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Garnet solid state electrolytes have been considered as potential candidates to enable next generation all solid state batteries (ASSBs). To facilitate the practical application of ASSBs, a high room temperature ionic conductivity and a low interfacial resistance between solid state electrolyte and electrodes are essential. In this work, we report a study of cerium doped $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ through X-ray pair distribution function analysis, impedance spectroscopy and electrochemical testing. The successful cerium incorporation was confirmed by both X-ray diffraction refinement and X-ray pair distribution function analysis, showing the formation of an extensive solid solution. The local bond distances for Ce and Nb on the octahedral site were determined using X-ray pair distribution function analysis, illustrating the longer bond distances around Ce. This Ce doping strategy was shown to give a significant enhancement in conductivity ($1.4 \times 10^{-4} \text{ S cm}^{-1}$ for $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$, which represents one of the highest conductivities for a garnet with less than 6 Li) as well as a dramatically decreased interfacial resistance ($488 \Omega \text{ cm}^2$ for $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$). In order to demonstrate the potential of this doped system for use in ASSBs, the long term cycling of a Li//garnet//Li symmetric cell over 380 h has been demonstrated.

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Introduction

State-of-the-art lithium ion batteries (LIBs) have dominated the energy storage market for more than 2 decades because of their long cycle life and high energy and power densities over other battery systems.^{1–3} Nowadays, the increasing demand in developing electrical vehicles (EVs) not only needs higher energy density, but more importantly, requires improved safety properties.⁴ The flammable organic solvent in LIBs can trigger safety issues, and so makes it a non-ideal system for EVs. Instead, all-solid-state batteries (ASSBs) using a non-flammable inorganic solid-state electrolyte are considered to be a promising candidate to address the safety issues for future EVs.⁵

Garnet lithium ion conductors have attracted increasing interest in the last decade due to their relatively good electrochemical stability against Li metal/cathode materials, and high room temperature lithium ionic conductivity compared to other solid state electrolyte materials.^{6–8} The ideal garnet framework has the chemical formula $\text{A}_3\text{B}_2\text{C}_3\text{O}_{12}$, where A, B

and C ions are located at eight, six and four oxygen coordinated sites, respectively. $\text{La}_3\text{M}_2\text{Li}_5\text{O}_{12}$ ($\text{M} = \text{Nb}, \text{Ta}$), was the first fast Li^+ ion conducting garnet reported by Thangadurai *et al.* in 2003.⁹ Detailed studies have revealed that the lithium content and the distribution of lithium in the structure are key to the diffusion pathway and resulting lithium ion conductivity in garnet materials.^{10,11} Li^+ ions occupy both 24d tetragonal and 96h/48g octahedral Wyckoff sites in cubic $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$, with lithium ions shifting from 24d tetragonal sites to 96h/48g octahedral sites as a function of increasing lithium content, thus enabling a reduction in the lithium–lithium interaction strain.^{6,7} The maximum lithium content in the structure has been found to be 7, leading to compositions such as $\text{Li}_7\text{La}_3\text{Zr}_2\text{O}_{12}$, which adopts tetragonal symmetry with lithium ions ordering in three fully occupied tetrahedral 8a and octahedral 16f and 32g sites; this Li ordering reduces its ionic conductivity.^{12–14} The highly conductive cubic phase can be stabilised through the creation of lithium vacancies, and the optimum conductivity is found for lithium contents of 6.4–6.6 per garnet formula unit.^{15–25}

In addition, to facilitate the practical application of ASSBs, issues such as high interfacial resistances between the electrode and electrolyte, and the lithium dendrite penetration problem within solid state electrolytes have attracted much attention in recent years.^{26–29} The interfacial impedance between the garnet

