

Cite this: *Energy Environ. Sci.*,  
2020, 13, 3663

## Framework for evaluating the performance limits of membraneless electrolyzers†

Xueqi Pang,<sup>a</sup> Jonathan T. Davis,<sup>a</sup> Albert D. Harvey III<sup>b</sup> and  
Daniel V. Esposito<sup>id</sup>\*<sup>a</sup>

Emerging membraneless electrolyzers offer an attractive approach to lowering the cost of hydrogen (H<sub>2</sub>) production from water electrolysis thanks to potential advantages in durability and manufacturability that are made possible by elimination of membranes or diaphragms from the device architecture. However, a fair comparison of the performance limits of membraneless electrolyzers to conventional electrolyzers is hindered by the early stage of research and absence of established design rules for the former. This task is made all the more difficult by the need to quantitatively describe multiphase flow between the electrodes in membraneless electrolyzers, which can have a huge impact on gas product purity. Using a parallel plate membraneless electrolyzer (PPME) as a model system, this study takes a combined experimental and modeling approach to explore its performance limits and quantitatively describe the trade-offs between efficiency, current density, electrode size, and product purity. Central to this work is the use of *in situ* high-speed videography (HSV) to monitor the width of H<sub>2</sub> bubble plumes produced downstream of parallel plate electrodes as a function of current density, electrode separation distance, and the Reynolds number (Re) associated with flowing 0.5 M H<sub>2</sub>SO<sub>4</sub> electrolyte. These measurements reveal that the HSV-derived dimensionless bubble plume width serves as an excellent descriptor for correlating the aforementioned operating conditions with H<sub>2</sub> crossover rates. These empirical relationships, combined with electrochemical engineering design principles, provide a valuable framework for exploring performance limits and guiding the design of optimized membraneless electrolyzers. Our analysis shows that the efficiencies and current densities of optimized PPMEs constrained to H<sub>2</sub> crossover rates of 1% can exceed those of conventional alkaline electrolyzers but are lower than the efficiencies and current densities achieved by zero-gap polymer electrolyte membrane (PEM) electrolyzers.

Received 17th July 2020,  
Accepted 7th September 2020

DOI: 10.1039/d0ee02268c

rsc.li/ees

### Broader context

Water electrolysis powered by solar- or wind-derived electricity is a promising technology for the production of H<sub>2</sub> in a renewable energy future. However, less than 5% of H<sub>2</sub> generated in the world today comes from water electrolysis because of its high cost compared to producing H<sub>2</sub> from steam methane reforming. Decreasing costs of electricity from solar and wind generators are helping to reduce this gap, but the lower capacity factor of electrolyzers operating primarily or exclusively using electricity from intermittent renewable sources makes it critically important to reduce the capital costs of electrolyzers. Membraneless electrolyzers developed in recent years represent a promising approach to decrease capital costs by reducing the number of cell components while presenting an opportunity to improve stack lifetime by omitting a membrane or diaphragm divider. Motivated by the need to better understand the performance limits of membraneless electrolyzers compared to conventional designs, this study describes and demonstrates a combined experimental and modelling framework to investigate the trade-offs between the efficiency, current density, product purity, and electrode size in model parallel plate membraneless electrolyzers. This framework shows that the performance limits of a model parallel plate membraneless electrolyzer can exceed the typical performance of conventional alkaline electrolyzers.

<sup>a</sup> Columbia University in the City of New York, Department of Chemical Engineering, Columbia Electrochemical Energy Center, Lenfest Center for Sustainable Energy, New York, NY 10027, USA. E-mail: de2300@columbia.edu; Tel: +1-212-854-2648

<sup>b</sup> Shell International Exploration & Production, Inc., Houston, TX 77082, USA

† Electronic supplementary information (ESI) available. See DOI: 10.1039/d0ee02268c

### Introduction

The continued growth in global solar photovoltaic and wind installations is expected to be one of the most important enablers for reducing anthropogenic carbon dioxide emissions,<sup>1</sup> but the intermittency of these energy sources raises significant challenges

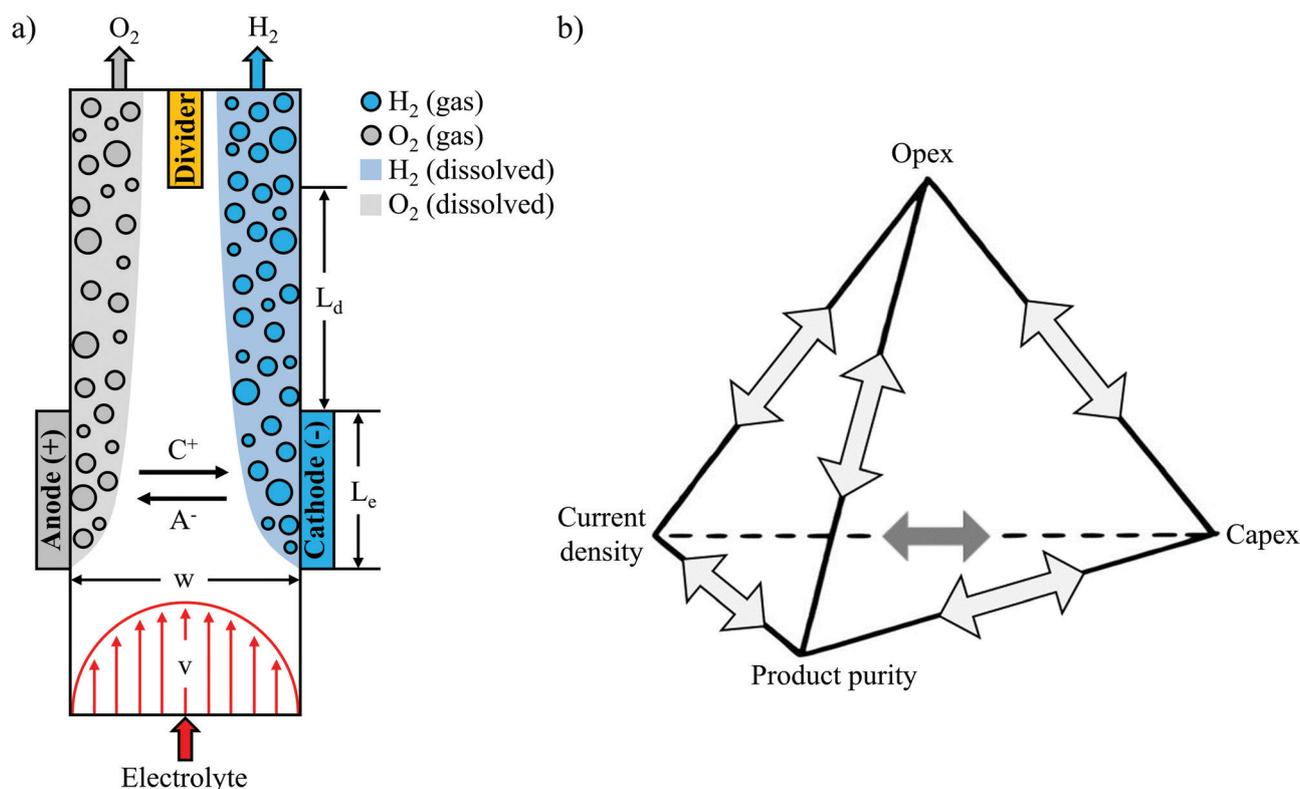


to integrate them into today's energy system.<sup>2</sup> Electrochemical devices called electrolyzers offer an attractive means of absorbing large amounts of excess renewable electricity by using it to convert water into a storable, carbon-free fuel in the form of H<sub>2</sub>.<sup>3–5</sup> However, the cost of splitting water to produce H<sub>2</sub> with conventional electrolyzers is currently much higher than steam methane reforming.<sup>6</sup>

Presently, there are two types of electrolyzers that dominate the market: alkaline electrolyzers and PEM electrolyzers.<sup>7</sup> As an alternative, membraneless electrolyzers have potential to decrease material and assembly costs thanks to simpler device designs that can be made from fewer components.<sup>8–10</sup> The absence of a membrane or diaphragm divider also creates an opportunity for stable operation in harsh environments, while the relative simplicity of membraneless electrolyzers also makes them attractive as a platform for studying multi-scale electrochemical phenomena or carrying out fundamental electrocatalysis research. Membraneless electrolyzers can be generally grouped into two categories that are differentiated based on the type of electrodes they employ: (i) flow-through and (ii) flow-by. Flow-through electrodes allow flowing electrolyte to pass through the electrodes, while flow-by electrodes rely on electrolyte flow that is parallel to the surface of the electrodes. One of the most commonly explored membraneless electrolyzer designs is based on flow-by electrodes that are

positioned in a parallel plate geometry,<sup>9,11</sup> as illustrated in Fig. 1a. During water electrolysis, H<sub>2</sub> and oxygen (O<sub>2</sub>) bubbles generated at the electrode surfaces are swept downstream into separate effluent channels by the flowing electrolyte. The gas bubbles travel in close proximity to the channel walls because the velocity gradients associated with laminar flow create a net inertial lift force that directs them away from the centerline of the channel, a phenomena known as the Segré-Silberberg effect.<sup>12–14</sup> The vast majority of membraneless electrolyzer demonstrations based on flow-through electrodes similarly rely on flow-induced separation of product species,<sup>15–17</sup> although a recent study from our lab has also demonstrated the ability to achieve relatively pure product streams using only the buoyancy of the product gas bubbles.<sup>18</sup> Regardless of the mechanism of product separation, the absence of a physical barrier between electrodes in membraneless electrolyzers makes it especially important to quantitatively understand how bubble dynamics affect product purity for different cell geometries and operating conditions.

Experimental studies of membraneless electrolyzers have demonstrated current densities exceeding 3.5 A cm<sup>-2</sup>,<sup>10</sup> efficiencies above 58% (based on  $\Delta G^\circ = 237 \text{ kJ mol}^{-1} \text{ H}_2$ ) at 0.2 A cm<sup>-2</sup>,<sup>11</sup> H<sub>2</sub> crossover rates below 0.4%,<sup>9,10</sup> and electrode sizes up to 344 cm<sup>2</sup>,<sup>19</sup> but never at the same time. The reason is that these performance metrics are all closely coupled, with numerous



**Fig. 1** (a) Schematic top-view of the parallel plate membraneless electrolyzer (PPME). H<sub>2</sub> and O<sub>2</sub> are generated at the cathode and anode, respectively. A<sup>-</sup> and C<sup>+</sup> represent the anion and cation species in the electrolyte. (b) "Unfortunate tetrahedron" of membraneless electrolyzers highlighting trade-offs between key techno-economic performance metrics. Current density is synonymous with production rate, while product purity is a proxy for process safety.



trade-offs between efficiency, current density, product purity, and electrode size. It is relatively straight-forward to modify the design of a membraneless electrolyzer to maximize any one of these parameters, but doing so will invariably have an adverse effect on one or more of the other key performance metrics. Taking as an example the parallel plate membraneless electrolyzer (PPME), the electrolysis efficiency at a given current density can be easily increased by decreasing the separation distance between the two electrodes, but doing so will lead to a concomitant decrease in product purity or maximum electrode size. The relationships and associated trade-offs between strongly-coupled technoeconomic performance metrics have been graphically illustrated in the design of batteries in what has been referred to as the “unfortunate tetrahedron”.<sup>20</sup> The lines forming the edges of the tetrahedron represent the design rules governing the direct binary relationships between two metrics located at the tetrahedron vertices, while each face of the tetrahedron, termed a “triplet”, connects three key metrics that are strongly coupled with each other. Within the unfortunate tetrahedron for batteries, the four key performance metrics of operating expenditure (opex), capital expenditure (capex), power density, and energy density are positioned at the vertices of the tetrahedron. Extending this concept to electrolyzers, Fig. 1b presents an unfortunate tetrahedron for membraneless electrolyzers, where opex (related to efficiency), capex (related to electrode size and scalability), current density, and product purity (a proxy for safety) have been selected as four key performance metrics. The electrolyzer and battery tetrahedrons have a lot in common, sharing several performance metrics of interest that are either identical or very similar. These similarities reflect the fact that there are similar design constraints for these two electrochemical energy conversion devices, but there are differences as well. In particular, electrolyzers are open systems with reactants flowing in and energy carriers flowing out, whereas batteries tend to be closed systems which are themselves energy storage devices. Thus, we have replaced “energy density” with “product purity” in Fig. 1b, where the latter is a key performance metric for membraneless electrolyzers that is closely interconnected with the other three vertices in the tetrahedron. With an eye towards the ultimate goals of minimizing the levelized cost of hydrogen (LCOH) production and maximizing safety, it is essential that engineering design rules be developed that accurately describe the trade-offs between the key performance metrics. Although select binary relationships have been previously studied for various membraneless electrolyzer designs,<sup>9–11</sup> a wholistic modeling framework that describes the relationships between all of the key performance metrics shown in Fig. 1b has yet to be established.

