


 Cite this: *RSC Adv.*, 2020, 10, 2747

# Significant enhancement of dielectric permittivity and percolation behaviour of $\text{La}_{2-x}\text{Sr}_x\text{NiO}_4$ /poly(vinylidene fluoride) composites with different Sr doping concentrations

 Keerati Meeporn,<sup>a</sup> Narong Chanlek<sup>b</sup> and Prasit Thongbai \*<sup>cd</sup>

The percolation behaviour and dielectric properties of  $\text{La}_{2-x}\text{Sr}_x\text{NiO}_4$  (LSNO)/poly(vinylidene fluoride) (PVDF) composites with different Sr doping concentrations were investigated. The semiconducting LSNO filler particles with  $x = 0.2$  (LSNO-1) and  $x = 0.4$  (LSNO-2) were prepared using a chemical combustion method. The microstructures, thermal properties, and phase compositions of the polymer composites and filler particles were systematically investigated. The conductivity of the LSNO fillers increased with the Sr content and had an important impact on the dielectric properties of the LSNO/PVDF composites. The percolation threshold of the LSNO-2/PVDF composite was lower than that of the LSNO-1/PVDF composite. An ultra-high dielectric permittivity ( $\epsilon'$ ) of 3384.7 (at 1 kHz and room temperature), which was approximately 340 times higher than that of pure PVDF, was obtained for the LSNO-2/PVDF composite with a filler volume fraction of 25 vol%. The enhanced dielectric properties were attributed to interfacial polarisation at the semiconductor–insulator interface, a micro-capacitor model, and the intrinsically remarkable dielectric properties of the LSNO ceramic.

 Received 20th November 2019  
 Accepted 7th January 2020

DOI: 10.1039/c9ra09719h

[rsc.li/rsc-advances](http://rsc.li/rsc-advances)

## 1. Introduction

Polymer-based composite materials with high dielectric permittivity ( $\epsilon'$ ) values have attracted increasing research focus in recent years because they have several advantages over dielectric ceramics and potential applications in various fields, including microelectronics, pulsed power, bandpass filtering, power conditioning, sensing, and energy storage.<sup>1–7</sup> The dielectric permittivity of a polymer-based composite material can be dramatically enhanced by introducing conductive metal fillers (Ni, Al, and Ag) and carbon-based fillers (multiwalled carbon nanotubes, graphene, and carbon black) into the polymer matrix. A dramatically increased dielectric permittivity was obtained when the volume fraction of filler used was increased to a critical fraction, *i.e.*, a percolation threshold ( $f_c$ ), according to the percolation theory.<sup>8–12</sup> However, a conductor–polymer composite with a greatly enhanced dielectric permittivity near  $f_c$  still cannot be considered an effective material in practical applications, especially in capacitor

applications, because of its high conductivity ( $\sigma$ ) and high dielectric loss ( $\tan \delta$ ). Therefore, it is necessary to develop a new composite system. Semiconductor–polymer composite systems have been studied as promising candidates because of their lower conductivities and high dielectric permittivities when the filler content is close to the percolation threshold.

Recently, several semiconducting materials have been used with several kinds of polymer matrices.<sup>13–19</sup> Among these,  $\text{La}_{2-x}\text{Sr}_x\text{NiO}_4$  (LSNO) ceramic oxides have attracted considerable attention as alternative semiconducting fillers for polymer composite dielectric materials because of their semiconductive properties and remarkably high dielectric permittivities of  $10^5$  to  $10^6$ . They can also exhibit a very high dielectric permittivity of approximately  $10^4$  even in the gigahertz region. The extremely good dielectric properties of LSNO ceramics are attributed to an intrinsic effect, caused by polaron hopping between  $\text{Ni}^{2+}$  and  $\text{Ni}^{3+}$ .<sup>20–24</sup> A filler material with an extremely high dielectric permittivity as a result of an intrinsic property is more suitable for compositing than one with such a dielectric permittivity resulting from an extrinsic effect at grain boundaries. Therefore, LSNO ceramics are promising effective fillers for polymer-based composite systems. Furthermore, it has been reported that their conductivities depend on the Sr content.<sup>25</sup> Therefore, it is necessary to understand the dielectric properties and percolation thresholds of composites filled with LSNO ceramic fillers.

In this work, two semiconductor–polymer composite systems were fabricated, using  $\text{La}_{2-x}\text{Sr}_x\text{NiO}_4$  (LSNO) ceramics

<sup>a</sup>Materials Science and Nanotechnology Program, Faculty of Science, Khon Kaen University, Khon Kaen 40002, Thailand

<sup>b</sup>Synchrotron Light Research Institute (Public Organization), 111 University Avenue, Muang District, Nakhon Ratchasima 30000, Thailand

<sup>c</sup>Department of Physics, Faculty of Science, Khon Kaen University, Khon Kaen 40002, Thailand. E-mail: pthongbai@kku.ac.th

<sup>d</sup>Institute of Nanomaterials Research and Innovation for Energy (IN-RIE), NANOTEC–KKU RNN on Nanomaterials Research and Innovation for Energy, Khon Kaen University, Khon Kaen, 40002, Thailand



with  $x = 0.2$  (LSNO-1) and  $x = 0.4$  (LSNO-2) as fillers. Poly(vinylidene fluoride) (PVDF) polymer was selected as the polymer matrix. The dependencies of the dielectric properties and percolation threshold on the conductivities of the LSNO fillers for the two composite systems are discussed in detail. The dielectrics and percolation mechanisms of the LSNO/PVDF composites are explained by the effective medium percolation theory model (EMPT). PVDF was chosen as the matrix because of its superior mechanical, thermal, and dielectric properties.

## 2. Experimental details

The precursors used for synthesising the LSNO fillers included  $\text{Sr}(\text{NO}_3)_2$  (99.0% purity),  $\text{La}(\text{NO}_3)_3 \cdot 6\text{H}_2\text{O}$  (99.99% purity), and  $\text{CH}_4\text{N}_2\text{O}$  (99.0–100.5% purity), which were purchased from Sigma-Aldrich.  $\text{Ni}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$  (99.0% purity) and  $\text{C}_6\text{H}_8\text{O}_7 \cdot \text{H}_2\text{O}$  (99.5% purity) were purchased from Kanto Chemical and RCI Labscan, respectively. The chemicals employed for the PVDF-based composite fabrication consisted of dimethylformamide (DMF) and PVDF ( $M_w$  of  $\sim 534\,000$ ), which were purchased from RCI Labscan and Sigma-Aldrich, respectively.

### Synthesis of $\text{La}_{2-x}\text{Sr}_x\text{NiO}_4$ (LSNO) ceramic fillers

LSNO fillers with  $x = 0.2$  (LSNO-1) and  $x = 0.4$  (LSNO-2) were synthesised using a chemical combustion method to prepare nanocrystalline LSNO.<sup>23,24</sup> First, a  $\text{C}_6\text{H}_8\text{O}_7 \cdot \text{H}_2\text{O}$  solution was produced by dissolving  $\text{C}_6\text{H}_8\text{O}_7 \cdot \text{H}_2\text{O}$  in deionised water. Second, stoichiometric amounts of  $\text{Ni}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$ ,  $\text{Sr}(\text{NO}_3)_2$ , and  $\text{La}(\text{NO}_3)_3 \cdot 6\text{H}_2\text{O}$  were dissolved in the  $\text{C}_6\text{H}_8\text{O}_7 \cdot \text{H}_2\text{O}$  solution, after which  $\text{CH}_4\text{N}_2\text{O}$  was added and stirred constantly at room temperature (RT). After this, the temperature of the mixed solution was increased from RT to  $200\text{ }^\circ\text{C}$  until a green treacly gel was obtained. In the next step, this green gel was dehydrated at  $300\text{ }^\circ\text{C}$  for 30 min and then calcined at  $1000\text{ }^\circ\text{C}$  for 6 h to obtain the LSNO phase without impurities.

