


 Cite this: *RSC Adv.*, 2020, 10, 5328

# The influence of molecular vicinity (expressed in terms of dielectric constant) on the infrared spectra of embedded species in ices and solid matrices

Pilling S. \* and Bonfim V. S.

In this theoretical work we evaluate how the chemical environment influences some features presented in the infrared spectrum, such as band intensities and band location of embedded species in icy matrices. The calculations were performed employing the Polarized Continuum Model (PCM) approach with the second-order Møller–Plesset perturbation theory (MP2) level using the Gaussian 09 package. Here, we simulate the effects of molecular vicinity around embedded species in terms of the effects of the dielectric constant ( $\epsilon$ ) of the icy and solid samples. Gas phase calculation was also performed for comparison purpose. The investigated embedded single molecules were CO, CO<sub>2</sub>, CH<sub>4</sub>, NH<sub>3</sub>, SO<sub>2</sub>, HCOOH, CH<sub>3</sub>OH and also H<sub>2</sub>O. The results suggest that for most vibrational modes, the strengths of IR bands show an increase with  $\epsilon$ , which implies they also decrease with respect to porosity. The frequency shifts showed opposite behavior in relation to the band strengths, with few exceptions. A correlation between calculated band intensities with the band strengths  $A$  (taken from literature) was determined and described by a linear function  $I \sim 6 \times 10^{18} A$  [km mol<sup>-1</sup>], with  $A$  in unity of cm per molecule. In addition, an associative exponential function was adjusted to the studied dataset to characterize the evolution of frequency-shift and intensity-shift and band strength ratio as function of the dielectric constant. Since astrophysical ice mantles over cold dust grains can vastly vary in composition in space (having different dielectric constants) they are a challenge to be well characterized. Therefore, this work can help the astrochemistry community to better understand astrophysical ices and its observations in the infrared.

Received 3rd December 2019

Accepted 24th January 2020

DOI: 10.1039/c9ra10136e

[rsc.li/rsc-advances](http://rsc.li/rsc-advances)

## 1 Introduction

The molecular abundances in solid phase observed in cold astrophysical environments (*e.g.* dense molecular clouds, inner part of protoplanetary disks, obscured regions of young stellar objects (YSOs), comets, and frozen surfaces of solar system) are mainly determined by measuring the area of specific peaks in emission or absorption (the latter is more frequent) in the infrared (IR) spectra which are related to a given molecular vibrational modes.<sup>1–3</sup> However, the procedure requires an accurate value of a molecular (physicochemical) parameter called band strength  $A$  (sometimes also called the  $A$ -value in the literature) ordinarily measured in the laboratory during ice simulation at low temperatures.<sup>4–7</sup> Additionally, it has been reported in the literature that the molecular band strength and peak position in the IR spectrum may vary depending on the sample temperature and chemical environment.<sup>7–13</sup>

Considering that in cold space environments, the ices may contain different chemical compositions in a water-ice matrix (embedded molecules lie in different chemical environments),

different temperatures and porosity degree, the employment of laboratory values of band strength (and also peak locations) derived for pure samples, or for simple mixtures, can introduce errors in the determination of real molecular abundances and also molecular identification in astronomical observations. In attempt to better handle and quantify molecular abundances of frozen species in the IR spectra of cold astrophysical environments, we performed a theoretical/computational work to evaluate, through vibrational analysis from *ab initio* calculations, how the chemical environment influences features in the spectrum, specifically, the peak positions and their respective intensities (here described in terms of band strengths).

Here, we simulate the effects of chemical environment in single molecular species embedded in an icy matrix by considering the net effect of a surrounding average electric field, parameterized by the dielectric constant ( $\epsilon$ ) of the medium. This level of description was achieved employing Polarized Continuum Model (PCM),<sup>14</sup> which is implemented in several computational quantum chemistry packages.<sup>15</sup> This approach has been used to model the average icy matrix effect on the embedded molecule reactivity.<sup>16–20</sup> The studied embedded species in this manuscript were CO, CO<sub>2</sub>, CH<sub>4</sub>, NH<sub>3</sub>, SO<sub>2</sub>, HCOOH (“*trans*” conformer), CH<sub>3</sub>OH and H<sub>2</sub>O.

Universidade do Vale do Paraíba – UNIVAP, Laboratório de Astroquímica e Astrobiologia – LASA, Av. Shishima Hifumi, 2911, Urbanova, São José dos Campos, SP, Brazil. E-mail: [sergiopilling@yahoo.com.br](mailto:sergiopilling@yahoo.com.br)



Section 2 presents, with details, the employed computational methodology with the description of the approximations and considerations used in the work. The results are listed in Section 3 with focus on the changes in the infrared features of the embedded selected species as function of dielectric constant of icy matrix. Conclusion and final remarks are given in Section 4.

## 2 Methodology

In this work we investigate through theoretical/computational analysis how cold chemical environments (astrophysical ice analogs) influence changes (*e.g.* peak positions and intensities) in the IR spectrum of single embedded species. Here we investigate eight different molecules (polar and non-polar ones): CO, CO<sub>2</sub>, CH<sub>4</sub>, NH<sub>3</sub>, SO<sub>2</sub>, HCOOH, CH<sub>3</sub>OH and also H<sub>2</sub>O. Such species are observed in several space environments.<sup>21,22</sup>

The PCM approach is here implemented *via* its “integral equation formalism” version, and employing the Gaussian 09 program.<sup>15,23,24</sup> The PCM solvation method is a good approximation to simulate the implicit solvation (sometimes termed continuum solvation). As it is a method to represent solvent as a continuum medium instead of individual “explicit” solvent molecules, it has been employed for liquids representation in molecular dynamics as well as good for simulating amorphous ices.<sup>16,18–20</sup> In the case of crystalline ices (not studied in this work) the explicit solvation would be more appropriated due to the system symmetry around the embedded molecule (unit cell) which may play a major influence on the molecular vibrations. The cavity model used here was the default available in the Gaussian 09 program.<sup>24</sup>

All calculations have been performed at second-order Møller–Plesset perturbation theory (MP2) level, with the correlation consistent triple zeta basis set, or just cc-pVTZ.<sup>25</sup> The number of electrons of the studied system varies quite well among the molecules investigated here, and also a larger set of molecules on which calculations are in progress. Thus, for reasons of computational convenience, we have used the MP2/cc-pVTZ calculation level for all the electronic structure calculations of the current work. Every chemical species had the equilibrium geometry confirmed for each  $\epsilon$  value (varying from 1 to 180) by applying harmonic vibrational analysis, which has also provided the desired output data. Here,  $\epsilon = 1$  explicitly indicates molecules in gas-phase (vacuum), low values of  $\epsilon$  also indicate molecules embedded in “very cold” ices (porous or not), and  $\epsilon = 180$  indicates molecules embedded in “hot” water-rich ices (see Aragonés *et al.*<sup>27</sup>), as we will discuss further. The calculated molecular infrared spectra of the studied species in the current work, including the assignment of molecular vibration (vibration modes), band intensities as well as the band locations (frequency or wavenumber) were performed using the Gaussian 09 package.<sup>24</sup>

In the adopted effective-medium approach context, we can use the link between dielectric constant  $\epsilon$  and other physical-chemical properties of the bulk ice to probe how they influence the IR spectra of the embedded molecules. For practical means,

it may be assumed that variations in such properties, such as ice temperature, are dominated by changes in the solid water matrix, and consequently in the corresponding  $\epsilon$  values. Fig. 1 presents an illustration of the PCM approximation employed (individual molecules are replaced by an average field represented by the dielectric constant). As discussed previously in this approach the molecular environment (force field) is represented by its dielectric constant  $\epsilon$ . The amount of polarization which occurs when a certain voltage (or electric potential) is applied to an object influences the amount of electrical energy that is stored in the electric field. This is described by the dielectric constant of the material. As discussed by Silla *et al.*<sup>26</sup> and Aragonés *et al.*,<sup>27</sup> and references therein in both works, the dielectric constant (or relative permittivity) of materials is also very dependent on chemical composition (diversity of polarity in molecules), sample temperature and degree of compaction (or its porosity). This figure also illustrates the porous sample approximation employed in this work (solid matrix filled with holes containing vacuum) in which the value of dielectric constant is lower as the amount of pores increases.