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and electrode mainly arises from the poor contact in association with microscopic voids and grain boundaries of garnet, as well as an insulating Li_2CO_3 surface layer formed in air initiated by the proton/lithium exchange at the surface.^{30–33} The ionic transport was reported to be limited by the grain boundaries, which are heavily affected by the segregation of dopants and Li.³⁴ Mechanical polishing, carbon annealing or acid treatment were reported to remove this unwanted surface layer.^{35,36} Metal or non-metal coatings on garnet pellets, which consequently enables lithium alloying, have also been used to minimise the interfacial resistance between lithium metal and garnet, and to reduce the dendrite growth.^{37–40} Despite the potential for short-circuiting from electronic conduction, a mixed ion/electron-conductive interface was demonstrated to be beneficial for Li dendrite suppression.^{41,42} Another common strategy in this field is to utilise polymer–ceramic composite electrolytes, which combine good wetting properties of polymer electrolytes with the high ionic conductivity of ceramic electrolytes, forming a uniform interfacial contact with decreased interfacial resistance.^{43–45}

Ce^{3+} or Ce^{4+} incorporation on La^{3+} site in $\text{Li}_7\text{La}_3\text{Zr}(\text{Hf})_2\text{O}_{12}$ has been previously studied albeit showing a limited solid solution range for this site substitution, before the detection of impurity phases.^{46,47} In contrast, we showed in a previous study, that there was a greater degree of Ce^{4+} substitution possible on the Zr^{4+} site in $\text{Li}_7\text{La}_3\text{Zr}_2\text{O}_{12}$ leading to a reduction in the tetragonal distortion, which consequently increased the ionic conductivity, although due to the stoichiometric Li content in this system, the room temperature conductivity was just below the value of $10^{-4} \text{ S cm}^{-1}$, which is considered the minimum for applications.⁴⁸ In this paper, the possibility to replace Ce^{4+} on the Nb^{5+} site in $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ with the creation of excess Li^+ as the charge compensation mechanism has been examined for the first time, aiming to enhance the overall conductivity for these Ce doped garnets, in addition to providing the beneficial decrease in the interfacial impedance previously reported for Ce-doped $\text{Li}_7\text{La}_3\text{Zr}_2\text{O}_{12}$. We investigate effects on the local structure, electrical and electrochemical properties of Ce-doped $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ through combined X-ray pair distribution function analysis, impedance spectroscopy and electrochemical testing.

Experimental

Synthesis

Li_2CO_3 (99.9%, Alfa Aesar), CeO_2 (99.9%, Acros), La_2O_3 (99.9%, Sigma Aldrich) and Nb_2O_5 (99%, Sigma Aldrich) were used as starting materials. Li_2CO_3 was dried at 180 °C while La_2O_3 and Nb_2O_5 were dried at 900 °C before use. Samples of $\text{Li}_{5+x}\text{La}_3\text{Nb}_{2-x}\text{Ce}_x\text{O}_{12}$ with $x = 0, 0.25, 0.5, 0.75$ and 1 were prepared from intimately ground stoichiometric amounts of starting reagents which were heated initially to 650 °C for 12 hours at a rate of 5 °C min^{-1} . 10–15% excess Li_2CO_3 was then added to the precursor and the powder milled for 30 minutes using a Pulverisette 5 planetary ball mill. The mixture was pressed into a pellet and heated to 950–1000 °C for 12 hours at a rate of 5 °C min^{-1} to form the final product.

Characterisation

A Bruker D8 X-ray diffractometer (XRD) with a $\text{CuK}\alpha$ radiation and a linear position sensitive detector was used to collect X-ray diffraction data. Patterns were recorded over the 2θ range 15° to 80° with a 0.02° step size. Structural refinement was carried out with the GSAS suite of Rietveld refinement software using the XRD data.⁴⁹

Scanning electron microscopy (SEM, HITACHI TM4000plus) was employed to assess the microstructure. Bulk samples were polished and thermally etched at 90% of the sintering temperature for 0.5 h. The distribution of elements was probed with an energy dispersive X-ray spectroscopy (EDX) detector.

Pellets (9.8 mm diameter) were pressed and sintered at 1000–1050 °C for 4 hours (ramp rate of 5 °C min^{-1}) in a dry room to limit H^+/Li^+ exchange and prevent the decomposition of samples for impedance measurements. Mother powder was used to cover the pellets to prevent the Li loss and reduce reaction with the Al_2O_3 crucible. Au paste was painted on both sides of the pellet and heated at 850 °C for 1 hour in air. Impedance data were collected with a HP 4192 analyser over the frequency range of 1 to 10^7 Hz with 100 mV ac amplitude.