### Fabrication of LSNO-1/PVDF and LSNO-2/PVDF composites

Two PVDF-based composite systems using LSNO-1 and LSNO-2 as fillers were fabricated *via* a solution processing method similar to a previously reported technique.<sup>26</sup> First, the PVDF powder was dissolved in DMF solvent at RT until the solution

was transparent. Second, the LSNO-1 and LSNO-2 fillers were separately added to the PVDF solution at 5–25 vol% and continuously stirred for 4 h. To enhance the dispersion of the ceramic filler, each mixture was further ultrasonicated for 10 min. Finally, each composite solution was poured into a casting plate and dried in an oven at  $70\text{ }^\circ\text{C}$  for 6 h. Afterward, LSNO-1/PVDF and LSNO-2/PVDF composite films with thicknesses of 50–195  $\mu\text{m}$  were obtained. The LSNO-1/PVDF and LSNO-2/PVDF composites with filler volume fractions of 5, 10, 15, 20, and 25 vol% were called 5 vol% LSNO-1, 10 vol% LSNO-1, 15 vol% LSNO-1, 20 vol% LSNO-1, and 25 vol% LSNO-1, respectively, and 5 vol% LSNO-2, 10 vol% LSNO-2, 15 vol% LSNO-2, 20 vol% LSNO-2, and 25 vol% LSNO-2, respectively.

### Material characterisations

The morphologies and phase compositions of the LSNO fillers and PVDF-based composites were characterised using scanning electron microscopy (SEM; SEC, SNE4500M), transmission electron microscopy (TEM; FEI, TECNAI G2 20), and powder X-ray diffraction (XRD; PANalytical, EMPYREAN). The chemical-state information of the Ni in the LSNO fillers and thermal properties of the PVDF-based composites were evaluated using X-ray photoelectron spectroscopy (XPS; PHI5000 VersaProbe II, ULVAC-PHI) and thermogravimetric analysis (TGA; Pyris Diamond, PerkinElmer), respectively. An analysis using a KEY-SIGHT E4990A impedance analyser was performed to investigate the dielectric properties of the PVDF-based composites. Prior to the dielectric measurement, silver electrodes were formed on both sides of each composite sample using silver paste.

## 3. Results and discussion

The microstructures of the LSNO-1 and LSNO-2 fillers prepared *via* a chemical combustion method and calcined at  $1000\text{ }^\circ\text{C}$  for 6 h are shown in Fig. 1(a) and (b), respectively. The SEM and TEM images (insets) reveal that the LSNO-1 and LSNO-2 powders consisted of nearly spherical particles with average radii of approximately 140–160 nm. It was found that as a result of the surface free energy, the assemblage between neighbouring particles was significantly decreased through the close

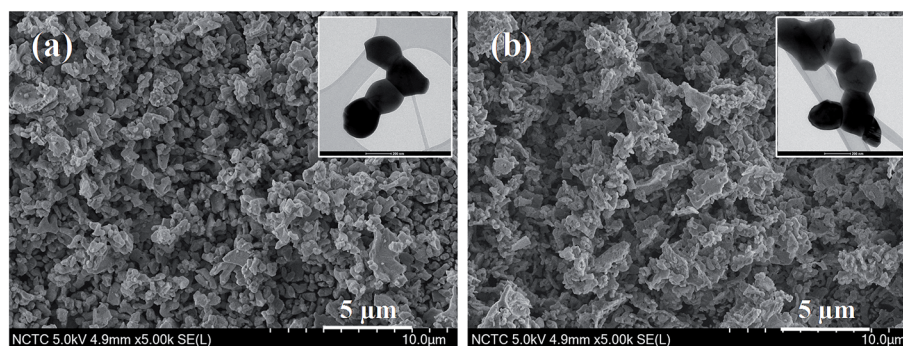


Fig. 1 SEM images of (a) LSNO-1 and (b) LSNO-2 fillers obtained using combustion method. The insets show TEM images of the LSNO-1 and LSNO-2 fillers.



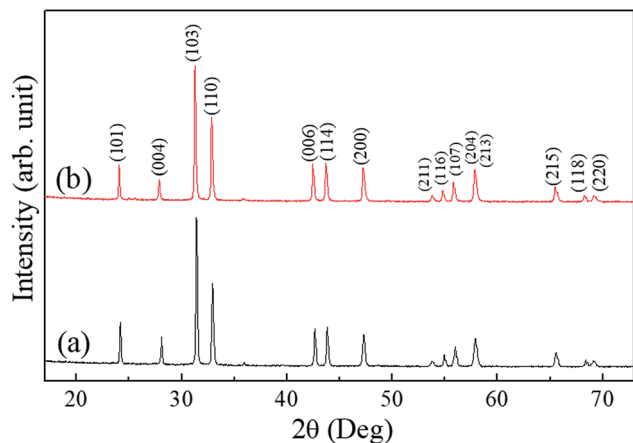


Fig. 2 XRD patterns of (a) LSNO-1 and (b) LSNO-2 fillers calcined at 1000 °C for 6 h.

packing of spherical particles.<sup>27</sup> The XRD patterns of the LSNO-1 and LSNO-2 fillers are given in Fig. 2. The peaks at  $2\theta$  in both samples interrelated to  $24^\circ$  (101),  $27^\circ$  (004),  $31^\circ$  (103),  $32^\circ$  (110),

$42^\circ$  (006),  $43^\circ$  (114),  $47^\circ$  (200),  $54^\circ$  (211),  $55^\circ$  (116),  $56^\circ$  (107),  $57^\circ$  (204),  $58^\circ$  (213),  $65^\circ$  (215),  $68^\circ$  (118), and  $69^\circ$  (220) can be assigned to LSNO with a tetragonal structure according to JCPDS no. 32-1241. The XRD patterns of the LSNO-1 and LSNO-2 ceramic particles are similar to those reported in the literature.<sup>24,33</sup> No visible diffraction peaks of impurity phases are observed, indicating that a chemical combustion method is an effective route for the preparation of  $\text{La}_{2-x}\text{Sr}_x\text{NiO}_4$  (LSNO) ceramic powders.

Fig. 3 shows macroscopic images of the flexibility of the 25 vol% LSNO-1 and 25 vol% LSNO-2 composite samples and SEM images of the fractured surfaces of the pure PVDF film and composites with the LSNO-1 and LSNO-2 fillers. As clearly seen in Fig. 3(b), the pure PVDF film prepared *via* solution casting is orderly, and a tight surface without any pores or impurities is observed. Fig. 3(c)–(f) show the fractured surfaces of the prepared 10 vol% LSNO-1, 10 vol% LSNO-2, 25 vol% LSNO-1, and 25 vol% LSNO-2 samples, respectively. The bright particles observed are the LSNO fillers, which are randomly dispersed across the PVDF matrix, as can be seen in Fig. 3(c) and

