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## FB-REDA: fragment-based decomposition analysis of the reorganization energy for organic semiconductors†

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We present a fragment-based decomposition analysis tool (FB-REDA) for the reorganisation energy ( $\lambda$ ). This tool delivers insights on how to rationally design low- $\lambda$  organic semiconductors. The contribution of the fragment vibrational modes to the reorganization energy is exploited to identify the individual contributions of the molecular building blocks. The usefulness of the approach is demonstrated by offering three strategies to reduce the reorganization energy of a promising dopant-free hole transport material (**TPA1PM**,  $\lambda = 213$  meV). A reduction of nearly 50% (**TPD3PM**,  $\lambda = 108$  meV) is achieved. The proposed design principles are likely transferable to other organic semiconductors exploiting common molecular building blocks.

## Introduction

Organic semiconductors based on small molecules are exploited in various optoelectronic applications owing to their well-defined chemical structures (monodispersity), low cost, flexibility and low-temperature and solution processability.<sup>1,2</sup> In the spirit of “molecular LEGO,”<sup>3,4</sup> their properties are easily tuneable with the charge mobility being the most relevant target<sup>5–11</sup> for rational design. Relevant examples include the incorporation of hole transport materials with high hole mobility into photovoltaics or organic light-emitting diodes with low hole–electron recombination rate.<sup>12–14</sup>

The reorganization energy ( $\lambda$ ) is one of the key parameters for determining the efficiency of charge transport in both the hopping and band transport (related to electron–phonon coupling term) regime.<sup>15,16</sup> Organic semiconductors with smaller reorganization energy (assuming the other parameters fixed) generally show a higher charge mobility.<sup>17</sup> Therefore, understanding the relationship between molecular structures and reorganization energy is the key to a successful rational design.<sup>2,18</sup>

A molecule’s reorganization energy is generally computed using Nelsen’s 4-point method or the normal mode (NM) analysis.<sup>17,19,20</sup> In order to assess the role played by each molecular building block, it is insightful to decompose the overall reorganization energy into the individual molecular fragment

contribution ( $\lambda_{\text{frag}}$ ). However, analysis tools giving direct access to the fragment contribution are still lacking. There are currently two existing analysis that serve to evaluate the atomic or fragment contributions to  $\lambda$  or the vibronic coupling constant ( $V$ ).

The first approach, developed by Shuai, Coropceanu *et al.* based on the NM analysis, expresses the displacements along the normal modes in terms of molecular internal coordinates.<sup>21,22</sup> The total reorganization energy is then equal to a sum of internal coordinate contributions (see eqn (S1) in the ESI†). Using this approach, the authors were able to demonstrate the influence of substituents and substitution sites on the overall reorganization energy of indolo [3,2-*b*]carbazole derivatives. While such an analysis tool is useful to identify the substituent leading to the lowest  $\lambda$ , it is less convenient to evaluate the fragment contributions to  $\lambda$  in a “molecular LEGO” framework. The reorganization energy in the internal coordinate representation contains cross terms involving two different internal coordinates, which are not necessarily small in comparison to the square terms (see eqn (S1) in the ESI†). However, there is no trivial way to divide these cross terms into internal coordinate contributions or to partition internal coordinates involving more than one fragment.

The second method developed by Sato *et al.*, is associated with the concept of atomic vibronic coupling constant (AVCC).<sup>23,24</sup> They investigated the chain-length dependence of the reorganization energy for oligofluorenes and oligothiophenes by evaluating the change in local degree of vibronic coupling when increasing the length of the oligomer. Because the sum of all AVCC is equal to the total vibronic coupling constant (see eqn (S2) in the ESI†), the fragment vibronic coupling constants can be defined as a sum over all the AVCC of the corresponding atoms within a fragment. However, given that the total reorganization energy is proportional to the square of the total vibronic coupling

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constant, the representation of  $\lambda$  in terms of AVCC inevitably leads to cross terms involving two different atomic vibronic constants (coupled contributions from different atoms).

This work proposes an alternative strategy to decompose the reorganization energy into local contributions. The approach is inspired by the idea of Huix-Rotllant *et al.*,<sup>25</sup> who decomposed the molecular normal modes in terms of local fragment modes. Instead of first decomposing  $\lambda$  into internal coordinates or atomic contributions, we avoid the generation of cross terms by computing the fragment contributions directly from the fragment modes. The equations and associated details are discussed in the Methodology section.

As a proof of concept, we use the proposed fragment reorganization decomposition analysis to identify organic hole transport materials (HTM) with low reorganization energy. We start with **TPA1PM**<sup>26</sup> that has proved to be an efficient dopant-free HTM good for defect-passivation. The molecule (shown in Fig. 1(a)) is composed of a triphenylamine core (TPA), a phenyl-substituted carbazole arm (PCZ) and methoxy (M) substituent groups. The chosen strategy for reducing the reorganization energy is established based on the fragment-based decomposition analysis. Three molecular design approaches were successively adopted and their effects on the total hole reorganization energy ( $\lambda_{\text{tot}}$ ) and the fragment contributions (divided into core ( $\lambda_{\text{core}}$ ), arm ( $\lambda_{\text{arm}}$ ) and substituent ( $\lambda_{\text{sub}}$ ) contributions) are discussed. The three strategies are: (1) introducing multi-arm, (2) non-covalent lock and (3) core rigidification. The usefulness of the computational approach is demonstrated by the reduction by half of the reorganization energy compared to **TPA1PM** (from 213 to 108 meV).

## Methodology

### Reorganization energy

There are two computational ways to evaluate the reorganization energy. The first method is known as Nelsen's 4-point method, where  $\lambda$  is calculated using four points on the adiabatic potential energy surfaces of the neutral and cation state of a molecule. The reorganization energy can be represented as:

$$\lambda = \lambda_N + \lambda_C = (E_{\text{cN}} - E_{\text{nN}}) + (E_{\text{nC}} - E_{\text{cC}}) \quad (1)$$

where  $\lambda_N$  and  $\lambda_C$  are the reorganization energy from neutral and cationic adiabatic potential energy surface. The "C" and "N" letters stand for cationic and neutral, respectively. Letters in lowercase denote the charge state of the molecule, while uppercase letters indicate the corresponding ground state molecular geometry of the charge state, as shown in Scheme 1. The second approach consists in performing a normal mode analysis (NM),<sup>28,29</sup> in which the reorganization energy is the sum of the potential energy of harmonic oscillators defined as:

$$\begin{aligned} \lambda = \lambda_N + \lambda_C &= \sum_i \lambda_i^N + \sum_j \lambda_j^C \\ &= \sum_i \frac{1}{2} k_i^N (\Delta Q_i^N)^2 + \sum_j \frac{1}{2} k_j^C (\Delta Q_j^C)^2 \end{aligned} \quad (2)$$

where  $\lambda_i^N/\lambda_j^C$  is the reorganization energy contribution from the  $i$ th/jth mode of the neutral/cationic ground state,  $k_i^N/k_j^C$  is the eigenvalue of the mass-weighted Hessian matrix for the neutral/cationic ground state (N/C).  $\Delta Q_i^N/\Delta Q_j^C$  are the projection of the

