrected local spin-density approximation (LSDA+U) was utilized for both structural relaxation and static calculation (U = 5 eV).^{31,48} The cutoff energy of the plane wave basis is set to

500 eV. Atomic structures are fully relaxed using the conjugate gradient method until the maximum force on each atom was less than 0.01 eV Å⁻¹ and the energy precision was set to 10^{-5} eV. Chemical bonding analyses were performed for La₆& [La@B₂₄]⁺ (3) using the adaptive natural density partitioning (AdNDP) approach at the PBE0 level^{49,50} and for 3D LaB₆ (7) crystal utilizing the solid-state AdNDP (SSAdNDP) method.^{51,52} Natural bonding orbital (NBO) analyses were performed on La₇B₂₄⁺ (3) using the NBO 6.0 program.⁵³ Born–Oppenheimer molecular dynamics (BOMD) simulations were performed on La₇B₂₄⁺ (3) for 30 ps at 300 K and 1000 K using the CP2K software package.⁵⁴

3. Results and discussions

3.1 Structures and stabilities

Starting from the experimentally observed $C_{2v} \operatorname{La_3B_{14}}^{-}(2)$ which possesses a tilted La-B₈-La-B₈-La inverse triple-decker structure with two conjoined B₈ rings sharing a B-B dumb-bell, we manually constructed the perfect cubic hepta-lanthanide-doped boron clusters O_h La₆&[La@B₂₄]⁺ (3, ¹A_{1g}) and O_h La₆&[La@B₂₄] $(4, {}^{2}A_{2g})$ from two C_{2v} La₃B₁₄⁻ (2) clusters which share two B-B dumb-bells on the edges and one La atom at the center, with two extra La atoms added on the top and bottom, respectively. With one La atom endohedrally coordinated at the center (ⓐ) and six La atoms exohedrally coordinated on the surface (&), $O_{\rm h}$ $La_{6} \& [La(a) B_{24}]^{+/0}$ (3/4) can be viewed as consisting of six equivalent $La_2B_8(1)$ inverse sandwiches embedded in a cube sharing twelve B-B dumb-bells on the edges and one La vertex at the center. Interestingly and encouragingly, as shown in Fig. S1 and S2,† extensive GM searches strongly suggest that the highly stable perfect cubic $La_6 \& [La(a)B_{24}]^+$ (3) and $La_6 \& [La(a)B_{24}]$ (4) with a B₂₄ framework be the well-defined GMs of the systems with the lowest vibrational frequencies of 111 cm⁻¹ and 112 cm⁻¹, respectively. La₆&[La@B₂₄]^{+/0} (3/4) possess the optimized La–B distances of $r_{\text{La-B}} = 2.94/2.95$ Å between the central La atom and its cubic B_{24} ligand and $r'_{La-B} = 2.70/2.69 \text{\AA}$ between the surface La atoms and their neighboring B₈ ligands on the surface, with the B–B distances of $r_{\rm B–B} = 1.67/1.67$ Å within the B–B dumb-bells on the edges and $r_{\rm B–B}^{'}=1.64/1.65{\rm \AA}$ within the B₃ triangles at the corners. The large HOMO-LUMO gap of $\Delta E_{gap} = 2.35$ eV calculated for La₆&[La@B₂₄]⁺ (3) well supports its high chemical stability. As shown in Fig. S1,† its second lowest-lying triplet C_1 La₆&[La@B₂₄]⁺ (³A) which is slightly distorted (with $r_{\text{La-B}} = 2.93-2.95$ Å and $r'_{\text{La}-\text{B}} = 2.70/2.71\text{\AA}$ due to Jahn-Teller effect lies 1.30 and 1.36 eV higher than the Oh GM in energy at PBE0 and TPSSH levels, respectively. Both the distorted cubic quintet and heptet isomers with the relative energies of 2.50 eV and 3.53 eV at PBE0, respectively, appear to be much less stable (Fig. S1⁺). Extensive molecular dynamics simulations indicate that La₆& $[La@B_{24}]^+$ (3) is also highly dynamically stable, with the calculated average root-mean-square-deviations of RMSD = 0.12 Å and maximum bond length deviations of MAXD = 0.38 Å at 1000 K. Similar situation exists in La₆&[La@B₂₄] (4). La₆& $[La@B_{24}]^{+/0}$ (3/4) are therefore the well-defined deep-lying GMs

of the systems highly stable both thermodynamically and dynamically.