A primary objective for this study was to develop such a modeling framework for the PPME shown in Fig. 1a. This particular electrolyzer was chosen as a basis for this study because its simple design is characterized by a small number of key geometric parameters: the distance between electrodes ( $w$ ), electrode length ( $L_e$ ), distance between the end of the electrodes and downstream divider ( $L_d$ ), and the height of the

electrodes ( $H$ ), which extends into the page. The geometric simplicity makes this design amenable to the application of established electrochemical modeling approaches, in addition to providing a convenient platform to study the multi-phase flow behavior of the plumes of product gases evolved from the electrodes using *in situ* high speed videography (HSV). Accurate descriptions of bubble plume dynamics and their dependencies on key operating parameters are essential for understanding how the product purity vertex in Fig. 1b is coupled to key performance metrics related to the other three vertices. This study builds off of several other recent studies that have used HSV to study bubble dynamics in membraneless electrolyzers. Notably, Hashemi *et al.* used *in situ* HSV to observe trends in bubble dynamics at different Reynolds numbers ( $Re$ ) and current densities,<sup>9,11</sup> while Gillespie *et al.* used *in situ* imaging to observe bubble build-up between flow-through electrodes—a phenomena termed void fracture—at low electrolyte flow rates.<sup>10</sup> More recently, our group has used *in situ* HSV to quantify bubble size distributions, gas evolution efficiencies, and current density distributions in a membraneless electrolyzer based on angled mesh flow-through electrodes.<sup>21</sup>

Expanding on the quantitative image analysis methods developed in prior work, the present study applies HSV to quantify the effective width of bubble plumes evolved downstream of parallel plate electrodes as a function of current density and  $Re$ . The fluid flow between the two electrodes is characterized for two different values of the channel width,  $w$ . As will be shown later in this article, the dimensionless bubble plume width, which is normalized by  $w$ , is found to be strongly correlated with  $H_2$  cross-over rates, and hence, product purity. After describing the experimental set-up and methodologies used to quantify bubble plume widths, correlations between plume widths, operating currents, and  $H_2$  crossover rates are presented. The correlations between these three parameters are then applied in conjunction with an electrolyzer model to map out the relationships between key performance metrics that are directly related to the vertices of the unfortunate tetrahedron in Fig. 1b. Finally, the modeled performance limits of the PPME based on this work are compared to the performance of conventional alkaline and PEM electrolyzers. Collectively, this work presents a combined experimental and modeling approach that can be broadly applied to establish the performance limits of membraneless electrolyzers and develop design rules for selecting a combination of geometric and operational parameters that will optimize the overall performance or cost of an electrolyzer.

## Materials and methods

### Chemicals

All solutions were prepared using 18.2 M $\Omega$  cm deionized water. Electrolytes were prepared using concentrated sulfuric acid (Certified ACS plus, Fisher Scientific). For every 100 mL electrolyte, one drop of Triton X-100 surfactant was added to the electrolyte in order to reduce the average bubble detachment



size from the electrodes, as is well known in literature.<sup>22,23</sup> Adding surfactant to the electrolyte is especially beneficial for smooth planar electrodes such as the Pt thin films used in this study, but can be avoided in commercial applications with higher surface area, textured electrodes that tend to make electrodes more hydrophilic and reduce the size of departing bubbles.<sup>24,25</sup> Electrolyte solutions were purged with argon (Ar) (PurityPlus, 99.998% purity). Green food coloring dye (McCormick) was injected into the entrance port of the flow cell to monitor streamlines at varying Re.

### Electrode fabrication

All electrodes were fabricated by using electron-beam evaporation to sequentially deposit 2 nm of titanium (Ti) (99.99%) and 50 nm platinum (Pt) (99.99%) onto Ti foil substrates (0.127 mm thick, annealed, 99% metals basis, Alfa Aesar). Electrodes were fabricated with a length of 0.90 cm and a width of 0.34 cm, resulting in an area of 0.30 cm<sup>2</sup>.

### Electrolyzer fabrication

Electrolysis cells were printed from polylactic acid (PLA) using a MakerGear M3-ID 3D printer. The rectangular fluidic channel within the electrolyzer was 145 mm long, 2 or 4 mm wide and 5 mm high, with a 29 mm long by 0.4 mm thick divider placed 12 mm downstream of the end of the electrodes. Two rectangular metal foils containing electrodes were bent to an angle of 135° at a length of 0.90 cm and slid into slits on both sides of the cell before being epoxied (ClearWeld, J.B. Weld) in place. Two glass microscope slides (76.2 mm × 25.4 mm × 1 mm) were then epoxied to the front and back of the device to allow backlighting during *in situ* HSV measurements. A computer aided design (CAD) rendering of the assembled cell is shown in Fig. S1 (ESI<sup>†</sup>), and CAD files for the two cells are freely available for download at echem.io.

### Device characterization and operation

Flexible silicone tubing (Masterflex L/S 24, Cole Parmer) was fitted to the electrolyzer inlet. 1/8" polyethylene straight connectors (Cole Parmer) were epoxied to the outlets of the device and flexible silicone tubing (Masterflex L/S 16, Cole Parmer) was fitted to the outlet connectors. The electrolyzer was then mounted vertically such that the inlet and outlet ports pointed downwards and upwards, respectively. The tubing connected to the inlet port was connected to a pulse dampener (Cole Parmer), which was in turn connected to a peristaltic pump (Cole Parmer). The end of the inlet tube was submerged in the 0.5 M sulfuric acid electrolyte, which was continuously purged with Ar gas for at least 90 min before each experiment. The outlet tubes were connected to a gas collection apparatus as shown in Fig. S2 (ESI<sup>†</sup>). The electrolyte flowing out of the collection apparatus was sent to a separate reservoir. This "once-through" set-up was chosen over a recirculating system to ensure that crossover rates were not artificially increased by dissolved gases passing through the system multiple times in the absence of a phase separator. Although the percentage of H<sub>2</sub> entering the dissolved phase will be increased with this

once-through set-up, it can be lowered by decreasing Re and/or increasing current density and electrode size (Fig. S3, ESI<sup>†</sup>).

All electrochemical measurements were carried out at room temperature using a Biologic SP-300 potentiostat/galvanostat. Electrochemical impedance spectroscopy (EIS) measurements were conducted at a constant applied cell voltage of 0.2 V vs. open circuit potential using an AC amplitude of 10 mV with a frequency range between 100 mHz and 200 kHz. Constant current electrolysis experiments were conducted for 30–90 minutes after purging the electrolyte feed reservoir for at least 90 minutes. HSVs of bubble dynamics within the fluidic channel were recorded using an Edgertronic SC1 high-speed camera operating at 1000 frames per second (fps) for Re up to 796 and 2000 fps for Re above 796 with a resolution of 288 × 1280 pixels. The camera was fitted with a FotodioX 52 mm reverse lens adaptor ring and a Nikon AF Nikkor 50 mm f/1.8D lens. The length of each HSV was 10 seconds.

### HSV processing

Still frame images from HSVs were cropped to an analysis area spanning the half channel width (0.5 $w$ ) on the cathode side of the cell. Cropped images started from the downstream edge of the electrode to the bottom of the divider separating the anode and cathode effluent streams. This analysis area located downstream of the electrodes was chosen because the total flux of bubbles entering and leaving the analyzed area remains constant, thereby simplifying the analysis for quantification of the average bubble plume width. The *imfindcircles* function within the MATLAB Image Processing Toolbox (R2019a) was then used to locate bubbles and estimate their radii.<sup>21</sup> In order to determine the two-dimensional coverage of bubbles in each image, the *bwareaopen* function in MATLAB was used to convert the greyscale raw image into a binary black and white image, with bubble size thresholds and bubble locations used to avoid underestimation in the bubble coverage caused by local bright spots from the bubbles. The 2D coverage of bubbles was then used to determine the local void fraction of bubbles within each 100 by 40 μm pixel in the image. This algorithm was applied to 100 frames of each video ( $\Delta t = 0.1$  s between consecutive frames), and the average void fraction values for each pixel were used to generate bubble void fraction maps.

### Finite element modeling

The potential field in the electrolyte was modeled using COMSOL 5.4 Electrodeposition Multiphysics package. The total solution resistance ( $R_s$ ) was computed from Ohm's law for different channel widths and electrode lengths with solution conductivities ( $\kappa$ ) of 0.226 S cm<sup>-1</sup> (corresponding to 0.5 M H<sub>2</sub>SO<sub>4</sub>) and 0.621 S cm<sup>-1</sup> (corresponding to 30 wt% KOH).

### Analysis of product gas composition

Gas chromatography (GC) measurements were performed using a model 8610C gas chromatograph (SRI Instruments) equipped with an 18" HayeSep D packed column, a 3' Molecular Sieve 5A packed column, and a thermal conductivity detector (TCD) using Ar as a carrier gas. The gaseous electrolysis products were



collected in separate vertically-oriented glass tubes that were sealed with rubber septa. Before electrolysis experiments, a syringe was used to draw fresh electrolyte upwards into each tube. 0.5–1 mL samples were taken from each tube using a gas-tight syringe (Hamilton) and manually injected into the GC. Product crossover was calculated by measuring the amount of hydrogen in the O<sub>2</sub> collection tube and dividing by the total amount of measured H<sub>2</sub>. Relationships between the % H<sub>2</sub> crossover and atomic % H<sub>2</sub> in the anode and cathode effluent streams are provided in Fig. S4 (ESI†).

## Results

### Description of the electrolyzer and experimental setup

As detailed in the methods section and illustrated in a simple flow diagram of the experimental set-up in Fig. S2 (ESI†), a peristaltic pump was used to pump 0.5 M H<sub>2</sub>SO<sub>4</sub> electrolyte through the vertically-oriented electrolyzer. During operation, the flowing electrolyte helps to facilitate removal of the electrolytically-generated H<sub>2</sub> and O<sub>2</sub> gas bubbles from the electrode surfaces before sweeping the “bubble plumes” downstream into the cathode and anode effluent channels, which are separated by an insulating baffle that is part of the cell body. After passing through the effluent channels, the two-phase mixture of gas bubbles and electrolyte flows into downstream glass collection tubes where the gas bubbles phase-separate and are captured in the headspace where they can be sampled for compositional analysis by GC. Although the majority of the sampled gas can be attributed to gas bubbles evolved from the upstream electrodes, some small fraction can also be assumed to originate from dissolved H<sub>2</sub> and O<sub>2</sub> species that are able to exchange between the liquid and gas phases inside the collection tubes. Given that the PPME is characterized by very large Péclet numbers (10<sup>4</sup>–10<sup>5</sup>) over the operating conditions for this study, H<sub>2(aq)</sub> and O<sub>2(aq)</sub> can be assumed to be co-located within the bubble plumes. Under the constraints of large Pe and laminar flow conditions, HSV-derived bubble plumes can thus be a good way to monitor the product species in both the gaseous and dissolved phases.

