Nanoscale Advances



PAPER

View Article Online
View Journal | View Issue



Cite this: Nanoscale Adv., 2020, 2, 5396

Received 28th July 2020 Accepted 4th October 2020

DOI: 10.1039/d0na00613k

rsc.li/nanoscale-advances

The dual interfacial modification of 2D g-C₃N₄ for high-efficiency and stable planar perovskite solar cells

Zhou Liu, ab Shuzhen Wu, ab Xiaojie Yang, ab Yijun Zhou, ab Jiaren Jin, ab Junmei Sun, c Li Zhao ab and Shimin Wang*ab

Carrier recombination and charge loss at the interfaces of perovskite layers have a significant influence on high-performance planar perovskite solar cells (PSCs). We employed two-dimensional graphitic carbon nitride (g- C_3N_4), which is a heat-resistant n-type semiconductor, to modify the electron-transport layer/perovskite and perovskite/hole-transport layer interfaces, respectively. g- C_3N_4 could passivate the surface trap states of the methylammonium lead iodide light absorber through the formation of a Lewis adduct between N and the under-coordinated Pb, and it could also remarkably reduce the grain boundaries between perovskite crystal particles. A maximum power conversion efficiency (PCE) of 19.67% ($V_{oc} = 1.14 \text{ V}$, $J_{sc} = 21.45 \text{ mA cm}^{-2}$, FF = 0.807) could be obtained from planar PSCs with long-term stability using dual-positioned g- C_3N_4 . Therefore, we consider that ultrathin semiconductor films with a Lewis base nature are suitable as dual-functional transport materials for devices. This work provides new guidance for dual-interfacial modification to improve the PCE and stability of devices.

1. Introduction

Organic–inorganic lead halide perovskite solar cells (PSCs) have attracted extensive attention because of their outstanding photovoltaic and optical properties, such as long charge carrier diffusion lengths, strong optical absorption, and low electronhole recombination rates. ¹⁻⁶ Therefore, the power conversion efficiency (PCE) of these devices has increased from 3.9% to 25.2% over the past few years. ⁷⁻¹¹ However, trap states have a great influence on PSC performance, which limits the ability of PSCs to obtain theoretical PCE of >30%. ¹²

Currently, four main defects exist on the surface of perovskite layer: (1) halide vacancies bringing about uncoordinated Pb²⁺ defects (2); Pb–I anti-site or halide excess; (3) cation vacancies (FA⁺ and MA⁺); (4) metallic lead.^{13–22} Energy disorder and obstruction of charge extraction are caused by the surface of perovskite layer defects as interfacial non-radiative recombination centers, which will result in a lower fill factor (FF), smaller open-circuit voltage (V_{oc}), and a sharply decreased PCE.^{23–25} Hence, interfacial modification is helpful to optimize perovskite films because it alters the interface of the perovskite

In this work, we used g-C₃N₄ to modify the interfaces of ETL/ perovskite and perovskite/HTL, respectively. A several layer-thick or monolayer-thick g-C₃N₄ have been shown to be outstanding two-dimensional (2D) nanomaterials in optoelectronics and electronics thanks to their excellent tunable optoelectronic properties and high chemical and thermal stability.^{37–43} Although g-C₃N₄ modification does not exhibit desirable band alignments among ETL, HTL, and MAPbI₃ perovskite absorbers at the two interfaces, an observably increased PCE of 19.69% for g-C₃N₄ dual-incorporated PSCs

layer to reduce defects.26 For example, Mohamadkhani et al. reported planar PSCs with CdS modification at the perovskite/ SnO₂ electron transport layer (ETL) interface, which achieved a PCE of 17.18%.27 Luo et al. showed devices with interfacial modification by thiourea to obtain a PCE of 19.18%.28 However, the complex defects in polycrystalline perovskite films (grain boundaries, pinholes, and crystal defects) are inevitable at the perovskite/hole transport layer (HTL) interface. Li et al. studied devices grown with the polymer poly(4-vinylpyridine) (PVP) at the perovskite/HTL interface, and an average PCE of 18.95% was obtained.29 Wang et al. introduced planar PSCs based on a multifunctional ammonium salt at the perovskite/HTL interface to yield an optimum PCE of 18.95% and a steady-state output PCE of 18.11%.30 Nevertheless, these methods solve only one problem, which is not conducive to gaining higher FF and short-circuit current density $(I_{sc})^{31-35}$ Thus, it is necessary to find a material to modify at ETL/perovskite and perovskite/HTL interfaces.36 Dual interfacial modification is conducive to reducing defects while passivating defects.