We initially constructed a 1D La_5B_{16} nanowire (P4/mmm) (Fig. S3[†]) by extending the finite $La_6\&[La@B_{24}]^{+/0}$ (3/4) clusters periodically in one direction (x), resulting in a 1D nanowire with one small imaginary phonon frequency at 49i cm⁻¹. This imaginary phonon frequency corresponds to a typical Peierls phase transition which leads to the slightly more stable 1D $La_{10}B_{32}$ nanowire (5, P4/mmm) (by 0.07 eV) when the unit cell is doubled in size, with the two La atoms inside slightly offcentered by 0.22 Å (Fig. 1). Similarly, 2D La₃B₁₀ (6) (P4/mmm) and 3D LaB₆ (7) ($Pm\bar{3}m$) networks can be obtained by expanding $La_{6} \left[La(B_{24})^{+/0} (3/4) \right]$ periodically in (x, y) and (x, y, z) directions, respectively (Fig. 1). The 1D (5), 2D (6), and 3D (7) networks possess the optimized lattice parameters of a = 8.66 Å, a = b =4.17 Å and a = b = c = 4.16 Å at PBE level, respectively. Our calculated lattice parameters and La–B distance ($r_{La-B} = 3.06$ Å) for 3D LaB₆ (7) agree well with the corresponding experimentally measured values of a = b = c = 4.16 Å and $r_{\text{La-B}} = 3.05$ Å.³¹ The newly obtained 2D La_3B_{10} (6) possesses the optimized La-B distances of $r_{\text{La}-B(x)} = r_{\text{La}-B(y)} = 3.05$ Å and $r_{\text{La}-B(z)} = 3.07$ Å between the central La atom and its B24 cubic ligand. Similar La-B distances exist in 1D $La_{10}B_{32}$ (5). These La-B distances turn out to be slightly longer than the corresponding La-B coordination bond lengths of $r_{\text{La-B}} = 2.94/2.95$ Å in La₆& $[La@B_{24}]^{+/0}$ (3/4) clusters. Interestingly, as shown in Fig. S4 and S5,[†] 1D La₁₀B₃₂ (5) has the lowest cohesive energy among the 1D structures obtained via extensive EOU global searches, while La_3B_{10} (6) is the lowest-lying structure in all the 2D La_3B_{10} conformations probed by extensive PSO global searches, with the second lowest-lying 1D and 2D structures lying 3.13 eV and 1.88 eV higher in cohesive energy per unit cell than $La_{10}B_{32}$ (5) and $La_{3}B_{10}$ (6), respectively. 1D $La_{10}B_{32}$ (5) and 2D $La_{3}B_{10}$ (6) are thus strongly favored in thermodynamics compared with their other low-lying counterparts.



Fig. 1 Optimized structures of inverse sandwich D_{8h} La₂&B₈ (1), inverse triple-decker C_{2v} La₃@B₁₄⁻ (2), cubic O_h La₆&[La@B₂₄]⁺ (3), cubic O_h La₆&[La@B₂₄] (4), 1D La₁₀B₃₂ nanowire (5), 2D La₃B₁₀ nanosheet (6), and 3D LaB₆ nanocrystal (7).

To evaluate the relative stability of 1D $La_{10}B_{32}$ (5), 2D La_3B_{10} (6), and 3D LaB_6 (7) networks, we computed their average cohesive energy per atom

$$E_{
m coh} = rac{n E_{
m La} + m E_{
m B} - E_{
m La_n B_m}}{n+m}$$

where E_{La} , E_{B} , and $E_{\text{La}_n B_m}$ are the total energies of a single La atom, a single B atom, and a unit cell of the La_nB_m lattices and *n* and *m* are the numbers of La and B atoms, respectively. La₁₀B₃₂ (5), La₃B₁₀ (6), and LaB₆ (7) appear to have the calculated cohesive energies of $E_{\text{coh}} = 6.33$, 6.50 and 6.95 eV per atom, respectively. These values appear to be systematically higher than the corresponding cohesive energy per atom of $E_{\text{coh}} =$ 5.65 eV calculated for neutral embryo La₆&[La@B₂₄] (4).