Electrolyzers with two different channel widths ( $w = 4$  mm and  $w = 2$  mm) and outfitted with identical 0.3 cm<sup>2</sup> platinized Ti foil electrodes were used for this study. 2-Electrode current density–voltage ( $iV$ ) curves for these cells were recorded in 0.5 M H<sub>2</sub>SO<sub>4</sub> for identical Reynolds number ( $Re = 547$ ) and are provided in Fig. S5 (ESI†). The onset voltage for both cells is around 2.0 V, which is 0.4 V higher than that reported in the same electrolyte for a similar parallel plate electrolyzer based on Ir-containing electrodes that have significantly lower overpotential for the OER.<sup>11</sup> As the voltage increases more than  $\approx 100$  mV beyond the onset voltage, the current density increases linearly with voltage. There is a significant increase in the slope of the  $iV$  curve ( $di/dV$ ) of the  $w = 2$  mm compared to the  $w = 4$  mm cell due to the lower solution resistance across the smaller channel gap, which was confirmed by EIS measurements (Fig. S6, ESI†). Solution resistance ( $R_s$ ) was found to

decrease from 3.6  $\Omega$  to 2.1  $\Omega$  when the channel width decreased from 4 mm to 2 mm, resulting in a decrease in the applied voltage of 140 mV at a current density of 200 mA cm<sup>-2</sup>.

### Flow characteristics of the electrolysis cell

Because parallel plate membraneless electrolyzers rely on flow-induced separation of product gases, it is essential to understand how the hydrodynamics of the cell vary as a function of electrolyte flow rate and the associated  $Re$ . It is particularly important to know what  $Re$  corresponds to a transition from laminar-to-turbulent flow since turbulent flow leads to rapid mixing of the electrolyte and the anode/cathode products.<sup>26</sup> In this work, the laminar-to-turbulent transition was monitored in the  $w = 4$  mm cell by visualizing changes in streakline characteristics as a function of  $Re$ . HSV was recorded while injecting colored dye into the inlet of the cell, with representative still frame images provided in Fig. 2a at four different  $Re$ . The images taken for  $Re = 131$  and  $Re = 1004$  contain vertical streaklines, consistent with laminar flow and the associated parabolic velocity profile. However, increasing  $Re$  above  $\approx 1200$  significantly distorts the streaklines and leads to small convective cells associated with unsteady flow. At a  $Re = 1211$ , transitional flow was observed, as evidenced by turbulent mixing of the dye in the center of the channel while narrow regions of laminar flow (light coloring) are still visible along the channel walls. These dye-free laminar regions are absent in the  $Re = 1419$  image, indicating that the flow is almost entirely turbulent under these conditions. This  $Re$  is significantly lower than a  $Re$  of 4000 that is typically associated with fully turbulent flow between perfectly flat parallel plates,<sup>27</sup> which we attribute to a combination of wall roughness and entrance effects associated with fluid flow transitioning from the circular inlet tube to the rectangular cell channel containing the electrodes.

In Fig. 2b, electrolyzer hydrodynamic characteristics were observed over the same range of  $Re$ , but instead of using colored dye to view streaklines, electrochemically-generated bubbles were monitored during constant current electrolysis experiments ( $i = 200$  mA cm<sup>-2</sup>). In these experiments, the bubbles serve as tracer particles to visualize flow characteristics, similar to particle image velocimetry experiments that are regularly used in the field of fluid dynamics.<sup>28</sup> Five representative images from HSVs are shown for different  $Re$  while running electrolysis, demonstrating how flowing electrolyte modulates the bubble plume width. H<sub>2</sub> bubbles were generated on the right (cathode), and O<sub>2</sub> bubbles were generated on the left (anode). In the absence of advection (Fig. 2b, “no flow”), bubble-induced convection leads to broad H<sub>2</sub> and O<sub>2</sub> bubble plumes that intersect before the bubbles reach the top of the electrodes, highlighting the necessity of using fluid flow to promote effective product separation in the parallel plate geometry. Even at the lowest flow rate used in this study, corresponding to  $Re = 131$ , product separation is observed to be significantly improved, although direct crossover of bubbles to the opposite effluent channel is still occasionally observed. As the flow rate increases further, the detached bubbles experience significantly larger forces from the fluid flow that push



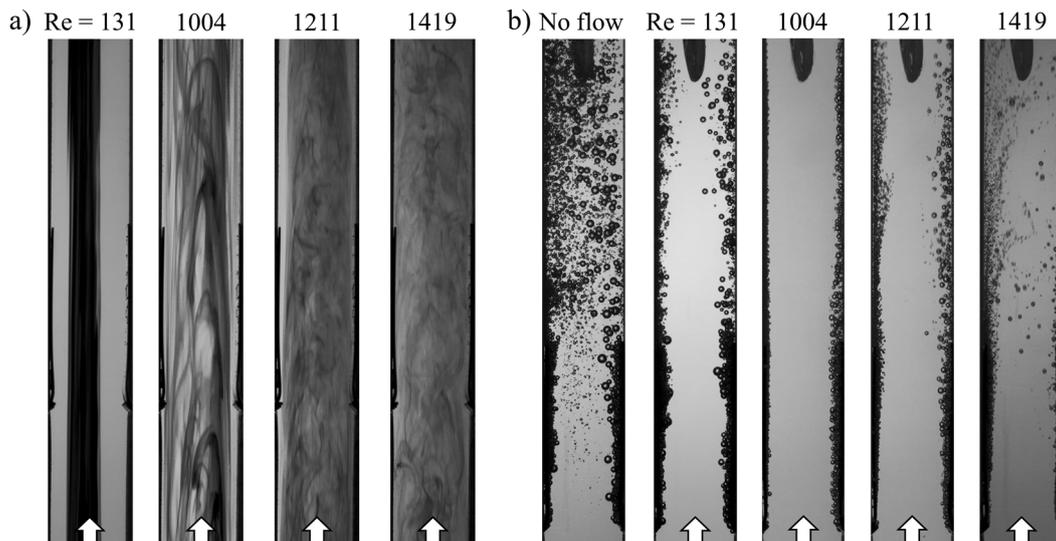


Fig. 2 (a) Flow visualization using dye injection at different  $Re$ , showing the flow transition from laminar to turbulent. (b) Representative still frame images from HSVs captured after 5 min of electrolysis showing bubbles evolving from electrodes, operated at  $200 \text{ mA cm}^{-2}$  in surfactant-containing  $0.5 \text{ M H}_2\text{SO}_4$  at different  $Re$ . The electrolyzer channel width is  $4 \text{ mm}$ , and the electrolyte flows from bottom to top in the image as indicated by the white arrows at the bottom of each image.

them closer to the channel wall from which they originated. At a flow rate corresponding to  $Re = 1004$ , the detached bubbles are effectively “pinned” to the channel wall, resulting in a very narrow bubble plume width and lack of any observed direct crossover events. However, increasing  $Re$  past the critical  $Re \approx 1200$  causes the bubble plume width to increase again, eventually leading to the reemergence of direct crossover events as the flow transitions to turbulent conditions. Increasing  $Re$  impacts not only the width of the bubble plume, but also the average bubble size, which is found to decrease monotonically with increasing  $Re$  (Fig. S7, ESI<sup>†</sup>), consistent with prior studies.<sup>9,11</sup> This trend can be more concisely attributed to the increasing ratio of the inertial force promoting bubble detachment relative to the surface tension force that promotes adhesion, a dimensionless number known as the Weber number ( $We$ ).

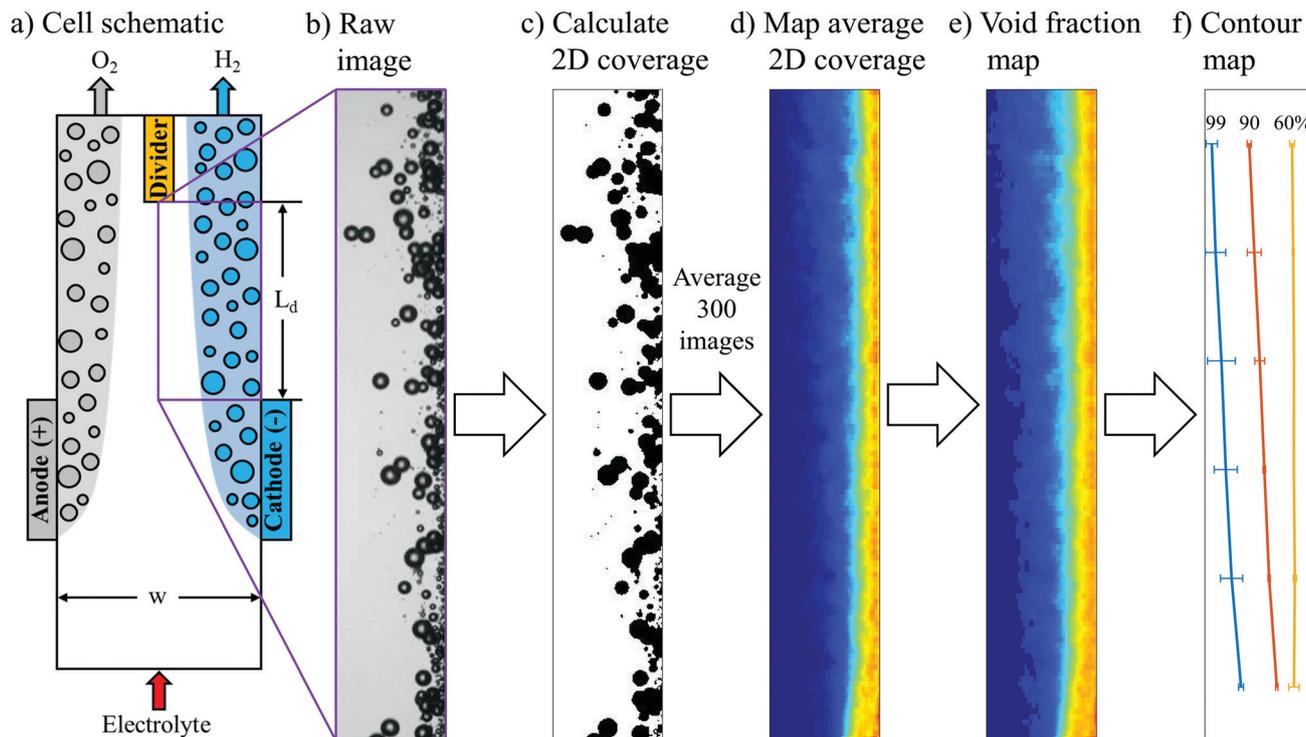
### Quantifying bubble plume width based on *in situ* high speed videography

A key premise of this study is that the effective width of bubble plumes generated by parallel plate electrodes strongly correlates with crossover rates, and can therefore be a useful descriptor for relating operating conditions to product purities. More specifically, it was our expectation that GC-determined crossover rates be very low ( $<1\%$ ) whenever the steady-state bubble plume width ( $w_p$ ) is significantly less than half of the electrolyzer channel width ( $0.5w$ ). Conversely, crossover should rapidly increase as  $w_p$  approaches ( $0.5w$ ). In this work, analysis was focused on the plume of  $\text{H}_2$  bubbles generated at the cathode, although similar analysis could also be applied to the  $\text{O}_2$  bubble plume. To quantify  $\text{H}_2$  bubble plume width, still frame images from HSVs were converted into bubble void fraction images. First, raw images (Fig. 3b) were cropped to an analysis area corresponding to the purple rectangle shown in Fig. 3a that is located immediately downstream of the

cathode and extending down to the divider separating the anode and cathode effluent streams. Next, the greyscale raw image was converted into a binary black and white image, and the bubbles were converted to black (Fig. 3c). Based on the black and white areas in the image, the two-dimensional (2D) coverage of bubbles within each  $100 \mu\text{m}$  by  $40 \mu\text{m}$  pixel in each image was calculated. Fig. 3d shows the 2D coverage map obtained by averaging 300 images from 3 HSVs for the same operating condition.