It is important to note that, a given dielectric constant might be associated with several types of ices/solid matrices (with different chemical composition, porosity degree and temperature). For this reason, we prefer to make all the calculations here in terms of the dielectric constant to make the manuscript as general as possible for the astrochemistry community. In future work we consider the explicit dependence of the dielectric constant with sample temperature and porosity and how this can influence spectral features of embedded molecules, following the directions already pointed out in the literature.<sup>17,28</sup>

Considering as an approximation that band area  $S$  is proportional to band intensity ( $S \propto I$ ), particularly good for symmetric band profiles, the relationship between the band intensity  $I$  and band strength ( $A$ -value) can be seen from the equation:

$$A = \frac{2.3}{N} S, \rightarrow A \propto \frac{2.3}{N} I \text{ [cm per molecule]} \quad (1)$$

where  $N$  is the column density (related with molecular abundance) of a given species, and  $S$  is the band area (in units of cm<sup>-1</sup>) of a specific IR vibrational mode.<sup>3,29,30</sup> The band strength is a molecular property for each vibrational mode (in the case of springs it would be equivalent to the spring displacement “ $x$ ”, as stated in Hooke’s law) and it is related with  $I$ , which in turn is the change in the system dipole moment generated by the displacement of atoms during vibration.

Assuming the change in the band area ( $S$ ) is very little affected by the enlargement or narrowing of the peak, eqn (1) states that the variations in the ratio between two different  $A$  values (*e.g.*  $A_1$  and  $A_2$ ) will be equivalent to the ratio between two different intensities  $I$  values (*e.g.*  $I_1$  and  $I_2$ ) when we compare systems in three different situations: (1) two different samples (1 and 2) containing the same amount of molecules or column density  $N$ ; (2) two different IR bands ( $i$  and  $j$ ) of the same molecular species; (3) the same species in two different environments (embedded into different media with different dielectric constants related with different temperatures,



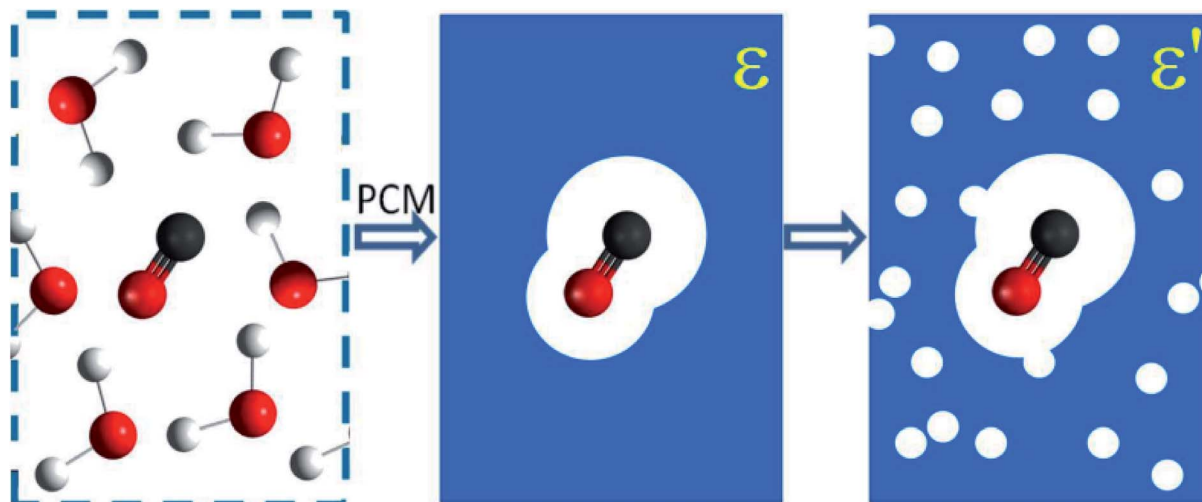


Fig. 1 Illustration of the PCM approximation and the porous sample approximation (solid filled with holes containing vacuum). The molecular environment is represented by its dielectric constant  $\epsilon$  and in this figure porous sample has a dielectric constant  $\epsilon' < \epsilon$ . See details in the text.

porosity or chemical composition). This methodology is very helpful since it avoids any systematic errors in the determination of individual values of band intensities due to the use of a given level of theory. Therefore, the prediction value for the ratio of two bands strengths (comparing two vibration modes of a single molecule or single vibration mode for same molecule in different environments) gets us as close as possible from connections with observational and/or experimental data.

It is important to note that, in an attempt not to be strongly affected by the level of calculation/theory, the current methodology focus on the relative values rather the absolute values. Here we always make a comparison (subtraction or a division) of calculated quantities with some reference calculated value and identify the main changes/differences from such comparison.

As it will be described in the following sections, in this manuscript we will consider the third situation listed above to describe, for each studied species, the dependence of IR band intensities (written in terms of  $I/I_0$ ) and IR bands strengths (written in terms of  $A/A_0$ ) as a function of different dielectric constants ( $\epsilon$ ).

### 3 Results and discussion

The evaluated IR spectrum features (peak location and intensities) have presented considerable variations with the systematic changes of  $\epsilon$ . As pointed out previously, a variation in the ratio of band intensities is proportional to the change on the ratio of band strengths for the same species in two different environments (for example, considering different dielectric constants). In this section we will present such variations for the studied embedded species (CO, CO<sub>2</sub>, CH<sub>4</sub>, NH<sub>3</sub>, SO<sub>2</sub>, HCOOH, CH<sub>3</sub>OH and also H<sub>2</sub>O).

As pointed out by Mate<sup>30</sup> the molecular band strengths are intrinsic properties of the material, and as the dielectric constants do, vary significantly with the temperature due to morphology and phase changes. Additionally, the band

strength values are also different in ice mixtures as compared with pure species, due to the interaction between the different kinds of molecules within the ice. The intensity variations are usually accompanied with frequency shifts of the IR absorptions. Both modifications must be taken into account for a proper description of the composition of ice samples. The chemical environments expressed in terms of the different type of neighbor molecules around a given embedded species and its influence on IR profiles were also discussed by Schmitt *et al.*<sup>31</sup> Similarly, the changes in the band strengths with chemical composition were extensively investigated experimentally in the literature.<sup>32–35</sup> As the surrounded species changes, the electric field and dielectric constant of the medium change, and this affects the molecular orbital and atom configuration of the embedded species (as well as electric dipole moments) which interact differently with the incoming electromagnetic field. These changes result, as a consequence, changes in the molecular band strength for a given energy transition (or vibration mode). For example, Bouwman *et al.*<sup>32</sup> show that most bands of water (stretch, bend and librations modes) decreases as the water species is surrounded by CO molecules. A similar behavior for the water bands was also observed by Öberg *et al.*<sup>33</sup> as the embedded water molecules are being surrounded by CO<sub>2</sub> molecules.

Table 1 lists the selected IR bands analyzed in this manuscript. The theoretical peak positions (in units of wavenumbers and micrometers) calculated for ice samples at 10 K (cold ice with  $\epsilon = 3$ ) and at gaseous-phase ( $\epsilon = 1$ ) are given. For comparison purpose, we also present some experimental values of band location and band strengths, taken from literature for pure ices at low temperature.<sup>6,8,36–39</sup>

As a summary of the results, for most vibration modes employed in this study, the band strengths have increased with  $\epsilon$ , which also implies they have decreased with respect to porosity. The frequency shifts (in wavenumbers) have shown an opposite behavior in relation to the band strengths, with few