^aHubei Collaborative Innovation Center for Advanced Organic Chemical Materials, Wuhan 430062, PR China. E-mail: zhaoli7376@163.com

^bMinistry-of-Education Key Laboratory for the Green Preparation and Application of Functional Materials, Faculty of Materials Science and Engineering, Hubei University, Wuhan 430062, P. R. China

^cCollege of Pharmacy and Biological Engineering, Chengdu University, Chengdu 610106, P. R. China

over the pristine device (PCE = 18.03%) has been obtained.⁴⁴ This result is principally because g-C₃N₄ can reduce trap density dramatically at ETL/perovskite and perovskite/HTL interfaces.

2. Experimental section

2.1 Materials

Paper

Perovskite (PbI₂, MAI) and HTL ((spiro-MeOTAD, lithium-bis(tri-fluoromethanesulfonyl)imide (Li-TFSI), and (4-tert-butylpyridine (tBP)) solutions were purchased from Xi'an Polymer Light Technology. Dimethyl sulfoxide (DMSO) (99.9%) and 4-hydroxybutyric acid lactone (DMF) (99.9%) were purchased from Sigma-Aldrich and Aladdin, respectively. SnCl₂·2H₂O (98%) and urea (99.9%) were purchased from Alfa Aesar.

2.2 Device fabrication

Preparation of g-C₃N₄. Urea was heated from 20 °C to 550 °C at a heating rate of 5 °C min⁻¹. Then, the light-yellow solid powder obtained was dispersed in isopropanol. Finally, a g-C₃N₄ precursor was obtained by ultrasonication for \sim 1 h.

The F-doped SnO₂ substrate was cleaned by using deionized water, isopropyl alcohol, ethanol, and isopropyl alcohol, respectively. The F-doped SnO2 substrate was irradiated by UVozone for 15 min before deposition of SnO₂ ETLs. The solution of SnO₂ precursor was spin-coated on glass at 4000 rpm for 30 s, and then annealed for 10 min at 180 °C. The g-C₃N₄ precursor was spin-coated on an SnO₂ film at 4000 rpm for 30 s, and then annealed for 5 min at 100 °C. A MAPbI₃ perovskite solution (1.3 M) was spin-coated on top of the g-C₃N₄ layer at 1000 rpm for 6 s and 4000 rpm for 30 s. During spinning, 100 µL of chlorobenzene was dropped onto the spinning perovskite film for 16 s. Then, the films were annealed for 10 min at 100 °C. The g-C₃N₄ precursor was spin-coated on top of the perovskite layer at 4000 rpm for 30 s, and then annealed for 5 min at 100 °C. Then, 20 μL of spiro-OMeTAD solution containing 36.1 mg of spiro-OMeTAD, 14.4 µL of t-BP, 9 µL of Li-TFSI solution (520 mg in acetonitrile) and 500 µL of chlorobenzene was spin-coated on the g-C₃N₄ layer at 4000 rpm for 30 s. Finally, Au of thickness 50 nm was deposited by thermal evaporation.

2.3 Device characterization

Absorbance was recorded with a UV-vis spectrophotometer (UV-3600; Shimadzu). The X-ray diffraction (XRD) patterns of samples and perovskite films were measured with an AXS D8 advance system (Bruker). The morphology of MAPbI₃ perovskite films was measured with a scanning electron microscope (500 series; Sigma). g-C₃N₄ particles were characterized *via* transmission electron microscopy (TEM) using a G2 F20 system (Tecnai). The incident photo-to-electron conversion efficiency (IPCE) of PSCs was analyzed using a 300 W xenon lamp (66984; Newport). The photocurrent–voltage (*J–V*) characteristics of PSCs were analyzed under simulated AM 1.5 G radiation (irradiance of 100 mW cm⁻²) using a solar simulator (91192-1000; Oriel) and a source meter (2400; Keithley). Electrochemical impedance spectroscopy (EIS) was employed using an electrochemical workstation (IM6; Zennium) over a frequency range of

100 mHz to 2 MHz with 10 mV AC amplitude at -1 V bias under simulated AM 1.5 G radiation (irradiance of 100 mW cm $^{-2}$). Devices were measured under ambient conditions [15% < relative humidity (RH) < 60%] each time (winter or summer). After measurements had been taken, devices were stored in a humidity-controlled dry room (15% < RH < 40%).

Results and discussion

A schematic diagram of g-C₃N₄ dual-modified PSCs is displayed as Fig. 1. We fabricated a typical planar PSC configuration of FTO/SnO₂/MAPbI₃/spiro-OMeTAD/Au. The other layers were coated by spin-coating except that Au was deposited by thermal evaporation. The upper and lower layers of perovskite were modified by multilayer g-C₃N₄. Lone-pair electrons on g-C₃N₄ could usefully coordinate unreacted lead ions at the perovskite/HTL interface, which could reduce the surface defects of the perovskite layer and improve the photoelectric performance of devices. Meanwhile, g-C₃N₄ could reduce the Gibbs free energy of the SnO₂ surface to serve as nucleation centers for perovskite growth and promote the nucleation process of perovskites at the ETL/perovskite interface. The porous structure of g-C₃N₄ also strengthened the contact between the perovskite layer and HTL and ETL, respectively.

