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## Composite formation in CdSe:Cu<sub>2</sub>Se nanocrystal films, charge transport characteristics and heterojunction performance†

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The use of nanocrystals as materials for potential technological applications depends on tailoring their properties through intentional doping with external impurities. We have used a new technique to synthesize nanocrystal thin films of CdSe:Cu<sub>2</sub>Se containing different weight percentages (wt%) of Cu<sub>2</sub>Se. The films were deposited on glass substrates at room temperature by co-evaporation of CdSe and Cu<sub>2</sub>Se powder in nitrogen gas at a pressure larger than that required for conventional thin film deposition. The films consisted of nanograins of CdSe doped with Cu<sub>2</sub>Se (i.e., nanograins of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se where x is the atom% of Cu<sub>2</sub>Se doped into CdSe) for lower wt% of Cu<sub>2</sub>Se, and nanocomposites of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se and Cu<sub>2</sub>Se for higher wt% of Cu<sub>2</sub>Se. An energy band diagram built using the Anderson model was used for discussing the heterojunction characteristics of the junction between nanograins of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se and Cu<sub>2</sub>Se. To investigate the usefulness of the nanocrystal thin films of CdSe:Cu<sub>2</sub>Se for practical applications, the *I*–*V* characteristics of p–p and p–n hetero-junctions formed by the films respectively with nanostructured films of similarly deposited Cu<sub>2</sub>Se and CdSe films were studied.

## Introduction

Semiconductor nanocrystals (NC) exhibit interesting size dependent properties which markedly deviate from their bulk properties, but additional control of these properties is required for their practical applications. The use of nanocrystals as materials for potential technological applications depends on tailoring their properties through intentional doping with external impurities. In the case of bulk semiconductors, under thermal equilibrium determined by Gibbs free energy and growth temperature, impurity atoms can be incorporated up to their solid solubility limit which can be as much as 50% for II–VI semiconductors. But the addition of only a few dopant atoms into nanomaterials, quantum dots (QDs), can make them heavily doped and doping foreign atoms into quantum dots is a challenging problem.<sup>1–5</sup> When doping of a semiconductor with another material is attempted, the dopant in excess of the quantity that gets incorporated into the semiconductor through doping would separate out as a second phase resulting in the formation of a composite of the doped semiconductor along with the dopant. The desirable properties of different nanoscale

building blocks can be combined through the synthesis of their nanocomposites which offer an opportunity for creating materials with tailored properties for the fabrication of devices for specific applications. Nanocomposites of n- and p-type semiconductors are of importance for technological applications in optoelectronic devices including photovoltaics.<sup>5</sup> Synthesis of nanocomposites such as metal–metal,<sup>6</sup> metal–polymer,<sup>7</sup> metal–dielectric,<sup>8</sup> dielectric–metal,<sup>9</sup> semiconductor–polymer,<sup>10</sup> semiconductor–dielectric,<sup>11</sup> and semiconductor–semiconductor<sup>12</sup> in different morphologies like powders, mesocrystals and films for specific applications such as solar cells, nanocomposites are required in the form of films it would be advantageous if the components of the nanocomposite are in the nanosize regime. These requirements necessitate that both the components of a two phase nanocomposite should be generated as nanoparticles which deposit simultaneously on substrates to form of nanostructured films with non-resistive inter-particle contacts. Nanu *et al.*<sup>13</sup> used atomic layer chemical vapour deposition to infiltrate CuInS<sub>2</sub> inside the pores of nanostructured TiO<sub>2</sub> to prepare nanocomposite of n-type TiO<sub>2</sub> and p-type CuInS<sub>2</sub>. Beal *et al.*<sup>14</sup> employed radio frequency sputtering to produce CdTe–ZnO nanocomposite thin films with varied semiconductor-phase extended structures and studied the optical absorption and the carrier transport behaviour of these films. Potter *et al.*<sup>15</sup> used a dual-source, sequential radio-frequency sputter deposition technique to study the spatial distribution of CdTe nanoparticles within ZnO thin-film matrix.

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Nanostructured films of doped semiconductors and nanocomposites have received much attention recently because of their unique electronic<sup>16–18</sup> and optical properties,<sup>19</sup> and important applications such as photochemical solar cells<sup>20</sup> transparent conductive films<sup>21</sup> laser diodes<sup>22</sup> and in magnetic storage devices.<sup>23–25</sup> CdSe is an important II–VI semiconductor and has been extensively investigated for studying quantum size effects in semiconductor quantum structures.<sup>26–28</sup> Zhang *et al.*<sup>29</sup> through theoretical studies predicted that CdSe can easily be n-doped and lightly p-doped. Cu<sub>2</sub>Se is a p-type semiconductor due to copper vacancies and has an indirect band gap of 1.1–1.27 eV<sup>30</sup> which makes it an absorber material and a direct bandgap between 2.0 and 2.3 eV which makes it a window material in solar cells.<sup>31–33</sup> P-type doping of CdSe with Cu<sub>2</sub>Se can tune the bandgap of CdSe for varying its optical response and this can improve the photocurrent conversion efficiency by providing extra holes. Simi *et al.*<sup>34</sup> reported the synthesis of type II core–shell quantum dots by a high temperature organometallic method. They inferred that the charge separation in the CdSe–Cu<sub>2</sub>Se core–shell system makes it a good choice for designing light harvesting devices. Xu *et al.*<sup>35</sup> synthesized (CuS)<sub>x</sub>:(ZnS)<sub>1-x</sub> ( $x = 1$  to  $x = 0$ ) nanocomposite system with the aim of maximizing both transparency and hole conductivity and found that the hole conductivity was linked to the proportion of the CuS phase.

In the present work, we have used a new technique to synthesize nanocrystal thin films of CdSe:Cu<sub>2</sub>Se containing different weight percentages (wt%) of Cu<sub>2</sub>Se. In this technique, the films were deposited on glass substrates at room temperature by co-evaporation of CdSe and Cu<sub>2</sub>Se powder in nitrogen gas using a vacuum coating unit at a pressure larger than that required for conventional thin film deposition. While the constituent phases of a composite can be precisely determined using X-ray diffraction, the substitutional incorporation of a dopant into the host lattice can also be known through X-ray diffraction since the intensity of a diffraction peak depends on the structure factor which depends on the distribution of atoms on the set of plane responsible for the peak. Because of the influence of deposition process, the synthesized films consisted of nanograins of CdSe doped with Cu<sub>2</sub>Se (*i.e.*, nanograins of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se where  $x$  is the atom% of Cu<sub>2</sub>Se doped into CdSe) for lower wt% of Cu<sub>2</sub>Se, and nanocomposites of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se and Cu<sub>2</sub>Se for higher wt% of Cu<sub>2</sub>Se. The electrical conductivity of the CdSe:Cu<sub>2</sub>Se thin films was found to be improved due to the p-type doping with Cu<sub>2</sub>Se. Energy band diagram built using the Anderson model considering the films to be made up of nanocomposites of nanograins of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se and Cu<sub>2</sub>Se for higher wt% of Cu<sub>2</sub>Se was used for discussing the heterojunction characteristics of the junction between nanograins of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se and Cu<sub>2</sub>Se, for the first time. To investigate the usefulness of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se for practical applications such as solar cells, the *I*–*V* characteristics of p–p and p–n heterojunctions formed by the films respectively with nanostructured films of similarly deposited Cu<sub>2</sub>Se and CdSe films were studied.

## Experimental

Nanocrystal thin films of CdSe:Cu<sub>2</sub>Se were deposited on glass substrates by thermal evaporation of CdSe and Cu<sub>2</sub>Se powders

for different wt% of Cu<sub>2</sub>Se. Glass substrates were cleaned with detergent, etched for 5 min in dilute nitric acid, degreased with methanol and ultrasonically cleaned with acetone and then dried in hot air before deposition of the films. The cleaned substrates were then dried in hot air. Depositions of the films were carried out not at high vacuum but at an appreciably high pressure of nitrogen gas (99.99%) of  $\sim 1 \times 10^{-4}$  mbar which was much larger than the vacuum required for thin film deposition. On evaporating CdSe and Cu<sub>2</sub>Se powders, supersaturation of the nitrogen gas with vapours of these materials produced nanoparticles which deposited on the substrates as nanostructured thin films of CdSe:Cu<sub>2</sub>Se. Nanostructured films were deposited for different wt% of 5, 10, 20 and 30% of Cu<sub>2</sub>Se. The film samples were assigned the sample codes CS5, CS10, CS20, and CS30 where the number after CS denotes the doping percent of Cu<sub>2</sub>Se.