Table 1 Molecules and IR bands studied in this work

Selected vibrational modes			MP2(fu)/cc-pVTZ (this work)				Experimental data (literature)	
			Cold ice (10 K; $\epsilon = 3$ )		Gas-phase (vacuum; $\epsilon = 1$ )		Cold pure ices ( $\epsilon = \text{unknown}$ )	
Molecule	Mode assignment	$\nu_{\#}$	$\nu$ ( $\text{cm}^{-1}$ )	$\lambda$ ( $\mu\text{m}$ )	$\nu$ ( $\text{cm}^{-1}$ )	$\lambda$ ( $\mu\text{m}$ )	$\nu$ ( $\text{cm}^{-1}$ )	$A$ (molecule per cm)
CO	Stretching	$\nu_1$	2145	4.661	2139	4.675	2139	$1.1 \times 10^{-17a}$
CO <sub>2</sub>	Bending	$\nu_2$	659	15.19	661	15.13	665	$1.2 \times 10^{-17a}$
	Asymmetric stretching	$\nu_3$	2427	4.120	2440	4.098	2342	$7.6 \times 10^{-17a}$
CH <sub>4</sub>	Asymmetric bending	$\nu_4$	1352	7.394	1358	7.364	1302	$8.0 \times 10^{-18a}$
	Asymmetric stretching	$\nu_3$	3214	3.112	3216	3.109	3010	$1.1 \times 10^{-17a}$
SO <sub>2</sub>	Bending	$\nu_2$	503	19.85	504	19.84	530	$4.8 \times 10^{-18b}$
	Symmetric stretching	$\nu_1$	1125	8.903	1120	8.929	1149	$2.2 \times 10^{-18c}$
	Asymmetric stretching	$\nu_3$	1333	7.495	1337	7.479	1335	$1.5 \times 10^{-17c}$
H <sub>2</sub> O	Bending	$\nu_2$	1645	6.080	1650	6.061	1659	$9.0 \times 10^{-18a}$
	Symmetric stretching	$\nu_1$	3864	2.588	3872	2.583	3298	$1.7 \times 10^{-16d}$
NH <sub>3</sub>	Asymmetric stretching	$\nu_3$	3977	2.515	3993	2.504	3257	$1.5 \times 10^{-16a}$
	Symmetric bending	$\nu_2$	1098	9.111	1077	9.285	1069	$1.6 \times 10^{-17a}$
	Asymmetric bending	$\nu_4$	1683	5.943	1691	5.914	1624	$4.7 \times 10^{-18e}$
HCOOH	Symmetric stretching	$\nu_1$	3533	2.831	3542	2.823	3375	$2.2 \times 10^{-17f}$
	Asymmetric stretching	$\nu_3$	3666	2.728	3680	2.717	3376	$2.3 \times 10^{-17a}$
	Out of plane O–H bond bending	$\nu_5$	685	14.67	692	14.45	929	$6.4 \times 10^{-17a}$
	Out of plane C–H bond bending	$\nu_8$	1141	8.774	1143	8.749	1064	$3.1 \times 10^{-19a}$
	Carbonyl group stretching	$\nu_3$	1819	5.514	1830	5.46	1708	$5.4 \times 10^{-17a}$
CH <sub>3</sub> OH	O–H bond stretching	$\nu_1$	3772	2.654	3781	2.64	2586	$1.4 \times 10^{-16a}$
	Methyl group torsion	$\nu_{12}$	306	32.70	309	32.36	694	$1.4 \times 10^{-17a}$
	C–O bond stretching	$\nu_8$	1067	9.368	1073	9.320	1031	$1.1 \times 10^{-17a}$
	C–H bond sym. stretching	$\nu_3$	3076	3.251	3070	3.257	2700–3600	$1.6 \times 10^{-16a}$
	C–H bond asym. stretching	$\nu_9$	3132	3.192	3124	3.20	2700–3600	$1.4 \times 10^{-16a}$

<sup>a</sup> At 25 K; Bouilloud *et al.*<sup>36</sup> <sup>b</sup> At 90 K; Khanna *et al.*<sup>37</sup> <sup>c</sup> At 16 K; Garozzo *et al.*<sup>38</sup> <sup>d</sup> At 10 K; Hudgins *et al.*<sup>8</sup> <sup>e</sup> At 10 K (presumably); Sandford and Allamandola.<sup>39</sup> <sup>f</sup> At 10 K; D'Hendecourt and Allamandola.<sup>6</sup>

exceptions. Moreover, all the trends here observed seem to be more pronounced in low temperature ices than in the higher temperature ones. Hopefully, this study will provide a better understanding of some astrophysical ices properties and their interaction with radiation, since from the IR spectral profile of diluted species, one could be able to infer physical and chemical properties of the surrounding medium. In the following section we go deeper into these findings for each studied species.

### 3.1 Influence of dielectric constant in infrared features (band location and intensity)

Fig. 2 presents the changes in the IR spectral features of the vibration modes of frozen CO, CO<sub>2</sub> and CH<sub>4</sub> molecules embedded in a solid matrix as a function of its dielectric constant  $\epsilon$  displayed in three different views: (a) frequencies shifts, (b) intensities shifts and, (c) band strength ratio. The calculated values for  $\nu_0$ ,  $I_0$ , and  $A_0$  were performed considering  $\epsilon_0 = 1$  (gas phase), as the reference. The results for these species clearly indicate that IR features seem to be much more dependent on the environment (expressed in terms of  $\epsilon$ ) for  $\epsilon < 20$  values. For dielectric constant higher than 50 it seems that band strength ratio  $A/A_0$ , as well as band location and band intensity lose their dependencies with dielectric constant.

The wavenumber of the CO vibrational mode showed an increase – blueshift in frequency – of up to  $12 \text{ cm}^{-1}$  ( $\sim 1\%$ ) over the studied interval. The variation calculated for the intensity, in turn, was numerically similar, around  $14 \text{ km mol}^{-1}$ , but much larger in relative terms, increasing more than 40% over  $I_0$  (or  $A_0$ ). Since CO has no other vibrational modes, those results shall be compared with CO<sub>2</sub> and CH<sub>4</sub> ones, the two apolar (non-polar) molecules also studied in this work. Comparing the frequency variations for the three molecules we see that only for CO it occurs in the same way as the matrix  $\epsilon$  changes. Despite this differentiation, there is an almost quantitative agreement with the laboratory data from Karssemeijer *et al.*<sup>40</sup> These authors reported a higher shift by  $12 \text{ cm}^{-1}$ , comparing the stretching of CO molecules trapped in H<sub>2</sub>O ice with that observed in a lower  $\epsilon$  (pure CO) chemical environment. On the other hand, there are also recent studies in which the authors argue there is an increase in the wavenumber also of CO embedded in a matrix formed from nonpolar molecules.<sup>41</sup> The results for CO<sub>2</sub> molecule are, on average, of the same range of frequency variation as those seen for CO, which does not hold for intensities, because mode  $\nu_3$  for CO<sub>2</sub> is much more intense. In addition, the wavelength variation profile thus agrees with that reported by Pilling *et al.*,<sup>3</sup> assuming that  $\epsilon$  has increased throughout his pure CO<sub>2</sub> ice irradiation experiment. The other nonpolar molecule, CH<sub>4</sub>, was studied by Boogert *et al.*,<sup>42</sup> in



