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Structural and thermodynamic study of Ca A- or Co B-site substituted $\text{SrFeO}_{3-\delta}$ perovskites for low temperature chemical looping applications†

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Perovskite-structured materials, owing to their chemical–physical properties and tuneable composition, have extended their range of applications to chemical looping processes, in which lattice oxygen provides the oxygen needed for chemical reactions omitting the use of co-fed gaseous oxidants. To optimise their oxygen donating behaviour to the specific application a fundamental understanding of the reduction/oxidation characteristics of perovskite structured oxides and their manipulation through the introduction of dopants is key. In this study, we investigate the structural and oxygen desorption/sorption properties of $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$ and $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ ($0 \leq x \leq 1$) to guide the design of more effective oxygen carriers for chemical looping applications at low temperatures (*i.e.* 400–600 °C). Ca A- or Co B-site substituted $\text{SrFeO}_{3-\delta}$ show an increased reducibility, resulting in a higher oxygen capacity at $T \leq 600$ °C when compared to the unsubstituted sample. The quantitative assessment of the thermodynamic properties (partial molar enthalpy and entropy of vacancy formation) confirms a reduced enthalpy of vacancy formation upon substitution in this temperature range (*i.e.* 400–600 °C). Among the examined samples, $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ exhibited the highest oxygen storage capacity (2.15 wt%) at 500 °C, complemented by excellent redox and structural stability over 100 cycles. The thermodynamic assessment, supported by *in situ* XRD measurements, revealed that the oxygen release occurs with a phase transition perovskite–brownmillerite below 770 °C, while the perovskite structure remains stable above 770 °C.

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1. Introduction

The increasing concern about climate change, primarily triggered by anthropogenic carbon dioxide (CO_2) emissions into the atmosphere, has stimulated research into finding solutions for reducing such emissions.^{1–3} Approaches to reduce the human CO_2 footprint include improving the energy efficiency of processes, transitioning to sustainable energy carriers and carbon dioxide capture and storage (CCS) processes. With regard to CCS processes, oxy-fuel combustion has received considerable attention in recent years.⁴ Here, a fuel is combusted in pure oxygen (O_2) (or mixtures of O_2 and CO_2) instead of air. The combustion product is a mixture of CO_2 and water, which can be readily separated. Owing to the absence of nitrogen (N_2) in the combustion process, also the formation of NO_x is largely eliminated.⁵ The main drawback of the conventional oxy-fuel processes is the high energy penalty associated with the production of pure O_2 , typically *via* cryogenic air separation (CAS).⁶

Although CAS is a very mature technology for the production of O_2 ,^{7–9} techno-economic studies of oxy-fuel combustion estimate that conventional CAS (specific power consumption of $0.24 \text{ kW h kgO}_2^{-1}$) consumes $\sim 10\%$ of the net oxy-fuel thermal power plant output and represents in addition $\sim 40\%$ of the total equipment costs.^{10–13} A potentially more cost-efficient method to produce pure O_2 is chemical looping air separation (CLAS), which dates back to the original Brin process.¹⁴ The CLAS process utilises metal oxides (referred to as oxygen carriers, OCs) capable of releasing gas phase O_2 at elevated temperatures (up to 1000 °C).^{9,15,16} The CLAS process consists of two steps that occur in a cyclic fashion: first, O_2 is released (uncoupling, reaction (1) by the metal oxide. Each metal oxide, Me_xO_y , possesses an equilibrium O_2 partial pressure ($p_{\text{O}_2,\text{eq}}$) for a given temperature at which the transition to a lower oxidation state, $\text{Me}_x\text{O}_{y-2}$, occurs. Gaseous O_2 is thus released when the O_2 partial pressure (p_{O_2}) in the given environment is lower than the respective thermodynamic equilibrium value (*i.e.* $p_{\text{O}_2,\text{eq}}$). The metal oxide is regenerated (reaction 2) at lower temperatures and/or higher p_{O_2} (typically using air).