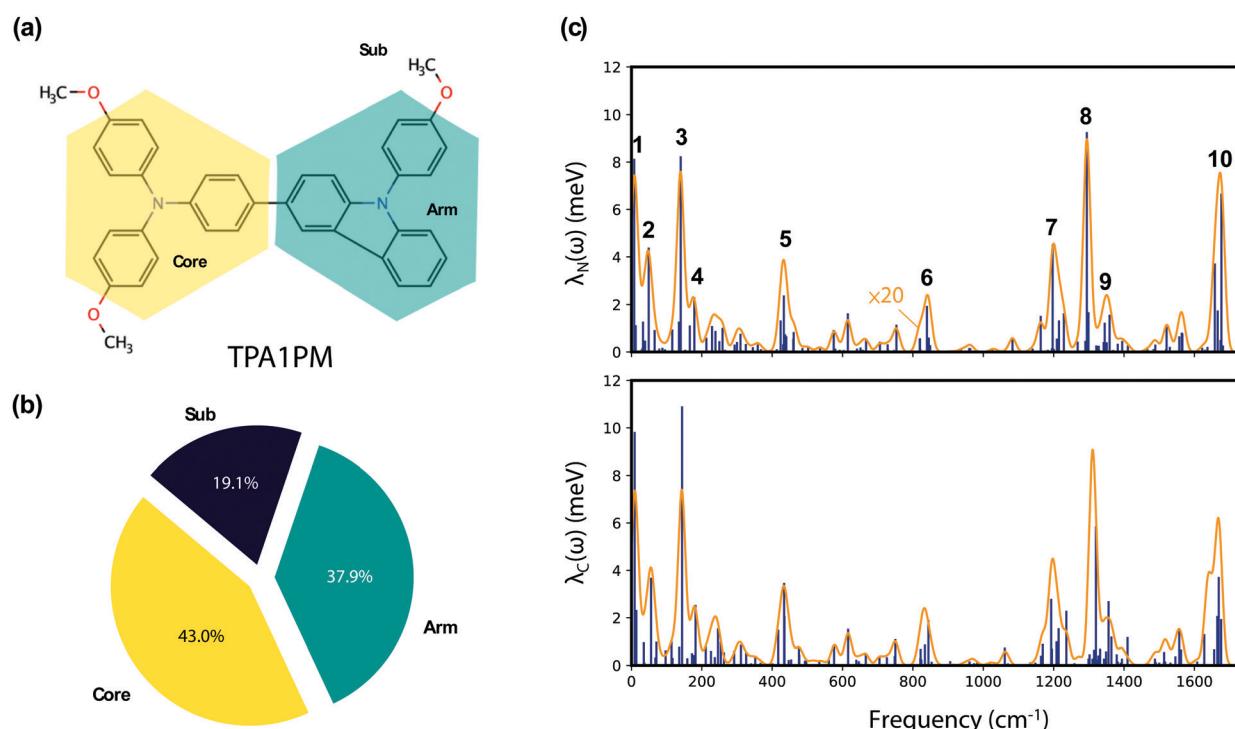
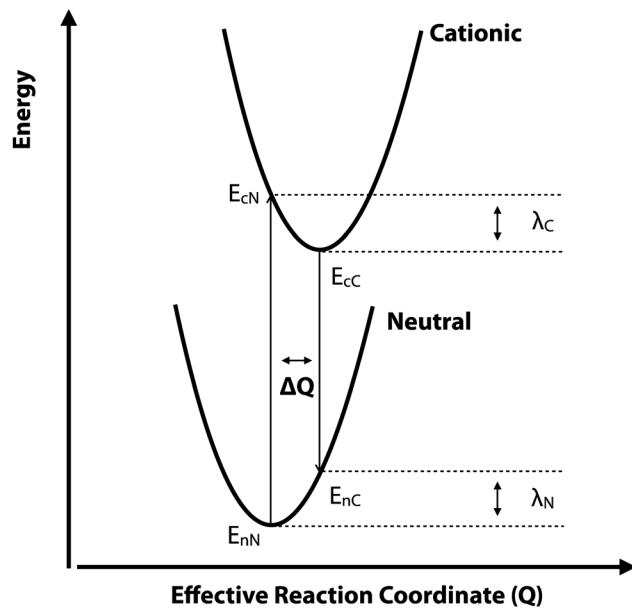


Fig. 1 (a) Illustration of the decomposition of **TPA1PM** into Core, Arm and Sub fragments. (b) The ratio of the fragment to the total reorganization energy for each fragment. (c) The spectral density plot, blue lines are  $\lambda_i^N/\lambda_j^C$  and orange lines are  $\lambda_N(\omega)$  and  $\lambda_C(\omega)$  with intensity multiplied by 20. The 10 most important peaks are indexed.





**Scheme 1** Sketch of the potential energy surface of the neutral and cationic molecular states. The relationship between the four points ( $E_{nN}$ ,  $E_{cN}$ ,  $E_{CC}$  and  $E_{nC}$ ),  $\lambda_N$ ,  $\lambda_C$  and  $\Delta Q$  on the potential energy surfaces is depicted.

eigenvector of the normal mode (N/C) onto the Cartesian coordinate difference between the cationic and neutral ground states, which is defined as:<sup>27</sup>

$$\begin{aligned}\Delta Q_i^N &= (C_i^N)^T m^{1/2} (R_C - R_N) \\ \Delta Q_j^C &= (C_j^C)^T m^{1/2} (R_N - R_C)\end{aligned}\quad (3)$$

where  $C_i^N/C_j^C$  is the eigenvector of the normal mode,  $m$  is the mass matrix and  $R_N/R_C$  is vector of the Cartesian coordinate of neutral/cation ground state geometry. Since large bending and torsional motions are present in the charge transfer process of the compounds investigated here, we use curvilinear coordinates to compute the  $\Delta Q_i^N/\Delta Q_j^C$  following the protocol proposed by Reimers.<sup>20</sup>

### Fragment contributions

Within the NM method, the reorganization energy can be further partitioned into individual fragment contributions *via* the fragment modes analysis. Following the protocol proposed by Huix-Rotllant and Ferré,<sup>25</sup> the modes for each fragment are obtained by first partitioning the total mass-weighted Hessian matrix  $H_{\text{tot}}$  into fragment submatrices  $H_f$ . The eigenvectors of the fragment modes  $C_f$  are then obtained by diagonalizing these submatrices. These eigenvectors form a complete orthonormal basis set to expand the total normal modes:

$$\begin{aligned}C_{\text{tot},i} &= \sum_j \sum_k d_{jk}^i C_{\text{frag},jk} \\ d_{jk}^i &= C_{\text{tot},i}^T C_{\text{frag},jk}\end{aligned}\quad (4)$$

where  $C_{\text{tot},i}$  is the eigenvector of the total normal mode  $i$ ,  $C_{\text{frag},jk}$  is the eigenvector of the mode  $j$  of the fragment  $k$  and  $d_{jk}^i$  is the

coefficient of expansion. The expansion into fragment modes gives access to the fragment contribution to  $\lambda$ .