Cell assembly and electrochemical test

For the cell tests, a $\text{Li}_{5.75}\text{La}_3\text{Zr}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ pellet (2 mm thickness) was sintered at 1000 °C for 12 h in a dry room. The $\text{Li} // \text{Li}_{5.75}\text{La}_3\text{Zr}_{1.25}\text{Ce}_{0.75}\text{O}_{12} // \text{Li}$ symmetric cell was hot pressed at 175 °C for 1 hour and assembled using a Swagelok cell in an Argon filled glove box. Electrochemical impedance spectroscopy (EIS) was performed over the frequency range from 0.1 to 10^7 Hz (100 mV ac applied potential) with a Solatron 1260 analyser. The Li plating/stripping performance was evaluated using a Bio-logic SP50 cell tester.

Pair distribution function data collection

Total scattering data were collected at the I15-1 XPDF beamline at Diamond Light Source, UK. Powdered samples were loaded into borosilicate capillaries (1.5 mm OD, 1.17 mm ID) and spun perpendicular to the beam during data collection to improve powder averaging. Scattering data were collected at an X-ray energy of 76.69 keV using a PerkinElmer XRD 4343 CT area detector placed ~200 mm from the sample. The 2-D data were corrected for polarization and flat-field, then integrated to 1-D using the DAWN package prior to processed the scattering range $0.7 \text{ \AA}^{-1} \leq Q \leq 25 \text{ \AA}^{-1}$ into PDFs using the GudrunX package.^{50,51} A modified Lorch Function ($\Delta_1 = 0.05 \text{ \AA}$) was applied to suppress spurious low- r features.⁵²

Results and discussion

Phase formation

Stoichiometric $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ (LLNO) can be indexed with an $Ia\bar{3}d$ cubic cell. In Kroger–Vink notation, the relationship for substitution of Nb^{5+} by Ce^{4+} is as follows:

$$[\text{Ce}'_{\text{Nb}}] = [\text{Li}_i]$$

Fig. 1 (a) XRD patterns of $\text{Li}_{5+x}\text{La}_3\text{Nb}_{2-x}\text{Ce}_x\text{O}_{12}$: $x = 0, 0.25, 0.5, 0.75, 1$. Triangle and rectangle labels represent CeO_2 and $\text{La}_2\text{CeLi}_2\text{O}_6$ respectively. (b) Variation of lattice parameter, a , and volume, V , of $\text{Li}_{5+x}\text{La}_3\text{Nb}_{2-x}\text{Ce}_x\text{O}_{12}$ with composition x .

As shown in Fig. 1a, XRD patterns of Ce-doped $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ with formulae $\text{Li}_{5+x}\text{La}_3\text{Nb}_{2-x}\text{Ce}_x\text{O}_{12}$ show peak shifts to lower 2θ angle with increasing Ce content, without the detection of any impurity phases up to $x = 0.75$. The cell parameter was shown to obey Vegard's law, increasing linearly as a function of x (Fig. 1b), which further confirms the successful incorporation of Ce into the structure due to the larger radius of Ce^{4+} to that of Nb^{5+} (0.87 and 0.64 ionic radius for Ce^{4+} and Nb^{5+} respectively in 6 oxygen coordinated octahedral sites). $\text{Li}_2\text{La}_2\text{CeO}_6$ and CeO_2 impurity phases were detected for compositions with $x = 1$, which is consistent with the leveling off of the cell parameter increase; from these data, a solid solution limitation around $x = 0.85$ is estimated. As shown in Fig. 2, EDX mapping of $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ shows a homogeneous distribution of La, Nb, Ce and O within the sample.