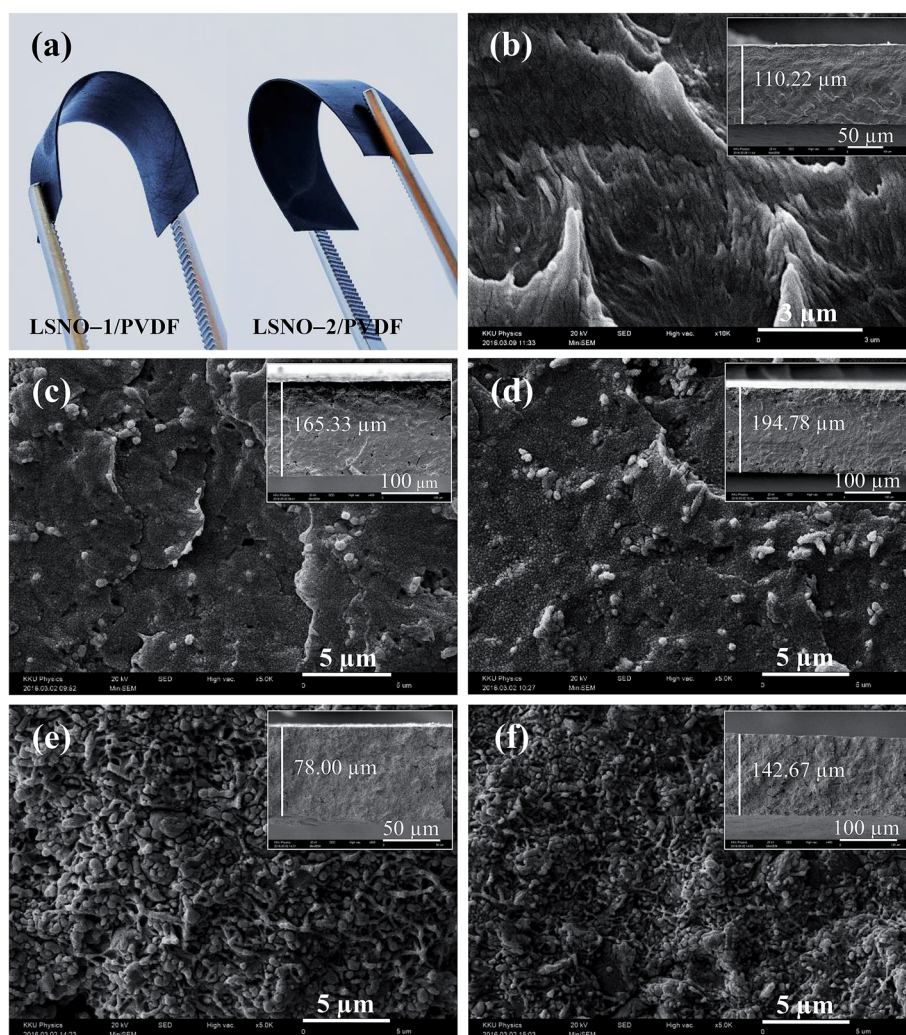


Fig. 3 Macroscopic images of (a) 25 vol% LSNO-1 and 25 vol% LSNO-2 composites. Fractured surfaces of (b) pure PVDF film, (c) 10 vol% LSNO-1, (d) 10 vol% LSNO-2, (e) 25 vol% LSNO-1, and (f) 25 vol% LSNO-2 composites.



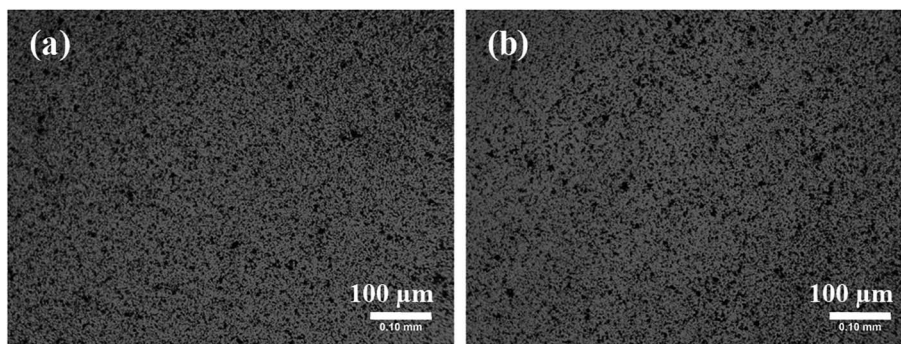


Fig. 4 Optical microscopy images showing dispersion states of LSNO-1 and LSNO-2 fillers in PVDF matrix: (a) 25 vol% LSNO-1 and (b) 25 vol% LSNO-2 composites.

(d). When the LSNO filler content increased to 25 vol%, as shown in Fig. 3(e) and (f), the fractured surfaces of the composites became rough, but the LSNO fillers were still well-dispersed. Interestingly, the composites with the 25 vol% filler concentrations still exhibited good mechanical flexibility, as seen in Fig. 3(a), which can be considered a beneficial effect of maintaining a well-dispersed LSNO filler distribution and using an effective preparation method. The overall dispersion states of the LSNO-1 and LSNO-2 fillers in a large area were also characterised using an optical microscope. Fig. 4(a) and (b) show the results of dispersion analyses of the 25 vol% LSNO-1 and 25 vol% LSNO-2 composites, respectively. Although it is observed that the LSNO fillers remained well-dispersed in the composites, small clusters of the LSNO fillers are also observed, as seen in the larger black spots. This may have resulted from the numerous particles and high surface free energy of the nanoparticles. It is important to note that the homogeneous dispersion of fillers can improve the dielectric properties of the composites.

The thermal stabilities of the LSNO-1/PVDF and LSNO-2/PVDF composites with various filler volume fractions are shown in Fig. 5(a) and (b), respectively. The TGA results for the LSNO/PVDF composites show the same trend, with the single state decomposition of the PVDF observed. The addition of the LSNO fillers caused an increase in the thermal stability of the PVDF. The temperature at a 25% weight loss ( $T_{25}$ ) increased from 450 °C for the pure PVDF film to 471, 480, 481, 482, and 484 °C for the 5 vol% LSNO-1, 10 vol% LSNO-1, 15 vol% LSNO-1, 20 vol% LSNO-1, and 25 vol% LSNO-1 composites and increases to 473, 480, 483, 484, and 484 °C for 5 vol% LSNO-2, 10 vol% LSNO-2, 15 vol% LSNO-2, 20 vol% LSNO-2, and 25 vol% LSNO-2 composites, respectively. These results indicate that the LSNO-1 and LSNO-2 fillers can slightly improve the thermal stability of the PVDF matrix.

Fig. 6(a) and (b) reveal the FT-IR spectra for the LSNO-1/PVDF and LSNO-2/PVDF composites with filler contents of 0 (pure PVDF) to 25 vol%, respectively. It was found that the LSNO/PVDF composites consisted of  $\alpha$ -phase and  $\beta$ -phase PVDF. The absorbance band at 766  $\text{cm}^{-1}$  was assigned to the nonpolar  $\alpha$ -phase,<sup>28</sup> and polar  $\beta$ -phase was observed at 840, 878, 1170, and 1231  $\text{cm}^{-1}$ . The absorbance bands at 1070 and 1402 were related to the

bending of C–C–C and the in-plane bending or scissoring of  $\text{CH}_2$ , respectively.<sup>27</sup> The absorbance bands at 766 and 840  $\text{cm}^{-1}$  corresponding to the  $\alpha$ -phase and  $\beta$ -phase, respectively, were utilised to calculate the content of the  $\beta$ -phase using eqn (1),<sup>28</sup>

$$F(\beta) = \frac{A_{\beta}}{(K_{\beta}/K_{\alpha})A_{\alpha} + A_{\beta}} \quad (1)$$

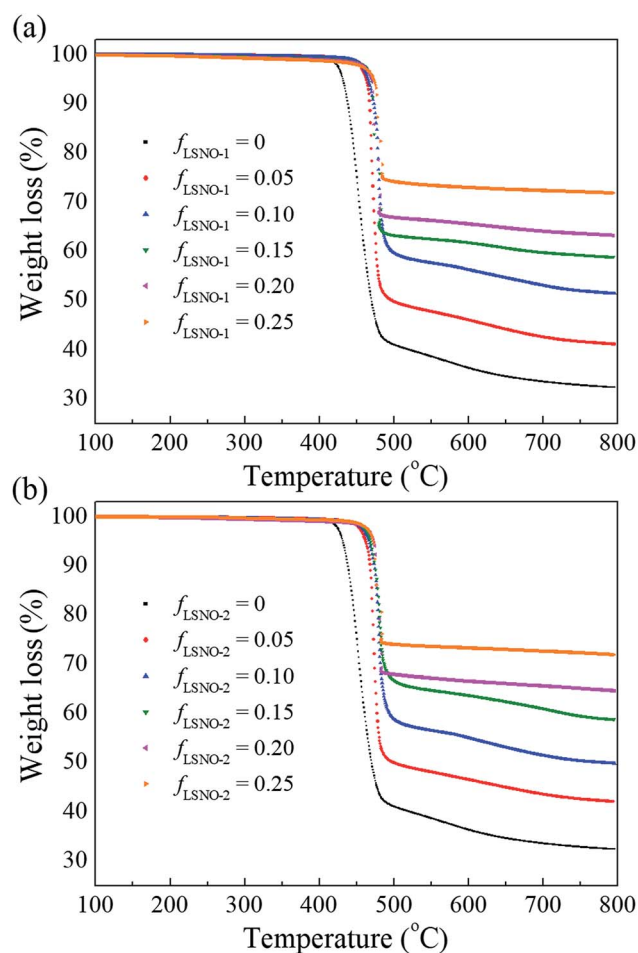


Fig. 5 TGA curves of (a) LSNO-1/PVDF and (b) LSNO-2/PVDF composites with filler content of  $f = 0$ –0.25.