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A series of experimental and techno-economic modelling studies have suggested that CLAS could produce O₂ at 40–70% lower costs than CAS.^{10,11,17,18} In fact, power costs for CLAS at an industrial scale have been estimated to be as low as 0.06 kW h kgO₂^{−1}, against the benchmark 0.24 kW h kgO₂^{−1} of a conventional cryogenic system.¹³

A key challenge for CLAS is finding OCs with suitable thermodynamic properties. The $p_{O_2,eq}$, the reaction enthalpy of a given OC and its oxygen storage capacity (OSC) largely determine its suitability for CLAS.¹⁹ For example, low values (< 1 wt%) of the OSC (defined as the mass of redox-active lattice O₂ of an OC divided by the total mass of the OC) imply that an OC would have to be regenerated (eqn (2)) more frequently in a given time compared to an OC with a high OSC. Further, the OC should have a stable reactivity over multiple redox cycles, besides being low cost and environmentally friendly. Typical OCs that have been demonstrated to be conceptually suitable for CLAS are based on the monometallic transition metal oxides CuO, Co₃O₄ or Mn₂O₃.^{20–23} Transition metal oxides often require high temperatures (up to 1000 °C) to release gaseous O₂. Under operation, metal oxides may thus reach temperatures close to their melting points, leading to sintering and in turn a decline in their redox activity and stability. To mitigate sintering, the metal oxide is often stabilized by an inert support, e.g. calcium aluminate.²⁴

To lower the energy costs of the CLAS processes further, it is desirable to identify metal oxides with suitable thermodynamic properties (reaction enthalpy and $p_{O_2,eq}$) that are capable of releasing O₂ at lower temperatures than typical CLOU temperatures, i.e. in the range 400–600 °C. However, operation at such temperatures may reduce the rates of the redox reactions, implying larger system volumes (and in increased capital costs) of the CLAS process. An interesting class of materials that has the potential to release O₂ at temperatures < 600 °C at a fast rate are perovskites.²⁵ Perovskites, with the general formula ABO_{3−δ} (where the A cation is typically a rare earth metal or alkaline earth metal, the B cation is a transition metal and delta is the oxygen non-stoichiometry), are non-stoichiometric compounds, with a non-stoichiometry δ that depends on temperature and p_{O_2} . Fast redox kinetics have been observed even at temperatures < 500 °C²⁶ due to high rates of lattice oxygen diffusion.²⁷ Perovskites can expand significantly the design space of OCs for chemical looping applications owing to their versatile composition and the ability to accommodate various A- and B-site cations.²⁵ This compositional flexibility allows the optimisation of their chemical-physical properties (i.e. thermodynamics, melting temperature, OSC). In fact, perovskite-type materials have been studied intensively in the past years for a wide range of applications, including thermochemical water splitting, SOFC and high temperature pressure swing adsorption (HT-PSA).^{27–31} The basic crystal structure of perovskites is cubic (*Pm3m*), although deviations in their composition (e.g. by cation substitution) can generate distortions of the ideal cubic lattice structure. While monometallic transition metal oxides donate lattice oxygen *via* a phase transition, in perovskites the release of gaseous O₂ can occur without any phase transition

and results in a change of the oxygen non-stoichiometry δ:

where $\Delta\delta = (\delta - \delta_0)$ can range from 0 to 0.5; $\Delta\delta = 0.5$ represents the maximal amount of O₂ that can theoretically be released while preserving the perovskite structure provided the initial non-stoichiometry δ_0 is zero.²⁷ When the complete reduction of the perovskite ABO₃ occurs, it transforms to its defect-ordered variant brownmillerite (ABO_{2.5}), with an orthorhombic basic crystal structure (*lbm2*).²⁷ Taking the non-stoichiometric oxide with the composition SrFeO₃ as an example, the maximal OSC due to the change in δ would be $0.5 \cdot M_{O_2}/M_{OC} = 4.17$ wt%, where M_{O_2} and M_{OC} are the molecular weight of O₂ and the OC, respectively. Hence, a small initial value of δ_0 is desired to ensure a high OSC.