For convenience, the following derivation is based on  $\lambda_N$ , but the conclusion is the same for  $\lambda_C$ . Inserting eqn (3) and (4) into eqn (2) leads to an expression for  $\lambda_N$  in which the diagonal terms comes from the same fragment mode ( $d_1^i \Delta q_{f,1}^N$ )<sup>2</sup> whereas the non-diagonal terms originates from two different fragment modes ( $d_1^i d_2^j \Delta q_{f,1}^N \Delta q_{f,2}^N$ ) (see eqn (S5) and (S6), ESI†). The difficulty associated with how to deal with the non-diagonal terms, is similar to the approaches mentioned in the introduction. To avoid the non-diagonal terms, we thus define the contribution of the mode  $j$  of the fragment  $k$  as follow:

$$\lambda_{jk}^i = w_{jk}^i \lambda^i = d_{jk}^{i,2} \lambda^i \quad (5)$$

where  $\lambda_{jk}^i$  is the reorganization energy of the normal mode  $j$  for fragment  $k$  and the total normal mode  $i$ ,  $\lambda^i$  is the reorganization energy arising from the total normal mode  $i$  and  $w_{jk}^i$  is the weight. According to this definition, we can further define the reorganization energy of fragment  $k$  arising from the total normal mode  $i$  ( $\lambda_k^i$ ) and the sum of  $\lambda_k^i$  over all total normal modes (fragment reorganization energy,  $\lambda_k$ ) as:

$$\begin{aligned}\lambda_k^i &= \sum_j \lambda_{jk}^i \\ \lambda_k &= \sum_i \lambda_k^i\end{aligned}\quad (6)$$

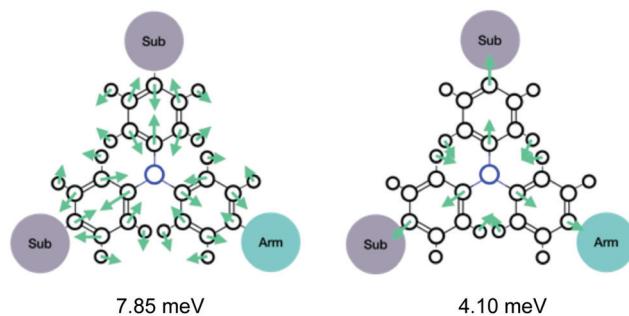
Since these fragment modes form a complete orthonormal basis, the following relation holds:

$$\begin{aligned}\sum_{jk} d_{jk}^{i,2} &= 1 \\ \sum_k \lambda_k^i &= \sum_i \lambda^i \sum_{jk} d_{jk}^{i,2} = \sum_i \lambda^i = \lambda_N\end{aligned}\quad (7)$$

The equation above shows that the sum of the reorganization energies of all fragments is equal to the total reorganization energy, justifying our definition of  $\lambda_k$ .

In addition to the fragment reorganization energy, we also define the fragment mode reorganization energy ( $\lambda_{jk}$ ), to identify crucial fragment modes:

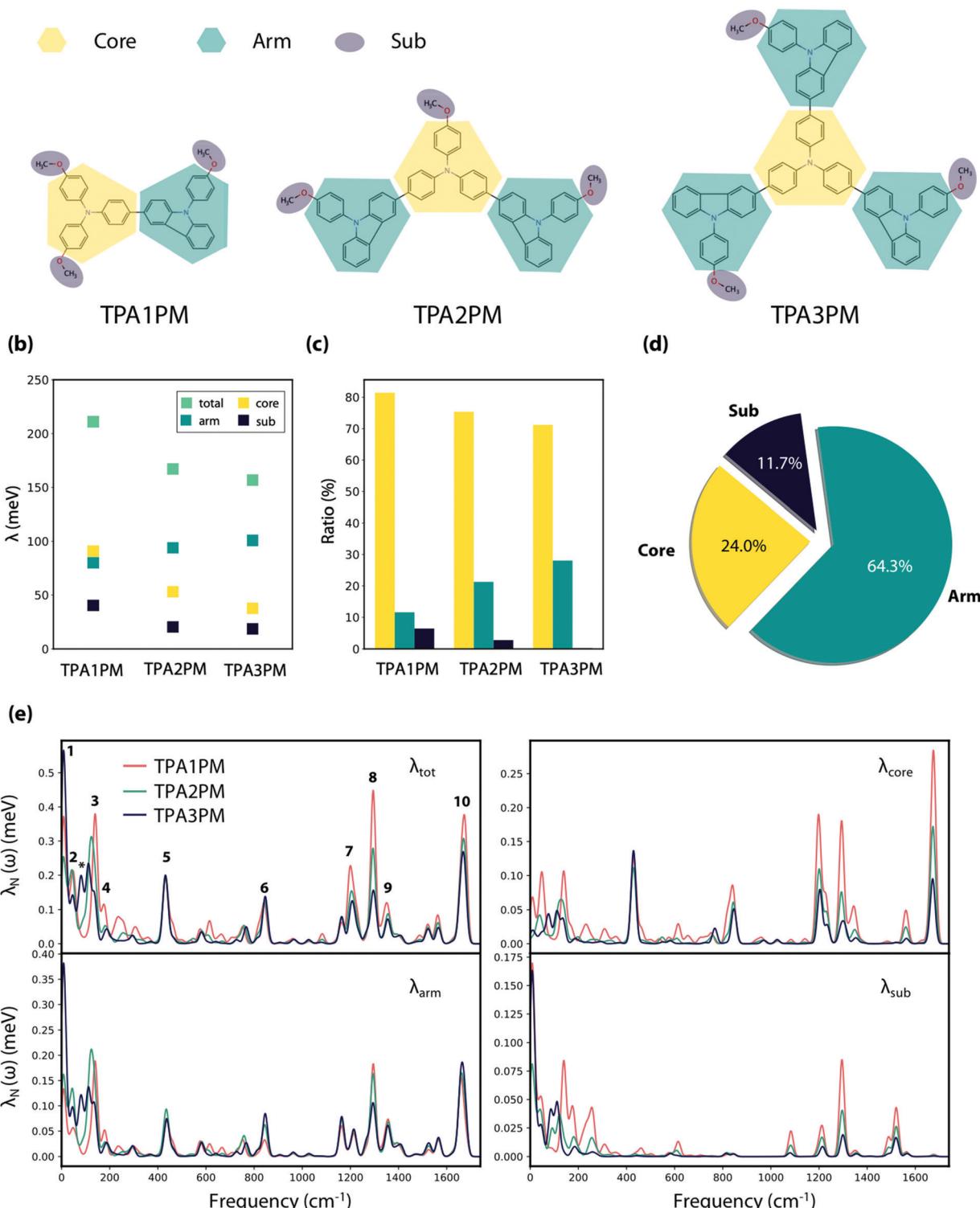
$$\lambda_{jk} = \sum_i \lambda_{jk}^i \quad (8)$$



**Fig. 2** Illustration of the two most contributing fragment modes and their corresponding fragment reorganization energies in TPA1PM.