The crystallinity, crystal structure and growth orientation of the film samples were investigated using X-ray diffraction (XRD). A Bruker AXS D8 advance X-ray diffractometer equipped with Kristalloflex 780, KF. 4KE X-ray source ( $\lambda \sim 1.5406 \text{ \AA}$ ) was used to record the XRD pattern of the samples. The average thickness of the film samples was measured using a Bruker Dektak-XT stylus profiler. Energy Dispersive Spectroscopy (EDS) measurements were performed using JEOL Model JED-2300 Energy Dispersive Spectrometer. AFM measurements were carried out on the film samples using a Digital Instruments Nanoscope-E and using a Si<sub>3</sub>N<sub>4</sub> 100  $\mu\text{m}$  cantilever having a force constant of 0.58 N m<sup>-1</sup> in contact mode. The optical transmission spectrum of the film sample was recorded using a JASCO V-650 double beam spectrophotometer at room temperature. Raman measurements on the sample were performed using Jobin Yvon Horiba Labram – HR 800 micro-Raman spectrometer with a He–Ne laser (633 nm) as the excitation source. XPS measurements were done using ESCA Instrument: VSW, UK with a Resolution:  $\sim 1$  eV, source: Al K alpha X-rays.

For DC and photoconductivity measurements, silver coplanar electrodes were painted at a separation of 2 mm on the surface of the films deposited on glass substrates. Electrical measurements on the samples were carried keeping the samples in a cell at a vacuum of  $\sim 0.001$  mbar. Ohmic nature of the contacts of the electrodes with the sample was verified from *I*–*V* measurements using a Keithley 2400 source meter. The resistance of the samples was measured using a standard DC two point probe method over the temperature range from 300 to 423 K with an applied voltage of 3 V corresponding to the ohmic region of the *I*–*V* plot. The dc conductivity ( $\sigma$ ) was determined using the relation,  $\sigma = \frac{1}{\rho}$ . The resistivity  $\rho$  was calculated as

$$\rho = \frac{Rbt}{l}$$
, where  $R$  is the resistance,  $b$  is the distance between the electrodes,  $t$  the thickness and  $l$  is the length of the electrodes. Hall measurements were carried out on the samples CS20 and CS30 using PPMS (Quantum Design DynaCool) (the samples CS5 and CS10 were too resistive for Hall measurements). After loading the sample of 4  $\times$  8 mm size on the sample holder, a linearity check was performed to confirm the ohmic behaviour



of the electrical contacts. To investigate the usefulness of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se for practical applications, the *I*-*V* characteristics of homo and heterojunctions formed by the films respectively with nanostructured thin films of similarly synthesized CdSe and Cu<sub>2</sub>Se films were measured.

## Results and discussion

The XRD patterns of nanocrystal thin films of pure CdSe, and of CdSe:Cu<sub>2</sub>Se containing different Cu<sub>2</sub>Se wt% of 5, 10, 20 and 30% are shown in Fig. 1. The EDS spectrum of CdSe:Cu<sub>2</sub>Se nanostructured thin film sample containing 10 wt% of Cu<sub>2</sub>Se is shown in Fig. S1† and the atom% (at%) of Cd and Cu in the film samples are given in Table S1.† The at% of Cu in the film samples CS5, CS10, CS20 and CS30 were ~4, ~9, ~18 and ~30 respectively. The XRD pattern of the sample CS5 shows a strong peak at 2 $\theta$  value of 25.4° corresponding to the (002) reflection and weak peaks at 2 $\theta$  values of 24° and 42° corresponding to the (100) and (110) reflections of CdSe of hexagonal phase (ICDD file no. 01-071-4772). The XRD pattern shows no peaks corresponding to Cu<sub>2</sub>Se even though EDS analysis results showed the presence of Cu<sub>2</sub>Se in the sample leading to the inference that Cu<sub>2</sub>Se got doped into CdSe lattice through replacement of Cd by Cu. The XRD pattern of the sample containing 10 wt% of Cu<sub>2</sub>Se (CS10) shown in Fig. 1(c) is identical with the pattern of CS5 in Fig. 1(b) indicating that the sample CS10 also does not contain Cu<sub>2</sub>Se as a separate phase but it got incorporated into CdSe through doping. In the XRD patterns of the samples CS20 in Fig. 1(d), in addition to the (002), (100) and (110) peaks of CdSe, a weak peak like feature at 2 $\theta$  value of 26.78° is observed, while in the XRD pattern of the sample CS30 in Fig. 1(e) a weak peak is observed at 2 $\theta$  value of 26.78° in addition to the CdSe peaks. The peak at 26.78° was identified as the (111) reflection of Cu<sub>2</sub>Se (ICDD file no. 03-065-2982) showing that Cu<sub>2</sub>Se existed as

a separate phase in the samples CS20 and CS30. The intensity of this peak increased in the pattern of CS30 in Fig. 1(e) compared to that in Fig. 1(d) of the sample CS20. From the XRD results it was inferred that the samples CS20 and CS30 are CdSe:Cu<sub>2</sub>Se nanocomposites with some percentage of Cu<sub>2</sub>Se doped into the CdSe lattice.

From the XRD results, it is possible to derive information on the nature of incorporation of Cu<sub>2</sub>Se into the nanocrystal thin films of CdSe:Cu<sub>2</sub>Se samples. The substitutional incorporation of dopant atoms in a host material would affect the atomic scattering factor and the structure factor causing a change in the relative intensities of the diffraction peaks.<sup>16</sup> From Fig. 2(a), it is clear that the intensity of the (002) peak of CdSe decreased with increase in the concentration of Cu<sub>2</sub>Se incorporated in the samples. The decrease in intensity of (002) plane is due to the incorporation of Cu<sub>2</sub>Se into the CdSe:Cu<sub>2</sub>Se samples. Table 1 shows that the thickness of the film samples increased with an increase in wt% of Cu<sub>2</sub>Se. Hence, the decrease in the intensity of the (002) peak in the XRD pattern is not due to decrease in the mass of the material scattering the X-rays but is due to decrease in the scattering power of the material of the samples with the increase in the wt% of Cu<sub>2</sub>Se. The atomic scattering factor is equal to the number of electrons in each atom and is 48 and 29 for Cd and Cu, respectively. Hence, the replacement of Cd<sup>2+</sup> by two numbers of Cu<sup>1+</sup> may be expected to decrease the intensity of a particular XRD peak depending on the structure factor.<sup>36,37</sup> The intensity of (hkl) peak in the XRD pattern of a material is represented by the structure factor (*F*) given by:

$$F_{hkl} = \sum_{N=1}^3 f_N \exp[2\pi i(hu_N + kv_N + lw_N)]$$

where *u<sub>N</sub>*, *v<sub>N</sub>* and *w<sub>N</sub>* are the fractional coordinates of the atoms and *N* represents the different types of atoms (*i.e.*, Cd, Cu and Se) in a unit cell of CdSe doped with Cu<sub>2</sub>Se. *f<sub>1</sub>*, *f<sub>2</sub>* and *f<sub>3</sub>* are the atomic scattering factors of Cd, Cu and Se respectively. It was found that the intensity of the (002) peak of CdSe decreased with increase in the wt% of Cu<sub>2</sub>Se incorporated into the sample (Fig. 2(a)) indicating that the incorporation of Cu<sub>2</sub>Se was through substitutional doping. From the XRD patterns in Fig. 1, it is clear that the intensity of the (111) peak of Cu<sub>2</sub>Se increased with increase in the concentration of Cu<sub>2</sub>Se from 20 wt% to 30 wt%. This revealed the existence of Cu<sub>2</sub>Se as a separate phase in the samples CS20 and CS30 and hence the samples CS20 and CS30 are nanocomposite films.

The (002) peaks of all the samples are enlarged and shown together in Fig. S2.† Fig. S2† shows that the (002) peak shifted towards lower 2 $\theta$  values for lower Cu<sub>2</sub>Se doping wt% and then shifted to larger 2 $\theta$  values for higher concentrations of Cu<sub>2</sub>Se. The shift towards lower 2 $\theta$  values is attributed to the shrinkage of CdSe crystal lattice due to the substitution of Cd<sup>2+</sup> ions of the larger ionic radius (0.078 nm) by Cu<sup>1+</sup> ions of the smaller radius (0.046 nm). The shifting towards higher 2 $\theta$  values for the highest concentration of Cu<sub>2</sub>Se may be a result of composite formation. The average crystallite size of the CdSe grains in the samples was calculated from the full-width at half-maximum of the (002) diffraction peaks using the Debye–Scherrer formula,<sup>38</sup>

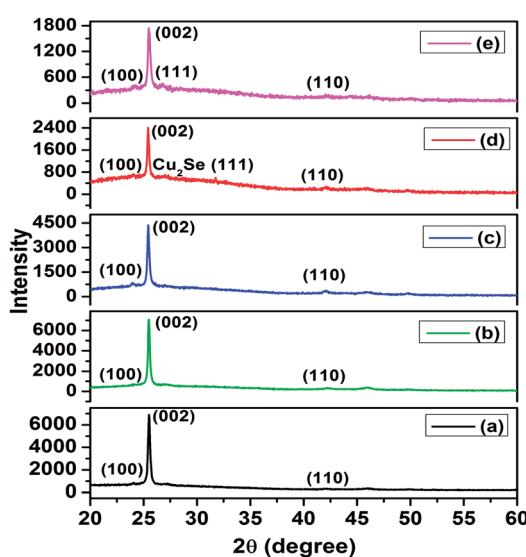


Fig. 1 XRD patterns of nanocrystal thin films of (a) pure CdSe, and of CdSe:Cu<sub>2</sub>Se containing different concentrations of Cu<sub>2</sub>Se: (b) 5 wt%, (c) 10 wt%, (d) 20 wt% and (e) 30 wt%.