Because  $\text{H}_2$  crossover depends on the net volume of  $\text{H}_2$  bubbles that cross over to the anode effluent stream, it was necessary to convert the 2D projected area of bubbles into three-dimensional (3D) bubble void fractions. The projected area of a bubble is easily converted into volume by multiplying the area of each bubble by  $(4/3)$  times the radius of the bubble. However, this method can result in significant error for dense bubble plumes, for which overlapping bubbles in the 2D projection leads to underestimation of the bubble void fraction. To more accurately convert local 2D bubble coverage to 3D void fraction, a correlation between these two parameters was first obtained from computer-generated 3D images of bubbles created within a control volume characterized by the same depth as the electrolysis cells. This was done by first randomly positioning bubbles within the control volume that were characterized by a bubble size distribution corresponding to that which was experimentally-observed during electrolysis at  $200 \text{ mA cm}^{-2}$  and  $Re = 796$  (Fig. S8, ESI<sup>†</sup>). One of such images is shown in Fig. S9a (ESI<sup>†</sup>) for a void fraction of  $0.032$ , which was calculated using the radii of these bubbles. 2D projections of bubbles located in 3D spaces for different bubble coverages are shown in Fig. S9b (ESI<sup>†</sup>). The 2D coverage of bubbles was computed for different bubble void fractions, with the resulting coverage-void fraction correlation provided in the blue curve in Fig. S9c (ESI<sup>†</sup>) for a cell depth of  $5 \text{ mm}$ . If the bubble overlap is





**Fig. 3** (a) Schematic of the PPME. The purple box indicates the region of the electrolyzer that was the focus of the *in situ* imaging work carried out in this study. (b) Representative raw still frame image of H<sub>2</sub> bubbles flowing upwards from the cathode during electrolysis at 200 mA cm<sup>-2</sup> and Re = 131. (c) Bubbles were converted to black after converting the greyscale raw image into a binary black and white image. (d) 2D map of bubble coverage obtained by averaging 300 still frame images from 3 HSVs for the same operating condition. (e) Time averaged void fraction map converted from the average 2D coverage map. For (d) and (e), the warmer colors indicate higher bubble coverages and void fractions, respectively. (f) Contour map with contour lines corresponding to the locations where 99%, 90% and 60% of H<sub>2</sub> within the plume is located to the right of the line.

neglected, a linear correlation will be obtained. With this linear relationship, void fraction is underestimated for bubble coverage higher than 0.2 (Fig. S9c, ESI†). In the following analyses, this true relationship was used to convert the HSV-measured 2D bubble coverages into local bubble void fractions. As an example, the 2D coverage map (Fig. 3d) was converted to a void fraction map (Fig. 3e) using this correlation. In order to quantify the bubble plume width, a contour map was generated to determine the distance from the wall at a given *y*-location that contains *X*% of the total detected H<sub>2</sub> bubble volume to its right. The analysis area was separated into six bins spanning the length of the channel from the end of the electrode to the downstream divider, and values of *X*% (*X* = 99, 90, 60) were calculated for each bin. A contour map generated from the void fraction image in Fig. 3e is shown in Fig. 3f.

Fig. 4a–g show the resulting void fraction images of the downstream H<sub>2</sub> bubble plume at different Re under steady-state electrolysis at 200 mA cm<sup>-2</sup>. Qualitatively, the H<sub>2</sub> bubble plume width is seen to decrease as Re increases from 131 to 1004, consistent with the raw HSV images in Fig. 2b where it was noted that the increased inertial lift force at high Re pins the bubble plume closer to the channel wall. For Re > 1000, the plume width expands with increasing flow rates as the fluid flow transitions to the turbulent regime. It should also be noted that the average bubble void fraction within the plume decreases at these high Re, which can be explained by

expansion of the plume width and a decreased residence time for bubbles flowing through the analysis area at higher fluid velocities.

A contour map generated from the void fraction image in Fig. 4a is shown in Fig. 4h, while the contour maps corresponding to the void fraction images in Fig. 4b–g are provided in Fig. S10 (ESI†). The *X* = 90% contour line value determined for the top of the image analysis area, corresponding to the bottom of the baffle divider, was taken to be representative of the H<sub>2</sub> bubble plume width in the following analyses.

### Correlating HSV-derived bubble plume width to H<sub>2</sub> crossover and operating parameters

Fig. 5a shows the dependency of the GC-measured H<sub>2</sub> crossover percent and HSV-determined H<sub>2</sub> bubble plume width (90%) on Re for a constant current density of 200 mA cm<sup>-2</sup>. Both curves are u-shaped, and can be broken up into three different regimes. At low Re, the shear stress is too low to push all the bubbles closely against the wall, which gives rise to wider bubble plume and higher crossover rate. For intermediate Re between roughly 300 and 1000, bubble plume widths are relatively narrow, and H<sub>2</sub> crossover rates are < 0.57%, with a minimum crossover rate of 0.12% for Re = 1004. Both crossover and bubble plume width are low in this flow regime due to the elevated shear force that helps to push the bubble plumes against the wall. Hashemi *et al.* also obtained the same trend



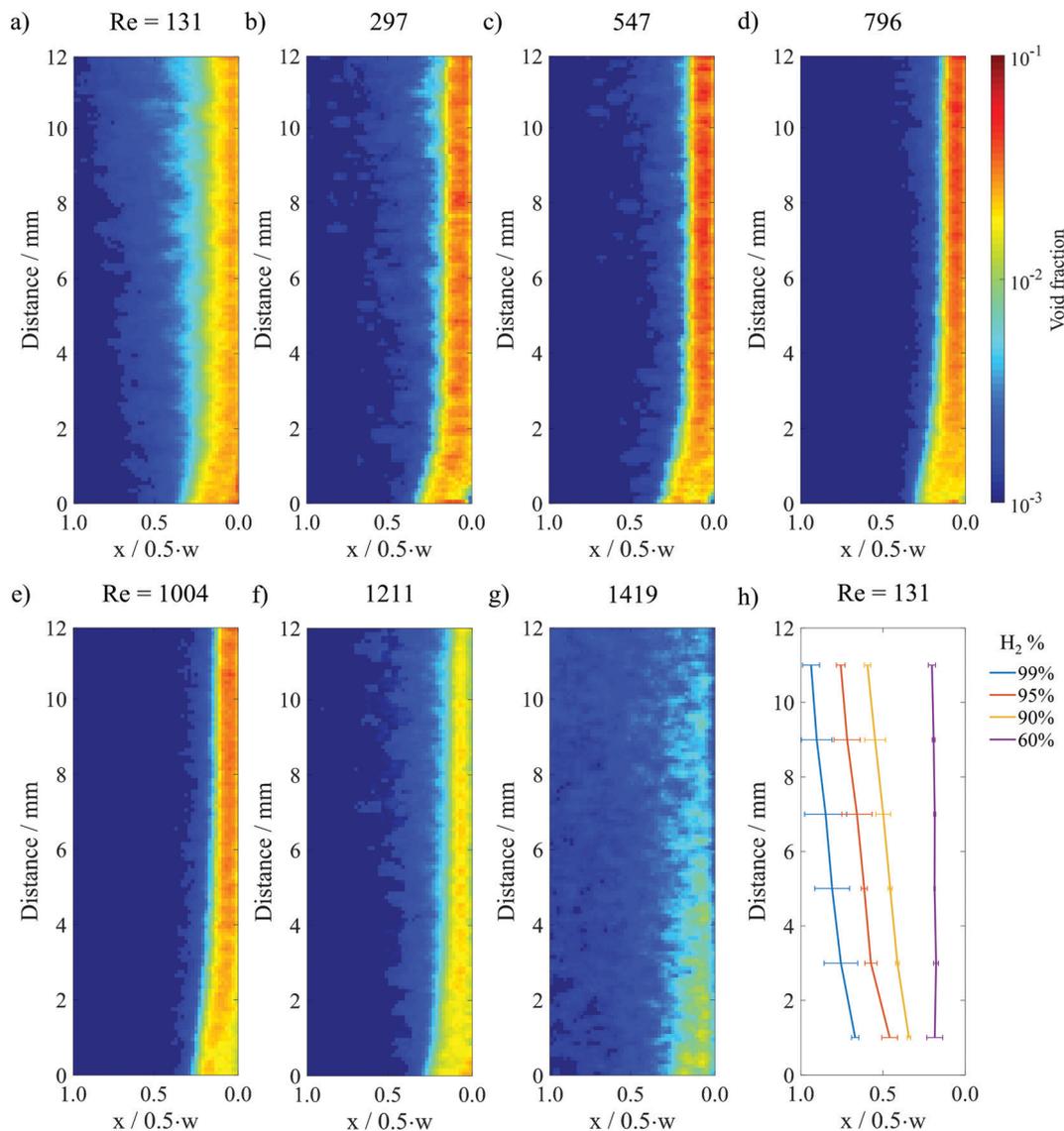


Fig. 4 (a)–(g) Time averaged void fraction images for a constant current density of  $200 \text{ mA cm}^{-2}$  and different Re in surfactant-containing  $0.5 \text{ M H}_2\text{SO}_4$ . 300 still frame images from 3 HSVs were analyzed to generate (a), (d), (e) and (f). 100 images from 1 HSV were analyzed to generate (b), (c) and (g). The image analysis area is marked by the purple box in the electrolyzer schematic provided in Fig. 3a. (h) Contour lines from the Re = 131 image that correspond to the locations where 99%, 95%, 90% and 60% of H<sub>2</sub> within the plume is located to the right of the line.

for H<sub>2</sub> crossover in the laminar flow regime up to  $\text{Re} = 312$ ,<sup>9,11</sup> although that study didn't explore Re above 312. When Re increases to 1211, the transition to turbulent flow characteristics shown in Fig. 2a results in significant increases in the bubble plume width, although the H<sub>2</sub> crossover rate still remains relatively low at 0.66%. We attribute the large variation in the HSV-determined plume width at  $\text{Re} = 1211$  to the stochastic nature of the transitional flow regime, while the crossover rate might remain low thanks to the presence of laminar flow characteristics close to the channel walls, as evidenced in Fig. 2a. At  $\text{Re} > 1211$  both H<sub>2</sub> crossover and bubble plume width rapidly increase as fluid flow becomes more unsteady.

Fig. 5b contains the same plot as Fig. 5a but generated from experiments carried out during electrolysis at a current density

of  $400 \text{ mA cm}^{-2}$ . Because of the higher rate of bubble generation at the elevated current density, notable increases in  $w_p$  and H<sub>2</sub> crossover rates are observed at all Re, but are especially pronounced for the low Re (131) and transitional flow ( $\text{Re} = 1211$ ) conditions. Fig. 5b indicates that the range of Re that allows for low H<sub>2</sub> crossover shrinks as current density increases. Similar trends can be found in Fig. S11 and S12 (ESI<sup>†</sup>) for contour line values at 99%, 95% and 60%. To achieve the lowest H<sub>2</sub> crossover, the operating Re should be increased to a value that is as high as possible without introducing flow instabilities into the system. Although operating at higher Re will require extra pumping power, the parasitic energy loss due to pumping is generally only a small fraction of the energy content of the generated H<sub>2</sub> for the cells studied in this work. As shown in Fig. S13 (ESI<sup>†</sup>), the ratio of required pumping power



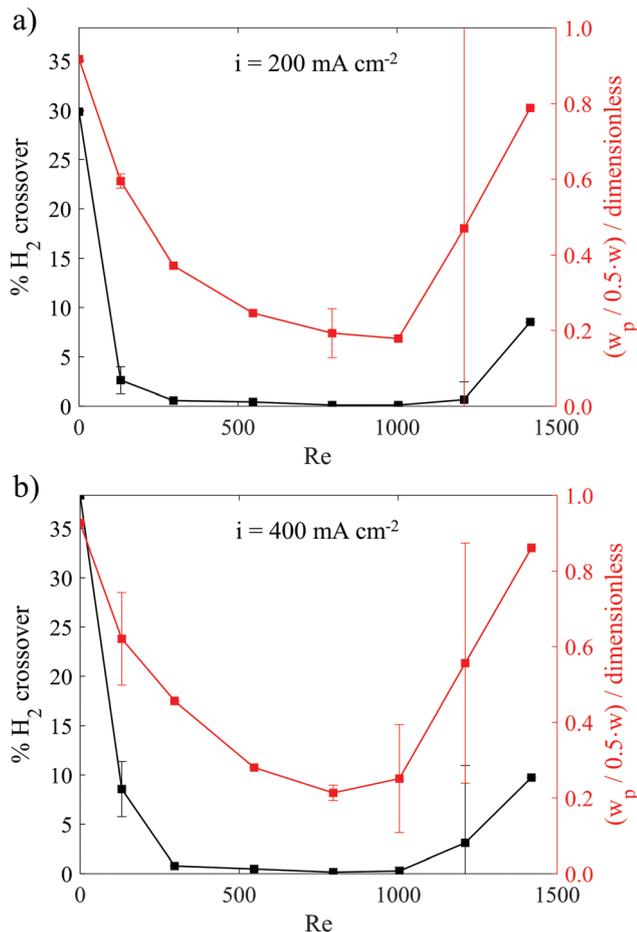


Fig. 5 H<sub>2</sub> crossover and normalized H<sub>2</sub> bubble plume width at different Re for electrolysis at a constant current density of (a) 200 mA cm<sup>-2</sup> and (b) 400 mA cm<sup>-2</sup>. The reported normalized bubble plume widths correspond to the 90% contour lines taken from contour plots such as Fig. 4h at the *y*-position corresponding to the start of the downstream divider that separates the anode and cathode effluent streams.