Density functional theory (DFT) calculations have shown that the variations in A- and B-site cations may induce a change in the oxidation state of the B-site transition metal and enhance the O₂ donation properties of the material, by reducing the energy of formation of oxygen vacancies.^{29,32} For example, Mishra *et al.*³³ performed DFT studies on the effect of dopants such as Sr and Fe in CaMnO₃ and found that the dopants enhance the O₂ donation properties of the material through a distortion of the lattice structure and changes in the oxygen coordination.

From an application point of view, OCs must possess excellent structural stability and reactivity. CaMnO₃, for example, is one of the most extensively investigated perovskites²⁵ but in spite of a reasonably high OSC of ~1 wt%³⁴ it undergoes an irreversible phase transformation to Ca₂MnO₄ and CaMn₂O₄ when exposed to repeated cycles of air and N₂ at high temperatures (i.e. $T = 1170$ °C), as these two latter phases are thermodynamically more stable than the parent CaMnO₃ under the given operating conditions.³⁵ The phase transition from perovskite to spinel is accompanied by a gradual decrease of the OSC, which makes such compositions unattractive for industrial applications.

Among the large number of possible non-stoichiometric oxides, strontium and iron-based perovskites have been identified as suitable OCs for low-temperature applications.^{19,36,37} For example, Miura *et al.* found that A-site substitution of the SrFeO₃ system by Ca improves the OSC when operating cyclically between 5% O₂ and air, identifying Sr_{0.76}Ca_{0.24}FeO_{3−δ} with an OSC of 1.1 wt% at 550 °C as an attractive OC (the OSC of the undoped SrFeO_{3−δ} was only 0.6 wt%).²⁸ In addition, substitution of the B-site of SrFeO₃ with Co was investigated by Ikeda *et al.* and they found that the substitution of 85% of Fe with Co allows the release of O₂ at temperatures as low as 300 °C instead of 600 °C for the unsubstituted SrFeO_{3−δ}.²⁹ Dou *et al.*³⁸ recently showed that the A- and B-site co-substitution in SrFeO₃ of Sr and Fe by Ca and Co, respectively, lowers the activation energy of oxygen diffusion and surface oxygen exchange i.e. faster redox kinetics compared to the unsubstituted perovskite. In particular, they found that Sr_{0.8}Ca_{0.2}Fe_{0.4}Co_{0.6}O₃ releases 1.2 wt% at 400 °C, even under when the p_{O_2} was relatively high (i.e. 0.05 atm).

As outlined above, the thermodynamic properties of the OC play an important role in the design of the CLAS process and the identification of the optimal working conditions.¹⁹ Although $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$ and $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ perovskites have been shown to possess promising chemical-physical properties, thermodynamic data for such materials is scarce or non-existent albeit critical for process modelling and further material optimisation.

Motivated by the promising results reported by Miura *et al.*²⁸ and Ikeda *et al.*,²⁹ we focused in our study on the thermodynamic and structural analysis of Ca A-site or Co B-site substituted SrFeO_3 ($\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$ and $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$, $0 \leq x \leq 1$) aiming at elucidating changes in the chemical-physical properties (*i.e.* OSC, thermodynamics, structural stability) upon substitution. This work shall provide guidelines for the design of more effective OCs for chemical looping applications at low temperatures (*i.e.* 400 to 600 °C). The structure of the substituted perovskites was characterised by X-ray powder diffraction (XRD) and their OSC and redox activity using a thermogravimetric analyser (TGA) under a wide range of oxygen partial pressures, $p\text{O}_2$, and temperatures, T , allowing us to calculate thermodynamic properties such as the molar enthalpy and entropy of vacancy formation. Finally, the long-term reactivity and chemical stability of the OC was investigated over multiple redox cycles *via* thermogravimetric, fixed bed and *in situ* XRD measurements.

In general, the OSC and reducibility vary distinctly with the composition. Among the Ca-substituted perovskites, $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ had the highest OSC (2.15 wt%) at temperatures as low as 500 °C, while among the Co-substituted perovskites, $\text{SrFe}_{0.2}\text{Co}_{0.8}\text{O}_{3-\delta}$ was individuated as the most promising OC for CLAS application due to its ability to release a high amount of O_2 (OSC of 1.1 wt%) at 500 °C in atmospheres containing as much as 5% O_2 .