XRD data of $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ and $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ were used to carry out structure refinement in order to confirm the incorporation of Ce on the Nb site. The U_{iso} parameters of all atoms were initially set to a typical value of 0.01 Å. It was not possible to refine the Li positions and occupancies due to the

Table 1 Refined structural parameters for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ using XRD data

Atom	x	y	z	Mult.	Occupancy	$u_{\text{iso}} \times 100$ (Å ²)
La1	0.125	0	0.25	24	1	0.6(1)
Nb1	0	0	0	16	1	0.2(1)
O1	0.2844	0.1044	0.1996	96	1	0.7(3)
	(5)	(6)	(6)			
Li1	0.25	0.875	0	24	0.836	1
Li2	0.125	0.6798	0.5702	48	0.110	1
Li3	0.095	0.6818	0.5778	96	0.152	1

$a = 12.7949(1)$ Å, $V = 2094.67(4)$ Å³, $\chi^2 = 5.266$, $R_{\text{wp}} = 11.49\%$, $R_p = 9.13\%$.

Fig. 2 EDX elemental mapping of $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$.

Table 2 Refined structural parameters for $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ using XRD data

Atom	x	y	z	Mult.	Occupancy	$u_{\text{iso}} \times 100 (\text{\AA}^2)$
La1	0.125	0	0.25	24	1	1
Nb1	0	0	0	16	0.61(1)	1
Ce1	0	0	0	16	0.39(1)	1
O1	0.2874 (4)	0.0980 (5)	0.1984 (5)	96	1	1.2(3)
Li1	0.25	0.875	0	24	0.961	1
Li2	0.125	0.6798	0.5702	48	0.127	1
Li3	0.095	0.6818	0.5778	96	0.175	1

$$a = 12.9761(1) \text{ \AA}, V = 2184.90(7) \text{ \AA}^3, \chi^2 = 2.823, R_{\text{wp}} = 8.18, R_p = 6.69\%.$$

poor sensitivity of Li with X-rays in the presence of heavier elements; Li positions were not refined and Li occupancies for the doped samples were scaled up from published occupancies of the undoped phase to ensure the charge-neutrality of the resulting structure. First the scale factor, background (6 terms

Table 3 Bond lengths for Nb1(Ce1)–O of $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ and $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ from refinement results

Composition	Bond	Length (\AA)
$\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$	Nb1–O1 ($\times 6$)	2.020(8)
$\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$	Nb1/Ce1–O1 ($\times 6$)	2.138(6)

of shifted Chebyshev function) and lattice parameters were refined. Then, peak profile parameters and 2θ zero error were refined before fixing them at the converged value. Finally, the atomic coordinates of O1, U_{isos} of La1, Nb1 were refined in sequence followed by refining together in the last step. Final refined parameters for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ are summarised in Table 1.

A similar structure refinement for $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ was performed. In this case, a constraint of the same atomic coordinates, the same U_{isos} and full occupancy between Nb1 and Ce1 was made initially. This refinement gave site occupancies of 0.61(1) for Nb1 and 0.39(1) for Ce1 which is consistent with the expected composition within 2 esds. Final refined parameters for $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ are shown in Table 2. Observed, calculated and difference profiles from Rietveld refinement for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ and $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ are shown in Fig. 3. Selected bond lengths for both $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ and $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ are shown in Table 3. The data show that all Nb1/Ce1–O1 bonds expand in the Ce doped sample (bond distance = 2.138 \AA) compared to undoped LLNO (bond distance = 2.020 \AA), consistent with the larger size of Ce^{4+} versus Nb^{5+} .

Pair distribution function (PDF)

X-ray PDF data for undoped $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ and Ce-doped $\text{Li}_{5.5}\text{La}_3\text{Nb}_{1.5}\text{Ce}_{0.5}\text{O}_{12}$ and $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ are shown in Fig. 4. The peaks in the low- r region of the PDF were assigned based on simulations of the partial PDFs of the undoped struc-

Fig. 3 Observed, calculated and difference profiles from Rietveld refinement for (a) $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ and (b) $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ using XRD data.