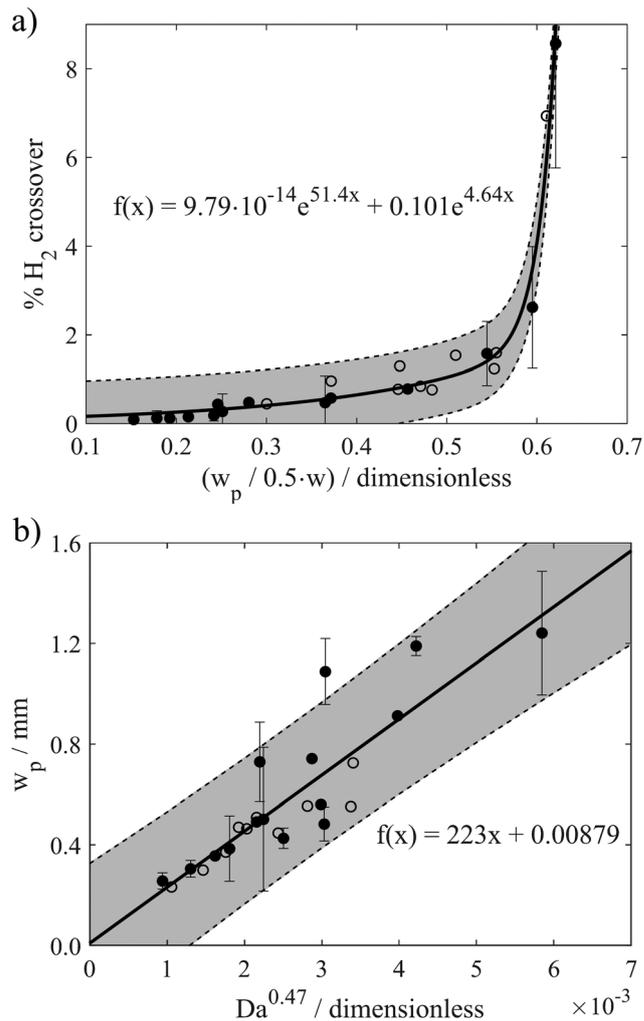


Fig. 6 (a) H<sub>2</sub> crossover as a function of normalized H<sub>2</sub> bubble plume width for 2 different channel widths and all operating conditions in the laminar flow regime with an exponential fitting. (b) H<sub>2</sub> bubble plume width as a function of  $Da^{0.47}$  for different operating conditions and cell geometries. The reported values of  $w_p$  correspond to the 90% contour lines at the *y*-position corresponding to the start of the downstream divider that separates the anode and cathode effluent streams. Data points were obtained for measurements recorded with cells based on  $w = 4$  mm (solid circles) and  $w = 2$  mm (open circles). The solid lines are best fit trendlines obtained using the functional form shown in each sub-figure. The grey regions surrounding the trendlines represent 95% prediction bounds. Error bars were determined by calculating the 95% confidence interval based on 3 repeated experiments.

to the power associated with hydrogen generation ( $P_{H_2} = (1.23 \text{ V}) \times (\text{current})$ ) is estimated to be 5% for the  $w = 4$  mm cell being operated at 400 mA cm<sup>-2</sup> and  $Re = 796$  using a pump efficiency of 75%. The relative pumping power can be minimized by increasing the current density, increasing the electrode size, or decreasing Re, but could become significant for channels characterized by long lengths and narrow channel widths that facilitate larger pressure drop across the cell.

To more explicitly view the relationship between crossover and plume width, Fig. 6a presents the H<sub>2</sub> crossover percent as a function of the dimensionless H<sub>2</sub> bubble plume width ( $X = 90\%$ ) for two different channel widths ( $w = 2$  mm and  $w = 4$  mm) and all operating conditions in the laminar flow regime. The trendline was obtained by fitting a two-term exponential model with equation shown in the subfigure, and the grey regions surrounding the trendline represent 95% prediction bounds for the fitted coefficient. The prediction bounds were calculated using the Curve Fitting Toolbox™ in MATLAB using 95% observation intervals and non-simultaneous bounds. This

correlation contains a total of 24 data points, which were recorded for various combinations of Re and  $i$  ranging between  $131 < Re < 1004$  and  $50 \text{ mA cm}^{-2} < i < 600 \text{ mA cm}^{-2}$ , with the exact combinations of  $w$ ,  $i$ , and Re for every data point provided in Fig. S14 (ESI†). As shown in Fig. S15 (ESI†), data points obtained at very high Re ( $Re > 1004$ ) or in the absence of advection do not agree with the same trendline; hence, the correlation provided in Fig. 6a should only apply to the laminar flow regime. For unsteady fluid flow ( $Re > 1211$ ) or when fluid flow rates are so low that the Segré-Silberberg effect is no longer present, this correlation is not applicable.



For normalized bubble plume widths less than  $\approx 0.5$ , Fig. 6a shows that  $H_2$  crossover rates are low and vary linearly with the normalized plume width. All of the data points obtained within this region correspond to high  $Re$  and/or low  $i$ , and good linearity is also maintained when plume widths are obtained from the  $X = 99\%$ ,  $95\%$  and  $60\%$  contour lines (Fig. S16, ESI<sup>†</sup>). Within this linear portion of the curve, the slope of  $H_2$  crossover *versus* the normalized  $H_2$  bubble plume width is very shallow because the small increases in plume width resulting from operation at higher current densities or lower  $Re$  are insufficient to induce a significant fraction of bubbles to cross the channel centerline. However, the  $H_2$  crossover percent shoots upwards when the  $X = 90\%$  dimensionless plume width ( $w_{p,90\%}/(0.5w)$ ) increases above  $\approx 0.6$ . This portion of the curve is largely populated by data points corresponding to low  $Re$ , and small increases in the plume width lead to large increases in crossover as the outer edge of the plume begins to extend out over the centerline. Non-linear correlations between  $H_2$  crossover and dimensionless plume widths are also observed in Fig. S17 (ESI<sup>†</sup>) based on analysis of the  $X = 99\%$  and  $X = 95\%$  contour lines, but the sharp upticks in crossover rates occur at larger values of ( $w_p/(0.5w)$ ), as expected. For  $X = 99\%$ ,  $H_2$  crossover rises sharply when the normalized  $H_2$  bubble plume width approaches 1, since in this case the  $X = 99\%$  contour corresponds more precisely to the outer edge of the plume. Importantly, all data points obtained for operation of the electrolyzers with two different channel widths in the laminar flow regime collapse onto the same curve, indicating that the dimensionless bubble plume width is a good descriptor that can be used to predict  $H_2$  crossover rates across a wide range of operating conditions for the PPME.

Having established that the bubble plume width is an appropriate descriptor for crossover, it is next necessary to describe how this descriptor depends on cell geometry and key operating conditions. Only then can one map out the performance limits describing the trade-offs between efficiency, current density, electrode size, and product purity. For this study, an additional empirical relationship was used to relate  $w_p$  to key geometric parameters ( $w$ ,  $L_e$ ) and operating conditions ( $i$ , velocity ( $U$ )) that are expected to influence  $w_p$ . According to Buckingham Pi Theorem,<sup>29</sup> this system of 5 variables containing 4 fundamental units (length, time, electric current, amount of substance) can be described with one dimensionless number, or pi grouping. In other words, there should be a universal relationship by which  $w_p$  can be computed from one dimensionless number:

$$w_p = \pi_i^\alpha \quad (1)$$

The dimensionless number relevant to the parallel plate electrolyzer system that can be obtained from the independent variables identified above is the Damkohler number ( $Da$ ), which represents the ratio of the reaction rate to the convective mass transfer rate and is given by eqn (2) for the parallel plate electrolyzer configuration:

$$Da = \frac{i \cdot L_e \cdot H}{C_{\text{solvent}} \cdot U \cdot w \cdot h} \quad (2)$$

where  $n$  is the number of electrons for HER,  $F$  is Faraday's constant,  $C_{\text{solvent}}$  is the concentration of the solvent ( $55.5 \text{ M}$ ) and  $h$  is the cell depth ( $5 \text{ mm}$ ).

Using HSV-derived bubble plume widths at different  $Re$ ,  $w$ , and  $i$ , a best fitting procedure was used to obtain a value of the constant in eqn (1) to be  $\alpha = 0.47$ , with the best fit curve provided in Fig. 6b. This relationship between  $w_p$  and  $Da^\alpha$  was obtained for 25 different data points recorded for two different channel widths and  $Re \leq 1004$ . Based on eqn (2),  $Da$  increases with increasing current and/or decreasing flow rate or  $Re$ . Since the bubble plume width also increases with current and decreases with flow rate, consistent with the HSV observations of Fig. 5,  $w_p$  increases with  $Da$ . In this linear relationship between  $w_p$  and  $Da^{0.47}$ , the lowest flow rate ( $Re = 131$ ) data points are the ones contributing the most scatter (Fig. S18, ESI<sup>†</sup>), which we attribute to a small pulsation effect from the peristaltic pump at the lowest flow rate. Using the empirical correlation shown in Fig. 6b,  $w_p$  can be predicted for any combination of electrode size ( $L_e \times H$ ), current density ( $i$ ), and  $Re$  so long as the electrolyzer is operating in the laminar flow regime. By extension, the predicted value of  $w_p$  obtained from this correlation can then be used in the first correlation shown in Fig. 6a to predict the percent  $H_2$  crossover at different current densities, channel widths, and electrode sizes. Although these empirical correlations were obtained from *in situ* HSV monitoring of bubble plumes in this study, an opportunity for future research efforts is to use computational fluid dynamics (CFD) to generate similar correlations between plume width, cell geometry, operating current density, and electrolyte flow rates.