(a)

## Introducing multiarm



**Fig. 3** (a) The **TPA1PM**, **TPA2PM** and **TPA3PM** molecules and their fragments. (b) Total reorganization energy, fragment reorganization energy, (c) contribution to the HOMO from the NAO analysis and (d) the ratio of  $\lambda_{\text{frag}}$  to  $\lambda_{\text{tot}}$  of each fragment. (e) Spectral density plots for  $\lambda_{\text{frag}}$  and  $\lambda_{\text{tot}}$  for **TPA1PM**, **TPA2PM** and **TPA3PM**.



In the following sections, we demonstrate the usefulness of the analysis in terms of  $\lambda_{\text{frag}}$ , while targeting the identification of low- $\lambda$  organic molecules.

### Computational details

Geometry optimization, electronic structure and frequency computations of the neutral and cationic molecules were performed at the B3LYP<sup>30,31</sup>D3BJ<sup>32</sup> level in conjunction with the 6-31G(d,p) basis set. All DFT computations were performed with Gaussian16.<sup>33</sup> The NM analysis was performed using our in-house code, DUSHIN<sup>20</sup> and Gaussian16 in tandem.

To evaluate the validity of harmonic approximation, the reorganization energies were computed using both the 4-point method and the NM analysis (Fig. S1, ESI<sup>†</sup>). Our results show that the reorganization energies obtained from these two methods are very similar, justifying the use of the harmonic approximation.

## Results and discussion

### Analysis of TPA1PM

We first compute the reorganization energy of the parent **TPA1PM** molecule (211 meV) prior to its decomposition into three fragments: (1) a TPA core (Core), (2) a PCZ arm (Arm), and (3) three methoxy substituents (Sub), as shown in Fig. 1(a). The fragment reorganization energy (in meV) of the core (91), Arm (80) and Sub (40) are obtained *via* eqn (6). The  $\lambda_{\text{frag}}/\lambda_{\text{tot}}$  ratio for each fragment is represented in the pie chart (Fig. 1(b)). The core contributes the most (43.0%) to total reorganization energy. The contribution from the arm is slightly less significant (37.9%), with Sub contributing the least (19.1%).

The normal modes that are dominating the reorganization energy are plotted in Fig. 1(c) with  $\lambda_i^N/\lambda_j^C$  representing the contribution for each normal mode *i* and *j*. Given the large number of normal modes in each system, we define the spectral density  $\lambda_{N/C}(\omega)$  to facilitate the analysis:<sup>34,35</sup>

$$\lambda_N(\omega) = \sum_i \lambda_i^N \delta(\omega - \omega_i) \quad (9)$$

$$\lambda_C(\omega) = \sum_j \lambda_j^C \delta(\omega - \omega_j)$$

where  $\omega$  is the vibrational frequency and the Dirac  $\delta$  function is approximated by a Gaussian broadening distribution with  $\sigma = 10 \text{ cm}^{-1}$  (Fig. 1(c)). Because  $\lambda_N(\omega)$  and  $\lambda_C(\omega)$  exhibit the same features and that the magnitude of  $\lambda_N$  and  $\lambda_C$  are very close (104 and 107 meV respectively) the discussion on the spectral density will focus on the neutral case  $\lambda_N(\omega)$ .

The 10 most contributing peaks in  $\lambda_N(\omega)$  are shown in Fig. 1(c) and the composition in terms of their fragment modes is shown in Table S1 (ESI<sup>†</sup>). Overall, several fragment modes recurrently contribute to these 10 important peaks and those with  $\lambda_{jk} > 1.2 \text{ meV}$  ( $10 \text{ cm}^{-1}$ ) are listed in Table S2 (ESI<sup>†</sup>). More specifically, the two most relevant fragment modes correspond to in-plane bond-stretching and angle-bending in the Core (Fig. 2). Considering the fact that the Core is the largest contributor to the reorganization energy, an efficient strategy

to reduce  $\lambda$  consist in spreading the HOMO away from the TPA core. As shown in our former work,<sup>17</sup> this objective can be achieved by introducing additional PCZ arms.

### Introducing multi-arm

The first strategy to reduce the reorganization energy of **TPA1PM** is thus to increase the number of PCZ arms, as exemplified in Fig. 3(a). Given that the methoxy groups serve as a Lewis base passivating defects in the perovskite layer,<sup>26</sup> the number of methoxy substituents are kept constant. They are inserted either directly on the core or on the arms depending on their number.

The design strategy is efficient as the total reorganization energy decreases by up to  $\sim 20\%$  with increasing the number of PCZ arms, from 213, 173 to 156 meV (Fig. 3(b)). As expected, the highest occupied molecular orbitals (HOMO) spread through the PCZ arms (Fig. S2, ESI<sup>†</sup>). The contribution from Core and Sub to the HOMO (evaluated using natural atomic orbital analysis (NAO) Fig. 3(c)) decreases from  $\sim 80\%$  to  $\sim 72\%$  for Core and from  $\sim 7\%$  to  $\sim 0\%$  for Sub respectively, while that of Arm increases. These trends are in line with the decrease of  $\lambda_{\text{core}}$  and  $\lambda_{\text{sub}}$ , which contrasts with the  $\lambda_{\text{arm}}$  increase from 80 to 101 meV. Overall, it is clear that the redistribution of the HOMO is the dominant factor affecting  $\lambda$  and that the largest contribution comes from the arms Fig. 3(d).

The spectral density of the Total, Core, Arm and Sub contribution plotted in Fig. 3(e) illustrates how the number of arms alter the fragment contribution. The overall features of  $\lambda_{\text{tot}}(\omega)$  are very similar for the three molecules with the exception of an additional peak

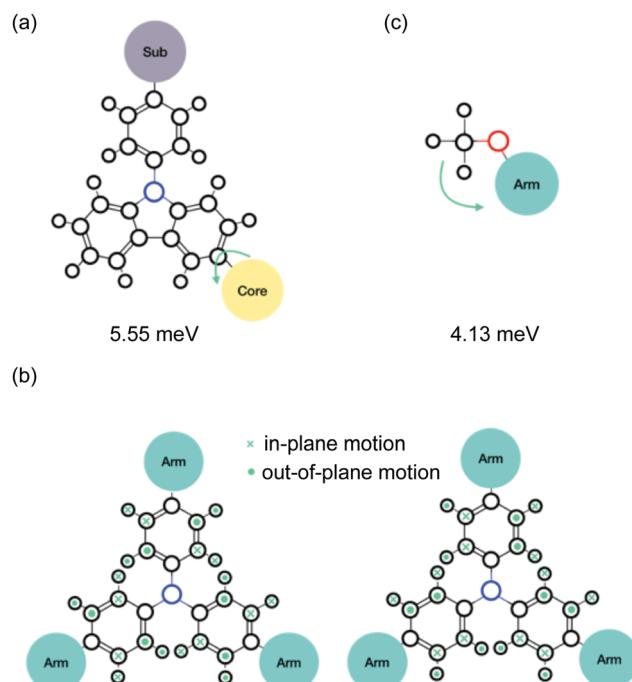


Fig. 4 Illustration of (a) the inter-fragment rotational modes present in the 1st peak and (b) the two dominant Core fragment modes present in the 5th and the 6th peak of **TPA3PM**. (a and c) The two most contributing fragment modes in **TPA3PM**.



(marked with \*) appearing in the low frequency range of the **TPA3PM** spectrum, (composition shown in Table S3, ESI<sup>†</sup>).