The XRD pattern of samples obtained after urea had been treated at high temperature is shown in Fig. 2(a). The intense diffraction peak at 27.5° was assigned to the (002) crystal plane of g-C₃N₄ (JCPDS-87-1526).⁴⁵ g-C₃N₄ was a typical polymer semiconductor in which the C-N atoms were sp² hybrid orbitals to form a highly delocalized conjugated system and electrondonor group. Covalent bonds were formed by the 6p empty orbit in under-coordinated Pb2+ and the lone pair electrons on C-N atoms, which helped to reduce the surface defects of the perovskite layer. A lower concentration of defects could increase the transportation and extraction of charge carriers. Fourier transform infrared (FTIR) spectroscopy of g-C₃N₄ was used to analyze absorption bonds. Peaks at ~ 3363 cm⁻¹ and \sim 3362 cm⁻¹ represent -NH₂ bending vibrations (Fig. 2(b)). Peaks at \sim 800 cm⁻¹ and 1200–1650 cm⁻¹ represent the stretching vibration of the triazine units and C-N heterocycle. g-C₃N₄ had several N-H groups, which could passivate the defects of the perovskite layer effectively and improve the PCE of PSCs. The optical changes of MAPbI3 perovskite films with and without g-C₃N₄ dual-positioned at both interfaces were evaluated (Fig. 2(c)). Dual interfacial modification (g-C₃N₄) could help the perovskite films to gain absorbance in the absorption

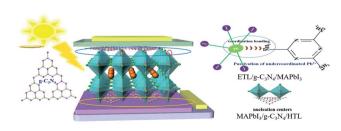


Fig. 1 g-C₃N₄ dual-modified PSCs (schematic diagram).

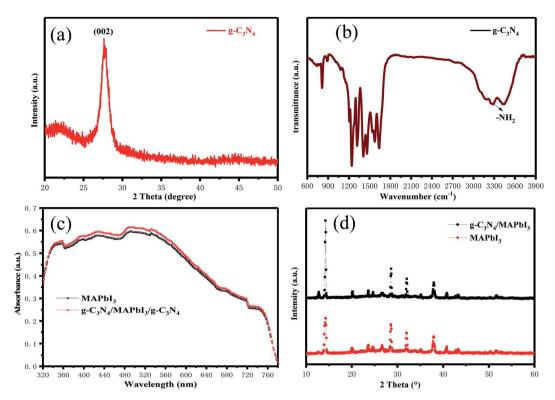


Fig. 2 (a) The XRD pattern of $g-C_3N_4$ film. (b) The FTIR spectrum of $g-C_3N_4$ film. (c) UV-vis absorption spectra of $SnO_2/MAPbI_3$ and $SnO_2/g-C_3N_4/MAPbI_3/g-C_3N_4$ films. (d) XRD patterns of MAPbI $_3$ and $g-C_3N_4/MAPbI_3$ films.

range 350–750 nm. This result showed that the perovskite layer had good crystallinity, which would enhance $J_{\rm sc}$ of PSCs. We also obtained the XRD of the samples of SnO₂ ETLs interface-modified with and without g-C₃N₄ (Fig. 2(d)). The characteristic peaks of MAPbI₃ at 14.5° and 28.5° corresponded to the

(110) and (220) crystal planes, respectively. The enhanced diffraction peak intensities showed that perovskite crystallization improved after g-C₃N₄ modified SnO₂ ETLs.

Fig. 3(a) shows the TEM image of pure $g\text{-}C_3N_4$ particles: a multi-layer sheet-like morphology with smooth surface was

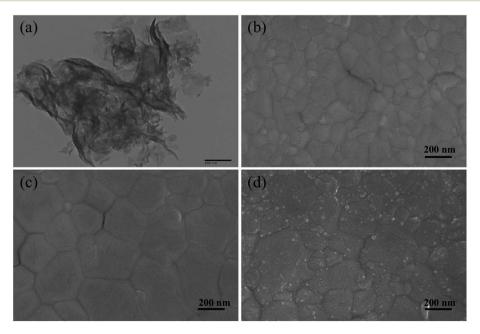


Fig. 3 (a) A TEM image of pure $g-C_3N_4$. (b-d) Top-view SEM images of planar PSCs with and without $g-C_3N_4$ -modified SnO₂ ETLs and a planar PSC with $g-C_3N_4$ simultaneously modifying the SnO₂ ETL and perovskite-layer (50 000×).