2. Experimental methods

Synthesis of the oxygen carriers

$\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$ and $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ ($0 \leq x \leq 1$) materials were synthesised by dissolving stoichiometric amounts of $\text{Sr}(\text{NO}_3)_2$ (Sigma Aldrich, ACS reagent, 99.9%), $\text{Fe}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$ (Sigma Aldrich, ACS reagent, $\geq 98\%$) and $\text{Ca}(\text{NO}_3)_2 \cdot 4\text{H}_2\text{O}$ (Sigma Aldrich, 99%) or $\text{Co}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$ (Sigma Aldrich, 99%) in deionised water together with citric acid (3:1 molar ratio of citric acid to metal ions) and ethylene glycol (1.2:1 molar ratio of ethylene glycol to citric acid). The solution was heated to 80 °C and held at this temperature until a gel formed. The gel was dried at 130 °C for 10 h in an oven and subsequently crushed to a fine powder. The powder was calcined in air in two steps: the first calcination step was performed with a heating rate of 3 °C min^{-1} and holding at 500 °C for 7 h. The second calcination step was performed with a heating rate of 5 °C min^{-1} and holding at 1000 °C (for $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$) or 1200 °C (for $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$) for 10 h. Calcination temperatures higher than the selected ones resulted in partially molten samples.

Characterisation

All gases used in this study were provided by Carbogas from cylinders (grade 5.0) and thus contained small quantities of impurities, such as ~ 10 ppm O_2 in N_2 . The flow rates of gas were measured and controlled at normal temperature and pressure (NTP) using mass flow controllers (Bronkhorst EL-Flow).

X-ray powder diffraction (XRD). X-ray diffractograms of the materials were collected using a PANalytical Empyrean X-ray powder diffractometer equipped with a X'Celerator Scientific ultra-fast line detector and Bragg-Brentano HD incident beam optics using Cu K α radiation (45 kV and 40 mA). The scans were collected in the 2θ range of 10–90° with a scanning speed of 0.022° s^{-1} and a step size of 0.016°. *In situ* XRD measurements were performed in the same diffractometer using an Anton Paar XRK 900 reactor chamber (Macor sample holder) in a reflection geometry, with gas flow through the sample and uniform temperature control. Each sample was heated up to 900 °C under a flow of air and then investigated under three different $p\text{O}_2$ obtained by using either pure N_2 (200 mL min^{-1}) or gas mixtures of N_2 and air (first 200 mL min^{-1} N_2 and 4 mL min^{-1} air, then 200 mL min^{-1} N_2 and 8 mL min^{-1} air), followed by re-oxidation in air (200 mL min^{-1}). The resulting $p\text{O}_2$ from the mixtures of N_2 and air were 7.09×10^{-5} bar, 7.87×10^{-3} bar, 8.64×10^{-3} bar and 2.09×10^{-1} bar, as measured by means of an oxygen sensor (SETNAG MicroPoas) located at the outlet of the instrument when no sample was loaded. The same procedure was repeated at 800, 750, 700, 650, 600, 550 and 450 °C. The scan time of one diffractogram over the 2θ range of 31–35°, to study the evolution of the main perovskite peak (110), was ~ 75 s.

Inductively coupled plasma – optical emission spectroscopy (ICP-OES). The actual extent of the substitution of the A- or B-site of the perovskite with Ca or Co, respectively, was quantified by ICP-OES using an Agilent 5100 VDV. A multi-element standard (Sigma Aldrich) was utilised for the calibration.

Nitrogen physisorption. A Quantachrome NOVA 4000e N_2 adsorption analyser was used to determine the pore volume of the OCs used in the fixed bed experiments. Each sample was degassed at 250 °C for 2 h prior to the measurements. The BJH model was used to calculate the pore volume of the materials.