### Predicting the performance limits of parallel plate membraneless electrolyzers

**Model description.** In this section, we combine a simple one-dimensional (1D) model describing the electrolyzer current-voltage ( $I$ - $V$ ) characteristics with the empirical correlations developed in the previous section to evaluate the trade-offs between the four key performance metrics shown in the unfortunate tetrahedron of Fig. 1b (opex, capex, current density, and product purity). Briefly, the 1D model employed in this study is used to predict electrolyzer efficiency as a function of operating current density for a given electrode size, channel width, and electrolyte conductivity. The electrolyzer efficiency ( $\eta_E$ ) is calculated as a voltage efficiency and given by the ratio of the reversible cell potential for the water splitting reaction at room temperature ( $|\Delta E^\circ| = 1.23 \text{ V}$ , based on  $\Delta G^\circ = 237 \text{ kJ mol}^{-1} H_2$ ) to the total cell voltage required to sustain water electrolysis at a given current density,  $\Delta V(i)$ :

$$\eta_E = \frac{|\Delta E^\circ|}{\Delta V(i)} \quad (3)$$

$\Delta V(i)$  is given by the sum of  $|\Delta E^\circ|$  and the absolute value of overpotential losses:

$$\Delta V = |\Delta E^\circ| + |\eta_{\text{HER}}| + \eta_{\text{OER}} + \eta_{\Omega} + \eta_{\text{conc}} \quad (4)$$



where  $\eta_{\text{HER}}$  and  $\eta_{\text{OER}}$  are the kinetic overpotentials associated with the HER and OER, respectively. The fourth term on the right side of eqn (4),  $\eta_{\Omega}$ , gives the ohmic overpotential losses due to ionic transport through the electrolyte and electron transport through the external circuit. Ohmic drops associated with the external circuit were neglected in this analysis since electronic conductivity is often very large compared to electrolyte conductivity ( $\kappa$ ). The final term in eqn (4),  $\eta_{\text{conc}}$ , is the concentration overpotential associated with concentration gradients at the cathode and/or anode. The following analysis assumes that  $\eta_{\text{conc}}$  is negligible, which is usually reasonable for water electrolysis in strong acids and bases. For this study,  $\eta_{\text{HER}}$  and  $\eta_{\text{OER}}$  were modeled using the Tafel equations:

$$\eta_{\text{HER}} = \beta_{\text{HER}} \log\left(\frac{i}{i_{0,\text{HER}}}\right) \quad (5)$$

$$\eta_{\text{OER}} = \beta_{\text{OER}} \log\left(\frac{i}{i_{0,\text{OER}}}\right) \quad (6)$$

where  $\beta$  is the Tafel slope,  $i$  is the current density, and  $i_0$  is the exchange current density. In the following analysis of PPME performance, kinetic overpotential losses were modeled based on Tafel parameters reported in literature for state-of-the-art electrocatalysts. Kinetic parameters for Pt HER electrocatalyst and iridium oxide (IrO<sub>2</sub>) OER electrocatalyst were used for acidic environment, while Tafel parameters for nickel (Ni) HER electrocatalyst and (nickel, iron) oxyhydroxide ((Ni,Fe)OOH) OER electrocatalyst were used for alkaline environment. In acid, values of  $\beta = 32 \text{ mV dec}^{-1}$  and  $i_0 = 1.3 \times 10^{-3} \text{ A cm}^{-2}$  were used for the HER electrocatalyst,<sup>30</sup> while values of  $\beta = 43 \text{ mV dec}^{-1}$  and  $i_0 = 7.6 \times 10^{-10} \text{ A cm}^{-2}$  were used for the OER electrocatalyst.<sup>31</sup> In alkaline media, values of  $\beta = 97.6 \text{ mV dec}^{-1}$  and  $i_0 = 7.0 \times 10^{-5} \text{ A cm}^{-2}$  were used for the HER electrocatalyst,<sup>32</sup> while values of  $\beta = 41.5 \text{ mV dec}^{-1}$  and  $i_0 = 9.9 \times 10^{-6} \text{ A cm}^{-2}$  were used for the OER electrocatalyst.<sup>33</sup> These kinetic parameters were normalized by the electrochemical surface area and were assumed to remain constant within the current density range investigated and independent of the fluid flow rate. To increase the accuracy of predicting kinetic overpotential losses over a wide range of operating conditions, especially for materials that are known to exhibit multiple Tafel slopes, more advanced kinetic and transport models are required.<sup>34</sup>  $\eta_{\Omega}$  associated with ion transport through the electrolyte was calculated using Ohm's law:

$$\eta_{\Omega} = I \cdot R_s \quad (7)$$

where  $R_s$  is the solution resistance and the total current  $I$  is obtained as the product of the current density and the electrode area ( $A_e$ ). For an electrolyzer with infinitely large parallel plate electrodes and uniform electrolyte composition between them,  $R_s$  is given by:

$$R_s = \frac{\rho \cdot w}{A_e} \quad (8)$$

where  $\rho$  is the resistivity of the electrolyte, which is inversely proportional to  $\kappa$ . For very large electrodes, the ionic current flowing between the two electrodes is uniform and follows

conduction paths that are perpendicular to the electrode surfaces. However, electrodes with lengths that are similar to or less than the characteristic length for ion transport (*i.e.*  $w$ ) can exhibit large edge effects whereby potential fields and the associated ion conduction pathways extend out beyond the edge of the electrodes such that the total resistance is less than that given by eqn (8). In this study, finite element simulations were carried out in COMSOL to account for these edge effects and compute  $R_s$  for electrodes with varying electrode length ( $L_e$ ) and cells with varying  $w$ . Fig. S19 (ESI†) shows the computed values of  $R_s$  for four different values of  $L_e$  as a function of  $w$  for constant  $H = 3.34 \text{ mm}$  and  $\kappa = 0.226 \text{ S cm}^{-1}$ . As expected,  $R_s$  is small and varies linearly with  $w$  for large values of  $L_e$ , but varies logarithmically with  $w$  for small  $L_e$  due to edge effects. It should be noted that the presence of gas bubbles between the electrodes can also decrease the effective solution conductivity,<sup>35</sup> but was neglected in this study since calculations based on the Bruggeman correlation show that the additional ohmic drop due to bubbles is very small compared to  $\eta_{\Omega}$  in the bubble-free solution (SI Section S20, ESI†).

Using experimentally measured solution resistance, Tafel parameters for Pt HER electrocatalyst<sup>30</sup> and fitted Tafel parameters for Pt OER electrocatalyst, this simple 1D model was found to give good agreement with experimental current density–voltage curves recorded for the PPME with Pt electrodes in 0.5 M H<sub>2</sub>SO<sub>4</sub>, as seen in Fig. S21 (ESI†).

**Analyzing trade-offs between current density, efficiency, and product purity.** To explore the performance limits of a parallel plate electrolyzer, the 1D electrolyzer model described above was used to predict electrolyzer efficiency as a function of current density and  $w$  for Pt and IrO<sub>2</sub> electrodes in acidic environment and for Ni and (Ni,Fe)OOH electrodes under alkaline condition, having the same dimensions as those used in the experimental work ( $L_e = 0.9 \text{ cm}$  and  $H = 0.334 \text{ cm}$ ). The results of these calculations are shown in Fig. 7a and b for electrolyte conductivities 0.226 S cm<sup>-1</sup> and 0.621 S cm<sup>-1</sup>, corresponding to 0.5 M H<sub>2</sub>SO<sub>4</sub> and 30 wt% KOH, respectively. For reference, 30 wt% KOH is typically the electrolyte used in conventional alkaline electrolyzers,<sup>36</sup> while the pH of 0.5 M H<sub>2</sub>SO<sub>4</sub> is consistent with that of Nafion solid electrolyte membranes used in conventional PEM electrolyzers.<sup>37</sup> Additionally, the HSV-derived empirical correlations shown in Fig. 6 were used to predict H<sub>2</sub> crossover rates at each current density and channel width. For this analysis, electrolyte flow rates were chosen to correspond to  $Re = 796$ , which was shown in this study to result in the narrowest bubble plume width, and hence, the lowest crossover rate at most current densities. The results of these calculations are provided as black curves in Fig. 7, which show the combinations of  $i$  and  $w$  that give constant H<sub>2</sub> crossover rates of 4%, 1%, and 0.5%. For reference, a H<sub>2</sub> crossover rate of 2.06% translates to 4 % H<sub>2</sub> in the anode effluent stream, the lower flammability limit for H<sub>2</sub> in pure O<sub>2</sub>, for the case where the molar O<sub>2</sub> crossover rate is exactly one half that of the H<sub>2</sub> crossover rate (Fig. S4a, ESI†).

As expected, the highest efficiencies are predicted at low current densities and small  $w$ , which minimize kinetic and



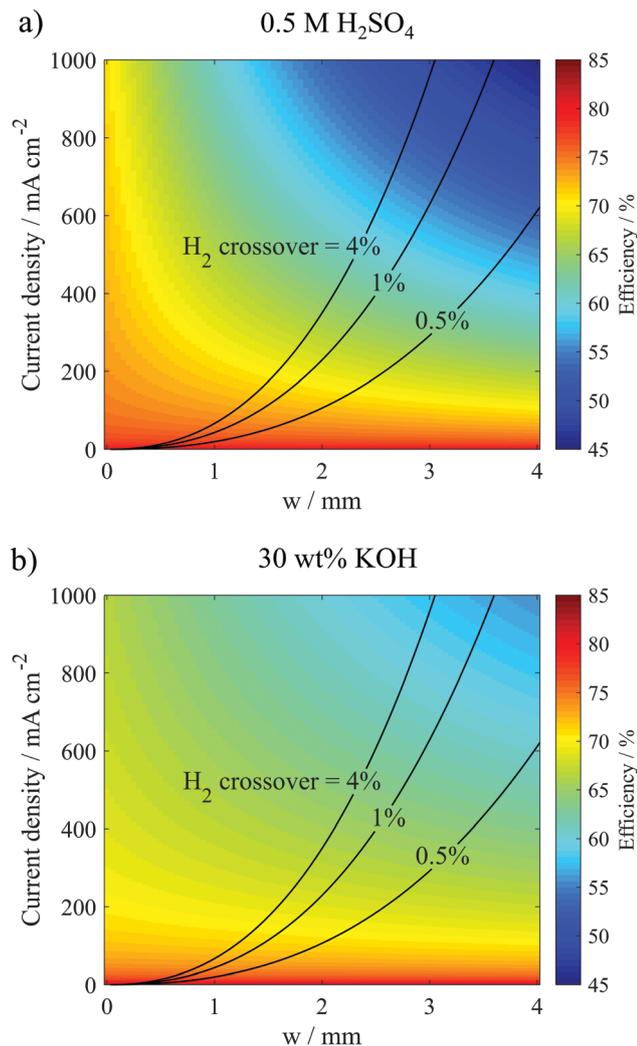


Fig. 7 Relation between current density, channel width and electrolysis efficiency for a PPME ( $L_e = 9$  mm,  $H = 3.34$  mm,  $L_d = 12$  mm) operated in (a) 0.5 M  $H_2SO_4$  and (b) 30 wt% KOH. The contour lines show different predicted  $H_2$  crossover rates based on the empirical correlations between bubble plume width, operating current density, and cross-over rates. Efficiencies are reported on the basis of  $\Delta G^\circ = 237$  kJ mol $^{-1}$   $H_2$  and were modeled based on Ohm's law and Tafel equations.