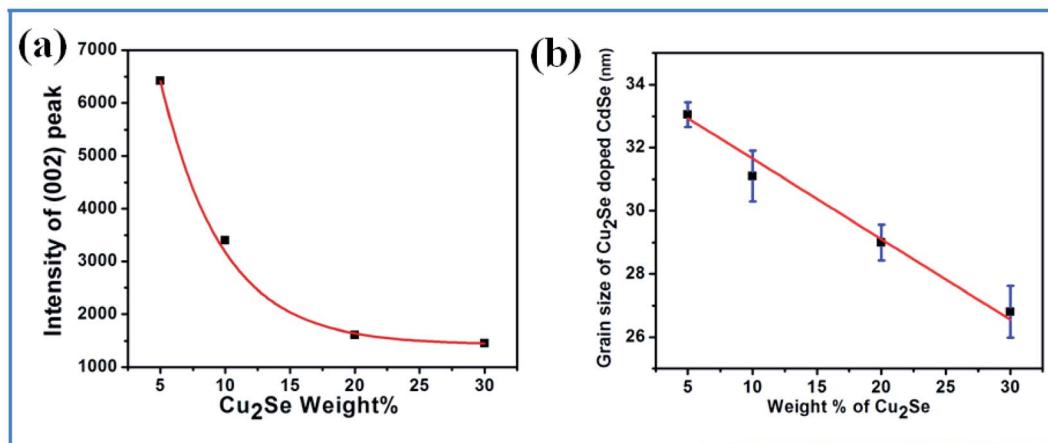


Fig. 2 (a) Plot of XRD intensity of (002) peak of CdSe with Cu<sub>2</sub>Se wt%; (b) plot of variation of grain size of CdSe in the nanocrystal thin films of CdSe:Cu<sub>2</sub>Se with wt% of Cu<sub>2</sub>Se.

$$D = \frac{k\lambda}{\beta \cos \theta}$$

where  $D$  is the mean grain size,  $k$  is a geometric factor ( $=0.89$ ),  $\lambda$  is the X-ray wavelength (1.5406 Å for Cu-K $\alpha$ ),  $\beta$  is the FWHM of diffraction peak, and  $\theta$  is the Bragg diffraction angle. The sizes ( $D$ ) of the CdSe grains in the nanocrystal thin films of CdSe:Cu<sub>2</sub>Se were obtained as 33.0, 31.0, 29.0 and 26.8 nm for the film samples containing 5, 10, 20 and 30 wt% of Cu<sub>2</sub>Se respectively (Table 1). In a recent report, Manu and Khadar<sup>16</sup> observed a decrease in the size of Fe doped nanocrystals of TiO<sub>2</sub> with the increase in the doping% of Fe and this observation was explained based on self purification mechanism prevalent in nanoparticles. Due to the high local energy associated with the dopant, nanoparticles always have a tendency to expel impurity atoms because of the thermodynamic requirement of minimising the energy of the system, which is the self-purification mechanism.<sup>16,39</sup> As a result, the concentration of dopant atoms near the grain boundaries would always be greater than that near the core of the nanocrystals. For the growth of nanoparticles of the host material, only the host atoms on the surface of the grains of the films can act as growth sites. With increase in the wt% of Cu<sub>2</sub>Se in the nanocrystal thin films of CdSe:Cu<sub>2</sub>Se, the number of the dopant copper ions near the surface of the host nanograins increases due to self purification mechanism and as a result the total number of available growth sites for the host Cd ions decreases leading to a reduction in the size of the CdSe nanograins with increase in the wt% of Cu<sub>2</sub>Se (Fig. 2(b)).

The SAED patterns (Fig. 3) of the sample CS5 containing 5 wt% of Cu<sub>2</sub>Se show rings corresponding to (103), (300), (310) and (323), and that of sample CS20 containing 20 wt% of Cu<sub>2</sub>Se show rings corresponding (103), (300), (310), (321) and (323) reflections of CdSe, respectively. The SAED pattern of the sample CS20 show ring corresponding to (640) reflection of Cu<sub>2</sub>Se indicating that the sample is nanocomposite. The spotty nature of the SAED pattern is an indication of the nanocrystalline nature of the samples.

X-ray photoelectron spectroscopy (XPS) measurements were carried out on CdSe:Cu<sub>2</sub>Se nanostructured film samples CS5 and CS20 containing 5 and 20 wt% Cu<sub>2</sub>Se to determine the valence states of the elements present in the samples. The XPS spectra of CS5 and CS20 in Fig. S3<sup>†</sup> show peaks corresponding to the elements C, Cd, Cu and Se elements. The C 1s peak located at 284.6 eV was taken as the internal standard for all the other elements and their peak positions were normalized with respect to the C 1s peak.<sup>40</sup> The binding energies corresponding to the peaks Cd 3d<sub>5/2</sub>, Cd 3d<sub>3/2</sub>, Cu 2p<sub>3/2</sub>, and Se 3d<sub>5/2</sub> obtained from the XPS spectra are 402.8, 411, 934 and 53.8 eV respectively.

The presence of doped Cu ions in samples CS5 and CS20 was investigated using narrow scan Cu 2p XPS data shown in Fig. 4. XPS plots of Cu 2p region show peaks at binding energies of 932.3 eV and 952.1 eV respectively corresponding to Cu 2p<sub>3/2</sub> and Cu 2p<sub>1/2</sub> for both the samples CS5 and CS20. The intensities of Cu 2p<sub>3/2</sub> and the Cu 2p<sub>1/2</sub> peaks increased as the Cu<sub>2</sub>Se wt%

Table 1 Composition, grain size and band gap energy of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se

Sl no.	Sample code	Composition (CdSe : Cu <sub>2</sub> Se)	Thickness (nm)	Grain size of Cu <sub>2</sub> Se doped CdSe nanoparticles (nm)	Bandgap energy (eV)	
					$E_{g1}$ (CdSe)	$E_{g2}$ (Cu <sub>2</sub> Se)
1	CS5	95 : 5	104	33.0 $\pm$ 0.4	—	1.85 $\pm$ 0.01
2	CS10	90 : 10	110	31.0 $\pm$ 0.8	1.77 $\pm$ 0.01	1.99 $\pm$ 0.04
3	CS20	80 : 20	162	29.0 $\pm$ 0.6	1.80 $\pm$ 0.01	2.07 $\pm$ 0.02
4	CS30	70 : 30	223	26.8 $\pm$ 0.8	1.90 $\pm$ 0.03	2.08 $\pm$ 0.03

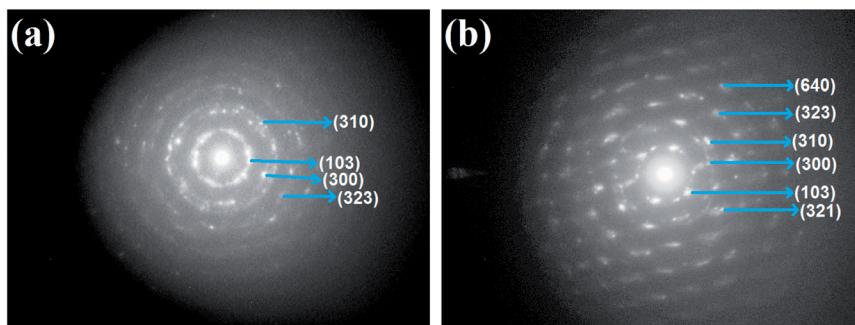


Fig. 3 SAED patterns of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se containing different concentrations of Cu<sub>2</sub>Se: (a) 5 wt%, (b) 20 wt%.

increased from 5% to 20%, but the peak positions did not change. The observed spin-orbit splitting was 19.8 eV which closely agreed with the reported value of 19.6 eV for Cu<sup>1+</sup> oxidation state in Cu<sub>2</sub>Se nanoparticle.<sup>41,42</sup> Satellite peaks at 940 and 945 eV of Cu<sup>2+</sup> oxidation state reported in the literature were absent in the spectra of the samples CS5 and CS20.<sup>43-46</sup> The absence of satellite peaks of Cu<sup>2+</sup> ruled out the possibility for +2 oxidation state of Cu in the present samples.