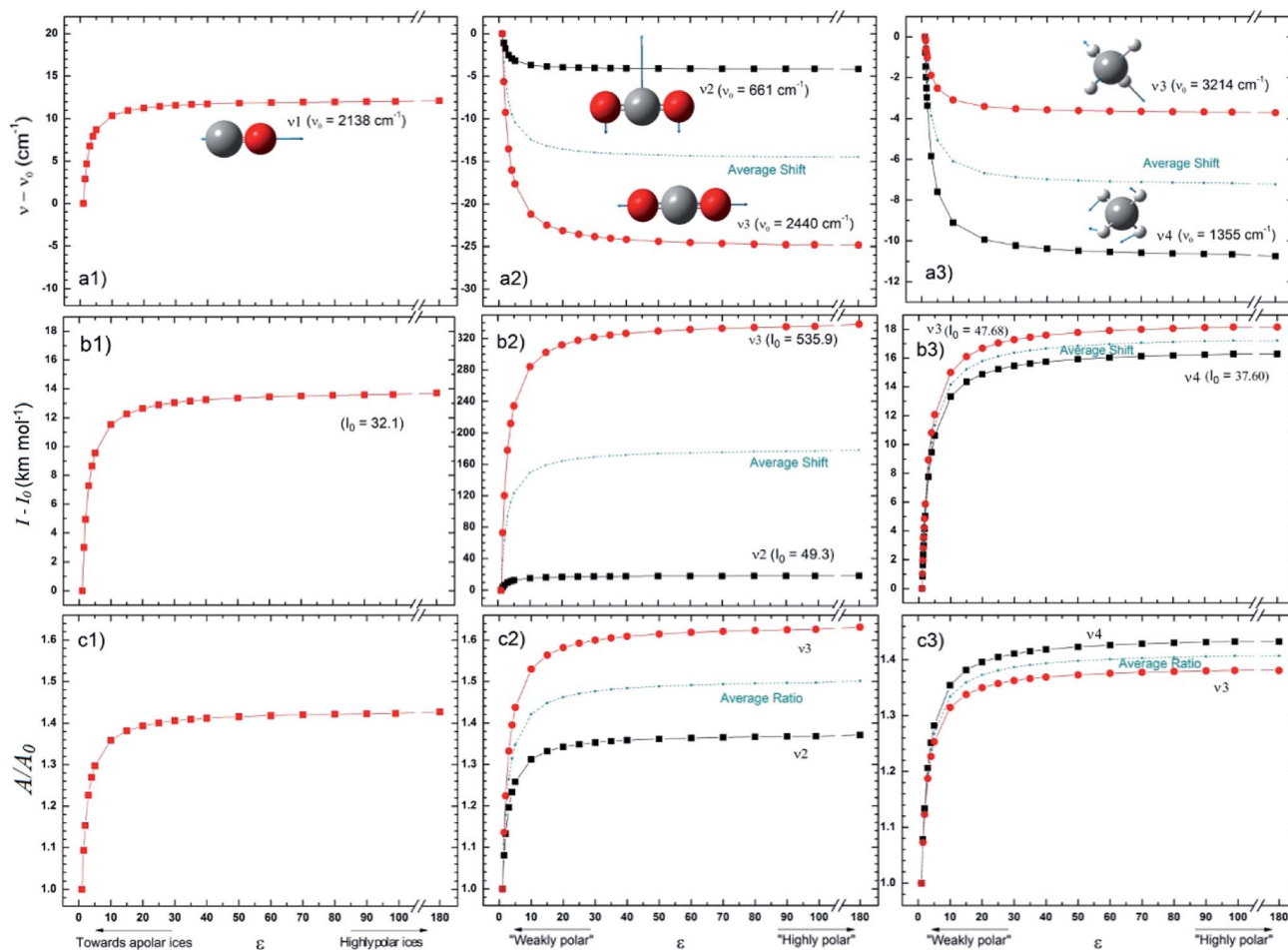


Fig. 2 Changes in the IR spectral features of the vibrational mode of frozen CO, CO<sub>2</sub> and CH<sub>4</sub> molecules embedded in solid matrix as a function of its dielectric constant  $\epsilon$  displayed in three different views: (a) frequencies shifts; (b) intensities shifts and (c) band strength ratio. The calculated values for  $\nu_0$ ,  $I_0$ , and  $A_0$  were performed considering  $\epsilon_0 = 1$  (gas phase), as the reference. See details in the text.

which the authors compared the characteristics of the  $\nu_4$  mode for some different chemical environments. The same authors found no direct correlation between band position and dipole moment (or polarity) of the predominant molecule in the solid matrix. Nevertheless, they found that the wavenumber of  $\nu_4$  is lower in solid methane ( $\epsilon = 1.766$ ) than in gas, which agrees with the trend observed from this work calculations.

The changes in the IR spectral features of the vibrational mode of frozen H<sub>2</sub>O, SO<sub>2</sub>, and NH<sub>3</sub> molecules embedded in solid matrix as a function of its dielectric constant is displayed in Fig. 3 at three different views: (a) frequencies shifts, (b) intensities shifts and, (c) band strength ratio. The calculated values for  $\nu_0$ ,  $I_0$ , and  $A_0$  were performed considering  $\epsilon_0 = 1$  (gas phase), as the reference. In the case of water species ice we observed that all 3 studied vibrational modes decreases as the value of  $\epsilon$  increases (more dramatic for  $\nu_3$  vibrational mode). Curiously, some IR bands of SO<sub>2</sub> and NH<sub>3</sub> have different behavior as the dielectric constant increases, some bands presenting a blueshift and other ones a redshift. Additionally, some ammonia bands seem to be highly affected by the changes in  $\epsilon$  (mainly between  $5 < \epsilon < 10$ ), for example, the  $\nu_1$  mode had a shift

of about  $35 \text{ cm}^{-1}$  and the  $\nu_3$  mode an opposite shift of roughly  $-22 \text{ cm}^{-1}$  when compared with nominal values at gas phase. For  $\epsilon > 50$  it seems that the band center location is no more affected by the dielectric constant.

For the species shown in Fig. 3 the band intensities have presented and increase as function of dielectric constant (with the exception of  $\nu_1$  of NH<sub>3</sub>). Curiously, some bands are much more sensitive to the dielectric constant (mainly between  $5 < \epsilon < 10$ ) than others as the case of  $\nu_3$  of SO<sub>2</sub> (with an outstanding increase of roughly  $120 \text{ km mol}^{-1}$ ) and  $\nu_2$  of NH<sub>3</sub> (increase of roughly  $55 \text{ km mol}^{-1}$ ). These two bands, have also present the larger values of intensities calculated in the gaseous-phase and despite obvious explanation this might have some influence. The bands strengths ratio  $A/A_0$  (where  $A_0$  indicates the value calculated at gas-phase) presented in the panels c, indicate an clear increase with the dielectric constant (mainly below  $\epsilon < 20$ ) for the 3 frozen species H<sub>2</sub>O, SO<sub>2</sub>, and NH<sub>3</sub> (with the exception of the  $\nu_1$  NH<sub>3</sub> vibration mode). As observed in Fig. 2, for  $\epsilon > 50$  (for highly polar ice matrix) it seems that the bands strengths ratio  $A/A_0$  lose its dependency with dielectric constant.



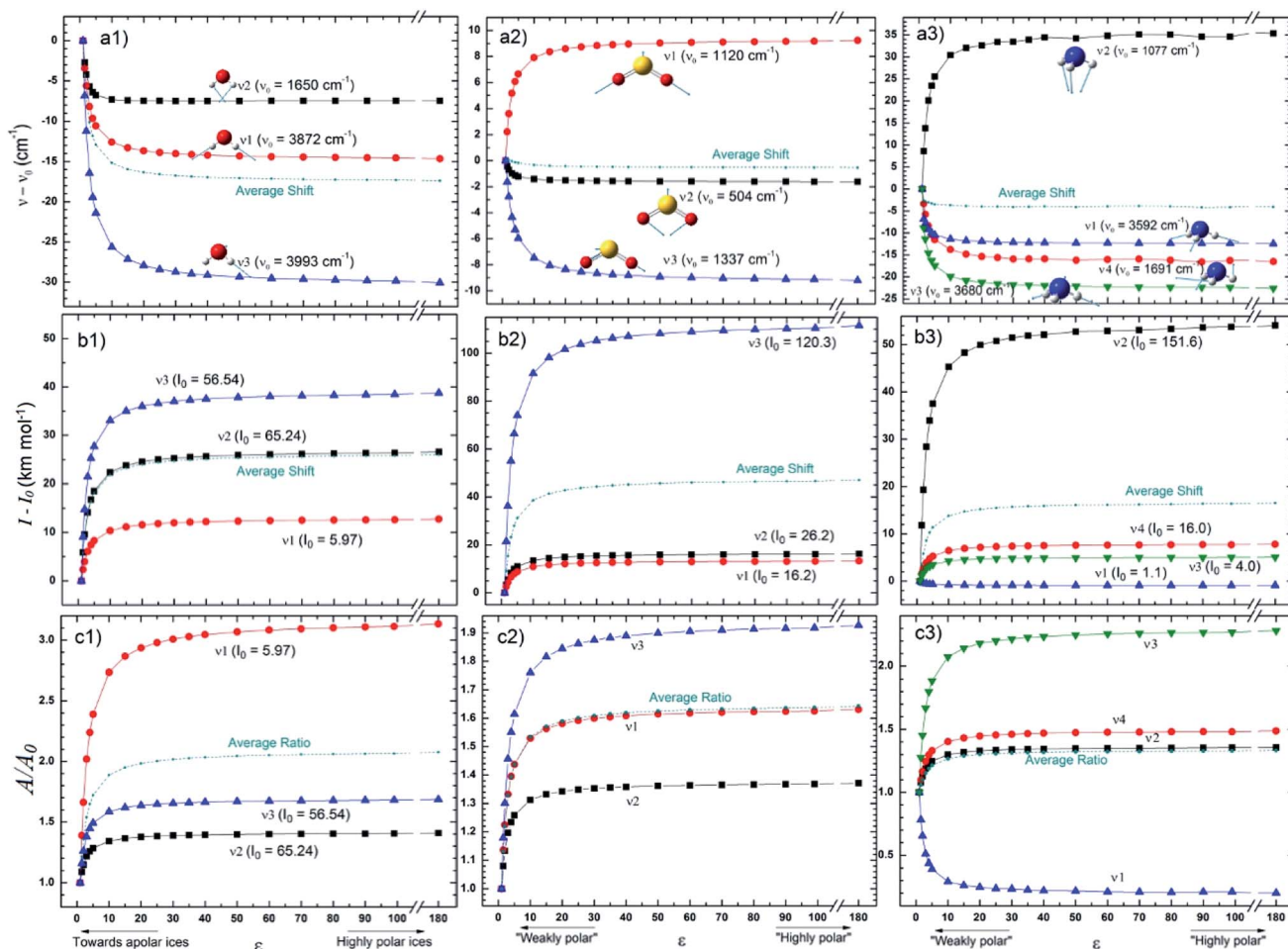


Fig. 3 Changes in the IR spectral features of three vibrational modes of frozen  $\text{H}_2\text{O}$ ,  $\text{SO}_2$  and  $\text{NH}_3$  molecules embedded in solid matrix as a function of its dielectric constant. (a) Frequencies shifts; (b) intensities shifts and (c) band strength ratio. The calculated values for  $\nu_0$ ,  $I_0$ , and  $A_0$  were performed considering  $\epsilon_0 = 1$  (gas phase), as the reference. See details in the text.