observed.46 We adopted dynamic spin-coating during the experiment, which was good for thinner and smoother g-C₃N₄ films. The top-view SEM images exhibited in Fig. 3(b) and (c) demonstrate that MAPbI₃ perovskite film grown on SnO₂/g-C₃N₄ exhibited a larger average particle size than that of perovskite on pristine SnO₂. The enlarged grains could be attributed to g-C₃N₄ with hydrophobic and low surface-restraining perovskite nucleation. SEM revealed that g-C₃N₄ obviously reduce the number of surface defects and grain boundaries. The size of layer-built g-C₃N₄ materials covered on perovskite films was from several nanometers to ~10 nm (Fig. 3(d)). Defects of perovskite grain boundaries are usually caused by loss of intergranular atoms, disordered arrangement or bond distortion, which leads to accumulation, serious trapping, and recombination of charge carriers. Small-scale g-C₂N₄ nanoparticles can effectively passivate defects at grain boundaries of the MAPbI₃ film surface by forming strong hydrogen bonds between MA⁺ and N atoms, which enhances charge extraction and suppresses carrier recombination.

Fig. 4(a) demonstrates the photovoltaic characteristics of typical PSCs (0.06 cm⁻²) and one with a g-C₃N₄-modified MAPbI₃ perovskite absorber at the two interfaces under AM 1.5 G irradiation at 100 mW cm⁻². The control device showed a PCE of 18.03% with $V_{\rm oc}$ of 1.12 V, $J_{\rm sc}$ of 20.3 mA cm⁻² and FF of 0.793. By contrast, the champion device with g-C₃N₄ dual positions exhibited a $J_{\rm sc}$ of 21.45 mA cm $^{-2}$, $V_{\rm oc}$ of 1.14 V, and FF of 0.807 and, thus, a much higher PCE of 19.67%. Dual-interface treatment led to an increase in $V_{\rm oc}$ from 1.12 to 1.14 V, of $J_{\rm sc}$ from $20.30 \text{ to } 21.45 \text{ mA cm}^{-2}$, and of FF from 0.793 to 0.807. The increase in Voc and FF could have been because g-C₃N₄ modified at the perovskite/HTL interface reduced the surface defects of perovskite films. g-C₃N₄ modified at the ETL/perovskite interface increased the absorption of the perovskite layer, which resulted in an increase in J_{sc} . The external quantum efficiency (EQE) spectra of devices and the one with g-C₃N₄ dualpositioned at both interfaces are sketched in Fig. 4(b). The EQE responses of the two champion devices exhibited high photon current conversion efficiency (75-90%). The integrated current value counted from the IPCE spectra was 20.59 and 21.17 mA cm⁻², respectively. Compared with the pristine device,

the IPCE spectra of the devices with g- C_3N_4 dual-incorporation increased significantly. These consequences were consistent with the corresponding best short-circuit current density, which demonstrated that the emission of our solar simulator was in accordance with the spectrum. IPCE = LHE $\times \varphi_{\rm inj} \times \eta_{\rm cc}$, where LHE is the light-harvesting efficiency, $\varphi_{\rm inj}$ is the electroninjection efficiency, and $\eta_{\rm cc}$ is the charge-collection efficiency. Hence, a better device with g- C_3N_4 dual-incorporation will have stronger light absorption and charge transport.

Devices with and without g-C₃N₄ dual-positioned at both interfaces and statistical analyses of the key photovoltaic parameters (PCE, FF, V_{oc} , and J_{sc}) are shown in the inset of Fig. 5(a)-(d). Pristine PSCs displayed a mean PCE of 17.46 \pm 0.57%, mean $V_{\rm oc}$ of 1.11 \pm 0.02 V, mean $J_{\rm sc}$ of 20.33 \pm 0.32 mA cm⁻², and mean FF of 0.772 \pm 0.021. The photovoltaic parameters of the g-C₃N₄ dual-modified PSCs were all higher than those of pristine devices because the PSCs could bring about an enhanced mean PCE of 19.1 \pm 0.68%, mean V_{oc} of 1.13 \pm 0.02 V, mean $J_{\rm sc}$ of 21.31 \pm 0.32 mA cm⁻², and mean FF of 0.793 \pm 0.024. The photovoltaic parameters of the devices and the one with g-C₃N₄ dual modifications are summarized in Table 1. Based on the analyses mentioned above, we think that ultrathin semiconductor films based on the property of Lewis bases are suitable for dual-modification at the two interfaces of the perovskite layer. Consequently, the device with g-C₃N₄ dualincorporation at both interfaces performed best.