The AFM images of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se containing different Cu<sub>2</sub>Se wt% of 5, 10, 20 and 30% are shown in Fig. 5. The AFM image of pure CdSe film is given as Fig. 1 in ref. 28. The surface morphology of the films was found to depend strongly on the Cu<sub>2</sub>Se wt%.

The AFM images in Fig. 5 show that particles in each sample were of almost the same size and shape and the films exhibited a tendency to form uniformly sized and shaped particles for all concentration of Cu<sub>2</sub>Se. The particles in the AFM images showed a clear decrease in size with an increase in the Cu<sub>2</sub>Se wt%. The average sizes of the particles seen in the AFM images were measured to be 83, 76, 40 and 37 nm for the samples CS5, CS10, CS20 and CS30 respectively which were larger compared to the grain sizes obtained from XRD analysis (Table 1) using Scherrer equation leading to the inference that

the grains in the nanocrystal thin films of CdSe:Cu<sub>2</sub>Se were aggregates of primary grain of size ~33 to ~26 nm (Table 1).

The optical transmission spectra of the nanocrystal thin films of CdSe:Cu<sub>2</sub>Se containing different Cu<sub>2</sub>Se wt% of 5, 10, 20 and 30 are shown in Fig. S4.<sup>†</sup> The optical transmission spectrum of pure CdSe film is given in Fig. 2(b) of ref. 28. The band gap energies of the film samples were determined from  $\alpha^2$  versus  $h\nu$  plots (Tauc plot) shown in Fig. 6 where  $\alpha$  is the optical absorption coefficient,  $h$  is the Planck's constant and  $\nu$  is the light frequency. The Tauc plots of all the samples except that of CS5 gave two band gap energy values  $E_{g1}$  and  $E_{g2}$ .<sup>47</sup> The band gap values of the samples CS5, CS10, CS20 and CS30 are given in Table 1.  $E_{g1}$  is close to the band gap of CdSe (1.74 eV), and it increases with increase in the wt% of Cu<sub>2</sub>Se in the thin film samples, while the value of  $E_{g2}$  is close to the band gap of Cu<sub>2</sub>Se and it increases with increase in the wt% of Cu<sub>2</sub>Se. The existence of two band gaps for these samples containing higher wt% of Cu<sub>2</sub>Se indicated that these CdSe:Cu<sub>2</sub>Se samples were nanocomposites.

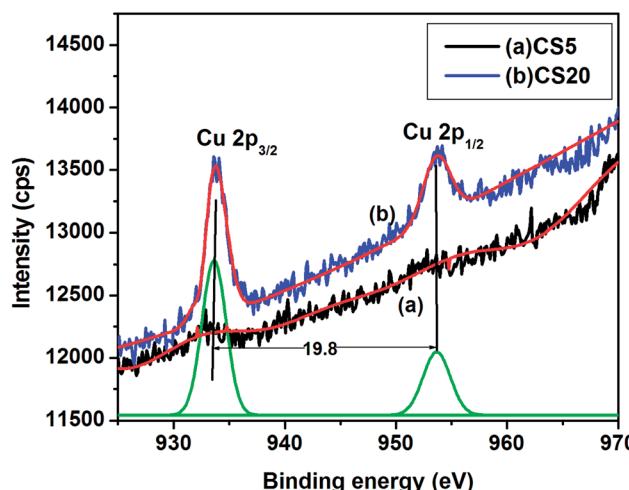


Fig. 4 XPS plots of Cu 2p spectrum of CdSe:Cu<sub>2</sub>Se nanostructured thin films: (a) sample CS5, and (b) sample CS20.

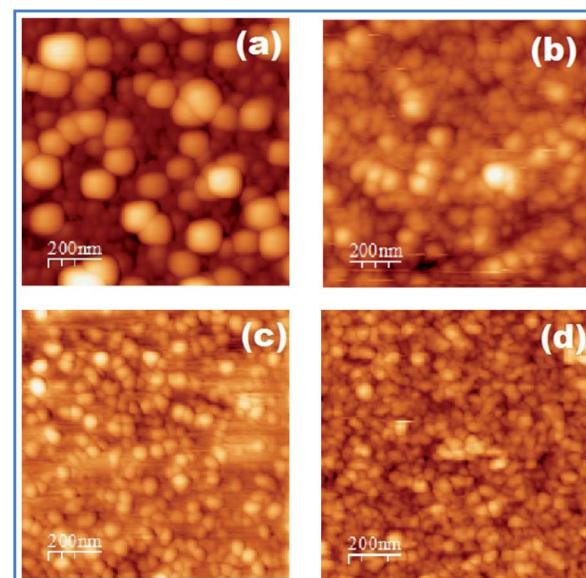


Fig. 5 AFM images of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se containing different concentrations of Cu<sub>2</sub>Se: (a) 5 wt%, (b) 10 wt%, (c) 20 wt% and (d) 30 wt%.



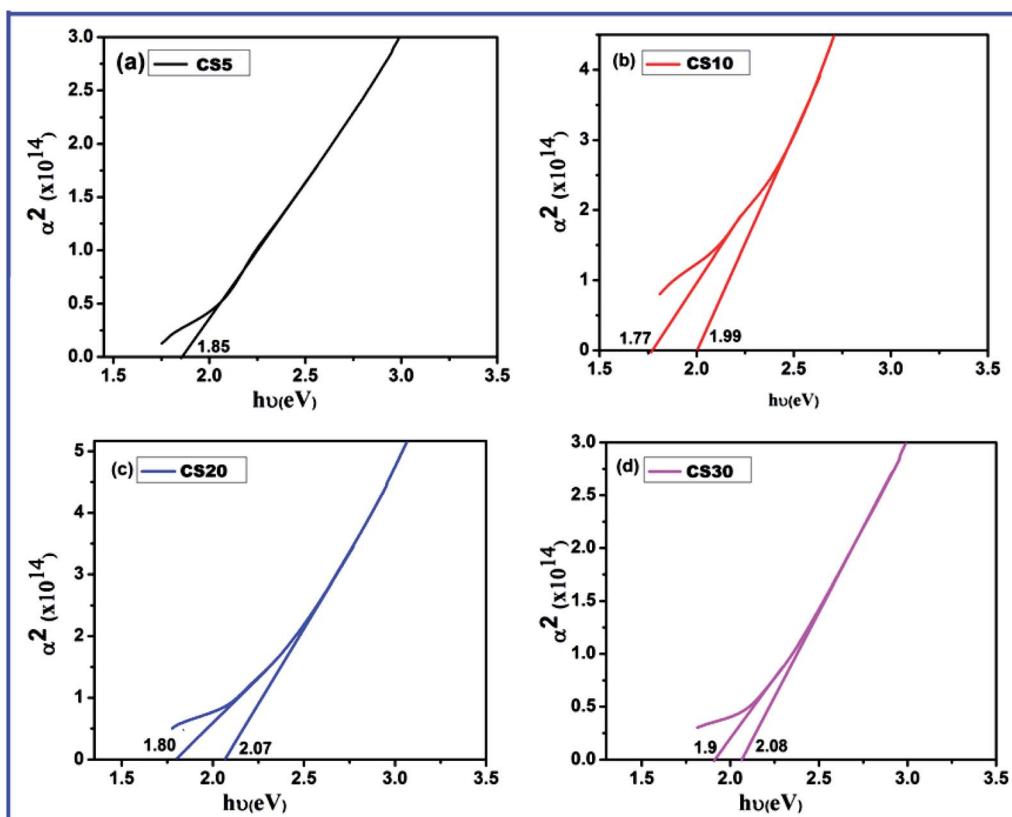


Fig. 6 Plot of  $\alpha^2$  against  $h\nu$  for nanocrystal thin films of CdSe:Cu<sub>2</sub>Se containing different concentrations of Cu<sub>2</sub>Se: (a) 5 wt%, (b) 10 wt%, (c) 20 wt% and (d) 30 wt%.