Fig. 4 presents the changes in the IR spectral features of the vibrational modes of frozen  $\text{HCOOH}$  and  $\text{CH}_3\text{OH}$  molecules embedded in icy matrix as a function of its dielectric constant  $\epsilon$  displayed in three different views: (a) frequencies shifts, (b) intensities shifts and, (c) band strength ratio. All the studied frequencies of  $\text{HCOOH}$  have presented a redshift (in comparison with the reference values at gas-phase) as the increasing of dielectric constant. For  $\text{CH}_3\text{OH}$ , only the frequencies  $\nu_{12}$  and  $\nu_8$  have presented such behavior, the frequencies  $\nu_3$  and  $\nu_9$  have presented a blueshift (in comparison with the reference values at gas-phase) with the increasing of dielectric constant. The two molecules have the intensities of its infrared band enhanced when the species are embedded in icy matrix in respect with the values at vacuum. Such behavior also suggests the enhancement of band strength as the dielectric constant increases, as it is observed in the panel c1 and c2, for  $\text{HCOOH}$  and  $\text{CH}_3\text{OH}$ , respectively.

Fig. 5 presents average changes in the IR spectral features of vibrational modes of studied frozen molecules embedded in icy matrix as a function of its dielectric constant  $\epsilon$ . Panels a, b and c indicates frequencies shifts, intensities shifts and band

strength ratio, respectively. The calculation considered  $\epsilon_0 = 1$  (gas phase) for  $\nu_0$ ,  $I_0$ , and  $A_0$ . Table 2 presents the employed parameters for the associative exponential that best fit the average dataset. From Fig. 5a we observe that the shift in frequencies due to the presence of dielectric medium was more enhanced for CO molecule (average blueshift effect as  $\epsilon$  increases) and  $\text{HCOOH}$  molecule (average redshift effect as  $\epsilon$  increases). In addition, two groups of molecules are identified: (i) the species that blueshift its vibration frequencies as  $\epsilon$  increases such as CO and  $\text{CH}_3\text{OH}$ , and (ii) the species the redshift its vibration frequencies as  $\epsilon$  increases such as  $\text{SO}_2$  (slightly affected),  $\text{NH}_3$ ,  $\text{CH}_4$ ,  $\text{CO}_2$ ,  $\text{H}_2\text{O}$  and  $\text{HCOOH}$ . Such blueshift/redshift behavior as dielectric constant increases is still a mystery and might be related with cavity effects of PCM ice model employed. Fig. 5b shows the influence of dielectric constant on band intensities for the studied species. We observe clearly that the presence of medium increase the band intensities and this dependence with dielectric constant could be described by an associative exponential function. The most affected species was  $\text{CO}_2$  and the less affected was the embedded CO species.



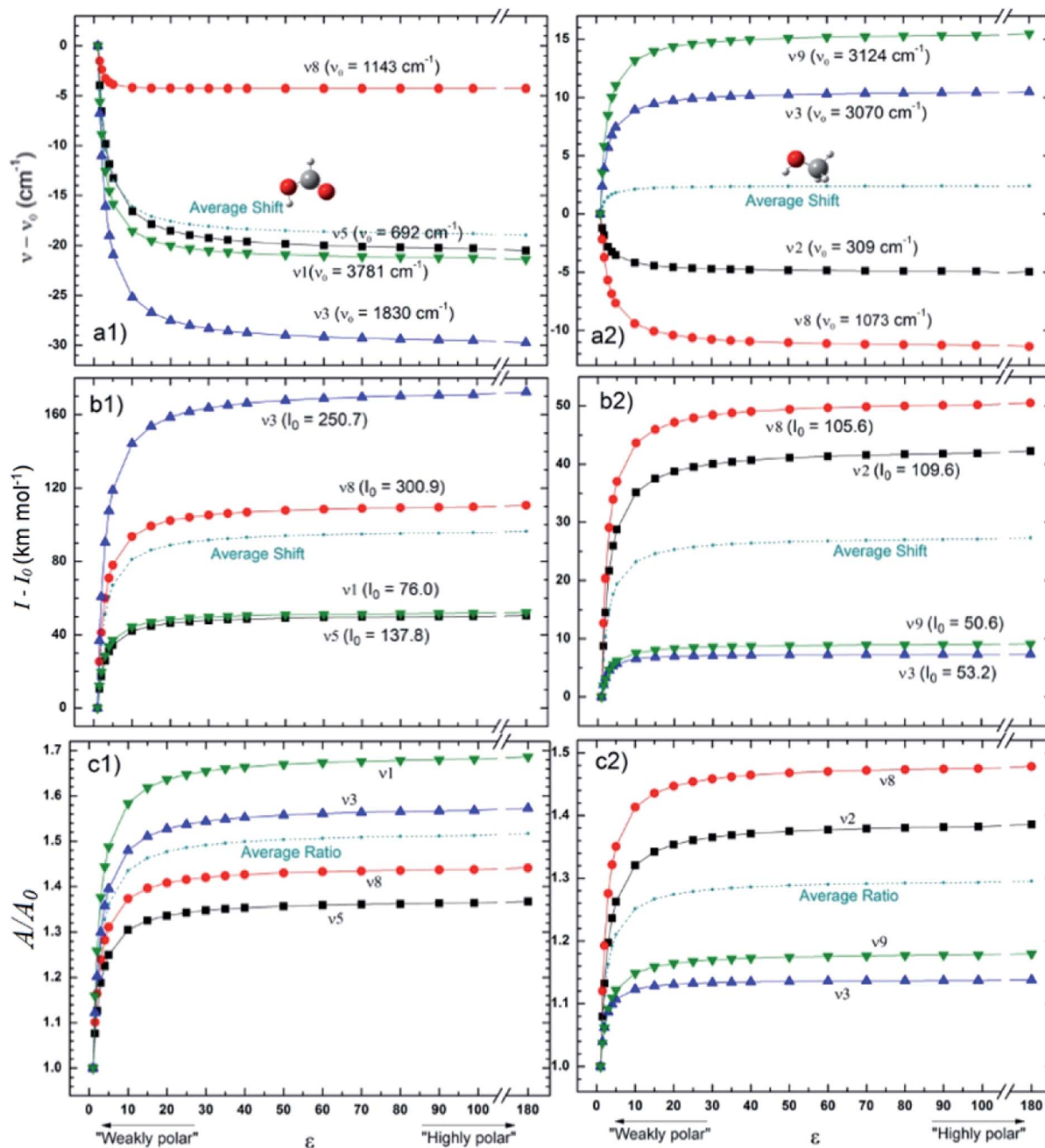


Fig. 4 Changes in the IR spectral features of vibrational modes of frozen HCOOH and CH<sub>3</sub>OH molecules embedded in icy matrix as a function of its dielectric constant  $\epsilon$ . The calculated values for  $\nu_0$ ,  $I_0$ , and  $A_0$  were performed considering  $\epsilon_0 = 1$  (gas phase), as the reference. See details in the text.

Curiously the average intensity of infrared bands of CO<sub>2</sub> and HCOOH presents the higher enhancement with the increases of dielectric constant (indicating that such species might be easier to be detected in astrophysical environments in highly polar ices than in apolar ices). Fig. 5c shows the dependence on band strength ratio  $A/A_0$  as a function of the dielectric constant. The most affected species was the H<sub>2</sub>O embedded species and the less affected was the embedded CH<sub>3</sub>OH species. Considering the entire studied dataset, the band strength (the average value) of embedded species at

dielectric constant of  $\epsilon = 10$  is 40% higher than the value calculated for the gaseous species or at  $\epsilon \sim 1$ . All panels in the Fig. 5 suggest that variation  $\epsilon$  in apolar ices induce much more changes than variation of  $\epsilon$  in polar ices (compared with the IR features in the gas-phase). This reinforces the concern on the effects of medium in the infrared properties of embedded species in space environments since its contain a variety of frozen species at different temperatures and porosity degrees.