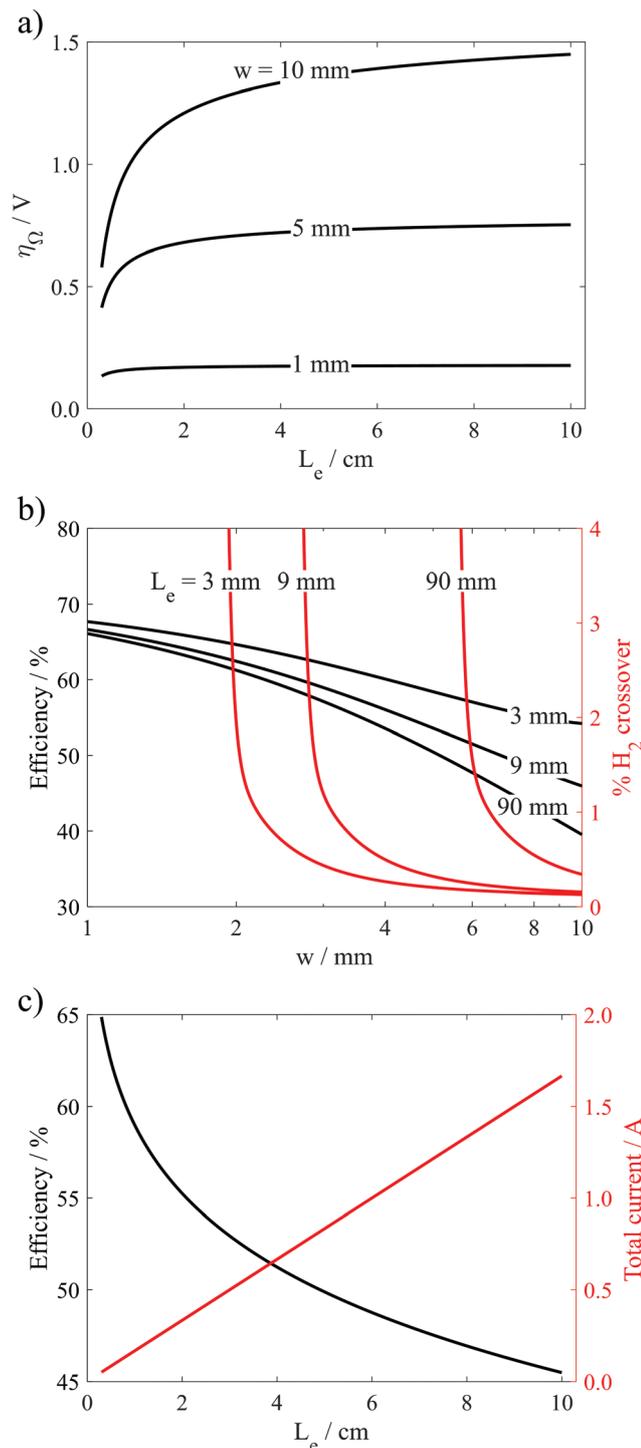
ohmic overpotential losses, respectively. Importantly, Fig. 7a shows that  $w$  needs to be less than 1.8 mm to achieve electrolyzer efficiencies  $> 65\%$  (wrt  $\Delta G^\circ = 237$  kJ mol $^{-1}$   $H_2$ ) at current densities  $\geq 400$  mA cm $^{-2}$  in the 0.5 M  $H_2SO_4$  electrolyte. In 30 wt% KOH, where the electrolyte conductivity is 1.75 times higher than that in the 0.5 M  $H_2SO_4$  electrolyte,  $w$  can be increased to much larger values to achieve the same efficiency and current density. For example,  $w$  only needs to be less than 2.6 mm to achieve electrolyzer efficiencies  $> 65\%$  at current densities  $\geq 400$  mA cm $^{-2}$ . Clearly, the higher conductivity electrolyte can enable operation at much higher current density and efficiency. However, these high efficiency/high current density regions in Fig. 7 also correspond to the highest crossover rates, clearly illustrating the trade-offs between current density, efficiency, and product purity. From a safety standpoint,

it is imperative that the composition of the product gases does not exceed the upper (94%  $H_2$  in  $O_2$ ) or lower (4%  $H_2$  in  $O_2$ ) flammability limits.<sup>38</sup> Thus, Fig. 7 allows one to predict the highest current density that can be safely achieved while maintaining some minimum desired electrolysis efficiency. For example, Fig. 7b predicts that the parallel plate electrolyzer with  $L_e = 0.9$  cm and  $H = 0.334$  cm operating at a desired electrolyzer efficiency of 60% can operate at current densities up to 795 mA cm $^{-2}$  in 30 wt% KOH before  $H_2$  crossover increases above 4%, and this maximum current density can be achieved at  $w = 2.8$  mm.

**Analyzing trade-offs between electrode size, efficiency, and product purity.** Another important parameter in electrolyzer design that can have a large impact on capex is the electrode size. Along with the operating current density, the electrode size also determines the throughput or total production rate of  $H_2$  per cell. In general, it is desirable to maximize throughput per cell and increase the scale of individual electrodes, the latter of which is often beneficial from a manufacturing standpoint.<sup>39</sup> Using the same modeling framework employed to generate Fig. 7, the effect of electrode length ( $L_e$ ) on the efficiency/cross-over trade-off was also investigated. These calculations were run for different values of electrode length ( $L_e$ ) between 0.3 cm and 10 cm and various electrode gap distances ranging from 1 mm to 10 mm for a constant electrode height of  $H = 3.34$  mm. It should be noted that the depth of the electrode can also be increased in order to increase the electrode size, and bubble plume width won't be affected by the electrode depth. Hence, increasing the depth of electrode shouldn't affect  $H_2$  purity based on the correlation in Fig. 6a, but can further increase the electrode size, and hence increase total current/throughput. The electrodes were again assumed to have kinetics consistent with Pt and  $IrO_2$  in acid, and the electrolyte conductivity was set to  $\kappa = 0.226$  S cm $^{-1}$ , corresponding to the electrolyte conductivity associated with 0.5 M  $H_2SO_4$ . In all cases, the current density was set to  $i = 500$  mA cm $^{-2}$ .

Results from these calculations are summarized in Fig. 8, which reveal that the electrode length has two primary influences on the performance limits of the PPME. First, increasing  $L_e$  reduces the "edge effects" associated with ion transport around the ends of the electrodes, leading to increases in the ohmic overpotential losses. Fig. 8a shows that this effect is amplified at large electrode gap distances where ohmic losses can dominate overall cell performance. Consistent with these observations, the black curves in Fig. 8b show that  $\eta_E$  decreases with increasing  $L_e$ , and that  $\eta_E$  becomes most sensitive to  $L_e$  at large  $w$ . The second influence of  $L_e$  on electrolyzer performance is that increasing  $L_e$  leads to increased crossover and reduced product purity. In Fig. 8b (red curves), the %  $H_2$  crossover was predicted for different  $L_e$  and  $w$  as described in the previous section, again assuming a constant  $Re = 796$ . For all three values of  $L_e$ , low  $H_2$  crossover rates are predicted for very large  $w$ , but each curve exhibits a critical value of  $w$  below which crossover increases exponentially. This critical  $w$  increases with  $L_e$ , which is explained by the fact that longer electrodes generate more  $H_2$ , and therefore wider bubble plumes at the





**Fig. 8** (a) Ohmic losses as a function of electrode length for different channel widths for a constant current density of  $500 \text{ mA cm}^{-2}$  and an electrolyte of  $0.5 \text{ M H}_2\text{SO}_4$ . (b) Efficiency and  $\text{H}_2$  crossover as a function of channel width for different electrode lengths for a PPME ( $H = 3.34 \text{ mm}$ ,  $L_d = 12 \text{ mm}$ ) operated in  $0.5 \text{ M H}_2\text{SO}_4$  and a constant current density of  $500 \text{ mA cm}^{-2}$ . (c) Efficiency and total current as a function of electrode length for constant product purity ( $1\% \text{ H}_2$  crossover) and current density ( $500 \text{ mA cm}^{-2}$ ) in  $0.5 \text{ M H}_2\text{SO}_4$  electrolyte.

downstream edge of the electrodes. In order to maintain a target crossover rate, a longer electrode length will require a

wider channel width to accommodate the wider bubble plume width. As shown in Fig. 8c for a constant predicted  $\text{H}_2$  crossover rate of  $1\%$ ,  $\eta_E$  decreases sharply with increasing  $L_e$ , but the throughput/total current increases linearly with electrode length because of the increasing electrode size. As a result, the maximum efficiency associated with safe operation of the electrolyzer at a constant current density decreases with increasing  $L_e$ , highlighting the efficiency/scalability/product purity trade-off that is commonly observed for membraneless electrolyzers.

## Discussion

By using a combination of *in situ* high speed videography and electrochemical engineering principles, this work has analyzed the trade-offs between efficiency, current density, electrode size, and product purity for a simple parallel plate membraneless electrolyzer. By establishing quantitative relationships between these key performance metrics and geometric design parameters, we have created a framework for exploring the performance limits of this electrolyzer architecture. Optimizing any one of these performance metrics is trivial, but the overall performance of the electrolyzer should be viewed in the context of balancing the trade-offs between the performance metrics illustrated in Fig. 1b, with the relative importance of each metric weighted according to the particular application for which the electrolyzer is being designed. Nonetheless, it is valuable to compare the predicted multi-metric performance limits of the parallel plate electrolyzer under favorable conditions to the typical performance metrics of commercial alkaline and PEM electrolyzers, which are presented in Table 1. As a basis for this comparison, we have considered the predicted performance limits presented in Fig. 7b for a membraneless electrolyzer with the same electrode dimensions presented in this study ( $L_e = 9 \text{ mm}$ ,  $H = 3.34 \text{ mm}$ ) operating in  $30 \text{ wt}\% \text{ KOH}$ . If the target  $\text{H}_2$  crossover rate is set to  $0.5\%$  (corresponding to  $\text{H}_2$  purity  $\geq 99.75\%$ ), the PPME is predicted to be capable of current densities between  $0.20\text{--}0.61 \text{ A cm}^{-2}$  with corresponding efficiencies of  $68.4\text{--}60.4\%$  (based on  $\Delta G^\circ = 237 \text{ kJ mol}^{-1} \text{ H}_2$ ), respectively. If the  $\text{H}_2$  cross-over constraint is relaxed to  $1\%$  cross-over (corresponding to  $\text{H}_2$  purity  $\geq 99.5\%$ ), the current density of the electrolyzer can approach  $0.73 \text{ A cm}^{-2}$  while maintaining a  $60\%$  efficiency, where this current density is  $1.8$  times as high as the upper limit for conventional alkaline electrolyzers at comparable efficiency. Meanwhile, a comparison of achievable current density and efficiency ranges for the

**Table 1** State-of-the-art for the performance metrics of commercial alkaline and PEM electrolyzers<sup>4,7,40</sup>

Electrolyzer type	Current density ( $\text{A cm}^{-2}$ )	Efficiency (% wrt $\Delta G^\circ$ )	$\text{H}_2$ purity (%)
Conventional alkaline	0.2–0.4	51.5–68.1	> 99.5
PEM	0.6–2.0	55.6–68.1	> 99.999



membraneless electrolyzer at safe cross-over levels are always expected to be inferior to the PEM electrolyzer.

The prediction that the H<sub>2</sub> purity/efficiency/current density performance limits of the PPME can exceed those of conventional alkaline electrolyzers is explained primarily by the fact that the former can achieve lower solution resistance thanks to the absence of a diaphragm divider and ability to place the electrodes closer together. Even when the electrodes in the PPME are increased to  $L_e = 9$  cm, the PPME still maintains a slight advantage, being predicted to achieve efficiencies of 61.6% and 58.5% with 1% H<sub>2</sub> crossover while operating at 200 and 400 mA cm<sup>-2</sup>, respectively. It should be noted that the optimal space between electrodes for an alkaline electrolyzer operating at a given current density is determined to be that which maximizes cell efficiency,<sup>41</sup> while the optimal electrode spacing for a PPME is the smallest spacing that still enables the required product purity to be achieved. Although it was not considered explicitly in this study, the forced advection used in the PPME is also expected to reduce bubble-induced losses that are known to be especially problematic for conventional alkaline electrolyzers at high current densities.<sup>42,43</sup> Meanwhile, the unfavorable current density/efficiency performance limits of the PPME compared to a PEM electrolyzer can be attributed to the intrinsically lower total electrolyte resistance associated with the “zero-gap” nature of the latter. However, this analysis does not take into consideration the capital costs associated with the electrolyzer materials and costs of construction. In general, the relative simplicity of membraneless electrolyzer architectures based on a limited number of component types and manufacturing steps creates an opportunity to reduce manufacturing costs compared to a PEM electrolyzer, although this advantage may be less significant when compared to alkaline electrolyzers. The size of electrodes is another important consideration for manufacturing, keeping in mind that cell areas for commercial alkaline and PEM electrolyzers have been reported to be up to 4 m<sup>2</sup> and 3 m<sup>2</sup>, respectively.<sup>40</sup> Ultimately, engineering analysis of electrolyzers, such as that presented in this work, needs to be combined with manufacturing cost models and technoeconomic models of entire electrolyzer systems to predict the levelized cost of hydrogen production under specific electricity pricing schemes. On average, 40–50% of the total capital cost of a commercial electrolyzer system is associated with the cost of balance of system (BOS) components.<sup>8</sup> Thus, it is important to consider not only the cost and performance of the electrolyzer stack itself, but also how the electrolyzer performance impacts the cost, sizing, and performance of BOS components within the larger system.