Micro-Raman spectra of the nanocrystal thin films of pure CdSe, and of CdSe:Cu<sub>2</sub>Se samples containing different wt% of Cu<sub>2</sub>Se are shown in Fig. S5.† The spectrum of the sample CS5 shows a strong and broad peak centered at 206 cm<sup>-1</sup> and a weak peak centered at 412 cm<sup>-1</sup>. The peak at 206 cm<sup>-1</sup> was attributed to the longitudinal optical phonon (1LO), and the weak one at 412 cm<sup>-1</sup> was attributed to the overtone (2LO) mode of CdSe. The very weak peak at 619 cm<sup>-1</sup> was attributable to the 3LO mode. As the wt% of the Cu<sub>2</sub>Se increased, no considerable change was observed in the frequencies of the peaks. But, a weak peak was observed at 260 cm<sup>-1</sup> in the spectrum of sample CS20 containing 20 wt% of Cu<sub>2</sub>Se which was attributed to Cu–Se vibration characteristic of binary copper selenides and which agreed well with the known Cu<sub>2</sub>Se Raman peak.<sup>48,49</sup> With an increase in the wt% of Cu<sub>2</sub>Se from 20% to 30% the intensity of this peak increased very much.

Analysis of the intensities of the peaks in the spectra in Fig. 7 showed a variation in the intensity ratios of the LO overtone to the fundamental. The intensity ratio of the LO overtone to the fundamental can be used as a direct measure for the exciton–LO phonon coupling strength. Fig. 7 shows the variation of the intensity ratio of 2LO to 1LO modes of CdSe in nanocrystal thin films of CdSe:Cu<sub>2</sub>Se as a function of grain size. The intensity ratio  $I_{2\text{LO}}/I_{1\text{LO}}$  for the films containing Cu<sub>2</sub>Se wt% of 5, 10, 20 and 30 were obtained as 0.32, 0.31, 0.3 and 0.29 respectively. For CdSe nanocrystals, Anne Myers Kelley reported values for 2LO/

1LO ratio in the range of 0.2–0.5 implying a weaker electron–phonon coupling in nanoparticles than that in the bulk for which this ratio was reported to be greater than unity.<sup>50</sup> Baranov *et al.*<sup>51</sup> reported a value of  $\approx 0.2$  for electron–phonon coupling strength in colloidal core/shell CdSe/ZnS quantum dots which was in the same order of magnitude as previously reported for

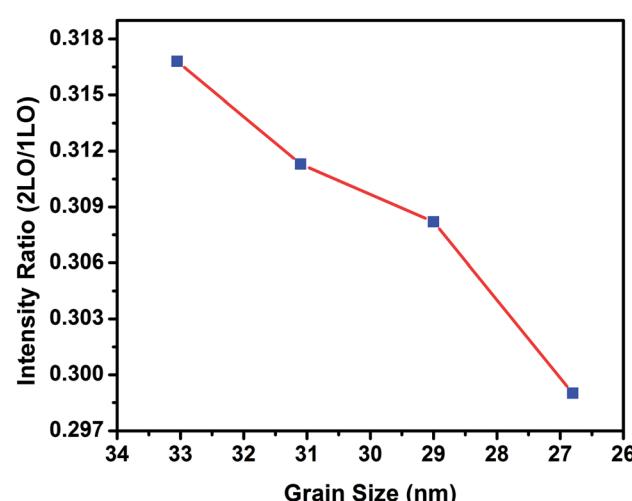


Fig. 7 Variation of the intensity ratio of 2LO to LO Raman lines of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se with grain sizes of CdSe.

bare CdSe quantum dots.<sup>52,53</sup> From Fig. 7 it can be observed that the intensity ratio of 2LO to LO decreased with the decrease in grain size. The electron–phonon coupling in nanoparticles is weakened with decrease in size due to the decrease in the density of states for both the electrons and the phonons, and also due to the increased overlap between the electron and the hole wave functions.<sup>54,55</sup>

Fig. S6† shows the DC electrical conductivity ( $\sigma$ ) of nanocrystal thin films of pure CdSe, and of CdSe:Cu<sub>2</sub>Se containing different wt% of Cu<sub>2</sub>Se. It was observed that the conductivity  $\sigma$  increased with increase in temperature, which confirmed the semiconducting nature of the nanostructured films. At room temperature, conductivity had the values  $2.63 \times 10^{-3}$ ,  $1.12 \times 10^{-2}$ ,  $1.94 \times 10^{-2}$  and  $2.72 \times 10^{-2} \Omega^{-1} \text{ cm}^{-1}$  for the films containing Cu<sub>2</sub>Se wt% of 5, 10, 20 and 30 respectively. The conductivity of the samples was modelled based on the Arrhenius equation:

$$\sigma(T) = \sigma_0 \exp\left(\frac{-E_A}{kT}\right)$$

where  $\sigma_0$  is the pre-exponential factor,  $k$  is the Boltzmann constant and  $T$  is the temperature and  $E_A$  the activation energy for conduction. Arrhenius plots of DC conductivity of the film samples are shown in Fig. 8. The determined values of activation energy ( $E_A$ ) were 0.2849, 0.1486, 0.1117 and 0.0897 eV respectively for the samples containing Cu<sub>2</sub>Se wt% of 5, 10, 20 and 30 respectively showing a decrease with increase in wt% of Cu<sub>2</sub>Se in the films. Since the grain size of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se samples decreased with increase in the wt% of Cu<sub>2</sub>Se (Table 1), it was found that the activation energy decreased with the decrease in the grain size of the nanostructured films. As concluded from the results of XRD analysis, the present samples consisted of nanoparticles of CdSe doped with Cu<sub>2</sub>Se and nanoparticles of Cu<sub>2</sub>Se. The decrease in the value of activation energy in nanocrystal thin films of CdSe:Cu<sub>2</sub>Se could be due to the creation of conduction paths of lower resistance enabling easy charge transfer between CdSe

nanocrystals doped with Cu<sub>2</sub>Se in the samples CS5 and CS10, and between the nanoparticles of CdSe doped with Cu<sub>2</sub>Se and nanoparticles of Cu<sub>2</sub>Se in the composite film samples CS20 and CS30.

Hall measurements could not be done on the CdSe:Cu<sub>2</sub>Se samples containing lower concentrations of Cu<sub>2</sub>Se since they were more resistive. Hall measurements carried at room temperature (308 K) and at 380 K were carried on the CdSe:Cu<sub>2</sub>Se film samples CS20 and CS30 containing Cu<sub>2</sub>Se concentrations of 20 and 30 wt% (Table 2). The positive Hall coefficient confirmed that both the film samples CS20 and CS30 were p-type semiconductors and that holes were the major charge carriers in these samples.<sup>56</sup> The carrier mobility decreased and carrier concentration increased with increase in the Cu<sub>2</sub>Se wt% and with temperature.

XRD and EDS results were utilized to determine the approximate value of  $x$ , where  $x$  is the at% of Cu<sub>2</sub>Se doped into CdSe. XRD patterns of the samples CS5 and CS10 in Fig. 1 did not show peaks of Cu<sub>2</sub>Se while the EDS analysis indicated the presence of respectively 3.87 and 9.24 at% of Cu (Table S1†) in these samples. Hence it was inferred that Cu<sub>2</sub>Se almost completely got substitutionally incorporated into the lattice of CdSe in the samples CS5 and CS10 and only a small percentage of Cu<sub>2</sub>Se below the level detectable through XRD remained in the samples as free Cu<sub>2</sub>Se. XRD patterns of the samples CS20 and CS30 showed the peaks of Cu<sub>2</sub>Se in addition to the peaks of CdSe. The at% of Cu<sub>2</sub>Se substitutionally incorporated into CdSe lattice in the samples CS20 and CS30 should be close to 10. Hence, we made the assumption that respectively 8 and 10 at% of Cu<sub>2</sub>Se got substitutionally incorporated into the CdSe lattice in the samples CS20 and CS30, and the remaining quantity existed as free Cu<sub>2</sub>Se nanograins in the samples. Hence the samples CS20 and CS30 actually consisted of nanograins of CdSe doped with Cu<sub>2</sub>Se (*i.e.*, nanograins of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se where  $x$  is the at% of Cu<sub>2</sub>Se doped into CdSe. Cu<sub>2(x)</sub> means that for any value  $x$ , there be  $x$  Cu ions in the +1 oxidation state) and nanograins of Cu<sub>2</sub>Se. In the nanograins of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se, the concentration of Cu<sub>2</sub>Se was assumed to be more near the grain boundaries of these nanograins due to self purification mechanism.<sup>16,39</sup> Hence CS20 and CS30 were considered as composites of nanograins of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se and Cu<sub>2</sub>Se, a model of which is shown in Fig. 9.