The 1st peak originates from three degenerated (by symmetry) Arm fragment modes, featuring inter-fragment rotation between

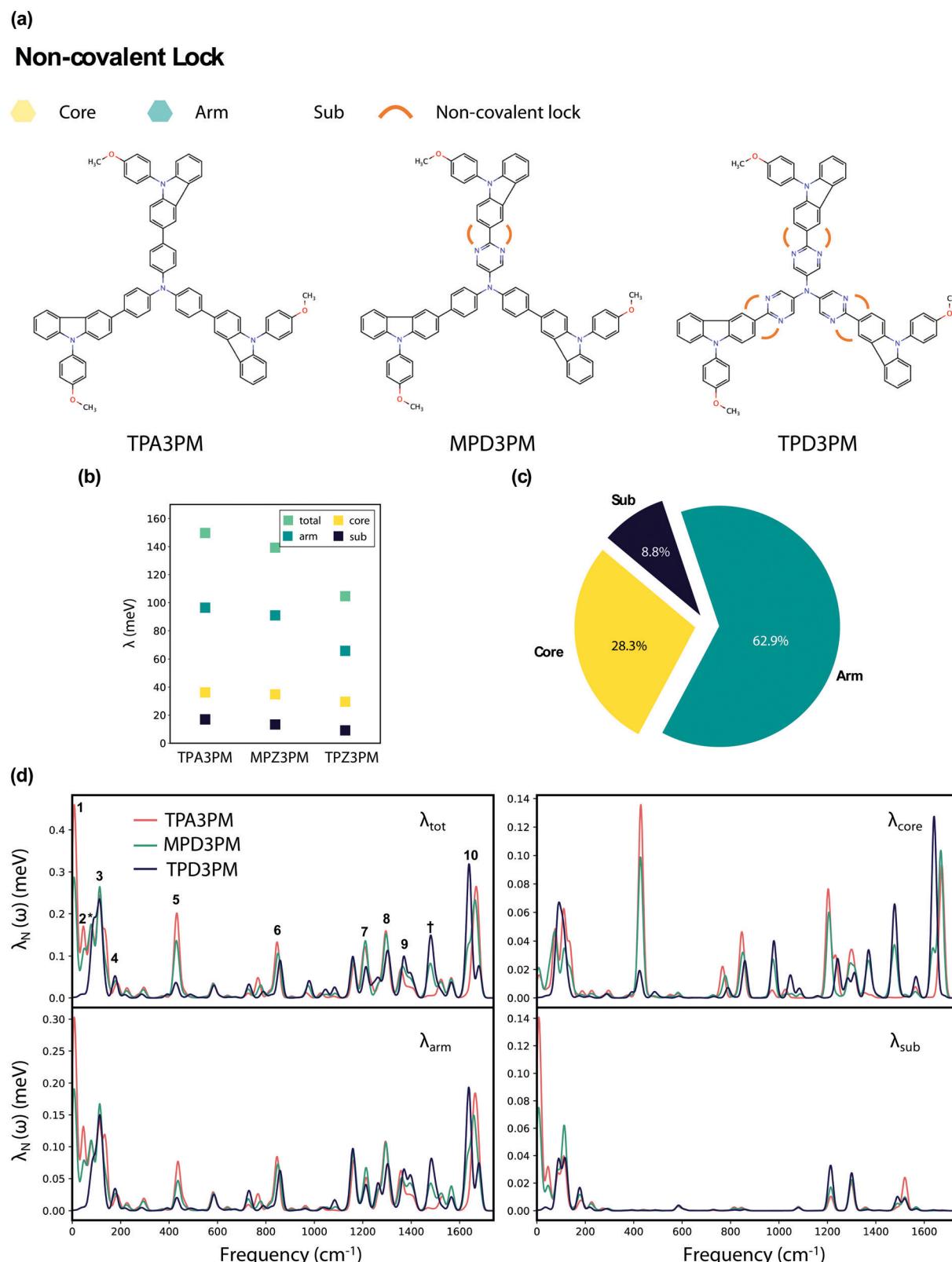


Fig. 5 (a) The **TPA3PM**, **MPD3PM** and **TPD3PM** molecules and their fragments. (b) Total reorganization energy, fragment reorganization energy, (c) the ratio of  $\lambda_{\text{frag}}$  to  $\lambda_{\text{tot}}$  of each fragment, and (d) spectral density plots for  $\lambda_{\text{frag}}$  and  $\lambda_{\text{tot}}$  for **TPA3PM**, **MPD3PM** and **TPD3PM**.



each Arm and the TPA Core (see Fig. 4(a)). As the number of arms increases, a larger number of “arm” atoms contribute to the normal mode, leading to a larger  $\Delta Q_i^N$ . Yet, the eigen frequency of this mode changes negligibly upon the addition of arms, leading to a larger  $\lambda_i^N$  in line with eqn (2). Except for the first peak, the intensity of nearly all peaks decreases as the number of arm increases. The 5th and 6th peaks that remain unchanged between **TPA1PM** and **TPA3PM** constitute another exception.

Unsurprisingly, the decrease of most peaks in  $\lambda_{\text{tot}}(\omega)$  comes from the fact that the relevant normal modes involve the Core and Sub fragments whose contribution is reduced. The main fragment modes present in the unchanged 5th and 6th peaks feature an out-of-plane torsion of benzene rings (Fig. 4(b)) that negligibly change across the **TPA1PM** to **TPA3PM** series. The multi-arm strategy is thus unable to efficiently reduce the vibronic coupling caused by this modes.

The fragment modes with  $\lambda_{jk} > 1.2$  meV ( $10 \text{ cm}^{-1}$ ) for **TPA3PM** (given in Table S4, ESI<sup>†</sup>) show that the two most important contributions comes from the Arm and Sub fragments and correspond to inter-fragment rotational motions (Fig. 4(a and c)). Interestingly, both these modes are present in the 1st peak *via* their coupling with the Arm and Core rotations around the C-C bridging bond. These results suggest that freezing the relative rotational motion between the Core and the Arm can further reduce  $\lambda$  by elimination of the 1st peak. Taking **TPA3PM** as our new starting point, the second strategy is thus to quench this inter-fragment rotational motion by introducing “non-covalent locks” that force the planarity between the benzene in the TPA Core and the carbazole in the PCZ Arm.<sup>36</sup> We expect that locking the out-of-plane mode will partially eliminate the contribution from the 1st peak (see next section).

An alternative strategy consists in reducing the contribution from the 5th and 6th peaks that are not affected by the number of arms. As discussed previously, the corresponding Core fragment modes correspond to an out-of-plane torsion in three benzene groups (Fig. 4(b)). This type of vibrational motion can be significantly restricted by introducing covalent bond between nearby benzene moieties. An example would be to substitute TPA by PCZ or indolo[3,2,1-*j*]-carbazole moiety (ICZ). This alternative will be discussed in the “Core Rigidification” section.

### Non-covalent Lock

The non-covalent locks are introduced to restrict the inter-moiety rotational motion between the Core and Arm in order to further reduce  $\lambda$ . In Fig. 5(a), the benzene rings of the core are replaced by 1 or 3 pyrimidine rings able to form 2 pairs of N···H noncovalent interaction for each pyrimidine-carbazole pair.