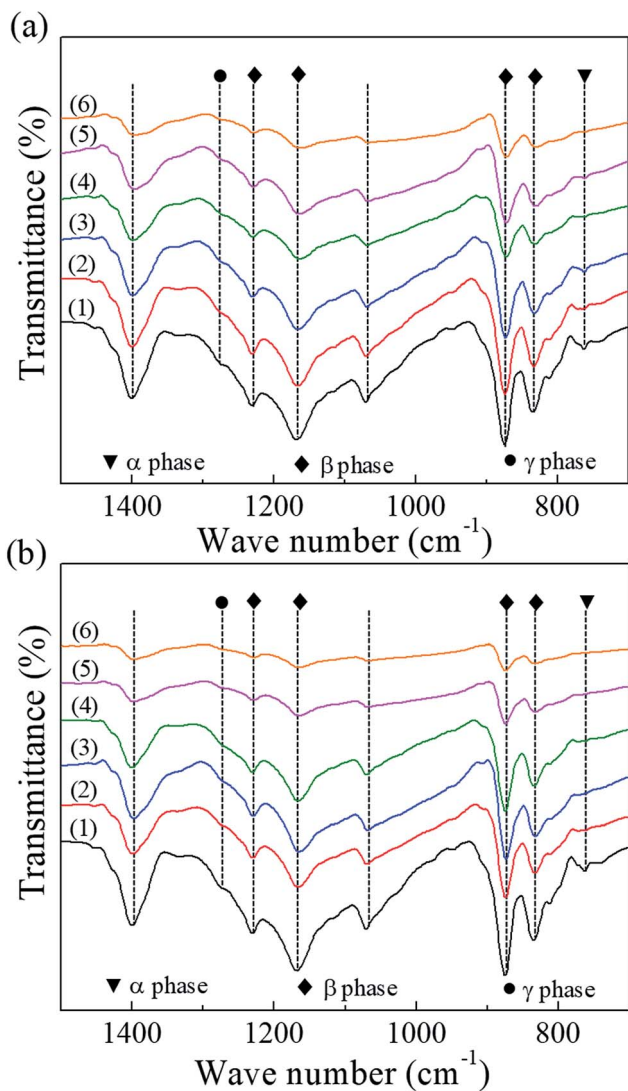


Fig. 6 FT-IR spectra of (a) LSNO-1/PVDF composites with LSNO-1 contents of (1) 0 vol%, (2) 5 vol%, (3) 10 vol%, (4) 15 vol%, (5) 20 vol%, and (6) 25 vol%, and (b) LSNO-2/PVDF composites with LSNO-2 contents of (1) 0 vol%, (2) 5 vol%, (3) 10 vol%, (4) 15 vol%, (5) 20 vol%, and (6) 25 vol%.

where  $F(\beta)$  is the relative fraction of the  $\beta$ -phase;  $A_\alpha$  and  $A_\beta$  are the absorbances at 766 and 840  $\text{cm}^{-1}$ , respectively; and  $K_\alpha$  and  $K_\beta$  correspond to the absorbance coefficients at the respective wavenumbers, which are  $6.1 \times 10^4$  and  $7.7 \times 10^4 \text{ cm}^2 \text{ mol}^{-1}$ , respectively. The  $F(\beta)$  values of the LSNO-1/PVDF composite samples with filler contents of 0, 5, 10, 15, 20, and 25 vol% are 77, 71, 63, 62, 56, and 54%, while  $F(\beta)$  for the LSNO-2/PVDF composite samples are 77, 66, 62, 61, 56, and 61%, respectively. This result shows that the crystallisation of PVDF into the  $\beta$ -phase was not promoted by the addition of the LSNO-1 and LSNO-2 fillers. Generally, incorporation of negative charge fillers into the PVDF polymer can increase the  $\beta$ -phase content.<sup>29</sup> Thus, due to a positive surface charge of the LSNO-1 and LSNO-2 fillers, the increase in  $\beta$ -phase was not active. Nevertheless, due to a high permittivity of the LSNO-1 and LSNO-2 fillers, the dielectric permittivity of the composites is expected to be increased by addition of the LSNO-1 and LSNO-2 particles.

In order to describe the relationships between the conductivities of the LSNO fillers with  $x = 0.2$  (LSNO-1) and  $x = 0.4$  (LSNO-2) and the percolation threshold, the conductivities of the LSNO fillers were estimated using the XPS analysis results to reveal the co-occurrence of  $\text{Ni}^{2+}$  and  $\text{Ni}^{3+}$  in the LSNO powder samples. The substitution of  $\text{La}^{3+}$  by  $\text{Sr}^{2+}$  and increase in the Sr content caused a decrease in the Ni-O1 bond length. This decrease was simultaneous with the cumulative oxidation of  $\text{Ni}^{2+}$  to  $\text{Ni}^{3+}$  and resulted in the enhancement of the electrical conductivity of the LSNO ceramic.<sup>25</sup> Fig. 7(a) and (b) show the XPS profiles of the LSNO-1 and LSNO-2 fillers, respectively. It was found that the  $\text{Ni}^{3+}/\text{Ni}^{2+}$  ratio for LSNO-2 was greater than that for LSNO-1. Therefore, the conductivity of LSNO-2 was higher than that of the LSNO-1 ceramic because of the increase in the charge carriers in the conduction band upon hole injection, which was simultaneous with the Sr chemical doping. Besides a high permittivity of the fillers, the conductivity of the LSNO-1 and LSNO-2 fillers also has a significant influence on the dielectric properties of the composites. A higher conductivity should result in a higher dielectric response in the composites.

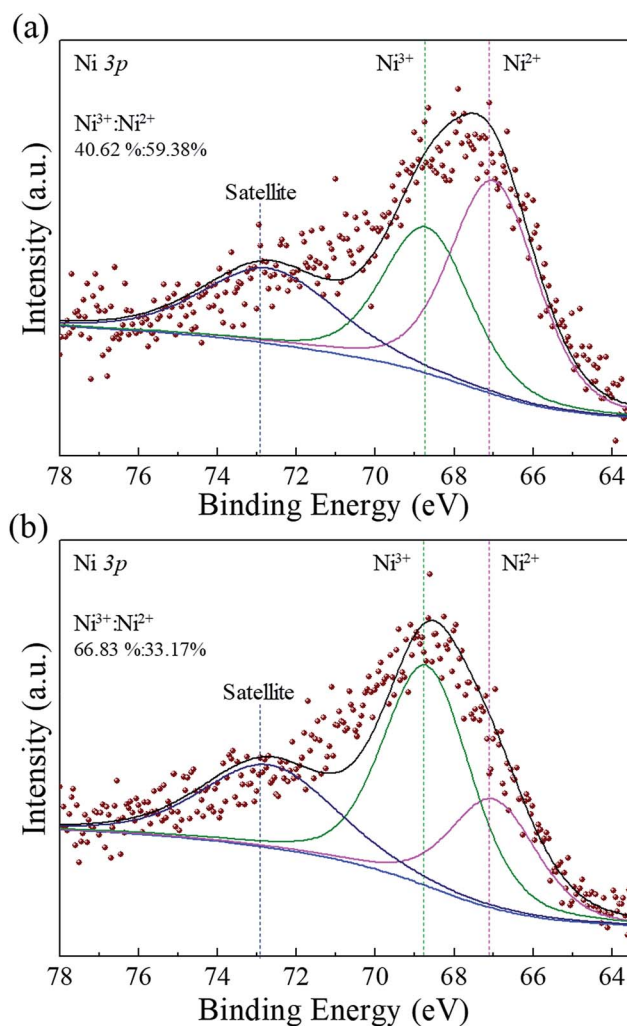


Fig. 7 XPS profiles of (a) LSNO-1 and (b) LSNO-2 fillers prepared via a chemical combustion method and calcined at 1000 °C for 6 h.