When a semiconductor heterojunction is formed between two dissimilar crystalline semiconductor materials with unequal band gaps, discontinuities between the valence band maxima or conduction band minima are produced at their interface. The discontinuities act as barriers to electrical transport across the interface.<sup>57,58</sup> The semiconductor device performance mainly depends on the valence band offsets ( $\Delta E_V$ ) and the conduction band offsets ( $\Delta E_C$ ). For equilibrium at the junction, the Fermi levels of the two semiconductors must attain same energy by a transfer of an electron from the semiconductor 1 of higher Fermi level to the semiconductor 2 to lower Fermi level. This result in a partial depletion of the electrons near the junction in semiconductor 1 and therefore a bending upward of the band edges (Fig. 10(b)). There is also a corresponding redistribution of charge in semiconductor 2

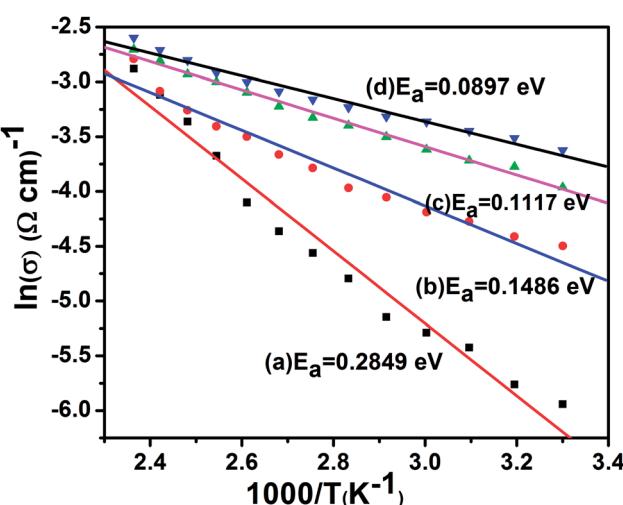
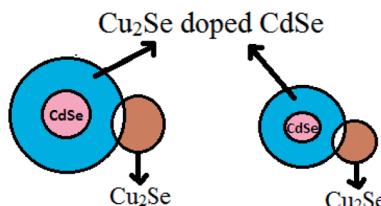


Fig. 8 Arrhenius plot of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se containing different concentrations of Cu<sub>2</sub>Se: (a) 5 wt%, (b) 10 wt%, (c) 20 wt% and (d) 30 wt%.



Table 2 Hall coefficient, carrier concentration and mobility of nanocrystal thin films of CdSe:Cu<sub>2</sub>Se

Sample code	CdSe : Cu <sub>2</sub> Se ratio in the samples	Grain size (nm)	Temperature (K)	Carrier concentration (cm <sup>-3</sup> )	Hall coefficient $R_H$ (m <sup>3</sup> C <sup>-1</sup> )	Mobility ( $\mu$ ) ( $\times 10^{-3}$ ) cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup>
CS20	80 : 20	29	308	$0.13 \times 10^{20}$	$+4.81 \times 10^{-7}$	9.15
			380	$0.08 \times 10^{21}$	$+8.10 \times 10^{-8}$	3.68
CS30	70 : 30	26.8	308	$0.39 \times 10^{20}$	$+1.58 \times 10^{-7}$	4.21
			380	$0.18 \times 10^{21}$	$+3.56 \times 10^{-8}$	1.76

Fig. 9 Model of the Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se:Cu<sub>2</sub>Se nanocomposite.

and a bending downward of the band edges. This is the model developed by Anderson for current flow in p-n (n-p) semiconductor heterojunctions.<sup>59</sup> Energy band diagram built using the Anderson model for semiconductor heterojunction between Cu<sub>2</sub>Se doped CdSe (*i.e.*, Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se) and Cu<sub>2</sub>Se in the present study is shown in Fig. 10(b). The band offsets of heterojunctions in an ideal case where no potential is created at the interface can be calculated using electron affinity model.<sup>58</sup>

The conduction band offset  $\Delta E_C$  for an electron at the heterojunction between Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se and Cu<sub>2</sub>Se is

$$\Delta E_C = \chi_{\text{Cd}_{1-x}\text{Cu}_{2(x)}\text{Se}} - \chi_{\text{Cu}_2\text{Se}} \quad (1)$$

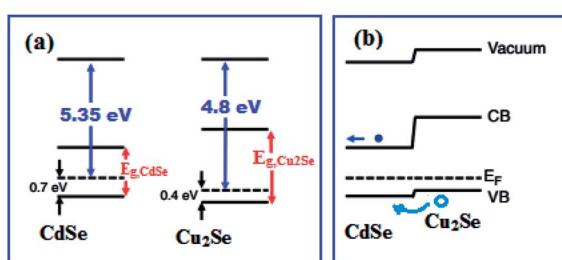
The valence band offset  $\Delta E_V$  for holes is

$$\Delta E_V = (\chi_{\text{Cd}_{1-x}\text{Cu}_{2(x)}\text{Se}} + E_g, \text{Cd}_{1-x}\text{Cu}_{2(x)}\text{Se}) - (\chi_{\text{Cu}_2\text{Se}} + E_g, \text{Cu}_2\text{Se}) \quad (2)$$

where  $\chi_{\text{Cd}_{1-x}\text{Cu}_{2(x)}\text{Se}}$  is the electron affinity and  $E_g, \text{Cd}_{1-x}\text{Cu}_{2(x)}\text{Se}$  is the band gap of Cu<sub>2</sub>Se doped CdSe (*i.e.*, Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se), and  $\chi_{\text{Cu}_2\text{Se}}$  is the electron affinity and  $E_g, \text{Cu}_2\text{Se}$  is the band gap of Cu<sub>2</sub>Se.

The electron affinity of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se determined using the equation,

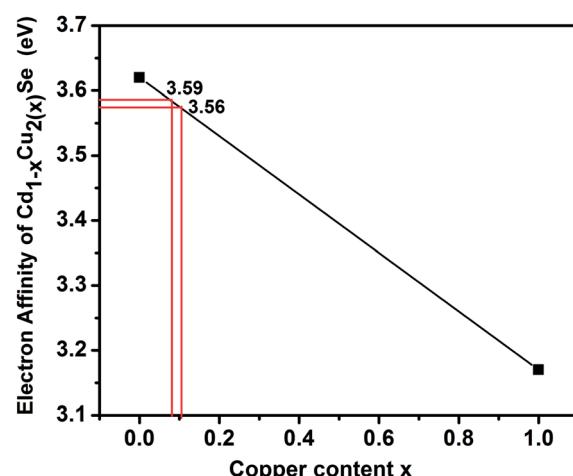
$$\chi_{\text{Cd}_{1-x}\text{Cu}_{2(x)}\text{Se}} = \Phi_{\text{Cd}_{1-x}\text{Cu}_{2(x)}\text{Se}} - (E_g, \text{Cd}_{1-x}\text{Cu}_{2(x)}\text{Se} - E_F, \text{Cd}_{1-x}\text{Cu}_{2(x)}\text{Se}) \quad (3)$$

Fig. 10 (a) Energy levels of the CdSe and Cu<sub>2</sub>Se, and (b) Anderson model energy band diagram of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se:Cu<sub>2</sub>Se heterojunction.

The electron affinity of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se is assumed to be linearly dependant on  $x$  and lie between 3.62 eV for nano CdSe determined using eqn (3) for  $x = 0$  and 3.17 eV for nano Cu<sub>2</sub>Se determined using eqn (3) for  $x = 1$ .

Fig. 11 depicts the electron affinity of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se in the Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se:Cu<sub>2</sub>Se nanocomposite as a function Cu at%  $x$ . The electron affinity  $\chi_{\text{Cd}_{0.8}\text{Cu}_{2(0.2)}\text{Se}}$  determined from the plot at  $x = 0.08$  is 3.59 eV for the sample CS20. The value of  $E_g, \text{Cd}_{0.8}\text{Cu}_{2(0.2)}\text{Se}$  was taken as 1.80 eV (Table 1) and the value of  $E_g, \text{Cu}_{2(0.2)}\text{Se}$  was taken as 2.07 eV (Table 1) for the sample CS20 to determine  $\Delta E_C$  and  $\Delta E_V$ . The electron affinity  $\chi_{\text{Cd}_{0.7}\text{Cu}_{2(0.3)}\text{Se}}$  determined from the plot at  $x = 0.1$  is 3.56 eV for the sample CS30. The value of  $E_g, \text{Cd}_{0.7}\text{Cu}_{2(0.3)}\text{Se}$  was taken as 1.90 eV (Table 1) and the value of  $E_g, \text{Cu}_{2(0.3)}\text{Se}$  was taken as 2.08 eV (Table 1) for the sample CS30 to determine  $\Delta E_C$  and  $\Delta E_V$ . The values of the  $\Delta E_C$  and  $\Delta E_V$  determined using the eqn (1) and (2) are respectively 0.46 eV and 0.19 eV for the sample CS20, and 0.44 eV and 0.26 eV for the sample CS30. The lower value of  $\Delta E_V$  created an easier pathway for hole movement from Cu<sub>2</sub>Se to CdSe. The applied voltage further reduced the valence band offset increasing the conductivity of the films.<sup>60</sup> The hole density is larger for the sample CS30 than that for the sample CS20 leading to an increased hole current for the sample CS30.