The total reorganization energy decreases with increasing the number of pyrimidine in the Core (**MPD3PM** and **TPD3PM**), from 156, 144 to 108 meV (Fig. 5(b)). A further  $\sim 31\%$  reduction in reorganization energy is achieved, demonstrating the efficiency of this strategy. The overall fragment contributions,  $\lambda_{\text{core}}$ ,  $\lambda_{\text{arm}}$  and  $\lambda_{\text{sub}}$  decrease with increasing the number of pyrimidine. By restricting the contribution to the first peak, one expects the reduction in  $\lambda_{\text{arm}}$  and  $\lambda_{\text{sub}}$ . However,  $\lambda_{\text{core}}$  also decreases from 36, 35 to 30 meV. Since all the fragment reorganization energies

decrease, the  $\lambda_{\text{frag}}$  to  $\lambda_{\text{tot}}$  ratio for each fragment of **TPD3PM** does not change significantly compared to **TPA3PM** (Fig. 5(c)). The arms are still the largest contributors (62.9%) followed by Core (28.3%) and Sub (8.8%).

The spectral density in Fig. 5(d) provides several explanations to the observed trends: first, the 1st peak indeed decreases with increasing the number of pyrimidine and almost disappear in the **TPD3PM** spectra of  $\lambda_{\text{tot}}(\omega)$ ,  $\lambda_{\text{arm}}(\omega)$  and  $\lambda_{\text{sub}}(\omega)$ . Interestingly, the 2nd peak, although less intense, follows a similar trend. The fragment mode analysis indicates that the 2nd peak contains an arm-based fragment modes associated with inter-moiety bending motions (see Table S3, ESI<sup>†</sup>). Those are also being restricted by the non-covalent locker.

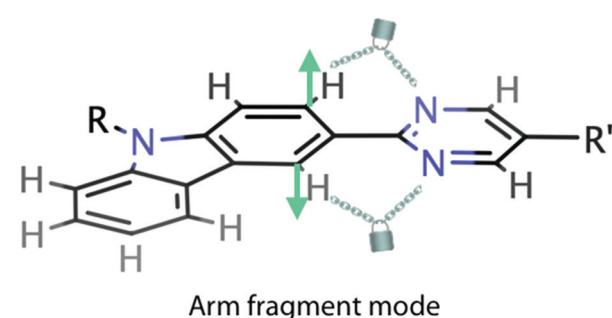
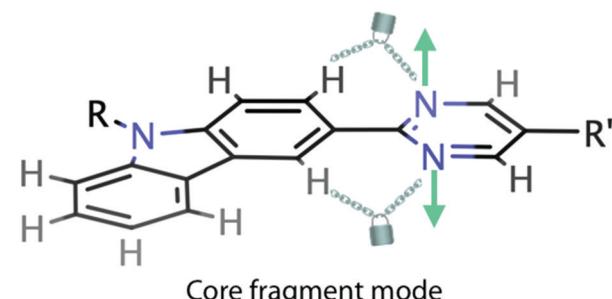


Fig. 6 The fragment modes of **TPD3PM** features out-of-plane torsions involving atoms in the non-covalent interaction pairs (only the motion of non-H atoms near the non-covalent interaction pairs is shown with the green arrows).

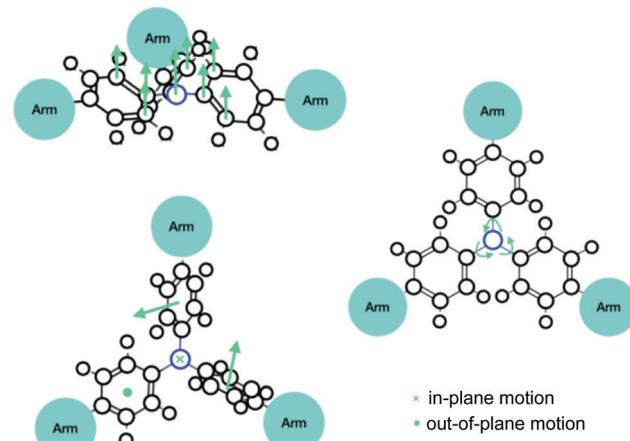


Fig. 7 The main Core fragment modes of **TPA3PM** present in the \* and the 3rd peaks.



Finally, the unexpected decrease in  $\lambda_{\text{core}}$  is shown to essentially arises from the significant reduction of the 5th peak (in  $\lambda_{\text{core}}(\omega)$ ). As mentioned earlier, the 5th peak of **TPA3PM** is dominated by the

benzene out-of-plane torsion. When the pyrimidine rings are introduced, the motion is attenuated by the intramolecular non-covalent interactions, (Fig. 6). The 5th peak of  $\lambda_{\text{arm}}(\omega)$  follows a