A photoluminescence (PL) quenching experiment was used to analyze charge-transfer kinetics. More efficient PL-quenching of perovskite films upon g- C_3N_4 dual-incorporation than that of pristine perovskite films are depicted in Fig. 6(a). PL quenching was due to enhancement in extraction of charge carriers across the interface at the ETL/perovskite and perovskite/HTL interfaces with g- C_3N_4 dual positions. Enhanced interfacial contact conductivity can lead to a reduction in charge accumulation and can improve the injection and transportation of electrons. The charge-transport behavior of devices with and without g- C_3N_4 dual-incorporation was studied further using time-resolved photoluminescence (TR-PL) spectra (Fig. 6(b)). A bi-exponential decay function could fit the PL decay time, and corresponded to a slow radiative decay and a fast radiative decay

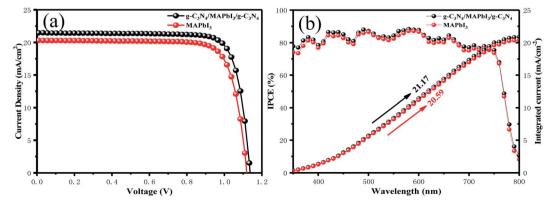


Fig. 4 (a) J-V curves of a pristine planar PSC and a planar PSC with $g-C_3N_4$ simultaneously modifying the SnO₂ ETL and perovskite layer. (b) IPCE spectra of a pristine device and a device with $g-C_3N_4$ -modified interfaces (ETL/ $g-C_3N_4$ /perovskite/ $g-C_3N_4$ /spiro-OMeTAD).

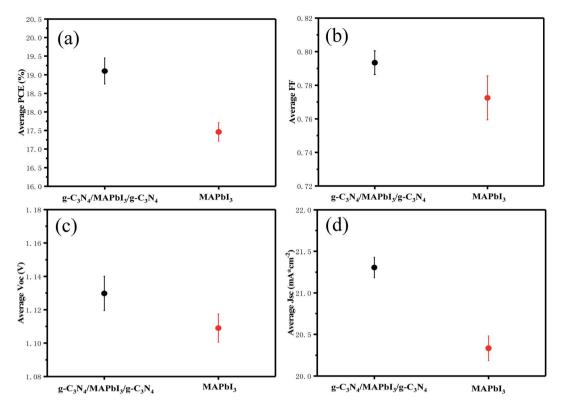


Fig. 5 Statistical analyses of (a) PCE, (b) FF, (c) V_{oc} , and (d) J_{sc} efficiency values, using 10 devices to calculate each value.

 $\textbf{Table 1} \quad \text{Summary of photovoltaic parameters of pristine planar PSCs and the one with $g-C_3N_4$-modified SnO_2 ETL and the perovskite layer and $g-C_3N_4$-modified SnO_2 ETL and $g-C_3N_4$-modified $g-C_$

PSC	$V_{\rm oc}$ (V)	$J_{ m sc}$ (mA cm $^{-2}$)	FF	PCE (%)	Highest PCE (%)
Standard Modified	$\begin{array}{c} \textbf{1.11} \pm \textbf{0.02} \\ \textbf{1.13} \pm \textbf{0.02} \end{array}$	$\begin{array}{c} 20.33 \pm 0.32 \\ 21.31 \pm 0.32 \end{array}$	$\begin{array}{c} 0.772 \pm 0.021 \\ 0.793 \pm 0.024 \end{array}$	$17.46 \pm 0.57 \\ 19.1 \pm 0.68$	18.03 19.67

in perovskite materials. Table 2 lists all the important parameters. For the pristine sample, the fast decay time (τ_1) was 116.9 ns, and the slow decay time (τ_2) was 31.47 ns, with an amplitude $\tau_{\rm ave}$ ($\tau_{\rm ave} = \Sigma A_i \tau_i^2 / \Sigma A_i \tau_i$, where A_1 and A_2 are pre-exponential factors) of 56.06 ns. Compared with the pristine perovskite, τ_1

was 99.5 ns, and τ_2 was 28.01 ns, and resulted in an amplitude $\tau_{\rm ave}$ of 46 ns. The average lifetime of the g-C₃N₄ dual-incorporated device decreased sharply, which is beneficial for suppressing charge recombination and increasing the number of photon-induced electrons and hole transfer at the ETL/

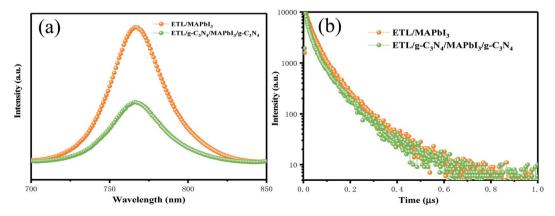


Fig. 6 (a) Steady-state PL spectra of PSCs with and without $g-C_3N_4$ dual-incorporation. (b) Time-resolved PL spectra of PSCs with and without $g-C_3N_4$ dual-incorporation.

Table 2 Parameters from the time-resolved PL spectra

PSC	τ_1 (ns)	τ_2 (ns)	τ _{ave} (ns)
$SnO_2/MAPbI_3 \\ SnO_2/g\text{-}C_3N_4/MAPbI_3/g\text{-}C_3N_4$	116.9	31.47	56.06
	99.5	28.01	46

perovskite and perovskite/HTL interfaces, respectively. Hence, FF and $J_{\rm sc}$ were enhanced markedly, a finding that was consistent with the J–V and IPCE (Fig. 4(a) and (b)) measurement.