3.2 Electronic structures and bonding patterns

The high stabilities of these cubic lanthanide boride nanostructures can be traced back to their unique electronic structures and bonding patterns. Detailed NBO analyses on closedshell $La_6 \& [La \textcircled{B}_{24}]^+$ (3) show that the central La atom possesses the natural atomic charge of $q_{\text{La}} = -0.50|e|$, electronic configuration of La[Xe]4f^{0.18}5d^{3.04}6s^{0.24}, and total Wiberg bond order of $WBI_{La} = 5.26$, while the six surface La atoms have the natural atomic charge of $q_{La} = +1.43 |e|$, electronic configuration of La[Xe]4f^{0.15}5d^{1.41}6s^{0.10}, and total Wiberg bond order of $WBI_{La} = 2.86$. Obviously, both the central and surface La atoms in La₆&[La@B₂₄]⁺ (3) donate their $6s^2$ electrons almost completely to the B₂₄ cubic ligand and there exists an effective La 5d \leftarrow B 2p back-donation from the cubic B₂₄ ligand to the central La atom. Especially, the central La atom with the 5d occupation of 5d^{3.04} has a much higher total La-B coordination bond order (5.26) than that of the partially coordinated surface La atoms (2.86) with the 5d occupation of 5d^{1.41}. Obvious La-B coordination interaction with the average La-B bond order of $WBI_{La-B} = 0.22$ exists between the central La and its 24 equivalent B ligands in La₆& $[La@B_{24}]^+$ (3). The La coordination centers in inverse sandwich $La_2\&B_8$ (1) and $La_3@B_{14}^-$ (2) possess the net atomic charges of $q_{\text{La}} = +1.29, +1.39/1.17|e|$, electronic configurations of La[Xe] 4f^{0.17}5d^{1.49}6s^{0.10}, La[Xe]4f^{0.13/0.16}5d^{1.37/1.63}6s^{0.11/0.11}, and total Wiberg bond orders of WBILa = 2.87, 2.88/3.08, repectively, similar to the situation in the six surface La atoms in La₆& $[La@B_{24}]^+$ (3). Detailed molecular orbital (MO) analyses indicate that all the occupied orbitals of the closed-shell $La_6 \& [La@B_{24}]^+(3)$ are effective bonding MOs, while its lowest unoccupied MO (LUMO, (a_{2g})) is basically an anti-bonding orbital mainly originated from the 5d atomic orbitals of the six surface La atoms (Fig. S6[†]). With a valence electron occupying the non-degenerate LUMO (a_{2g}) of $La_6 \& [La@B_{24}]^+$ (3), neutral $La_6 \& [La@B_{24}]$ (4) possesses a singly occupied highest occupied MO (SOMO (a_{2g})). Such an electron occupation helps maintain the Oh symmetry of the open-shell neutral which possesses the same geometry as $La_{6}\&[La@B_{24}]^{+}(3).$

To better understand the bonding nature and stabilization mechanism, we performed a detailed AdNDP bonding analysis on the closed-shell $La_6\&[La@B_{24}]^+$ (3) which recovers both the localized and delocalized bonds of the system. As clearly

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indicated in Fig. 2a, $La_6\&[La@B_{24}]^+$ (3) possesses 12 equivalent 2c-2e B-B σ bonds on twelve B-B dumb-bells with the occupation number of ON = 1.78|*e*| and 8 equivalent 3c-2e σ bonds over eight B₃ triangles at the corners of the cube with ON = 1.83|*e*|, forming the σ -skeleton on the cubic B₂₄ ligand. The remaining 26 delocalized bonds are mainly responsible for the La-B coordination interactions in the complex, including 6 10c-2e B₈(σ)-La₂(d_{σ}) bonds, 12 10c-2e B₈(σ)-La₂(d_{π}) bonds, and 6 10c-2e B₈(π)-La₂(d_{δ}) bonds with ON = 1.72-1.90|*e*| on six equivalent conjoined La₂B₈(1) inverse sandwiches embedded in the cube and 2 31c-2e B₂₄(π)-La₇(d_{σ}) bonds totally delocalized

over the whole molecule with ON = 2.00|e| (Fig. 2a, see Fig. S7[†] for more details). Overall, cubic La₆&[La@B₂₄]⁺ (3) has similar (d–p) σ , (d–p) π , and (d–p) δ coordination interactions with that of observed inverse sandwich La₂B₈ (1) and triple-decker La₃B₁₄⁻ (2).^{19–21} The bonding pattern presented above provides further evidence that all the occupied MOs of La₆&[La@B₂₄]⁺ (3) are intrinsically bonding MOs (Fig. S6[†]), while all its antibonding MOs remain empty. Such a unique electron occupation renders an exceptionally high stability to the monocation.

We also performed a detailed SSAdNDP analysis on the periodical 3D LaB₆ (7) cubic lattice in Fig. 2b. There are 3 equivalent 2c-2e B-B σ bonds on three B-B dumb-bells in the cubic unit cell with ON = 1.80|*e*| which correspond to the 12 2c-2e B-B σ bonds on the B₂₄ ligand in La₆&[La@B₂₄]⁺ (3). The remaining 8 equivalent 4c-2c B₃(σ)-La(d_{σ}) bonds between the central La atom and the eight B₃ triangles at the corners with ON = 1.77|*e*| represent mainly electrostatic interactions between the trivalent La center and B₂₄ cubic framework, with weak but discernible coordination contributions from La (5d) \leftarrow B (2p) back-donations (Fig. 2b). Similar bonding patterns exist in 1D La₁₀B₃₂ (5) and 2D La₃B₁₀ (6). Detailed Bader charge analyses⁵⁵ indicate that the central La atom in LaB₆ (7) carries the net atomic charges of $q_{La} = +1.71|e|$, with its B neighbors possessing the average atomic charge of $q_B = -0.29|e|$. The

formation of rare earth (RE) boride crystals REB₆ in infinite lattices involve an obvious electron transfer from the trivalent RE centers to the boron sublattice. 1D (5), 2D (6), and 3D (7) nanocrystals with the elongated La–B distances of $r_{\text{La-B}(z)} =$ 3.05-3.07 Å possess obviously weaker La–B covalent coordination interactions than that in La₆&[La@B₂₄]^{+/0} (3/4) clusters which have the La–B coordination distances of $r_{\text{La–B}} = 2.94-2.95$ Å. The size of the trivalent RE atoms, rather than their electron configuration, is the main factor responsible for their boride structures.⁵⁶