The dielectric permittivity ( $\epsilon'$ ), loss tangent ( $\tan \delta$ ), and conductivity ( $\sigma$ ) values at 1 kHz for the LSNO-1/PVDF and LSNO-2/PVDF composites as functions of the filler volume fraction are shown in Fig. 8(a) and (b), respectively. The  $\epsilon'$  values for the 5 vol% LSNO-1, 10 vol% LSNO-1, 15 vol% LSNO-1, 20 vol% LSNO-1, and 25 vol% LSNO-1 composites were 21.2, 34.8, 54.9, 178.0, and 579.6, while the  $\epsilon'$  values for the 5 vol% LSNO-2, 10 vol% LSNO-2, 15 vol% LSNO-2, 20 vol% LSNO-2, and 25 vol% LSNO-2 composites were 25.4, 38.0, 109.1, 1407.2, and 3384.7, respectively. Fig. 8(a) clearly shows that the LSNO-2/PVDF composite had an abrupt change in  $\epsilon'$  when the LSNO-2 concentration became  $\geq 20$  vol%. Furthermore, the  $\tan \delta$  and  $\sigma$  values followed a similar trend with an increase in the volume fraction of LSNO-2, as illustrated in Fig. 8(b). In contrast, both the dielectric properties and conductivity of the LSNO-1/PVDF composite slowly increased over the whole filler loading range studied (0–25 vol%). This could indicate the formation of a percolation network in the LSNO-1/PVDF composite at an LSNO-1 concentration of  $\geq 25$  vol%. Therefore, the percolation threshold ( $f_c$ ) of the LSNO-2/PVDF composite was lower than that of the LSNO-1/PVDF composite. The two kinds of LSNO/PVDF composites exhibited improved and superior  $\epsilon'$  values compared to the pure PVDF ( $\epsilon' = 10$ ), and the  $\epsilon'$  values increased with the addition of the LSNO fillers. These enhanced  $\epsilon'$  values were related to interfacial polarisation or the Maxwell–Wagner–Sillars effect,<sup>4,5,17</sup> which is produced by the accumulation of

charge carriers at semiconductor–insulator interfaces under an applied electric field. As shown by the previously mentioned XPS results, the higher conductivity of LSNO-2 compared to the LSNO-1 filler resulted in the LSNO-2/PVDF composite possessing higher  $\epsilon'$  values for all of the filler concentrations compared to the LSNO-1/PVDF composite. This can be described as strong interfacial polarisation caused by the numerous charge carriers in the LSNO-2 filler.

To understand the effects of the filler content and interface on the dielectric properties of the LSNO/PVDF composites, different filler volume fractions were chosen based on the interparticle distances. The interparticle distances ( $d$ ) in the LSNO-1/PVDF and LSNO-2/PVDF composites can be calculated using the following derived equation:<sup>30</sup>

$$d = r \left[ \left( \frac{4\pi}{3f} \right)^{\frac{1}{3}} - 2 \right] \quad (2)$$

where  $f$  is the volume fraction of the LSNO filler, and  $r = 141$  and  $154$  nm are taken as the average radii of the LSNO-1 and LSNO-2 particles, respectively. The variations in the  $d$  and  $\epsilon'$  values as a function of LSNO filler volume fraction are shown in Fig. 8(d). As the filler loading increased, the particles became compact, and the interparticle distances of the LSNO fillers decreased. At 20 vol% filler, the  $d$  values of the LSNO-1/PVDF and LSNO-2/PVDF composites were 107 and 116 nm, while the  $\epsilon'$  values were 178.0 and 1407.2, respectively. As a result, the dielectric

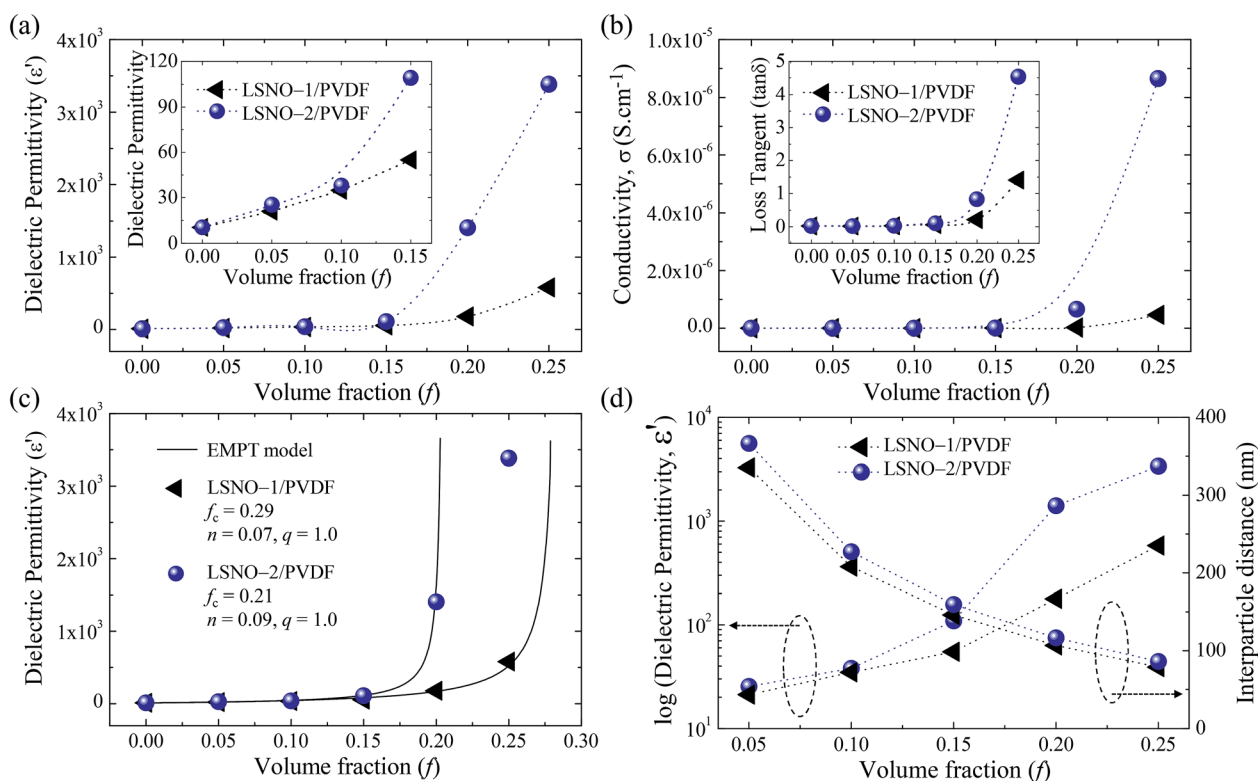


Fig. 8 Variation of (a) dielectric permittivity ( $\epsilon'$ ) and (b) conductivity as function of volume fraction of LSNO filler for LSNO/PVDF composites at 1 kHz and RT. (c) Experimental data for the  $\epsilon'$  values at 1 kHz and RT fitted by the EMPT model. (d) The relationship between the dielectric permittivity ( $\epsilon'$ ) and interparticle distances ( $d$ ) of LSNO fillers in the LSNO/PVDF composites. The insets of (a) and (b) show an enlarged scale in the range of  $f = 0$ – $0.15$  and the dependence of the loss tangent ( $\tan \delta$ ) of the LSNO/PVDF composites on the LSNO volume fraction, respectively.



properties and percolation thresholds of the LSNO/PVDF composites were likely due to the dominant effect of the filler conductivity.