For the interpretation of the effect of g-C₃N₄ dual-modified devices on charge recombination, we measured EIS within the frequency range of 10 mHz to 2 MHz at −1 V bias under dark conditions. The Nyquist plots of the devices with and without g-C₃N₄ dual-incorporation are shown in Fig. 7. The latter shows the equivalent circuit containing a recombination resistance $(R_{\rm rec})$, chemical capacitance (C), and series resistance $(R_{\rm s})$. A smaller R_s of 11.31 Ω for the device with g-C₃N₄ dual-positioned than that of the pristine device $(R_s = 15.26 \Omega)$ indicated enhanced conductivity. The devices had a much lower R_s with a relatively higher FF. A larger $R_{\rm rec}$ of 716.2 Ω for the device with g-C₃N₄ dual modifications than that of the pristine device (R_{rec}) = 413.8 Ω) showed more efficient charge dissociation and transport. As a consequence, the boosted charge transfer relieved the loss in interfacial charge, thereby contributing to increase V_{oc} ($V_{\text{oc}} = E_g/q - V_{\text{loss}}$) and J_{sc} .

Fig. 8 shows the steady-state output at the maximum power point (MPPT) under the constant bias voltages and AM 1.5 G illumination measured. Champion devices with and without g-C₃N₄ dual-incorporation were biased at 0.97 and 0.89 V, respectively. In a continuous measurement for 200 s, the g-C₃N₄ dual-incorporated device achieved a stabilized efficiency of 19.2%, which was superior to that of the pristine device (17.32%). High values of MPPT and biasing voltage could have been because g-C₃N₄ reduced the surface defects of perovskite films while passivating defects. The surface defects of the perovskite layer and trapped charge can give rise to the migration of iodide and ion mobilization within perovskite lattices,

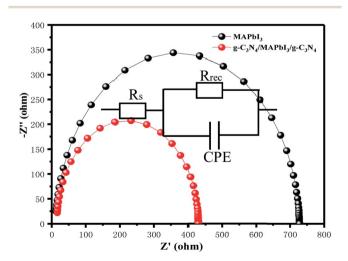


Fig. 7 Nyquist plots of PSCs with and without $g-C_3N_4$ -modified interfaces (ETL/ $g-C_3N_4$ /perovskite/ $g-C_3N_4$ /spiro-OMeTAD).

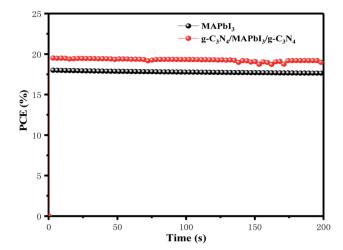


Fig. 8 Steady-state efficiencies of PSCs with and without $g-C_3N_4$ dual-incorporation measured at maximum power output.

which eventually affect PSC stability. Therefore, the highly stable PCE of PSCs explained how the g-C₃N₄ dual interface reduced defects significantly.

4. Conclusions

We demonstrated a simple and effective method to fabricate n-type g-C₃N₄ ultrathin films at HTL/perovskite and ETL/perovskite interfaces in planar PSCs to obtain high efficiency and long-term stability. FF, $V_{\rm oc}$, and $J_{\rm sc}$ values were enhanced to 0.807, 1.14 V, and 21.45 mA cm⁻², respectively, and enabled the PCE to increase to 19.67%. The better performance of the planar PSCs after g-C₃N₄ dual-incorporation was attributed mainly to two reasons: (i) g-C₃N₄ dual-incorporation enhanced the absorption of visible light by the perovskite layer; and (ii) the device with g-C₃N₄ dual-incorporation had a smaller $R_{\rm s}$, smaller $\tau_{\rm ave}$, and larger $R_{\rm rec}$. This phenomenon helped to accelerate the charge extraction and transfer rates. Furthermore, ultrathin semiconductor films with Lewis base character may effectively passivate defects at the dual interfaces of the perovskite layer to achieve better PSC photovoltaic performance and long-term stability.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This research was financially supported by the Key Program for Intergovernmental S&T Innovation Cooperation Projects of National Key R&D Program of China (2019YFE0107100), National Natural Science Foundation of China (5170021087), and China Postdoctoral Science.