Fig. 4 Pair distribution function data for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ (black), $\text{Li}_{5.5}\text{La}_3\text{Nb}_{1.5}\text{Ce}_{0.5}\text{O}_{12}$ (green) and $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ (red) focusing in the low- r region.

ture using the computer program PDFgui (Fig. S1†).⁵³ The single Nb–O peak at ~ 1.94 Å in undoped LLNO develops a shoulder at high- r with doping, which increases in intensity concurrent with Ce content. Gaussian peaks were fit to the observed PDF data using the computer program TOPAS (Fig. S2†) allowing for the extraction of the local structure

bond lengths given in Table 4.⁵⁴ In contrast to the long-range average bond length from Rietveld refinement which calculates a single average Nb,Ce–O bond length which increases from 2.020(8) Å to 2.138(6) Å on doping, the local structure information provided by the PDF allows us to separate the contributions of the individual bond distances. These results show

Table 4 Bond lengths for Nb–O, Ce–O, La–O, La–Nb,Ce, La–La of $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$, $\text{Li}_{5.5}\text{La}_3\text{Nb}_{1.5}\text{Ce}_{0.5}\text{O}_{12}$ and $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ from X-ray PDF data

Bond	$\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$	$\text{Li}_{5.5}\text{La}_3\text{Nb}_{1.5}\text{Ce}_{0.5}\text{O}_{12}$	$\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$
Nb–O	1.93(3) Å	1.95(2) Å	1.94(2) Å
Ce–O	—	2.16(4) Å	2.16(2) Å
La–O	2.474(5) Å	2.475(7) Å	2.476(7) Å
La–Nb,Ce	3.531(1) Å	3.572(1) Å	3.591(1) Å
La–La	3.904(1) Å	3.929(1) Å	3.942(2) Å

Fig. 5 Pair distribution function data for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ (black), $\text{Li}_{5.5}\text{La}_3\text{Nb}_{1.5}\text{Ce}_{0.5}\text{O}_{12}$ (green) and $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ (red), with the r axis normalized to the lattice parameter of $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ at (a) low r (b) high r .

Fig. 6 SEM images of (a) $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$, (b) $\text{Li}_{5.5}\text{La}_3\text{Nb}_{1.5}\text{Ce}_{0.5}\text{O}_{12}$ and (c) $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ pellets.

that the Nb–O bond remains constant at ~ 1.94 Å, while the Ce–O bond length of 2.16 Å can be extracted for the doped samples, which is consistent with the increase expected on moving to the larger Ce^{4+} ion on simple ionic radii arguments.

Fig. 5 shows the PDF data with the r axis for the doped samples normalised to the lattice parameter of the undoped LLNO; this normalisation highlights changes upon doping beyond simple expansion of the lattice. At low- r (Fig. 5a), peaks around 5.9 Å due to the second La–La and Nb,Ce–Nb,Ce coordination shell contract relative to the normalised lattice parameter, but the rest of the features remain unchanged. Neutron PDF data would be required to determine the full impact of doping on the Li and O environments (previous neutron PDF study shows lithium atoms within either tetrahedral or octahedral sites prefer to stay at the off-center positions in the undoped $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ ⁵⁵). However, it is clear that doping introduces more disorder into the system, evidenced by broader peaks at high r (Fig. 5b); this additional disorder is expected due to the significantly different sizes of Nb and Ce, and the additional lithium which requires occupation of more varied coordination environments.

Conductivity measurement

The morphology of undoped and cerium doped $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ are shown in Fig. 6. The relative density of garnet pellets increased as a function of Ce content (51.2% for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ and 75.1% for $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$). A typical dataset of $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ impedance spectroscopy is shown in Fig. 7. From the Z^* complex plane (Fig. 7a), a non-ideal semicircle was observed at higher frequency while a sharp spike is seen at lower frequency. To fit the circuit, a resistor R1 and a constant phase element (CPE) CPE1 in parallel were used to simulate the high frequency semicircle. Another CPE2 which is in series with the R1/CPE1 in the circuit was utilised to interpret the low frequency spike.

The corresponding spectroscopic C' plot (Fig. 7b) shows a higher frequency plateau with a capacitance of 8.2 pF cm^{-1} with an associated permittivity of 91 (calculated from $\epsilon'_{\infty} = C/\epsilon_0$, where ϵ_0 is the permittivity of free space with a value of $8.854 \times 10^{-14} \text{ F cm}^{-1}$), which is a typical value for the bulk response.⁵⁶ In addition, a low frequency plateau with a capacitance of $3 \text{ } \mu\text{F cm}^{-1}$ was observed, which is due to the double layer effect at the sample–electrode interface; hence consistent with Li^+ ion conduction.