To clarify the dielectric response and dependence of the percolation threshold on the filler conductivity for the LSNO-1/PVDF and LSNO-2/PVDF composites, comparisons were made between the theoretical approximations and experimental data at 1 kHz and RT. Fig. 8(c) shows the calculation results when using the EMPT model,<sup>31–33</sup>

$$\varepsilon_{\text{eff}} = \varepsilon_{\text{PVDF}} \left[ 1 + \frac{f_{\text{LSNO}}(\varepsilon_{\text{LSNO}} - \varepsilon_{\text{PVDF}})}{\varepsilon_{\text{PVDF}} + n(1 - f_{\text{LSNO}})(\varepsilon_{\text{LSNO}} - \varepsilon_{\text{PVDF}})} \right] \left| \frac{f_c - f}{f} \right|^{-q} \quad (3)$$

where  $\varepsilon_{\text{eff}}$  is the effective dielectric permittivity of the composite;  $f_{\text{LSNO}}$  is the filling factor, or the volume fraction of the LSNO filler;  $f_c$  is the percolation threshold; and  $\varepsilon_{\text{PVDF}} = 10.5$  and  $\varepsilon_{\text{LSNO}} = 2.2 \times 10^5$  are taken as the dielectric permittivities of the PVDF matrix and LSNO filler, respectively. Furthermore,  $n$  and  $q$  are the filler morphology fitting factor and a critical exponent of approximately one, respectively. The best fits of the dielectric permittivity data to the EMPT model yield  $f_c = (0.29, 0.21)$ ,  $q = (1.0, 1.0)$ , and  $n = (0.07, 0.09)$  for the LSNO-1/PVDF and LSNO-2/PVDF composites, respectively. The fitted  $n$  and  $q$  values correlate to the nearly spherical shape of the LSNO particles and the

universal value of 1.0, respectively.<sup>31</sup> The fitting results of the EMPT model confirm that the percolation threshold of the LSNO-2/PVDF composite is lower than that of the LSNO-1/PVDF composite. Moreover, the percolation thresholds of the two kinds of LSNO/PVDF composites are higher than that of a spherical-shape conductor-polymer composite ( $f_c = 0.16$ ).<sup>34</sup> This clarifies that the percolation threshold of the LSNO/PVDF composite system depends on the conductivity of the LSNO filler, which decreases as the conductivity of the filler increases, as reported for conductor-polymer composite systems.<sup>30</sup> According to the fitting of the experimental data to the EMPT model, it is likely that the dielectric responses in the LSNO/PVDF composites are influenced by more than one mechanism. The possible mechanisms can be described as follows.

A comparison shows that at a filler concentration of less than 30 vol%, various other ceramic/PVDF composite systems such as BaTiO<sub>3</sub>/PVDF, Ba(Fe<sub>1/2</sub>Nb<sub>1/2</sub>)O<sub>3</sub>/PVDF, CaCu<sub>3</sub>Ti<sub>4</sub>O<sub>12</sub>/PVDF, and (Li + Ti) co-doped NiO/PVDF had  $\varepsilon'$  values that were lower than 50 (at RT and 10<sup>2</sup> Hz).<sup>14,18,35–38</sup> Moreover, the  $\varepsilon'$  values of these composites were still lower than 100 even when the filler concentration was increased to 50 vol%. These results may be attributed to (i) the  $\varepsilon'$  values of the ceramics used in the composites not being high enough (10<sup>3</sup> to 10<sup>4</sup>),<sup>39,40</sup> and (ii) the  $\varepsilon'$  values of these ceramics generally being due to the grain

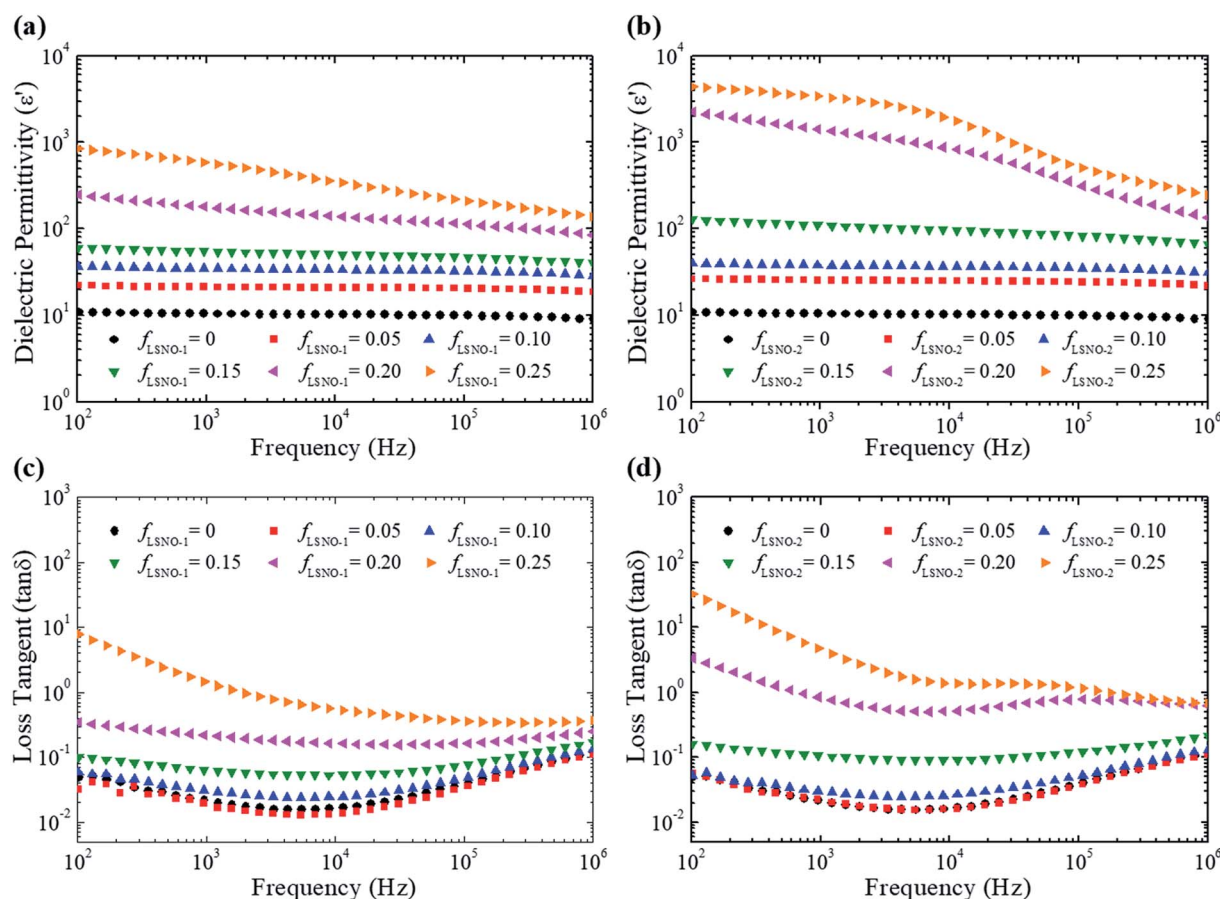


Fig. 9 Frequency dependence of (a and b) dielectric permittivity ( $\varepsilon'$ ) and (c and d) loss tangent ( $\tan \delta$ ) for LSNO-1/PVDF and LSNO-2/PVDF composites, respectively.