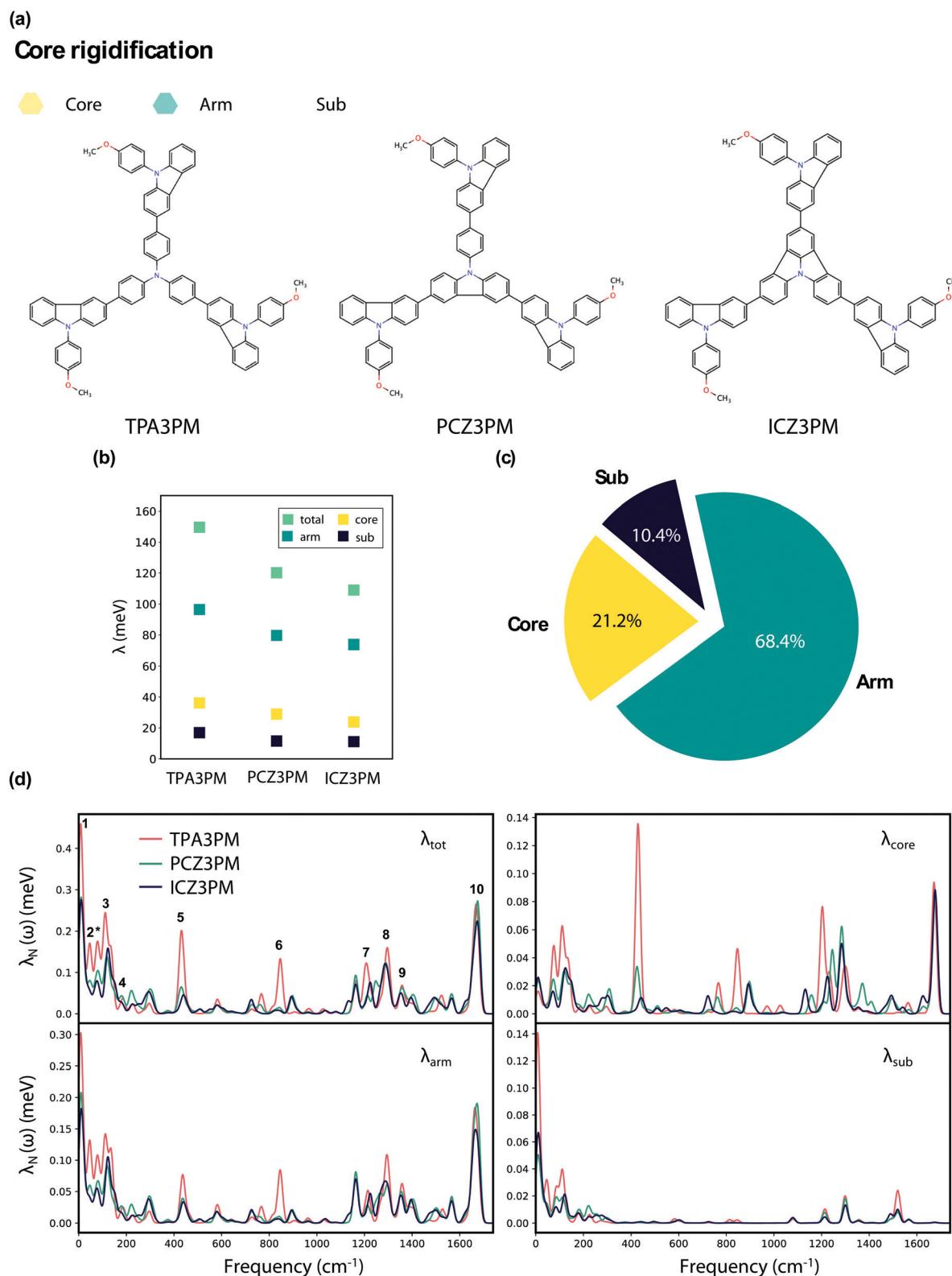


Fig. 8 (a) The **TPA3PM**, **PCZ3PM** and **ICZ3PM** molecules and their fragments. (b) Total reorganization energy, fragment reorganization energy, (c) the ratio of  $\lambda_{\text{frag}}$  to  $\lambda_{\text{tot}}$  of each fragment, and (d) spectral density plots for  $\lambda_{\text{frag}}$  and  $\lambda_{\text{tot}}$  for **TPA3PM**, **PCZ3PM** and **ICZ3PM**.



very similar trend to that of  $\lambda_{\text{core}}(\omega)$  owing to the key Arm fragment mode featuring analogue out-of-plane torsional motion of the carbazole in the PCZ Arm, as shown in Table S3 (ESI<sup>†</sup>). This movement is also restricted by the non-covalent locks (Fig. 6).

Overall, this section demonstrates that the non-covalent lockers reduce the contribution from the fragment modes dominated by out-of-plane inter-moieties motions involving the Core and Arm fragments. This strategy can be applied to other systems. For instance, Malagoli and Brédas showed that the reorganization energy of *N,N'*-diphenyl-*N,N'*-bis(3-methylphenyl)-(1,1'-biphenyl)-4,4'-diamine (TPD), a widely used hole transport material for organic light-emitting diodes, is more than twice of that of its monomer (TPA).<sup>37</sup> They found that  $\lambda$  of TPD is dominated by the central biphenyl fragment, which may be due to the huge change in the inter-ring dihedral angle between neutral and cation ground state geometry. We anticipate the reorganization energy of TPD to be largely reduced upon introduction of non-covalent locks between the two phenyl groups.

### Core rigidification

As mentioned in the “Multi-arm” section, the alternative strategy to reduce the reorganization energy of **TPA3PM** is to restrict the out-of-plane torsion of the benzene moieties in the TPA Core. To achieve this goal, the core can be “rigidified” by introducing covalent bonds between the neighboring benzene moieties transforming TPA into a PCZ (1 bond) or a ICZ Core (2 bonds) shown in Fig. 8(a).

Overall, the core rigidification is relatively efficient and leads to up to 28% reduction in total reorganization energy from 156 to 112 meV (Fig. 8(b)). As expected, the value of  $\lambda_{\text{core}}$  goes down but more surprisingly  $\lambda_{\text{arm}}$  and  $\lambda_{\text{sub}}$  also decreases upon replacement of TPA by ICZ. Given that the reorganization energy of every fragment decreases, the  $\lambda_{\text{frag}}$  to  $\lambda_{\text{tot}}$  ratios of **TPD3PM** is very similar to those of **TPA3PM** (Fig. 8(c)).

The spectral density of  $\lambda_{\text{tot}}(\omega)$  and  $\lambda_{\text{core}}(\omega)$  (Fig. 8(d)) indicates that the intensity of the 5th and 6th peaks decreases from TPA to ICZ. In addition, the \* and 3rd peaks in  $\lambda_{\text{core}}(\omega)$  decreases upon core-rigidification. Since their dominant Core modes involve relative motions between the TPA benzene rings (Fig. 7), it is not surprising that the contribution from these modes is attenuated by the core-rigidification.

The unexpected decrease in  $\lambda_{\text{arm}}$  and  $\lambda_{\text{sub}}$  originates from the reduced contributions in the low-frequency region,  $\omega < 200 \text{ cm}^{-1}$ . In the case of  $\lambda_{\text{arm}}(\omega)$ , this region features relative motions (rotational and translational) between Arms and Core shown in Table S3 (ESI<sup>†</sup>). Given that **ICZ3PM** undergoes less geometrical reorganization in relevant internal coordinates (defined in Fig. S3, ESI<sup>†</sup>) upon oxidation,  $\Delta Q$  and thus the reorganization energy are smaller (see Table S5, ESI<sup>†</sup>). The 1st peak in  $\lambda_{\text{sub}}(\omega)$  decreases upon binding covalently the benzene rings because of the smaller dihedral angle change between the Core and Arm in **ICZ3PM**, which cause the reduction of the contribution from the Sub rotational modes to which they are coupled (*vide supra*).

This section demonstrated that ICZ is a promising building block in organic semiconductors if the goal is to lower the reorganization energy. Interestingly, a recent work by Jiang,

Wang and collaborators concluded that perovskite solar cells exploiting ICZ-based HTMs exhibited outstanding performance.<sup>38</sup> As a building block, the use of the ICZ core is, however, scarce and our work calls for a more systematic analysis of the performance of HTMs for PSCs involving this core fragment.

## Conclusions

We developed a fragment-based decomposition analysis tool for the reorganization energy and demonstrated its usefulness by rationally designing low  $\lambda$  organic HTMs. Starting from the promising dopant-free **TPA1PM** HTM composed of triphenyl amine and carbazole building blocks, three design strategies were adopted to lower the reorganization energy of **TPA1M** ( $\lambda = 213 \text{ meV}$ ): multi-arm, non-covalent lock, and core rigidification. The introduction of multiple arms reduces  $\lambda$  by extending the  $\pi$ -conjugation away from the core. The lower  $\lambda$  of **TPA3PM** (156 meV) originates from diminishing the contributions from the TPA core that arises from two types of modes: (1) the Arm/ Core rotational motions around the C-C bridge and (2) the out-of-plane torsions of the core benzene rings. We attenuated the former by introducing non-covalent locks at the Core and Arm bridging point. The intramolecular non-covalent interactions constrain the reorganization of the nuclei, which allows for a ~31% reduction in the  $\lambda$  of **TPD3PM** (108 meV). On the other hand, the out-of-plane motion of the benzene rings is constrained by the insertion of covalent bonds between the benzene rings, which leads to a nearly 28% decrease in  $\lambda$  (**ICZ3PM**, 112 meV). Given the extensive use of triphenylamine and carbazole building blocks in organic semiconductors, we expect the two types of aforementioned fragment modes to largely contribute to their organization energies. For this reason the proposed strategies should be transferable to a variety of systems, such as, for instance, TPD (a common HTM for organic light emitting diodes). As such, we believe that the precise insight brought by our computational analysis tools has significant potential for facilitating and accelerating the search for novel low- $\lambda$  organic semiconductors.