Electron microscopy. Transmission electron microscopy (TEM) was performed on an FEI Talos F200X operated at 200 kV. The instrument was equipped with SuperX EDX comprising four SDD detectors. The morphology of the fresh OCs was investigated using scanning electron microscopy (FEI Quanta 200F).

Iodometric titration. The actual oxygen content of the as-synthesised OCs was determined by iodometric titration. An OC sample (250.0 mg) was dissolved in 50 mL HCl (6 M). A portion of the solution (10 mL) was placed in an Erlenmeyer flask and 1.00 g of potassium iodide (KI) was added. After a few minutes, the solution was diluted with DI water and the liberated iodine was titrated with sodium thiosulfate $\text{Na}_2\text{S}_2\text{O}_3$ (0.05 M). Sodium thiosulfate was always standardised against KIO_3 before each series of titration. The titration end-point was detected using a

starch-indicator, which was introduced into the solution before the titration.

Each titration was repeated at least twice and the values reported represent an average of the individual measurements. All reagents used were of analytical grade.

Thermogravimetric analyser (TGA). The OSC of the materials was determined using a TGA (Mettler Toledo TGA/DSC 3+). A small amount of the sample (~30 mg) was placed in an alumina crucible (70 μL), heated to 900 $^{\circ}\text{C}$ at a rate of 25 $^{\circ}\text{C min}^{-1}$ and held for 150 min. The total flow rate of gas was 150 mL min^{-1} and included a constant N_2 flow of 25 mL min^{-1} over the microbalance. During that isothermal period, the gas atmosphere was altered to three different $p\text{O}_2$ by using either pure N_2 (125 mL min^{-1}), a gas mixture of N_2 and air (60 mL min^{-1} and 65 mL min^{-1} , respectively) or only air (125 mL min^{-1}). The resulting $p\text{O}_2$ were 7.09×10^{-5} bar, 9.96×10^{-2} bar and 2.09×10^{-1} bar. The same procedure was repeated for temperatures ranging from 800 $^{\circ}\text{C}$ to 450 $^{\circ}\text{C}$ in decrements of 50 $^{\circ}\text{C}$. The temperature range in which the OCs release O_2 was determined through temperature-programmed reduction (TPR) experiments. In these experiments, the sample was heated to 1000 $^{\circ}\text{C}$ at a rate of 5 $^{\circ}\text{C min}^{-1}$ in N_2 and held for 30 min.

To determine the thermodynamic properties of the OCs, a high pressure thermobalance (LINSEIS L81PT High Pressure TGA-DTA/DSC) was used. Blank experiments were performed before the analysis of the material to correct for the effect of buoyancy on the balance. About 300 mg of sample, placed in an alumina crucible (500 μL), was used in each experiment. The sample sat on top of a layer of quartz wool to reduce the depth of the crucible and to minimise gas diffusion effects. First the sample was heated to 125 $^{\circ}\text{C}$ under a flow of synthetic air at a total flow rate of 1 L min^{-1} and held for 30 min to remove adsorbed water. The temperature program of each experiment consisted of four consecutive temperature segments: 800, 700, 600 and 500 $^{\circ}\text{C}$. At each temperature, the gas environment was varied to obtain seven different $p\text{O}_2$. At the beginning and at the end of each temperature segment, the $p\text{O}_2$ was identical to assess whether irreversible structural changes occurred. For each sample, experiments were performed at an absolute pressure of 1, 5 or 7 bar, covering a wide range of $p\text{O}_2$ (*i.e.* from 7.09×10^{-5} to 1.4 bar). Additionally, a constant N_2 flow of 1.5 L min^{-1} was used to purge the balance. The change in oxygen non-stoichiometry $\Delta\delta$ was calculated from the recorded weight changes as:

$$\Delta\delta = \Delta m_s \frac{M_{\text{OC}}}{M_{\text{O}_2}} \quad (4)$$

where Δm_s is the measured weight loss, calculated as the difference between the mass at the set point temperature (*i.e.* before switching to pure N_2) and the mass at the investigated $p\text{O}_2$.