boundary response.<sup>41</sup> As previously mentioned, the superior dielectric properties ( $\epsilon' \approx 10^5$  to  $10^6$ ), which could primarily be attributed to intrinsically small polaronic hopping inside the LSNO particles, could have resulted in the higher  $\epsilon'$  values of the LSNO/PVDF composites compared to those of other ceramic/PVDF composites. The  $\epsilon'$  and  $\tan \delta$  values of the LSNO-1/PVDF and LSNO-2/PVDF composites as a function of the filler content were measured over a frequency range of  $10^2$  to  $10^6$  Hz at RT, and the results are shown in Fig. 9(a–d). The  $\epsilon'$  values of both the LSNO-1/PVDF and LSNO-2/PVDF composites increased with the filler content, as seen in Fig. 9(a) and (b), respectively. The  $\epsilon'$  values of the LSNO/composites with the filler contents of 0 (pure PVDF) to 15 vol% were independent of the frequency in the range of  $10^2$  to  $10^5$  Hz, in contrast to the frequency dependence behaviour of the  $\epsilon'$  values for LSNO ceramics in this frequency range.<sup>23</sup> This can be interpreted to mean that the dielectric response of the LSNO/PVDF composites with filler contents below 20 vol% is primarily caused by the intrinsic  $\epsilon'$  values of the LSNO fillers and PVDF matrix. In contrast, at filler contents greater than 15 vol%, a slight decrease in the  $\epsilon'$  value with increasing frequency was observed. Moreover, when the loading of the LSNO filler was increased from 15 to 20 vol%, the  $\epsilon'$  values increase from 54.9 to 178.0 for the LSNO-1/PVDF composite and drastically increased from 109.1 to 1407.2 for the LSNO-2/PVDF composite. This clearly indicated that the effect of interfacial polarisation on the enhanced  $\epsilon'$  values was

dominant, especially for the LSNO-2/PVDF composite, as a result of their percolative nature. Such a drastic increase in the  $\epsilon'$  value for the LSNO-2/PVDF composite could also be rationalised by the micro-capacitor model,<sup>33,42</sup> which considers the consistency of a network of several micro-capacitors formed by a thin layer of PVDF polymer between two neighbouring LSNO-2 particles, where each of these micro-capacitors contributes to an abnormally high capacitance. Fig. 9(c) and (d) show the frequency dependences of the  $\tan \delta$  values of the LSNO-1/PVDF and LSNO-2/PVDF composites, respectively. All of the prepared composites with a filler content below 20 vol% showed similar frequency-dependent behaviours of  $\tan \delta$ . The apparent  $\tan \delta$  peak in the frequency range of  $10^4$  to  $10^5$  Hz corresponding to the step-like decrease in  $\epsilon'$  for the LSNO-2/PVDF composites with filler contents of 20–25 vol% was characteristic of the occurrence of a dielectric relaxation process. The high  $\epsilon'$  and  $\tan \delta$  values in the low-frequency range and their decrease with increasing frequency confirmed the existence of interfacial polarisation.

The temperature dependences of the dielectric properties of the LSNO-1/PVDF and LSNO-2/PVDF composites at 1 kHz as a function of the filler content were measured over a temperature range of  $-60$  °C to  $+100$  °C, and the results are shown in Fig. 10. The  $\epsilon'$  values of the LSNO/PVDF composites increased with the temperature, as seen in Fig. 10(a) and (b). A step increase in the  $\epsilon'$  value and the appearance of a  $\tan \delta$  peak

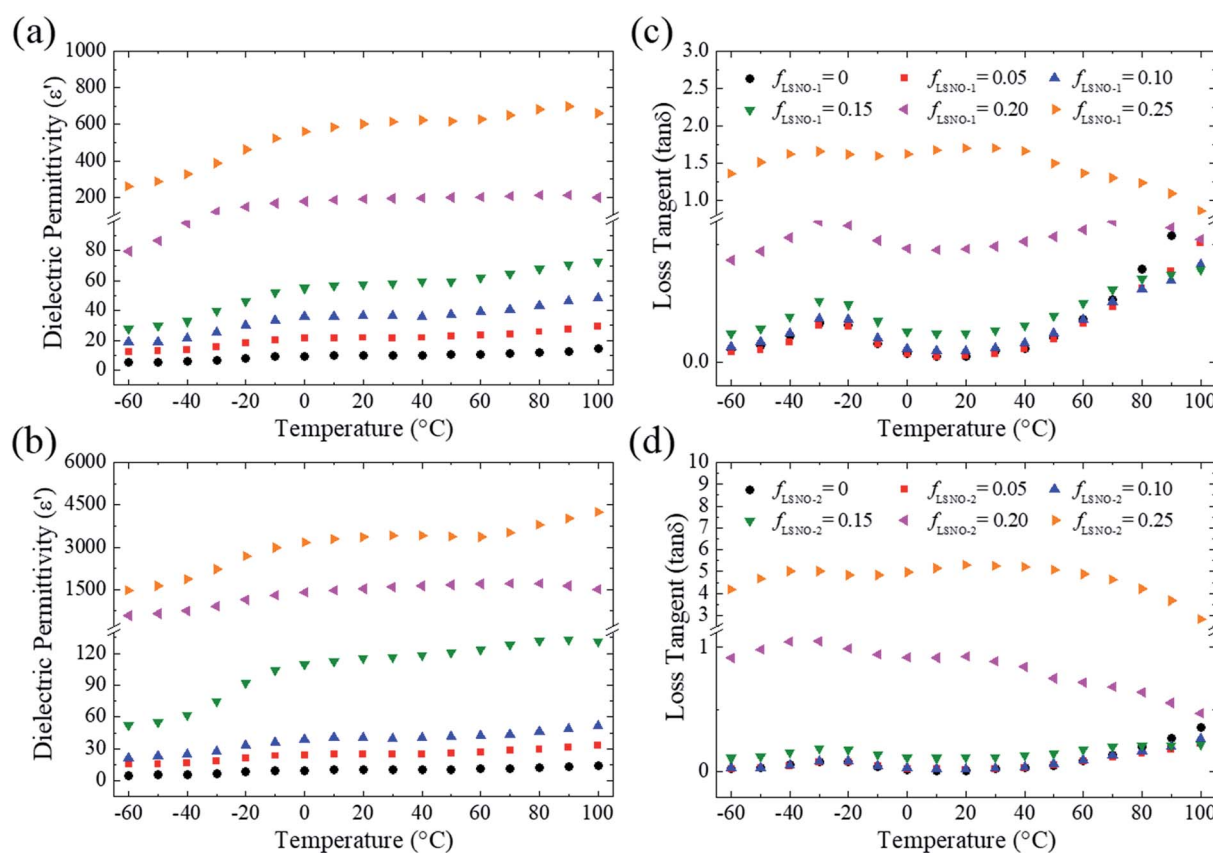


Fig. 10 Temperature dependences of (a and b) dielectric permittivity ( $\epsilon'$ ) and (c and d) loss tangent ( $\tan \delta$ ) at 1 kHz for LSNO-1/PVDF and LSNO-2/PVDF composites, respectively.



(Fig. 10(c) and (d)) in the low-temperature region correspond to the rotational motions of dipolar groups of the PVDF matrix, which result in the  $\beta$  relaxation. The enhanced  $\tan \delta$  values at the high-temperature region for the composites with the filler contents of 5–15 vol% indicate that the dominant polarisation at high temperature is the interfacial polarisation. The temperature dependence behaviour of the  $\tan \delta$  values for the composites with the filler contents of 20–25 vol% reflects the change in the PVDF matrix due to the addition of the LSNO fillers. The crystallisation of the PVDF into the  $\beta$ -phase was expected to decrease with an increase in the LSNO filler, and this was supported by the FT-IR results.

## 4. Conclusion

Promising semiconductor–polymer composites were successfully fabricated using a solution processing method. LSNO-1 and LSNO-2 particles were used as fillers to study the effects of the conductivity of the semiconducting LSNO filler on the dielectric properties and percolation threshold of PVDF-based composites. The results demonstrated that the conductivity of the LSNO filler had a significant impact on the physical properties of the LSNO/PVDF composites. The percolation thresholds of the LSNO/PVDF composites decreases as the conductivity of the LSNO increased. The LSNO-1/PVDF and LSNO-2/PVDF composites showed percolation thresholds of 0.29 and 0.21, respectively. Drastic changes in the dielectric properties and conductivity were observed only in the LSNO-2/PVDF composite, while both the dielectric properties and conductivity of the LSNO-1/PVDF composite increased slowly over the whole filler loading range studied (0–25 vol%). At a filler content of 20 vol%, the  $\epsilon'$  and  $\tan \delta$  values of the LSNO-1/PVDF composite were 178.0 and 0.22, while those of the LSNO-2/PVDF composite were 1407.2 and 0.83, respectively. Nevertheless, the LSNO/PVDF composites with the filler content of 25 vol% exhibited excellent flexibility. Hence, the LSNO ceramic proved to be an important semiconductive filler, and these semiconductor–polymer composites are potential candidates for dielectric charge storage capacitors and electromagnetic wave absorption applications.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

This work was financially supported by the Synchrotron Light Research Institute [Grant No. BRG6180003], Khon Kaen University [Grant No. BRG6180003], and the Thailand Research Fund (TRF) [Grant No. BRG6180003]. This work was partially supported by the Research Network NANOTEC (RNN) program of the National Nanotechnology Center (NANOTEC), NSTDA, Ministry of Science and Technology and Khon Kaen University, Thailand. K. M. would like to thank the Thailand Research Fund

under the Royal Golden Jubilee PhD Program [Grant No. PHD/0191/2556] for his PhD scholarship.