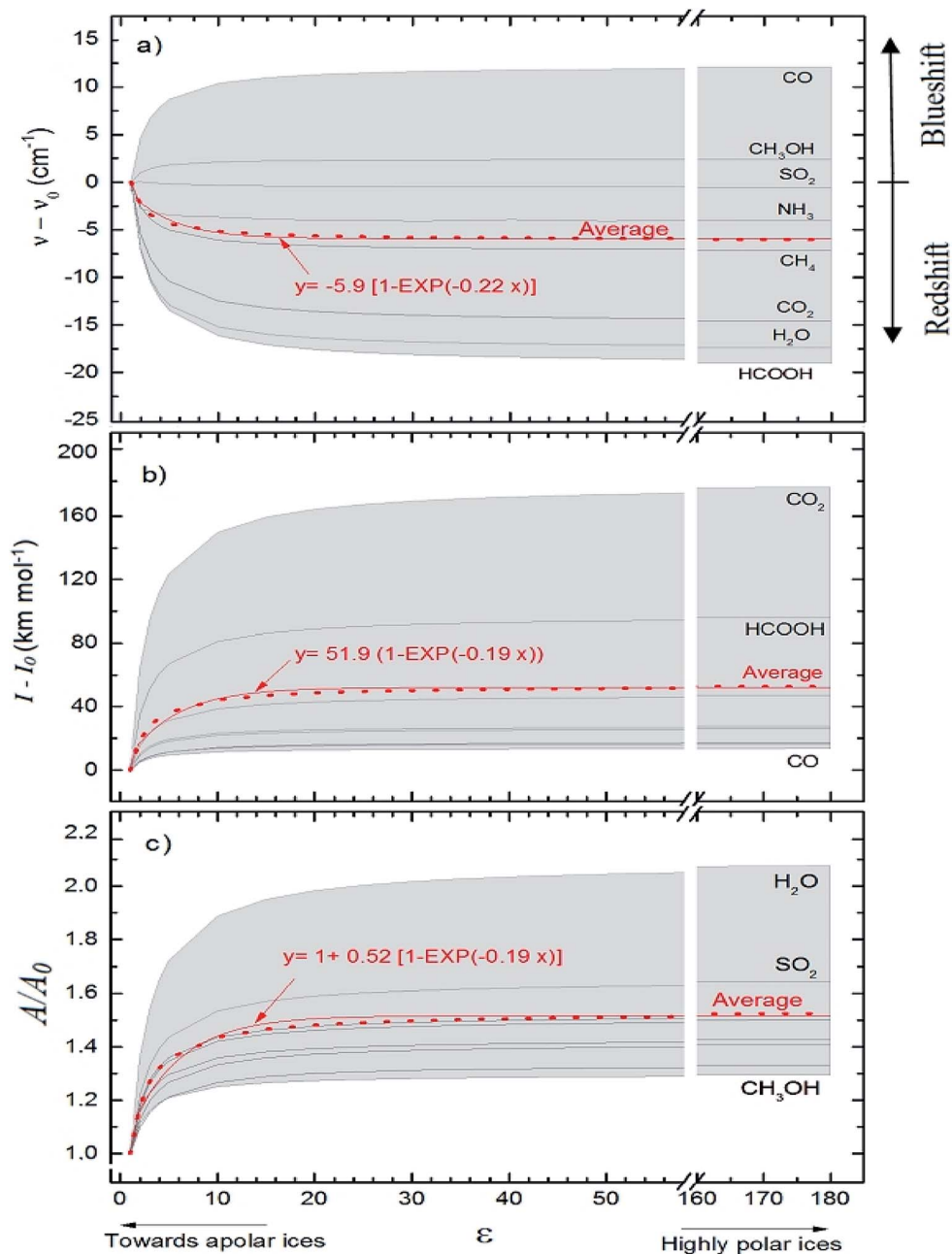


Fig. 5 Average changes in the IR spectral features of vibrational modes of studied frozen molecules embedded in icy matrix as a function of its dielectric constant  $\epsilon$ . Gas phase calculation is indicated by the data at  $\epsilon = 1$ . (a) Frequencies shifts; (b) intensities shifts and (c) band strength ratio (considering  $\epsilon_0 = 1$  (gas phase) for  $\nu_0$ ,  $I_0$ , and  $A_0$  calculations). The associative exponential that best fits the average dataset was also presented (see also Table 2).

### 3.2 Correlation between band strengths and band intensities

The band intensities observed in the infrared spectrum, as parameterized by the Lambert-beer law, is mainly related with the number of molecules (oscillators) in the line of sight and a quantum property related with the specific molecule (plus the vicinity) with take into an account how much the incoming radiation is affected by the molecular vibration in a given frequency (ruled by the band strength or  $A$ -value). In the current

calculation, the infrared intensity (absorption) was determined by the amount of variation of the electric dipole, of a specific vibration mode, of a single molecular species embedded in an electric field described by the given dielectric constant.

The correlation between calculated band intensities (calculated at epsilon equal to 3, expected for cold ices at  $\sim 10$  K)<sup>43,44</sup> with band strengths (taken from literature mostly at temperatures below 25 K) for the studied vibration modes (listed in Table 1) is shown in Fig. 6. The vertical lines over each black circle represents calculated values ranging from  $\epsilon = 1$  (for the



Table 2 Parameters employed for the associative exponential function that best fit average datasets in the Fig. 6

Fig. 6 <sup>a</sup>	Parameters employed in the equation $y = a + b[1 - \exp(-cx)]$		
	$a$	$b$	$c$
Panel a ( $\nu - \nu_0$ )	0	-5.9	0.22
Panel b ( $I - I_0$ )	0	51.9	0.19
Panel c ( $A/A_0$ )	1	0.52	0.19

<sup>a</sup> Considering  $\epsilon_0 = 1$  (gas phase) for  $\nu_0$ ,  $I_0$ , and  $A_0$  calculations.

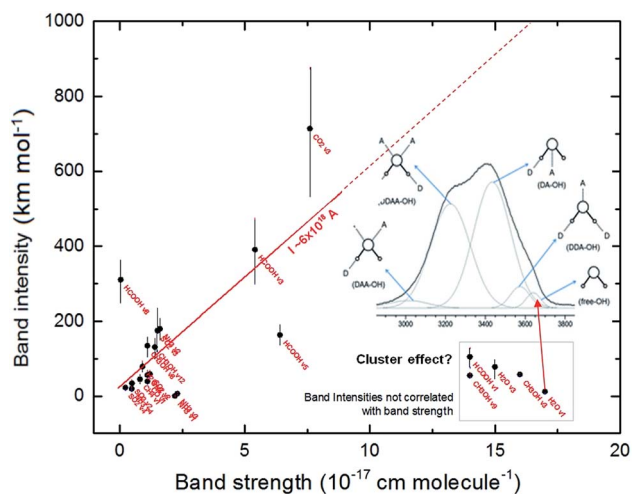


Fig. 6 Correlation between calculated band intensities (calculated at epsilon equal to 3) and band strengths (taken from literature) for the studied vibration modes. The vertical lines over each black circle represents calculated values ranging from  $\epsilon = 1$  (for the gas-phase) to 180 (higher value in the figure). The values inside the box in right-bottom corner indicate the vibration modes in which the calculated intensities do not have a correlation with band strengths and might be related with molecular cluster within the ice. Figure inset with band profile is taken from Choe *et al.*<sup>46</sup> and illustrates the influence of water molecular cluster in the  $\text{H}_2\text{O}$   $\nu_1$  band. See details in text.

gas-phase) to 180 (higher value in the figure). For the most part of the current studied vibration modes, the band intensities ( $I$ ) is correlated with the band strength ( $A$ ), in units of cm per molecule, by the linear function:

$$I \sim 6 \times 10^{18} A [\text{km mol}^{-1} \text{ or } \text{D}^2 \text{ \AA}^{-1} \text{ u}^{-1}] \quad (2)$$

However, for some vibration modes (e.g.  $\text{H}_2\text{O}$   $\nu_3$  and  $\nu_1$ ,  $\text{HCOOH}$   $\nu_1$  and  $\text{CH}_3\text{OH}$   $\nu_9$ ) the calculated intensities does not follow the linear tendency described by eqn (2). These vibration modes have, in the literature, large value of band strengths ( $\sim 1.5 \times 10^{-16}$  cm per molecule) and have presented very low values of calculated band intensities. One suggestion in attempt to explain such behavior might be the influence of the vibration of molecular clusters (molecules linked by hydrogen bonds) in the infrared band of ice associated with of such vibration mode. Such behavior might also be related with asymmetric band profiles of such vibration modes in the ices.<sup>45–48</sup> The inset figure,

with the spectrum profile, in the right part of Fig. 6 illustrates the influence of water molecular cluster in the  $\text{H}_2\text{O}$   $\nu_1$  band of typical water ice or liquid (taken from Choe *et al.*<sup>46</sup>).

The current investigation shows that the PCM approach can be valid for ices with low dielectric constants (apolar ices, very cold ices or for porous ices). Whereas, if there are other bonding considerations that may alter the IR intensities or  $A$ -values then the PCM approach may not be valid (the species shown in the gray box of Fig. 6 are examples of these). Bennett *et al.*<sup>51</sup> presented some examples of how PCM may be valid for ices of low dielectric values but did also note its failing for species such as  $\text{H}_2\text{O}$  in some cases.