The standard partial molar Gibbs free energy change for reduction (eqn (3)) is defined as:

$$\Delta\bar{g}_{\text{O}}^{\circ}(\delta, T) = -\frac{1}{2}RT \ln(p\text{O}_2/p^{\circ}) \quad (5)$$

where p° is the standard pressure (1 bar), T is the temperature (K) and R is the ideal gas constant (8.314 $\text{J K}^{-1} \text{mol}^{-1}$). $\Delta\bar{g}_{\text{O}}^{\circ}$ can be related to the standard partial molar enthalpy ($\Delta\bar{h}_{\text{O}}^{\circ}$) and entropy change ($\Delta\bar{s}_{\text{O}}^{\circ}$):

$$\Delta\bar{g}_{\text{O}}^{\circ}(\delta, T) = \Delta\bar{h}_{\text{O}}^{\circ}(\delta) - T\Delta\bar{s}_{\text{O}}^{\circ}(\delta) \quad (6)$$

Eqn (5) and (6) can be combined such that $\Delta\bar{h}_{\text{O}}^{\circ}$ and $\Delta\bar{s}_{\text{O}}^{\circ}$ are expressed as a function of $p\text{O}_2$ and temperature, for a constant oxygen non-stoichiometry δ :

$$-\ln(p\text{O}_2/p^{\circ}) = \frac{1}{T} \frac{2\Delta\bar{h}_{\text{O}}^{\circ}}{R} - \frac{2\Delta\bar{s}_{\text{O}}^{\circ}}{R} \Big|_{\delta=\text{const.}} \quad (7)$$

Accordingly, plotting the logarithm of $p\text{O}_2/p^{\circ}$ over the inverse of temperature (Arrhenius plot) for constant $\Delta\delta$ allows extracting the thermodynamic properties of the OCs (*i.e.* the standard partial molar enthalpy and entropy change) from the slope $2\Delta\bar{h}_{\text{O}}^{\circ}/R$ and intercept $2\Delta\bar{s}_{\text{O}}^{\circ}/R$.

Cyclic redox experiments. The cyclic redox stability of the OCs $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ and $\text{SrFe}_{0.2}\text{Co}_{0.8}\text{O}_{3-\delta}$, *i.e.* the two most promising materials of the Ca- or Co-substituted materials, respectively, were investigated also in a fixed bed reactor over 100 cycles. Here, a quartz tube reactor (ID 10 mm, 30 cm long) was loaded with 3 g of the OC (sieved to 180–250 μm), which was held in the centre of the tube by means of a plug of quartz wool, giving a bed of length 20 mm. The OC was heated to 600 $^{\circ}\text{C}$ in air (50 mL min^{-1}) by using a tubular furnace controlled *via* an N-type thermocouple. Each redox cycle consisted of 40 min reduction under a flow of N_2 (50 mL min^{-1}) followed by 20 min oxidation in air (50 mL min^{-1}). The fraction of O_2 in the gas stream at the outlet of the reactor was determined by means of a paramagnetic analyser (ABB EL3020, Magnos 206). A blank experiment was performed at 600 $^{\circ}\text{C}$ under a flow of N_2 (50 mL min^{-1}) using the same quartz tube reactor filled with an inert material (SiC) to account for mixing effects in the gas sampling system when switching between gases. The total amount of O_2 released by the OC in each cycle was quantified using eqn (8):

$$N_{\text{O}_2} = \int_{t_1}^{t_2} \frac{x_{\text{O}_2}}{1 - x_{\text{O}_2}} dt \dot{N}_{\text{N}_2} \quad (8)$$

where $\dot{N}_{\text{N}_2}$ is the total molar flowrate of N_2 , t_1 and t_2 are the time when the gas flow rate is switched from air to N_2 and when $x_{\text{O}_2} = 0$, respectively, and x_{O_2} is the mole fraction of O_2 measured at the outlet of the reactor when flowing only N_2 , after subtraction of the blank measurement.