## References

- 1 K. Zou, Y. Dan, Y. Yu, Y. Zhang, Q. Zhang, Y. Lu, H. Huang, X. Zhang and Y. He, *J. Mater. Chem. A*, 2019, **7**, 13473–13482.
- 2 X. Zhang, Y. Shen, B. Xu, Q. Zhang, L. Gu, J. Jiang, J. Ma, Y. Lin and C. W. Nan, *Adv. Mater.*, 2016, **28**, 2055–2061.
- 3 Prateek, D. Singh, N. Singh, A. Garg and R. K. Gupta, *Compos. Sci. Technol.*, 2019, **174**, 158–168.
- 4 F. Liu, Q. Li, J. Cui, Z. Li, G. Yang, Y. Liu, L. Dong, C. Xiong, H. Wang and Q. Wang, *Adv. Funct. Mater.*, 2017, **27**, 1606292.
- 5 M. Yang, C. Hu, H. Zhao, P. Haghi-Ashtiani, D. He, Y. Yang, J. Yuan and J. Bai, *Carbon*, 2018, **132**, 152–156.
- 6 Prateek, V. K. Thakur and R. K. Gupta, *Chem. Rev.*, 2016, **116**, 4260–4317.
- 7 X. Chen, X. Han and Q.-D. Shen, *Adv. Electron. Mater.*, 2017, **3**, 1600460.
- 8 S. Tu, Q. Jiang, X. Zhang and H. N. Alshareef, *ACS Nano*, 2018, **12**, 3369–3377.
- 9 S. Maiti, S. Suin, N. K. Shrivastava and B. B. Khatua, *J. Appl. Polym. Sci.*, 2013, **130**, 543–553.
- 10 L. Zhang, W. Wang, X. Wang, P. Bass and Z. Y. Cheng, *Appl. Phys. Lett.*, 2013, **103**, 232903.
- 11 Z. Chen, H. Li, G. Xie and K. Yang, *RSC Adv.*, 2018, **8**, 1–9.
- 12 L. Zhang, Z. Liu, X. Lu, G. Yang, X. Zhang and Z. Y. Cheng, *Nano Energy*, 2016, **26**, 550–557.
- 13 Z.-M. Dang, C.-W. Nan, D. Xie, Y.-H. Zhang and S. C. Tjong, *Appl. Phys. Lett.*, 2004, **85**, 97.
- 14 Z.-M. Dang, H.-P. Xu, D. Xie and L. Li, *Mater. Lett.*, 2007, **61**, 511–515.
- 15 W. Wu, X. Huang, S. Li, P. Jiang and T. Toshikatsu, *J. Phys. Chem. C*, 2012, **116**, 24887–24895.
- 16 C. Dagdeviren and M. Papila, *Polym. Compos.*, 2009, **31**, 1003–1010.
- 17 Y. Zhang, Y. Wang, Y. Deng, M. Li and J. Bai, *ACS Appl. Mater. Interfaces*, 2012, **4**, 65–68.
- 18 H. Tang, Z. Zhou, C. C. Bowland and H. A. Sodano, *Nano Energy*, 2015, **17**, 302–307.
- 19 Z. Zhang, Y. Gu, J. Bi, S. Wang, M. Li and Z. Zhang, *Mater. Lett.*, 2015, **160**, 16–19.
- 20 S. Krohns, P. Lunkenheimer, C. Kant, A. V. Pronin, H. B. Brom, A. A. Nugroho, M. Diantoro and A. Loidl, *Appl. Phys. Lett.*, 2009, **94**, 122903.
- 21 S. Krohns, P. Lunkenheimer, S. Meissner, A. Reller, B. Gleich, A. Rathgeber, T. Gaugler, H. U. Buhl, D. C. Sinclair and A. Loidl, *Nat. Mater.*, 2011, **10**, 899–901.
- 22 X. Q. Liu, S. Y. Wu and X. M. Chen, *J. Alloys Compd.*, 2010, **507**, 230–235.
- 23 K. Meeporn, N. Chanlek and P. Thongbai, *RSC Adv.*, 2016, **6**, 91377–91385.
- 24 K. Meeporn, T. Yamwong and P. Thongbai, *Jpn. J. Appl. Phys.*, 2014, **53**, 06JF01.
- 25 A. Aguadero, M. J. Escudero, M. Perez, J. A. Alonso, V. Pomjakushin and L. Daza, *Dalton Trans.*, 2006, 4377–4383, DOI: 10.1039/b606316k.



- 26 K. Meeporn, S. Maensiri and P. Thongbai, *Appl. Surf. Sci.*, 2016, **380**, 67–72.
- 27 B. Luo, X. Wang, Y. Wang and L. Li, *J. Mater. Chem. A*, 2014, **2**, 510–519.
- 28 P. Martins, A. C. Lopes and S. Lanceros-Mendez, *Prog. Polym. Sci.*, 2014, **39**, 683–706.
- 29 C. Ribeiro, C. M. Costa, D. M. Correia, J. Nunes-Pereira, J. Oliveira, P. Martins, R. Gonçalves, V. F. Cardoso and S. Lanceros-Méndez, *Nat. Protoc.*, 2018, **13**, 681.
- 30 K. S. Deepa, S. Kumari Nisha, P. Parameswaran, M. T. Sebastian and J. James, *Appl. Phys. Lett.*, 2009, **94**, 142902.
- 31 F. Fang, W. Yang, S. Yu, S. Luo and R. Sun, *Appl. Phys. Lett.*, 2014, **104**, 132909.
- 32 X. Chen, F. Liang, W. Lu, Z. Jin, Y. Zhao and M. Fu, *Polymers*, 2018, **10**, 1–10.
- 33 K. Meeporn, P. Thongbai, T. Yamwong and S. Maensiri, *RSC Adv.*, 2017, **7**, 17128–17136.
- 34 C. W. Nan, Y. Shen and J. Ma, *Annu. Rev. Mater. Res.*, 2010, **40**, 131–151.
- 35 W. Wan, J. Luo, C.-e. Huang, J. Yang, Y. Feng, W.-X. Yuan, Y. Ouyang, D. Chen and T. Qiu, *Ceram. Int.*, 2018, **44**, 5086–5092.
- 36 S. Luo, S. Yu, R. Sun and C. P. Wong, *ACS Appl. Mater. Interfaces*, 2014, **6**, 176–182.
- 37 Z. Wang, M. Fang, H. Li, Y. Wen, C. Wang and Y. Pu, *Compos. Sci. Technol.*, 2015, **117**, 410–416.
- 38 T. Zhou, J. W. Zha, R. Y. Cui, B. H. Fan, J. K. Yuan and Z. M. Dang, *ACS Appl. Mater. Interfaces*, 2011, **3**, 2184–2188.
- 39 P. Thomas, K. T. Varughese, K. Dwarakanath and K. B. R. Varma, *Compos. Sci. Technol.*, 2010, **70**, 539–545.
- 40 A. R. West, T. B. Adams, F. D. Morrison and D. C. Sinclair, *J. Eur. Ceram. Soc.*, 2004, **24**, 1439–1448.
- 41 X. Sun, J. Deng, S. Liu, T. Yan, B. Peng, W. Jia, Z. Mei, H. Su, L. Fang and L. Liu, *Appl. Phys. A: Mater. Sci. Process.*, 2016, **122**, 864.
- 42 P. W. Jaschin, R. Bhimireddi and K. B. R. Varma, *ACS Appl. Mater. Interfaces*, 2018, **10**, 27278–27286.