## Conflicts of interest

There are no conflicts to declare.

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## References

- 1 A. Mishra and P. Bäuerle, *Angew. Chem., Int. Ed.*, 2012, **51**(9), 2020–2067.
- 2 G. Gryn’Ova, K. H. Lin and C. Corminboeuf, *J. Am. Chem. Soc.*, 2018, **140**(48), 16370–16386.
- 3 C. Kunkel, C. Schober, J. T. Margraf, K. Reuter and H. Oberhofer, *Chem. Mater.*, 2019, **31**(3), 969–978.



4 S. Vela, C. Krüger and C. Corminboeuf, *Phys. Chem. Chem. Phys.*, 2019, **21**(37), 20782–20790.

5 K.-H. Lin, A. Prlj and C. Corminboeuf, *J. Phys. Chem. C*, 2017, **121**(39), 21729–21739.

6 Y. Yamashita, *Sci. Technol. Adv. Mater.*, 2009, **10**(2), 024313–024321.

7 H. Sun, J. Gerasimov, M. Berggren and S. Fabiano, *J. Mater. Chem. C*, 2018, **6**(44), 11778–11784.

8 B. Kippelen and J. L. Brédas, *Energy Environ. Sci.*, 2009, **2**(3), 251–261.

9 I. Yavuz, J. B. Lin and K. N. Houk, *Phys. Chem. Chem. Phys.*, 2019, **21**(2), 901–914.

10 N. R. Tummala, S. A. Elroby, S. G. Aziz, C. Risko, V. Coropceanu and J. L. Brédas, *J. Phys. Chem. C*, 2016, **120**(31), 17242–17250.

11 J. Guo, Y. Zhang, W. Cai, Z. Zhang, R. He, W. Shen and M. Li, *Mater. Chem. Phys.*, 2020, **240**, 122058.

12 L. Caliò, S. Kazim, M. Grätzel and S. Ahmad, *Angew. Chem., Int. Ed.*, 2016, **55**(47), 14522–14545.

13 S. Shahnawaz, S. Sudheendran Swayamprabha, M. R. Nagar, R. A. K. Yadav, S. Gull, D. K. Dubey and J. H. Jou, *J. Mater. Chem. C*, 2019, **7**(24), 7144–7158.

14 P. Agarwala and D. Kabra, *J. Mater. Chem. A*, 2017, **5**(4), 1348–1373.

15 C. Schober, K. Reuter and H. Oberhofer, *J. Phys. Chem. Lett.*, 2016, **7**(19), 3973–3977.

16 H. Oberhofer, K. Reuter and J. Blumberger, *Chem. Rev.*, 2017, **117**(15), 10319–10357.

17 K. H. Lin, A. Prlj, L. Yao, N. Drigo, H. H. Cho, M. K. Nazeeruddin, K. Sivula and C. Corminboeuf, *Chem. Mater.*, 2019, **31**(17), 6605–6614.

18 G. Gryn'ova and C. Corminboeuf, *J. Phys. Chem. Lett.*, 2016, **7**(24), 5198–5204.

19 S. F. Nelsen, S. C. Blackstock and Y. Kim, *J. Am. Chem. Soc.*, 1987, **109**(3), 677–682.

20 J. R. Reimers, *J. Chem. Phys.*, 2001, **115**(20), 9103–9109.

21 H. Geng, Y. Niu, Q. Peng, Z. Shuai, V. Coropceanu and J. L. Brédas, *J. Chem. Phys.*, 2011, **135**, 10.

22 Z. Shuai, H. Geng, W. Xu, Y. Liao and J. M. André, *Chem. Soc. Rev.*, 2014, **43**(8), 2662–2679.

23 T. Sato, M. Uejima, N. Iwahara, N. Haruta, K. Shizu and K. Tanaka, *J. Phys.: Conf. Ser.*, 2013, **428**, 1.

24 M. Uejima, T. Sato, K. Tanaka and H. Kaji, *Phys. Chem. Chem. Phys.*, 2013, **15**(33), 14006–14016.

25 M. Huix-Rotllant and N. Ferré, *J. Chem. Theory Comput.*, 2016, **12**(10), 4768–4777.

26 S. J. Park, S. Jeon, I. K. Lee, J. Zhang, H. Jeong, J. Y. Park, J. Bang, T. K. Ahn, H. W. Shin, B. G. Kim and H. J. Park, *J. Mater. Chem. A*, 2017, **5**(25), 13220–13227.

27 H. Geng, Y. Niu, Q. Peng, Z. Shuai, V. Coropceanu and J. L. Brédas, *J. Chem. Phys.*, 2011, **135**, 10.

28 O. Kwon, V. Coropceanu, N. E. Gruhn, J. C. Durivage, J. G. Laquindanum, H. E. Katz, J. Cornil and J. L. Brédas, *J. Chem. Phys.*, 2004, **120**(17), 8186–8194.

29 V. Coropceanu, M. Malagoli, D. A. da Silva Filho, N. E. Gruhn, T. G. Bill and J. L. Brédas, *Phys. Rev. Lett.*, 2002, **89**(27), 1–4.

30 P. J. Stephens, F. J. Devlin, C. F. Chabalowski and M. J. Frisch, *J. Phys. Chem.*, 1994, **98**(45), 11623–11627.

31 A. D. Becke, *J. Chem. Phys.*, 1993, **98**(7), 5648–5652.

32 S. Grimme, S. Ehrlich and L. Goerigk, *J. Comput. Chem.*, 2011, **32**(7), 1456–1465.

33 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, *Gaussian 16, Revision C.01*, Gaussian, Inc., Wallingford CT, 2016.

34 X. Xie, A. Santana-Bonilla, W. Fang, C. Liu, A. Troisi and H. Ma, *J. Chem. Theory Comput.*, 2019, **15**(6), 3721–3729.

35 J. Bednarska, R. Zaleśny, W. Bartkowiak, B. Ośmiałowski, M. Medved and D. Jacquemin, *J. Chem. Theory Comput.*, 2017, **13**(9), 4347–4356.

36 Y. Cheng, Y. Qi, Y. Tang, C. Zheng, Y. Wan, W. Huang and R. Chen, *J. Phys. Chem. Lett.*, 2016, **7**(18), 3609–3615.

37 M. Malagoli and J. L. Brédas, *Chem. Phys. Lett.*, 2000, **327**(1–2), 13–17.

38 X. J. Ma, X. D. Zhu, K. L. Wang, F. Igbari, Y. Yuan, Y. Zhang, C. H. Gao, Z. Q. Jiang, Z. K. Wang and L. S. Liao, *Nano Energy*, 2019, **63**, 103865.