References

1 Y. Wu, X. Li, S. Fu, L. Wan and J. Fang, *J. Mater. Chem. A.*, 2019, **14**, 8078–8084.

- 2 G. C. Xing, N. Mathews, S. Y. Sun, S. S. Lim, Y. M. Lam, M. Gratzel, S. Mhaisalkar and T. C. Sum, Science, 2013, **6156**, 344-347.
- 3 S. D. Stranks, G. E. Eperon, G. Grancini, C. Menelaou, M. J. P. Alcocer, T. Leijtens, L. M. Herz, A. Petrozza and H. J. Snaith, Science, 2013, 6156, 341-344.
- 4 M. Z. Liu, M. B. Johnston and H. J. Snaith, Nature, 2013, 7467, 395.
- 5 M. Saliba, T. Matsui, K. Domanski, J. Y. Seo, A. Ummadisingu, S. M. Zakeeruddin, J. P. Correa-Baena, W. R. Tress, A. Abate, A. Hagfeldt and M. Gratzel, Science, 2016, 63, 206-209.
- 6 Y. Li, L. Li, A. S. Yerramilli, Y. Chen, D. Fang, Y. Shen and T. L. Alford, Org. Electron., 2019, 73, 130-136.
- 7 X. Qi, G. Liu, D. Wang, N. Zhu, Y. Zhang, Z. Zhang, C. Wu, X. Li, W. Luo, Y. Li, H. Hu, Z. Chen, L. Xiao, S. Wang and B. Qu, Org. Electron., 2019, 74, 52-58.
- 8 H. Liu, Z. Chen, H. Wang, F. Ye, J. Ma, X. Zheng, P. Gui, L. Xiong, J. Wen and G. Fang, J. Mater. Chem. A, 2019, 17, 10636-10643.
- 9 C. Chen, Z. Song, C. Xiao, D. Zhao, N. Shrestha, C. Li, G. Yang, F. Yao, X. Zheng, R. J. Ellingson, C.-S. Jiang, M. Al-Jassim, K. Zhu, G. Fang and Y. Yan, Nano Energy, 2019, 4, 141-147.
- 10 G. Yang, H. Zhang, G. Li and G. Fang, Nano Energy, 2019, 63, 103835.
- 11 L. Wang, C. McCleese, A. Kovalsky, Y. Zhao and C. Burda, J. Am. Chem. Soc., 2014, 35, 12205-12213.
- 12 C. Chen, Y. Jiang, J. Guo, X. Wu, W. Zhang, S. Wu, X. Gao, X. Hu, Q. Wang, G. Zhou, Y. Chen, J. M. Liu, K. Kempa and J. Gao, Adv. Energy Mater., 2019, 30, 1900557.
- 13 A. Klasen, P. Baumli, Q. Sheng, E. Johannes, S. A. Bretschneider, I. M. Hermes, V. W. Bergmann, C. Gort, A. Axt, H. Kim, H. J. Butt, W. Tremel and R. Berger, J. Phys. Chem. Lett., 2019, 22, 13458-13466.
- 14 Z. Arain, C. Liu, Y. Yang, M. Mateen, Y. Ren, Y. Ding, X. Liu, Z. Ali, M. Kumar and S. Dai, Sci. China Mater., 2018, 2, 161-172.
- 15 B. Yang, C. C. Brown, J. Huang, L. Collins, X. Sang, R. R. Unocic, S. Jesse, S. V. Kalinin, A. Belianinov, J. Jakowski, D. B. Geohegan, B. G. Sumpter, K. Xiao and O. S. Ovchinnikova, Adv. Funct. Mater., 2017, 26, 1700749.
- 16 F. Wang, M. Yang, S. Yang, X. Qu, L. Yang, L. Fan, J. Yang and F. Rosei, Nano Energy, 2019, 67, 104224.
- 17 P. Schulz and D. Cahen, Chem. Rev., 2019, 5, 3349-3417.
- 18 P. Wang, Q. Jiang, Y. Zhao, Y. Chen, Z. Chu, X. Zhang, Y. Zhou and J. You, Sci. Bull., 2018, 11, 726-731.
- 19 Q. Jiang, Z. Chu, P. Wang, X. Yang, H. Liu, Y. Wang, Z. Yin, J. Wu, X. Zhang and J. You, Adv. Mater., 2017, 29, 1703852.
- 20 J. Li, T. Bu, Y. Liu, J. Zhou, J. Shi, Z. Ku, Y. Peng, J. Zhong, Y.-B. Cheng and F. Huang, *ChemSusChem*, 2018, 17, 2898–2903.
- 21 M. Wang, F. Cao, K. Deng and L. Li, Nano Energy, 2019, 63, 103867.
- 22 M. M. Tavakoli, M. Saliba and P. Yadav, Adv. Energy Mater., 2019, 4, 1900413.
- 23 Z. Liu, F. Cao, M. Wang, M. Wang and L. Li, Angew. Chem. Int. Ed., 2020, 10, 4161-4167.