Literature reports show that the electrical transport properties in semiconductor heterojunctions are affected by the band offsets at the junction.<sup>57</sup> Mews *et al.*<sup>61</sup> fabricated an amorphous/crystalline silicon heterojunction with n-type crystalline silicon

Fig. 11 Plot of the electron affinity of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se vs. Cu at% ( $x$ ) in Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se:Cu<sub>2</sub>Se nanocomposite thin films.

absorber and studied the hole transport with varying band offsets. Holes generated in the n-type crystalline silicon were directed towards the p-n junction and overcame the valence band offset ( $\Delta E_V$ ) to enter the amorphous –  $\text{SiO}_x\text{H}$  and the amorphous – Si:H layer. They suggested that two possible transport paths existed for holes. First, the holes overcame the band offset barrier by thermionic emission. Second, they tunnelled into the amorphous – Si:H passivation layer and then travelled by tunnel hopping in its valence band tail states. Tunnel hopping was expected to become increasingly important for increased  $\Delta E_V$ .<sup>61</sup> Alivov *et al.*<sup>62</sup> reported the *I*-*V* characteristics of n-ZnO/p-AlGaN heterojunction light-emitting diodes on 6H-SiC substrates. They concluded from the study that the predominant device current was hole injection from the p-type  $\text{Al}_{0.12}\text{Ga}_{0.88}\text{N}$  into the n-type ZnO region of the heterojunction. Dogan *et al.*<sup>63</sup> fabricated nanowire CdSe/Cu<sub>2</sub>Se heterojunctions by masking cation exchange *via* electron-beam irradiation. They performed electrical measurements on single CdSe nanowires before and after cation exchange and found that the partial conversion (CdSe  $\rightarrow$  Cu<sub>2</sub>Se) of the nanowire between the two electrodes can enhance its dark and photoconductance by several orders of magnitude. The low valence band offset is responsible for the high conductivity in CdSe/Cu<sub>2</sub>Se heterojunctions.<sup>63</sup> The explanation of easy hole transport leading to current in the present case of nanostructured CdSe:Cu<sub>2</sub>Se thin films are in agreement with the above interpretation of easy hole transport in the case of amorphous/crystalline silicon heterojunction, n-ZnO/p-AlGaN heterojunction and nanowire CdSe/Cu<sub>2</sub>Se heterojunction.

Current *versus* voltage (*I*-*V*) curves recorded for p-p heterojunctions formed between the CdSe:Cu<sub>2</sub>Se nanostructured films CS5, CS10, CS20 and CS30 deposited on ITO glass with similarly deposited Cu<sub>2</sub>Se film on ITO glass are shown in Fig. S7.† Current-voltage characteristic of pure CdSe film is given as Fig. S5 in ref. 28. The *I*-*V* curves are linear suggesting ohmic behaviour.

Fig. 12 shows *I*-*V* characteristics of the heterojunction formed by nanostructured CdSe:Cu<sub>2</sub>Se films CS5, CS10, CS20 and CS30 deposited on ITO coated glass with a similarly deposited ITO-CdSe film.<sup>64</sup> The *I*-*V* characteristic for CS5 is liner showing ohmic behaviour of the junction, while that for CS10 indicates the behaviour of a p-n junction with a small knee voltage of 0.2 volts. The *I*-*V* characteristics of CS20 and CS30 are similar to that of a standard p-n junction diode. Although conduction starts from the lowest applied voltages, the knee voltage can be considered as  $\sim 0.6$  V and the forward resistance for both these samples is  $\sim 100$   $\Omega$ . This resistance is not high even though the contact between the two films at the junction may not be perfect and can be modulated by the sequential deposition of the two films on a single substrate, and such junctions may find practical applications such as photovoltaic (PV) solar cells and light-emitting diodes (LEDs).<sup>65,66</sup> An ideal absorber for solar cells it should be preferably of p-type with a direct band gap of medium band gap energy with high optical absorption coefficient.<sup>67</sup> The CdSe:Cu<sub>2</sub>Se nanocomposite films of the present study satisfy these conditions and hence may be used as absorber material in solar cells.

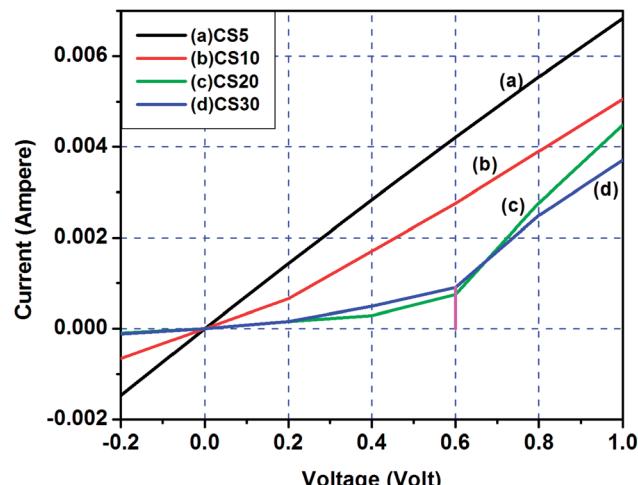


Fig. 12 *I*-*V* characteristics of the heterojunction formed by nanostructured CdSe:Cu<sub>2</sub>Se films CS5, CS10, CS20 and CS30 deposited on ITO coated glass with a similarly deposited ITO-CdSe film.

## Conclusions

Nanocrystal thin films of CdSe:Cu<sub>2</sub>Se containing different wt% of Cu<sub>2</sub>Se were synthesized by a new technique of co-evaporation of CdSe and Cu<sub>2</sub>Se powder in nitrogen gas at a pressure larger than that required for conventional thin film deposition. The films consisted of nanograins of CdSe doped with Cu<sub>2</sub>Se for lower wt% of Cu<sub>2</sub>Se, and nanocomposites of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se and Cu<sub>2</sub>Se for higher wt% of Cu<sub>2</sub>Se. Energy band diagram built using the Anderson model was used for discussing the heterojunction characteristics of the junction between nanograins of Cd<sub>1-x</sub>Cu<sub>2(x)</sub>Se and Cu<sub>2</sub>Se. To investigate the usefulness of the nanocrystal thin films of CdSe:Cu<sub>2</sub>Se for practical applications, the *I*-*V* characteristics of p-p and p-n heterojunctions formed by the films respectively with nanostructured films of similarly deposited Cu<sub>2</sub>Se and CdSe films were studied. The p-n heterojunctions studied may be useful for practical applications such as PV solar cells and LEDs.

## Conflicts of interest

There are no conflicts of interest to declare.

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## References

- 1 M. Shim and P. G. Sionnest, *Nature*, 2000, **407**, 981–983.
- 2 M. Mahima and R. Viswanatha, *RSC Adv.*, 2018, **8**, 22103–22112.



3 A. Sahu, M. S. Kang, A. Kompch, C. Christian Notthoff, A. W. Wills, D. Deng, M. Winterer, C. D. Frisbie and D. J. Norris, *Nano Lett.*, 2012, **12**, 2587–2594.

4 D. Mocatta, G. Cohen, J. Schattner, O. Millo, E. Rabani and U. Banin, *Science*, 2011, **332**, 77–81.

5 A. P. Litvin, I. V. Martynenko, F. P. Milton, A. V. Baranov, A. V. Fedorov and Y. K. Gun'ko, *J. Mater. Chem. A*, 2017, **5**, 13252–13275.

6 C. Ray and T. Pal, *J. Mater. Chem. A*, 2017, **5**, 9465–9487.

7 M. Karttunen, P. Ruuskanen, V. Pitkänen and W. M. Albers, *J. Electron. Mater.*, 2008, **37**, 951–954.

8 P. Zukowski, T. N. Koltunowicz, O. Boiko, V. Bondariev, K. Czarnacka, J. A. Fedotova, A. K. Fedotov and I. A. Svitov, *Vacuum*, 2015, **120**, 37–43.