### 3.3 Overview of changes in band strengths compared with typical of cold ice as function of degree of polarity

Fig. 7 presents the variation in the band strength of selected vibrational modes as a function dielectric constant normalized by its typical value for cold ice, characterized by a compact water-like matrix at 10 K ( $\epsilon_{\text{cold}} = 3$ ).<sup>28,43</sup> As two of the studied molecules, methanol and formic acid, have much more complex IR spectrum than the other ones, the intensities shown in this figure correspond to the four most intense modes in vacuum, while simpler molecules had all modes presented. A special exception was made for the ammonia molecule, because it showed one vibration mode ( $\sim 3337 \text{ cm}^{-1}$ ) with a decreasing behavior of  $A/A_{\text{cold}}$  with the enhancement of  $\epsilon/\epsilon_{\text{cold}}$ , contrary to the other values in the dataset, and so it was omitted in order to facilitate visualization of the whole set of modes behavior. In this figure we also present two different models, employing associative exponential functions, that best fit the calculated dataset of band strengths ratio. The apolar (and porous ices) are indicated by  $\epsilon/\epsilon_{\text{cold}} < 1$ , very polar ices are expected to be at  $\epsilon/\epsilon_{\text{cold}} < 5-10$ . We observe in this figure a decrease up to 20% for the

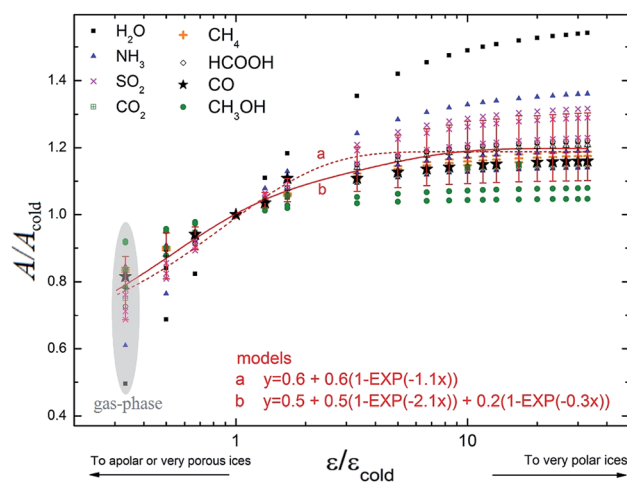


Fig. 7 Changes in the band strength of selected vibrational modes as a function dielectric constant normalized by its typical value for compact cold ices ( $\epsilon_{\text{cold}} = 3$ ). Apolar and porous ices are indicated by  $\epsilon/\epsilon_{\text{cold}} < 1$ , very polar ices are expected to be at  $\epsilon/\epsilon_{\text{cold}} < 5-10$ .  $A_{\text{cold}}$  indicated values calculated at  $\epsilon_{\text{cold}}$ . Two different models, employing associative exponential functions that best fits the studied dataset of band strengths ratio are also given.



values of band strength towards very porous ices (or highly apolar) and an increase of 20% for polar ices in comparison with the typical value of compact cold ice at 10 K.

The  $A/A_0$  value for the  $\text{H}_2\text{O}$  symmetric stretching mode varied by more than 50% with the change from  $\epsilon = 3.0$  (expected for cold ices around 10 K; Tsekouras *et al.*;<sup>43</sup> Ruepp<sup>44</sup>) to  $\epsilon = 180$  (expected for hot water-rich ices; Aragonés *et al.*<sup>27</sup>). Considering that the determination of molecular abundances is highly dependent on the band strength, it is therefore very important to better precise such value (and its dependence with temperature, porosity degree and chemical environment) to obtain more accurate molecular abundances of astrophysical ices.

Finally, we would like to point out that besides the employ of PCM be a good approximation to simulate the effects solid/icy matrices around embedded species, there are some works in the literature that addressed some deficiencies of the this approach towards predicting IR intensities.<sup>49,50</sup> Additionally, as pointed out by Bennett *et al.*<sup>51</sup> the PCM approach would likely provide good estimates of the band intensities (and band strengths) in ices with low dielectric constants (apolar ices, very cold ices or porous ices), however if there are other bonding considerations that may alter the IR intensities or  $A$ -values then the PCM approach may not be valid, for example when molecular clusters might exist within the ice.

Finally, following the ideas of Bonfim and Pilling,<sup>28</sup> the approach performed in the current work is likely to be better applicable to the very cold astrophysical ices where the dielectric constant is the dominant effect such as the case of apolar ices at very low temperatures (<50 K) and ices with some degree of porosity, and should be employed with caution for weakly-bounded very porous ices ( $\epsilon \sim 1$ ) and very polar ices ( $\epsilon \gg 10$ ).

## 4 Conclusions

In this work, we computationally simulate and investigate, through vibrational analysis, how the chemical environment influences some molecular properties in the IR spectra such as band position and intensities (linked with absorption cross section), of individual species embedded in the solid phase water-like matrix employing the Polarized Continuum Model (PCM) approach. The solid phase was simulated using different values of dielectric constant representing different types of matrix containing the studied molecules at PCM approach. The studied molecules employed in this research were CO,  $\text{CO}_2$ ,  $\text{CH}_4$ ,  $\text{NH}_3$ ,  $\text{SO}_2$ ,  $\text{HCOOH}$ ,  $\text{CH}_3\text{OH}$  and also  $\text{H}_2\text{O}$ .

Our main conclusions are listed below:

(i) In respect of frequency shift (in units of  $\text{cm}^{-1}$ ) two groups of molecules are identified: (1) molecular species that blueshift their vibration frequencies as  $\epsilon$  increases such as CO and  $\text{CH}_3\text{OH}$ , and (2) molecular species the redshift their vibration frequencies as  $\epsilon$  increases such as  $\text{SO}_2$  (slightly affected),  $\text{NH}_3$ ,  $\text{CH}_4$ ,  $\text{CO}_2$ ,  $\text{H}_2\text{O}$  and  $\text{HCOOH}$ . The shift in frequencies due to the presence of dielectric medium was more enhanced for CO molecule (average blueshift effect as  $\epsilon$  increases) and  $\text{HCOOH}$  molecule (average redshift effect as  $\epsilon$  increases).

(ii) We observe clearly that the presence of medium increases the band intensities and this dependence on dielectric constant

could be described by an associative exponential function. The most affected species was  $\text{CO}_2$  and the less affected was the embedded CO species. Curiously the average intensity of infrared bands of  $\text{CO}_2$  and  $\text{HCOOH}$  presents the higher enhancement with the increases of dielectric constant (indicating that such species might be easier to be detected in astrophysical environments in highly polar ices than in apolar ices).

(iii) In respect of the dependence on band strength ratio  $A/A_0$  as a function of dielectric constant we observe that the most affected species was the  $\text{H}_2\text{O}$  embedded species and the less affected was the embedded  $\text{CH}_3\text{OH}$  species. The  $A/A_0$  value for the  $\text{H}_2\text{O}$  symmetric stretching mode varied by more than 50% with the change from  $\epsilon = 3.0$  to  $\epsilon = 180$ . Considering the entire studied dataset, the band strength (the average value) of embedded species at dielectric constant of  $\epsilon = 10$  is 40% higher than the value calculated for the gaseous species or at  $\epsilon \sim 1$ . The results also suggest that variation  $\epsilon$  in apolar ices induce much more changes than variation of  $\epsilon$  in polar ices. This reinforces the concern on the effects of medium (icy matrix) in the infrared properties of embedded species in space environments since it contains a variety of frozen species at different temperatures and porosity degrees.

(iv) A correlation between calculated band intensities with the band strengths (taken from literature) were determined and described by a linear function  $I \sim 6 \times 10^{18} A [\text{km mol}^{-1}]$ , with  $A$  in unity of cm per molecule. In addition, an associative exponential functions were adjusted to the studied datasets to characterize the evolution of frequency-shift and intensity-shift and band strength ratio as function of dielectric constant. We hope the tendencies and behavior that we pointed out here in this manuscript might be better characterized in future works employing high precision quantum chemistry calculation methods.

Considering that the determination of molecular abundances is highly dependent on the band strength and that it can be affected by the dielectric constant of the medium (molecular environment), this study helps us to better characterize solid samples in the infrared, specially, for astrophysical ices monitored by infrared observations.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

The authors would like to thank the Brazilian agencies FAPESP (#JP2009/18304-0), CAPES and CNPq (#306145/2015-4; #302985/2018-2).