The calculated band structures and corresponding projected density of states (PDOS) of 1D $La_{10}B_{32}$ (5) and 2D $La_{3}B_{10}$ (6) nanostructures are depicted in Fig. 3. Notably, these nanostructures all appear to be metallic in nature. Both 2p orbitals from B atoms and 5d orbitals from La centers contribute to the calculated PDOS near the Fermi level, with La-5d orbitals making major contributions to the PDOS above the Fermi level, while B-2p orbitals dominating the PDOS below the Fermi level. Both spin unpolarized and spin-polarized computations were carried out to determine the ground states for these nanomaterials during the structural optimizations. 1D $La_{10}B_{32}$ (5) network turned out to be a ferromagnetic metal with the total magnetic moment of 2.03 $\mu_{\rm B}$ per unit cell which mainly originates from the surface La atoms (with each surface La carrying the magnetic moment of 0.12 $\mu_{\rm B}$, while each central La carrying the magnetic moment of only 0.008 $\mu_{\rm B}$). In contrast, both 2D $La_{3}B_{10}$ (6) and 3D LaB_{6} (7) are nonmagnetic in nature with the calculated magnetic moments of zero.

3.3 Simulated IR and Raman spectra of the $La_6\&[La@B_{24}]^+$ embryo

Infrared photodissociation (IR-PD) in combination with firstprinciples theory calculations has proven to be a powerful means to characterize novel gas-phase clusters.^{57,58} We computationally simulate the IR and Raman spectra of the closed-shell $La_6 \& [La @ B_{24}]^+$ (3) in Fig. 4a to facilitate its future experimental characterizations. The high-symmetry O_h $La_6 \& [La @ B_{24}]^+$





Fig. 3 Calculated band structures and projected densities of states (PDOS) of (a) 1D $La_{10}B_{32}$ (5) and (b) 2D La_3B_{10} (6) using the PBE functional.



Fig. 4 Simulated IR and Raman spectra of $O_h La_6 \vartheta [La@B_{24}]^+$ (3) at PBE0/6-311+G(d) level.

possesses relatively simple vibrational spectra, with sharp asymmetrical IR vibrations at 190 (t_{1u}), 301 (t_{1u}), 431 (t_{1u}), and 1215 (t_{1u}) cm⁻¹ and symmetrical Raman features at 187 (a_{1g}), 289 (t_{2g}), 497 (a_{1g}), 636 (t_{2g}), 1190 (t_{2g}) and 1298 (a_{1g}) cm⁻¹, respectively. The strong Raman peaks at 187 cm⁻¹ and 497 cm⁻¹ mainly represent the typical radial breathing modes (RBMs) (a_{1g}) of the six evenly distributed surface La atoms and the B₂₄ cubic framework, respectively. Such RBM spectral features can be used to characterize hollow boron structures.⁵⁹ As shown in Fig. S8,† similar but slightly more complicated IR and Raman spectra are obtained for the open-shell neutral La₆&[La@B₂₄] (4).

4. Conclusions

Based on extensive GM searches and first-principles theory calculations, we have predicted in this work the highly stable perfect cubic $O_{\rm h}$ La₆&[La@B₂₄]^{+/0} (3/4) which are the embryos of the metallic 1D La₁₀B₃₂ (5) nanowire, 2D La₃B₁₀ (6) nanosheet, and 3D LaB₆ (7) nanocrystal, building a bridge between gasphase lanthanide boride clusters and their low-dimensional nanomaterials. La₆&[La@B₂₄]^{+/4} (3/4) clusters and 1D (5), 2D (6), and 3D (7) nanocrystals exhibit strong similarities in structures, with the periodical nanostructures possessing obviously weaker La (5d) \leftarrow B (2p) back-donations than their cubic embryos. Using the commercially available $3D \text{ LaB}_6(7)$ as target, it is possible to synthesize the newly predicted 0D(3/4), 1D (5), and 2D (6) nanomaterials by laser ablations, chemical vaporizations, or arc-discharges. The construction strategy demonstrated in this work may be applicable to the design of more rare-earth boride nanomaterials.

Conflicts of interest

There are no conflicts to declare.

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