The spectroscopic Y' plot (Fig. 7c) shows a frequency-independent plateau at intermediate frequency, which is associated with the bulk conductivity. A dispersion at lower frequency was seen due to the blockage of lithium ions at the sample–electrode interface. At higher frequency, a curvature was detected related to the Jonscher's power law.⁵⁷

The Arrhenius plot (Fig. 8) shows the bulk conductivities of undoped and Ce-doped LLNO. Conductivities increased with increasing cerium content in agreement with more excess lithium in the structure. The maximum conductivity was found in $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ sample, which reached a value of $1.4 \times 10^{-4} \text{ S cm}^{-1}$ at 25 °C, which is $\sim$ two orders more

Fig. 7 Impedance data of $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$. (a) Complex Z^* plot (b) spectroscopic plot of C' (c) spectroscopic plot of Y' .

than that of LLNO, and represents one of the highest conductivities for a garnet with less than 6 Li. The activation energy also decreased from 0.41(1) eV for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ to 0.35(1) eV for $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$. This decrease in activation energy may be indicative of an opening of the bottleneck for Li

Fig. 8 Arrhenius conductivity plot for Ce-doped $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$.

migration between sites, due to the increase in cell size through Ce doping.

Electrochemical properties

The electrochemical impedance spectroscopy of the $\text{Li}/\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}/\text{Li}$ cell is shown in Fig. 9a. For circuit fitting, a parallel R1 and CPE1 representing a bulk semicircle in series with another parallel R2 and CPE2 related to the

Fig. 9 (a) Electrochemical impedance spectroscopy of $\text{Li}/\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}/\text{Li}$ cell. (b) Stripping/plating properties of $\text{Li}/\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}/\text{Li}$ cell at 55 °C.

interfacial semicircle were used. A significant interfacial resistance decrease from $3802 \Omega \text{ cm}^2$ for $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ to $488 \Omega \text{ cm}^2$ for $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ was observed (Fig. S3†), which is consistent with our previous work on Ce doped $\text{Li}_7\text{La}_3\text{Zr}_2\text{O}_{12}$, and represents a reduced interfacial resistance attributed to the formation of a mixed valence $\text{Ce}^{4+}/\text{Ce}^{3+}$ product at the electrode/lithium interface.⁴⁸ Long-term cycling properties of the $\text{Li}/\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}/\text{Li}$ cell at 55 °C are shown in Fig. 9b. A voltage plateau with a voltage overpotential of 251 mV at $50 \mu\text{A cm}^{-2}$ current density was observed; the long-term cycling stability over 380 h was shown.

Conclusion

Cerium was successfully doped into the $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ garnet structure through a conventional solid-state method. The lattice parameter increased linearly as a function of composition to $x \sim 0.85$, indicating the formation of an extensive solid solution range. XRD PDF analysis confirmed the incorporation of cerium, allowing the separate Nb/Ce–O bond distances to be determined, and thus illustrating the longer bond distances for Ce on the octahedral site to those of Nb. The conductivity increased by two orders of magnitude of $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$ compared to that of $\text{Li}_5\text{La}_3\text{Nb}_2\text{O}_{12}$ with a notable reduction in the activation energies for Ce-doped samples. The interfacial resistance also dramatically decreased to $488 \Omega \text{ cm}^2$ for $\text{Li}_{5.75}\text{La}_3\text{Nb}_{1.25}\text{Ce}_{0.75}\text{O}_{12}$, demonstrating that cerium doping benefits the interfacial issue. The long-term cycling stability of a $\text{Li}/\text{garnet}/\text{Li}$ symmetric cell over 380 h was demonstrated, thus illustrating the potential of this Ce doped garnet system. Further studies are required to characterise the formation and possible change of the interlayer between garnet and lithium metal.

Conflicts of interest

The authors declare no competing financial interest.

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