## References

- 1 A. C. A. Boogert and P. Ehrenfreund, *ASP Conf. Ser. 309, Astrophysics of Dust*, ed. A. N. Witt, G. C. Clayton and B. T. Draine, ASP, San Francisco, 2004, p. 547.



- 2 E. L. Gibb, D. C. B. Whittet and J. E. Chiar, *Astrophys. J.*, 2001, **558**, 702.
- 3 S. Pilling, E. Seperuelo Duarte, A. Domaracka, H. Rothard, P. Boduch and E. F. da Silveira, *Astron. Astrophys.*, 2010, **523**, A77.
- 4 P. A. Gerakines, W. A. Schutte, J. M. Greenberg and E. F. van Dishoeck, *Astron. Astrophys.*, 1995, **296**, 810.
- 5 O. Kerkhof, W. A. Schutte and P. Ehrenfreund, *Astron. Astrophys.*, 1999, **346**, 990.
- 6 L. B. D'Hendecourt and L. J. Allamandola, *Astron. Astrophys., Suppl. Ser.*, 1986, **64**, 453.
- 7 P. A. Gerakines, J. J. Bray, A. Davis and C. R. Richey, *Astrophys. J.*, 2005, **620**, 1140.
- 8 D. M. Hudgins, S. A. Sandford, L. J. Allamandola and A. G. G. M. Tielens, *Astrophys. J., Suppl. Ser.*, 1993, **86**, 713.
- 9 D. Fulvio, B. Sivaraman, G. A. Baratta, M. E. Palumbo and N. J. Mason, *Spectrochim. Acta, Part A*, 2009, **72**, 1007.
- 10 R. L. Hudson, P. A. Gerakines and R. F. Ferrante, *Spectrochim. Acta, Part A*, 2018, **193**, 33.
- 11 R. L. Hudson, R. F. Ferrante and M. H. Moore, *Icarus*, 2014, **276**, 287.
- 12 M. H. Moore, R. F. Ferrante, W. J. Moore and R. L. Hudson, *Astrophys. J., Suppl. Ser.*, 2010, **191**, 96.
- 13 K. I. Öberg, H. J. Fraser, A. C. A. Boogert, *et al.*, *Astron. Astrophys.*, 2007, **462**, 1187.
- 14 S. Miertuš, E. Scrocco and J. Tomasi, *Chem. Phys.*, 1981, **55**, 117.
- 15 E. Cancès, B. Mennucci and J. Tomasi, *J. Chem. Phys.*, 1997, **107**, 3032.
- 16 V. S. Bonfim, R. B. Castilho, L. Baptista and S. Pilling, *J. Phys.: Conf. Ser.*, 2015, **635**, 102012.
- 17 V. S. Bonfim, R. B. Castilho, L. Baptista and S. Pilling, *Phys. Chem. Chem. Phys.*, 2017, **19**, 26906.
- 18 J. Y. Park and D. E. Woon, *J. Phys. Chem. A*, 2004, **108**, 6589.
- 19 S. Pilling, L. Baptista, H. M. Boechat-Roberty and D. P. P. Andrade, *Astrobiology*, 2011, **11**, 883.
- 20 D. E. Woon and J.-Y. Park, *Astrophys. J.*, 2004, **607**, 342.
- 21 P. Ehrenfreund and S. B. Charnley, *Annu. Rev. Astron. Astrophys.*, 2000, **38**, 427.
- 22 P. Ehrenfreund and W. A. Schutte, Infrared Observations of Interstellar Ices, in *Astrochemistry: From Molecular Clouds to Planetary Systems, Proceedings of IAU Symposium 197*, ed. Y. C. Minh and E. F. van Dishoeck, 2000, p. 135.
- 23 R. Cammi and J. Tomasi, *J. Comput. Chem.*, 1995, **16**, 14498.
- 24 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria and M. A. Robb, *et al.*, *Gaussian 09 (Revision B.01)*, 2010.
- 25 T. H. Dunning Jr, *J. Chem. Phys.*, 1989, **90**, 1007.
- 26 E. Silla, *et al.*, Fundamental principles governing Solvents Use, in *Handbook of Solvents*, ed. G. Wypych, ChemTec Publishing, 2nd edn, 2014, pp. 11–72, ISBN 9781895198645, DOI: 10.1016/B978-1-895198-64-5.50004-0.
- 27 J. L. Aragonés, L. G. Macdowell and C. Vega, *J. Phys. Chem. A*, 2011, **115**, 5745.
- 28 V. S. Bonfim and S. Pilling, *Proc. Int. Astron. Union*, 2017, **13**(S332), 346–352.
- 29 S. Pilling, E. S. Duarte, A. Domaracka, H. Rothard, P. Boduch and E. F. da Silveira, *Phys. Chem. Chem. Phys.*, 2011, **13**, 15755.
- 30 B. Mate, in *Infrared optical constants and Band Strengths of ices, Laboratory Astrophysics, Astrophysics and Space Science Library 451*, ed. G. M. Muñoz-Caro and R. Escribano, Springer, 2018, pp. 71–86, DOI: 10.1007/978-3-319-90020-9\_5.
- 31 B. Schmitt, E. Quirico, F. Trotta and W. M. Grundy, in *Optical properties of ices from UV to infrared, Solar System Ices*, ed. B. Schmitt, *et al.*, Kluwer Academic Publishers, Netherlands, 1998, pp. 199–204.
- 32 J. Bouwman, L. Wiebke, *et al.*, *Astronomy & Astrophysics*, 2007, **476**, 995.
- 33 K. I. Öberg, H. J. Fraser, A. C. A. Boogert, *et al.*, *Astronomy & Astrophysics*, 2007, **462**, 1187.
- 34 R. Luna, G. Molpeceres, *et al.*, *Astronomy & Astrophysics*, 2018, **617**, A116.
- 35 S. E. Bisschop, *et al.*, *Astronomy & Astrophysics*, 2007, **470**, 749.
- 36 M. Bouilloud, N. Fray, Y. Bénilan, H. Cottin, M. C. Gazeau and A. Jolly, *Mon. Not. R. Astron. Soc.*, 2015, **451**, 2145.
- 37 R. K. Khanna, G. Zhao, M. J. Ospina and J. C. Pearl, *Spectrochim. Acta, Part A*, 1988, **44**, 581–586.
- 38 M. Garozzo, D. Fulvio, O. Gomis, M. E. Palumbo and G. Strazzulla, *Planet. Space Sci.*, 2008, **56**, 1300–1308.
- 39 S. A. Sandford and L. J. Allamandola, *Astrophys. J.*, 1993, **417**, 815.
- 40 L. J. Karssemeijer, S. Ioppolo, M. C. van Hemert, A. van der Avoird, M. A. Allodi, G. A. Blake and H. M. Cuppen, *Astrophys. J.*, 2014, **781**, 16.
- 41 A. C. A. Boogert, P. A. Gerakines and D. C. B. Whittet, *Annu. Rev. Astron. Astrophys.*, 2015, **53**, 541–581.
- 42 A. C. A. Boogert, W. A. Schutte, F. P. Helmich, A. G. G. M. Tielens and D. H. Wooden, *Astron. Astrophys.*, 1997, **317**, 929–941.
- 43 A. A. Tsekouras, M. J. Iedema and J. P. Cowin, *Phys. Rev. Lett.*, 1998, **80**, 5798.
- 44 R. Ruepp, *Physics and chemistry of ice. Papers presented at the Symposium on the Physics and Chemistry of Ice, Ottawa-Canada, 14–18 August 1972*, ed. E. Whalley, S. J. Jones and L. W. Gold, Royal Society of Canada, Ottawa, 1973, pp. 179–86, ASIN: B007ETWKN0.
- 45 J. O. Daldrop, M. Saita, M. Heyden, *et al.*, *Nat. Commun.*, 2018, **9**, 311.
- 46 C. Choe, J. Lademann and M. E. Darvin, *Analyst*, 2016, **141**, 6329.
- 47 D. A. Schmidt and K. Miki, *J. Phys. Chem. A*, 2007, **111**, 10119.
- 48 N. A. Chumaevskii and M. N. Rodnikova, *J. Mol. Liq.*, 2003, **106**, 167.
- 49 S. A. Katyuba, E. E. Zvereva and T. I. Burganov, *J. Phys. Chem. A*, 2013, **117**, 6664.
- 50 A. Gladney, C. Qin and H. Tamboue, *Open J. Phys. Chem.*, 2012, **2**, 117.
- 51 C. J. Bennett, P. Courtney and R. I. Kaiser, *Astrophys. J.*, 2014, **782**, 63.