It should be emphasized that the analysis provided in this article is specific to a parallel plate electrolyzer with electrode morphology and bubble dynamics consistent with those used in this study. Thus, the performance limits identified in this analysis are not true upper limits to performance. Several additional modifications are anticipated to push the performance bounds even further. For example, using more textured electrodes are anticipated to help (i) reduce kinetic overpotential losses thanks to the higher catalytic surface area and

(ii) facilitate smaller bubble departure diameters and the associated bubble size distribution associated with the bubble plumes evolved from the downstream electrodes.<sup>25</sup> Reducing the average bubble size is also expected to help decrease kinetic overpotential losses by decreasing the average bubble coverage on the electrode surface, and may also alter the plume width by enabling tighter packing of the smaller bubbles within the bubble plume. Anecdotally, we have observed that the occasional build-up and detachment of “big bubbles” with radii that are substantially larger than the average bubble radius can preferentially lead to cross-over events by promoting flow disturbances that lead to crossover of other smaller bubbles. A sequence of still frames capturing one such “big bubble” departure event is provided in Fig. S22 (ESI<sup>†</sup>), with the associated video provided as Video S1 (ESI<sup>†</sup>). It should also be noted that the product purity is underestimated in this study because the baffle separating the anode and cathode effluent streams was not located immediately at the end of the electrodes, which would correspond to  $L_d = 0$ . An optimized PPME design would have  $L_d = 0$  in order to minimize crossover. Collectively, changes in electrode morphology, entrance effects, divider position, or gas/electrolyte surface tension that affect bubble size distribution or dynamics can all be expected to lead to slightly different correlations than those provided in Fig. 6. This study has shown that *in situ* HSV offers an attractive means of generating these correlations. However, multi-phase computational fluid dynamics (CFD) represents a huge opportunity to generate these correlations without ever constructing a cell and for cell geometries that may not be amenable to *in situ* HSV.

In future research efforts, the use of *in situ* HSV and CFD, combined with electrochemical engineering design principles, represents a well-rounded toolbox for further exploring the performance limits of membraneless electrolyzers characterized by a far broader design space than that explored in this study. For example, membraneless electrolyzers based on flow-through electrodes offer an even richer design space that is characterized by more geometric parameters. As the number of design parameters increase and advances in additive manufacturing enable rapid prototyping of electrochemical reactors,<sup>44,45</sup> the automated discovery and exploration of new membraneless electrolyzer architectures with the help of artificial intelligence tools like machine learning<sup>46,47</sup> and genetic algorithms<sup>48</sup> also present new opportunities for advancing the performance limits of membraneless electrolyzers. Given the urgency of developing high performance, low-cost electrolysis technologies for a wide range of clean energy applications, the coordinated use of these emerging computational and experimental tools may play a critical role in accelerating the development of membraneless electrolyzer technologies to the point where they can make a meaningful impact in the near future.

## Conclusions

This study has demonstrated a method of using *in situ* HSV analysis to quantify H<sub>2</sub> bubble plume width at different Re,



current densities, and channel widths in parallel plate membraneless electrolyzers. The HSV-derived dimensionless bubble plume width is found to strongly correlate with H<sub>2</sub> crossover, and hence can be a useful descriptor for relating operating conditions and cell geometries to crossover rates and product purities. Combined with electrochemical engineering design principles, the trade-offs between efficiency, current density, electrode size, and product purity were investigated for membraneless electrolyzers. Furthermore, the performance limits were analyzed and compared with the typical operating performance of commercial water electrolyzers. This analysis shows that it is possible for membraneless electrolyzers based on the parallel plate design to exceed the performance of alkaline electrolyzers, but that they are unlikely to be able to match all of the performance metrics of PEM electrolyzers. Importantly, this work has presented a combined experimental and modeling framework that can be extended to guide the design of many other membraneless electrolyzer architectures and analyze their performance limits.

## Conflicts of interest

DVE is a co-founder of sHYp B.V.

## Acknowledgements

The authors acknowledge funding from Shell International Exploration & Production New Energies Research & Technology Dense Energy Carriers Program. The authors also acknowledge Ziyang Mao for his contributions to HSV analysis.

## References

- 1 D. Gielen, F. Boshell, D. Saygin, M. D. Bazilian, N. Wagner and R. Gorini, *Energy Strateg. Rev.*, 2019, **24**, 38–50.
- 2 M. Liserre, T. Sauter and J. Y. Hung, *IEEE Ind. Electron. Mag.*, 2010, **4**, 18–37.
- 3 K. Ayers, N. Danilovic, R. Ouimet, M. Carmo, B. Pivovar and M. Bornstein, *Annu. Rev. Chem. Biomol. Eng.*, 2019, **10**, 219–239.
- 4 M. Carmo and D. L. Fritz, *Int. J. Hydrogen Energy*, 2013, **8**, 4901–4934.
- 5 K. Zeng and D. Zhang, *Prog. Energy Combust. Sci.*, 2010, **36**, 307–326.
- 6 G. Maggio, A. Nicita and G. Squadrito, *Int. J. Hydrogen Energy*, 2019, **44**, 11371–11384.
- 7 C. Xiang, M. Papadantonakis and N. S. Lewis, *Mater. Horiz.*, 2016, **3**, 169–173.
- 8 D. V. Esposito, *Joule*, 2017, **1**, 651–658.
- 9 S. M. H. Hashemi, M. A. Modestino and D. Psaltis, *Energy Environ. Sci.*, 2015, **8**, 2003–2009.
- 10 M. I. Gillespie, F. Van Der Merwe and R. J. Kriek, *J. Power Sources*, 2015, **293**, 228–235.
- 11 S. M. H. Hashemi, P. Karnakov, P. Hadikhani, E. Chinello, S. Litvinov, C. Moser, P. Koumoutsakos and D. Psaltis, *Energy Environ. Sci.*, 2019, **12**, 1592–1604.
- 12 G. Segre and A. Silberg, *Nature*, 1961, **189**, 209–210.
- 13 H. Amini, W. Lee and D. Di Carlo, *Lab Chip*, 2014, **14**, 2739–2761.
- 14 P. Hadikhani, S. M. H. Hashemi, G. Balestra, L. Zhu, M. A. Modestino, F. Gallaire and D. Psaltis, *Lab Chip*, 2018, **18**, 1035–1046.
- 15 G. D. O'Neil, C. D. Christian, D. E. Brown and D. V. Esposito, *J. Electrochem. Soc.*, 2016, **163**, F3012–F3019.
- 16 M. M. Monroe, P. Lobaccaro, Y. Lum and J. W. Ager, *J. Phys. D: Appl. Phys.*, 2017, **50**, 154006.
- 17 I. Holmes-Gentle, F. Hoffmann, C. A. Mesa and K. Hellgardt, *Sustainable Energy Fuels*, 2017, **1**, 1184–1198.
- 18 J. T. Davis, J. Qi, X. Fan, J. C. Bui and D. V. Esposito, *Int. J. Hydrogen Energy*, 2018, **43**, 1224–1238.
- 19 M. I. Gillespie and R. J. Kriek, *J. Power Sources*, 2017, **372**, 252–259.
- 20 D. Steingart, The Unfortunate Tetrahedron, <https://medium.com/the-unfortunate-tetrahedron/the-unfortunate-tetrahedron-ce1e44d0b961>.
- 21 J. T. Davis, D. E. Brown, X. Pang and D. V. Esposito, *J. Electrochem. Soc.*, 2019, **166**, F312–F321.
- 22 H. Y. Lee, C. Barber and A. R. Minerick, *Electrophoresis*, 2014, **35**, 1782–1789.
- 23 D. Fernández, P. Maurer, M. Martine, J. M. D. Coey and M. E. Möbius, *Langmuir*, 2014, **30**, 13065–13074.
- 24 D. V. Esposito, Y. Lee, H. Yoon, P. M. Haney, N. Y. Labrador, T. P. Moffat, A. A. Talin and V. A. Szalai, *Sustainable Energy Fuels*, 2017, **1**, 154–173.
- 25 A. Angulo, P. van der Linde, H. Gardeniers, M. Modestino and D. Fernández Rivas, *Joule*, 2020, **4**, 555–579.
- 26 E. R. Choban, L. J. Markoski, A. Wieckowski and P. J. A. Kenis, *J. Power Sources*, 2004, **128**, 54–60.
- 27 N. de Nevers, *Fluid mechanics for chemical engineers*, McGraw-Hill, 3rd edn, 1999.
- 28 A. Taqieddin, R. Nazari, L. Rajic and A. Alshawabkeh, *J. Electrochem. Soc.*, 2017, **164**, E448–E459.
- 29 E. Buckingham, *Phys. Rev.*, 1914, **4**, 345.
- 30 H. Kita, S. Ye and Y. Gao, *J. Electroanal. Chem.*, 1992, **334**, 351–357.
- 31 Y. Zhao, N. M. Vargas-Barbosa, E. A. Hernandez-Pagan and T. E. Mallouk, *Small*, 2011, **7**, 2087–2093.
- 32 C. González-Buch, I. Herraiz-Cardona, E. Ortega, J. García-Antón and V. Pérez-Herranz, *J. Appl. Electrochem.*, 2016, **46**, 791–803.
- 33 H. Zhou, F. Yu, Q. Zhu, J. Sun, F. Qin, L. Yu, J. Bao, Y. Yu, S. Chen and Z. Ren, *Energy Environ. Sci.*, 2018, **11**, 2858–2864.
- 34 T. Shinagawa, A. T. Garcia-Esparza and K. Takanabe, *Sci. Rep.*, 2015, **5**, 1–21.
- 35 M. Philippe, H. Jérôme, B. Sebastien and P. Gérard, *Electrochim. Acta*, 2005, **51**, 1140–1156.
- 36 L. M. Gandía, R. Oroz, A. Ursúa, P. Sanchis and P. M. Diéguez, *Energy Fuels*, 2007, **21**, 1699–1706.



- 37 M. Paidar, V. Fateev and K. Bouzek, *Electrochim. Acta*, 2016, **209**, 737–756.
- 38 M. T. Janicke, H. Kestenbaum, U. Hagendorf, F. Schüth, M. Fichtner and K. Schubert, *J. Catal.*, 2000, **191**, 282–293.
- 39 E. A. K. Ayers and C. Capuano, *ECS Trans.*, 2012, **41**, 15–22.
- 40 H. A. Miller, K. Bouzek, J. Hnat, S. Loos, C. I. Bernäcker, T. Weißgärber, L. Röntzsch and J. Meier-Haack, *Sustainable Energy Fuels*, 2020, **4**, 2114–2133.
- 41 N. Nagai, M. Takeuchi, T. Kimura and T. Oka, *Int. J. Hydrogen Energy*, 2003, **28**, 35–41.
- 42 B. E. Bongenaar-Schlenter, L. J. J. Janssen, S. J. D. Van Stralen and E. Barendrecht, *J. Appl. Electrochem.*, 1985, **15**, 537–548.
- 43 F. Hine, *J. Electrochem. Soc.*, 1980, **127**, 292.
- 44 J. C. Ruiz-Morales, A. Tarancón, J. Canales-Vázquez, J. Méndez-Ramos, L. Hernández-Afonso, P. Acosta-Mora, J. R. Marín Rueda and R. Fernández-González, *Energy Environ. Sci.*, 2017, **10**, 846–859.
- 45 S. M. H. Hashemi, U. Babic, P. Hadikhani and D. Psaltis, *Curr. Opin. Electrochem.*, 2020, **20**, 54–59.
- 46 D. E. Blanco, B. Lee and M. A. Modestino, *Proc. Natl. Acad. Sci. U. S. A.*, 2019, **116**, 17683–17689.
- 47 L. Mao and L. Jackson, *Proceedings of the International MultiConference of Engineers and Computer Scientists*, 2016, vol. 1, p. 825.
- 48 K. Chen, S. Laghrouche and A. Djerdir, *Appl. Energy*, 2019, **252**, 113439.