9 M. Kumar, T. Kumar and D. K. Avasthi, *Scr. Mater.*, 2015, **105**, 46–49.

10 Z. B. Sun, X. Z. Dong, W. Q. Chen, S. Nakanishi, X. M. Duan and S. Kawata, *Adv. Mater.*, 2008, **20**, 914–919.

11 C. Dey, A. R. Molla, M. Goswami, G. P. Kothiyal and B. Karmakar, *J. Opt. Soc. Am. B*, 2014, **31**, 1761–1770.

12 W. Fan, Q. Zhang and Y. Wang, *Phys. Chem. Chem. Phys.*, 2013, **15**, 2632–2649.

13 M. Nanu, J. Schoonman and A. Goossens, *Adv. Funct. Mater.*, 2005, **15**, 95–100.

14 R. J. Beal, J. B. K. Kana and B. G. Potter Jr, *Appl. Phys. Lett.*, 2012, **101**, 031102.

15 B. G. Potter, R. J. Beal and C. G. Allen, *J. Appl. Phys.*, 2012, **111**, 034305.

16 S. Manu and M. A. Khadar, *J. Mater. Chem. C*, 2015, **3**, 1846–1853.

17 L. Maya, W. R. Allen, A. L. Glover and J. C. Mabon, *J. Vac. Sci. Technol., B: Microelectron. Nanometer Struct.–Process., Meas., Phenom.*, 1995, **13**, 361–365.

18 T. Onishi, E. Iwamura, K. Takagi and K. Yoshikawa, *J. Vac. Sci. Technol., A*, 1995, **14**, 2728–2735.

19 W. Zhang, Z. Liu, Z. Liu and J. Zhao, *Appl. Surf. Sci.*, 2012, **258**, 6103–6106.

20 A. Kongkanand, K. Tvrdy, K. Takechi, M. Kuno and P. V. Kamat, *J. Am. Chem. Soc.*, 2008, **130**, 4007–4015.

21 W. C. Ching and Y. C. Fu, *Nanoscale Res. Lett.*, 2013, **8**, 33.

22 S. Liu and X. Su, *Anal. Methods*, 2013, **5**, 4541–4548.

23 R. A. Andrievski and A. M. Gleze, *Scr. Mater.*, 2001, **44**, 1621–1624.

24 B. Cantor and C. M. Allen, *Scr. Mater.*, 2001, **44**, 2055–2059.

25 A. I. Gusev, *Phys.-Usp.*, 1998, **41**, 49–76.

26 P. D. Persans, A. Tu, Y. J. Wu and M. Lewis, *J. Opt. Soc. Am. B*, 1989, **6**, 818–823.

27 A. I. Ekimov, F. Hache, M. C. S. Klein, D. Ricard and C. Flytzanis, *J. Opt. Soc. Am. B*, 1993, **10**, 100–107.

28 N. S. Babu and M. A. Khadar, *Sol. Energy Mater. Sol. Cells*, 2018, **178**, 106–114.

29 W. Zhang, C. Wang, L. Zhang, X. Zhang, X. Liu, K. Tang and Y. Qian, *J. Solid State Chem.*, 2000, **151**, 241–244.

30 G. P. Sorokin, Y. M. Papshev and P. T. Oush, *Sov. Phys. Solid State*, 1966, **7**, 1810–1811.

31 H. Okimura, T. Matsumae and R. Makabe, *Thin Solid Films*, 1980, **71**, 53–59.

32 Al-Mamun, A. B. M. O. Islam and A. H. Bhuiyan, *J. Mater. Sci.: Mater. Electron.*, 2005, **16**, 263–268.

33 A. M. Hermann and L. Fabick, *J. Cryst. Growth*, 1983, **61**, 658–664.

34 N. J. Simi, R. Vinayakan and V. V. Ison, *RSC Adv.*, 2019, **9**, 15092–15098.

35 X. Xu, J. Bullock, L. T. Schelhas, E. Z. Stutz, J. J. Fonseca, M. Hettick, V. L. Pool, K. F. Tai, M. F. Toney, X. Fang, A. Javey, L. H. Wong and J. W. Ager, *Nano Lett.*, 2016, **16**, 1925–1932.

36 S. Bhattacharyya and A. Gedanken, *J. Phys. Chem. C*, 2008, **112**, 4517–4523.

37 A. Anu and M. A. Khadar, *SN Appl. Sci.*, 2019, **1**, 1057.

38 B. D. Cullity and S. R. Stock, *Elements of X-ray Diffraction*, Prentice Hall, New Jersey, 2001.

39 G. M. Dalpian and J. R. Chelikowsky, *Phys. Rev. Lett.*, 2006, **96**, 226802.

40 N. Gleason, J. Guevremont and F. Zaera, *J. Phys. Chem. B*, 2003, **107**, 11133–11141.

41 J. F. Moulder, W. F. Stickle, P. E. Sobol and K. D. Bomben, *Handbook of X-ray Photoelectron Spectroscopy*, Perkin-Elmer Corporation, Physical Electronics Division, Eden Prairie, 1995.

42 S. C. Riha, D. C. Johnson and A. L. Prieto, *J. Am. Chem. Soc.*, 2011, **133**, 1383–1390.

43 C. Ge, L. Liu, X. Yao, C. Tang, F. Gao and L. Dong, *Catal. Sci. Technol.*, 2013, **3**, 1547–1557.

44 L. Wang, J. R. Gaudet, W. Li and D. Weng, *J. Catal.*, 2013, **306**, 68–77.

45 R. Zhang, D. Shi, Y. Zhao, B. Chen, J. Xue, X. Liang and Z. Lei, *Catal. Today*, 2011, **175**, 26–33.

46 L. Liu, Z. Yao, B. Liu and L. Dong, *J. Catal.*, 2010, **275**, 45–60.

47 G. K. Larsen, B. C. Fitzmorris, C. Longo, J. Z. Zhang and Y. Zhao, *J. Mater. Chem.*, 2012, **22**, 14205–14218.

48 V. I. Roca, A. P. Rodríguez, A. R. Rodríguez and J. R. Morante, *J. Appl. Phys.*, 2007, **101**, 103517.

49 B. M. Sukarova, M. Najdoski, I. Grozdanov and C. J. Chunnillall, *J. Mol. Struct.*, 1997, **410–411**, 267–270.

50 A. M. Kelley, *J. Phys. Chem. Lett.*, 2010, **1**, 1296–1300.

51 A. V. Baranov, Y. P. Rakovich, J. F. Donegan, T. S. Perova, R. A. Moore, D. V. Talapin, A. L. Rogach, Y. Masumoto and I. Nabiev, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2003, **68**, 165306.

52 M. C. Klein, F. Hache, D. Ricard and C. Flytzanis, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1990, **42**, 11123–11132.

53 C. Lin, K. Gong, D. F. Kelley and A. M. Kelley, *ACS Nano*, 2015, **9**, 8131–8141.

54 M. Nirmal, C. B. Murray, D. J. Norris and M. G. Bawendi, *ZPhys-e.D: At., Mol. Clusters*, 1993, **26**, 361–363.

55 J. Z. Zhang, *Acc. Chem. Res.*, 1997, **30**, 423–429.

56 S. C. Riggs, R. D. McDonald, J. B. Kemper, Z. Stegen, G. S. Boebinger, F. F. Balakirev, Y. Kohama, A. Migliori, H. Chen, R. H. Liu and X. H. Chen, *J. Phys.: Condens. Matter*, 2009, **21**, 412201.

57 N. N. Halder, P. Biswas, B. Nagabhushan, S. Kundu, D. Biswas and P. Banerji, *J. Appl. Phys.*, 2014, **115**, 203719.



58 B. Roul, M. Kumar, M. K. Rajpalke, T. N. Bhat and S. B. Krupanidhi, *J. Phys. D: Appl. Phys.*, 2015, **48**, 423001.

59 A. G. Milnes and D. L. Feucht, *Heterojunctions and Metal Semiconductor Junctions*, Academic Press, New York, 1972.

60 R. Dalven, *Introduction to Applied Solid State Physics*, Springer Publication, Boston, 1990.

61 M. Mews, M. Liebhaber, B. Rech and L. Korte, *Appl. Phys. Lett.*, 2015, **107**, 013902.

62 Y. I. Alivov, E. V. Kalinina, A. E. Cherenkov, D. C. Look, B. M. Ataev, A. K. Omaev, M. V. Chukichev and D. M. Bagnall, *Appl. Phys. Lett.*, 2003, **83**, 4719–4721.

63 S. Dogan, S. Kudera, Z. Dang, F. Palazon, U. Petralanda, S. Artyukhin, L. D. Trizio, L. Manna and R. Krahne, *Nat. Commun.*, 2018, **9**, 505.

64 M. Zhou, H. Zhu, X. Wang, Y. Xu, Y. Tao, S. Hark, X. Xiao and Q. Li, *Chem. Mater.*, 2010, **22**, 64–69.

65 S. Menezes and A. Samantilleke, *Sci. Rep.*, 2018, **8**, 11350.

66 R. Graft, T. Fischer, A. Gray and S. Kennerly, *J. Appl. Phys.*, 1993, **74**, 5705–5712.

67 S. R. Kodigala, *Thin Films and Nanostructures*, Elsevier, UK, 2010.