- 24 W. Ke, C. Xiao, C. Wang, B. Saparov, H. S. Duan, D. Zhao, Z. Xiao, P. Schulz, S. P. Harvey, W. Liao, W. Meng, Y. Yu, A. J. Cimaroli, C. S. Jiang, K. Zhu, M. Al-Jassim, G. Fang, D. B. Mitzi and Y. Yan, Adv. Mater., 2016, 26, 5214-5221.
- 25 D. W. deQuilettes, K. Frohna, D. Emin, T. Kirchartz, V. Bulovic, D. S. Ginger and S. D. Stranks, Chem. Rev., 2019, 20, 11007-11019.
- 26 W. Cao, K. Lin, J. Li, L. Qiu, Y. Dong, J. Wang, D. Xia, R. Fan and Y. Yang, J. Mater. Chem. C, 2019, 40, 12717-12724.
- 27 F. Mohamadkhani, S. Javadpour and N. Taghavinia, Sol. Energy, 2019, 56, 647-653.
- 28 H. Luo, J. Wu, X. Liu, Y. Yang, Q. Liu, M. Zhang, P. Yuan, W. Sun, Z. Lan and J. Lin, ACS Appl. Energy Mater., 2018, 12, 6700-6706.
- 29 Z. Li, S. Wu, J. Zhang, Y. Yuan, Z. Wang and Z. Zhu, Sol RRL, 2019, 3, 1900413.
- 30 S. Wang, Y. Zhu, C. Wang and R. Ma, Small, 2019, 19, 11867-11876.
- 31 W. Hui, Y. Yang, Q. Xu, H. Gu, S. Feng, Z. Su, M. Zhang, J. Wang, X. Li, J. Fang, F. Xia, Y. Xia, Y. Chen, X. Gao and W. Huang, Adv. Mater., 2020, 4, 1906374.
- 32 K. Wang, X. Liu, R. Huang, C. Wu, D. Yang, X. Hu, X. Jiang, J. C. Duchamp, H. Dorn and S. Priya, ACS Energy Lett, 2019, 8, 1852-1861.
- 33 Q. Jiang, Y. Zhao, X. Zhang, X. Yang, Y. Chen, Z. Chu, Q. Ye, X. Li, Z. Yin and J. You, Nat. Photonics, 2019, 7, 460-466.
- 34 J. Chen, H. Dong, L. Zhang, J. Li, F. Jia, B. Jiao, J. Xu, X. Hou, J. Liu and Z. Wu, J. Mater. Chem. A., 2020, 5, 2644-2653.
- 35 S. S. Mali, J. V. Patil, H. Kim and C. K. Hong, Matter, 2019, 2, 452-464.
- 36 B. Li, M. Li, C. Fei, G. Cao and J. Tian, J. Mater. Chem. A., 2017, 46, 24168-24177.
- 37 L.-L. Jiang, Z.-K. Wang, M. Li, C.-C. Zhang, Q.-Q. Ye, K.-H. Hu, D.-Z. Lu, P.-F. Fang and L.-S. Liao, Am. J. Chem., 2018, 7, 1705875.
- 38 Y. Xiang, J. Yu, T. Zhang, Q. Wen, J. Zhuang and J. Guo, Mater. Sci. Semicond. Process., 2016, 15, 48-53.
- 39 Z.-L. Yang, Z.-Y. Zhang, W.-L. Fan, C. Hu, L. Zhang and J.-J. Qi, Sol. Energy, 2019, 7, 859-865.
- 40 A. B. Salyulev and A. M. Potapov, J. Phys. Chem. Lett., 2018, 3, 259-263.
- 41 P. Kumar, E. Vahidzadeh, U. K. Thakur, P. Kar, K. M. Alam, A. Goswami, N. Mahdi, K. Cui, G. M. Bernard, V. K. Michaelis and K. Shankar, J. Am. Chem. Soc., 2019, 13, 5415-5436.
- 42 L. Zhao, Z. Zhao, Y. Li, X. Chu, Z. Li, Y. Qu, L. Bai and L. Jing, Nanoscale, 2020, 18, 10010-10018.
- 43 K. Sun, J. Shen, Q. Liu, H. Tang, M. Zhang, S. Zulfiqar and C. Lei, Chin. J. Catal., 2020, 41, 72-81.
- 44 Y. Tian, L. Zhou, Q. Zhu, J. Lei, L. Wang, J. Zhang and Y. Liu, Nanoscale, 2019, 43, 20638-20647.
- 45 H. Tang, R. Wang, C. Zhao, Z. Chen, X. Yang, D. Bukhvalov, Z. Lin and Q. Liu, Chem. Eng. J., 2019, 1064-1075.
- 46 W. Liu, J. Shen, X. Yang, Q. Liu and H. Tang, Appl. Surf. Sci., 2018, 369-378.
- 47 J. Dagar, S. Castro-Hermosa, G. Lucarelli, F. Cacialli and T. M. Brown, Nano Energy, 2018, 32, 290-299.