3. Results and discussion

Characterisation of the oxygen carriers

Fig. 1 shows the XRD patterns of as-prepared $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ and $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$ at ambient temperature. The diffraction patterns of the samples with composition $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ ($0 \leq x \leq 0.8$) and $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$ showed good agreement with the reference perovskite cubic structure (PDF 39-0954)

($0 \leq x \leq 0.25$), and no impurity phases were detected. For Co B-site substitution and $x = 1$, additional diffraction peaks were observed corresponding to hexagonal $\text{SrCoO}_{2.52}$ (PDF 40-1018), in agreement with other works.^{29,39}

In the case of Ca A-site substitution, the samples with a composition $0.3 \leq x \leq 1$ revealed reflections due to further oxide phases that were identified as CaO (PDF 38-5590) and the brownmillerite phase $\text{Ca}_{1.5}\text{Sr}_{0.5}\text{Fe}_2\text{O}_5$ (PDF 46-0480). In an attempt to enhance the purity of the perovskites with Ca substitutions in the range $0.3 \leq x \leq 1$, higher calcination temperatures (*i.e.* 1100 °C or 1200 °C, for 10 h) were used, yet the same extent of impurity phases was observed, in agreement with previous reports.²⁸

High signal-to-noise X-ray diffractograms of some of the samples (Fig. S1, ESI†) were collected also at the ID31 beamline of the European Synchrotron Radiation Facility (ESRF), confirming the formation of a phase-pure perovskite structure only in the range $0 \leq x \leq 0.8$ for $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ and $0 \leq x \leq 0.25$ for $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$. For each sample in this range of substitution, the actual ratio of the single elements was determined by ICP-OES analysis. The values obtained are in good agreement with the expected values based on the amount of nitrates used in the preparation method, confirming the successful synthesis of phase-pure materials (Table S1 of the ESI†).

The microscopic distribution of the elements composing the OCs was probed by energy dispersive X-ray spectroscopy (EDS) coupled with scanning transmission electron microscopy (STEM). Fig. 1(c) shows a high-angle annular dark-field (HAADF) image and the corresponding EDS maps of Fe, Ca and Sr, confirming a homogeneous dispersion of all elements. Agglomerates or impurity phases were not observed. For compositions that, according to

XRD, do not yield a phase-pure perovskite phase ($\text{Sr}_{0.2}\text{Ca}_{0.8}\text{FeO}_{3-\delta}$, Fig. 2(d)), a non-homogenous distribution of the elements was observed, due to the formation of additional (segregated) mixed oxide phases.

The following investigations consider phase-pure perovskites only to exclude the influence of the impurity phases on the observed reduction/oxidation behaviour.

Oxygen desorption properties

When heating from ambient temperature to the set point temperature in a flow of air, all samples exhibited a significant weight loss that suggests their $p\text{O}_{2,\text{eq}}$ increases with temperature (Fig. S2, ESI†). The effect of the composition of the perovskite on their $p\text{O}_{2,\text{eq}}$ and thermodynamic properties is discussed in more detail in the next section.

Thermogravimetric analysis revealed a strong influence of the perovskite composition (for both A- and B-site substitution) on the amount of desorbed/sorbed O_2 in the temperature range 400–900 °C.

The measured oxygen storage capacity (OSC_m) for each temperature was evaluated as the weight change at equilibrium (eqn (9)) after changing the atmosphere from air ($p\text{O}_2 = 2.09 \times 10^{-1}$ bar) to N_2 ($p\text{O}_2 = 7.09 \times 10^{-5}$ bar), and is reported in Fig. 2(a) and (b) for $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ ($x = 0, 0.2, 0.4, 0.6, 0.8$) and $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$ ($x = 0.1, 0.2, 0.25$), respectively.

$$\text{OSC}_\text{m} = \frac{m_\text{air} - m_{\text{N}_2}}{m_{\text{N}_2}} \times 100 \quad (9)$$

where m_air and m_{N_2} are the weight of the OC measured at equilibrium in air and N_2 , respectively.

Fig. 1 XRD patterns of as-prepared (a) $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ ($x = 0, 0.2, 0.4, 0.6, 0.8, 1$) and (b) $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$ ($x = 0.1, 0.2, 0.25, 0.4, 0.6, 0.8, 1$). HAADF images and elemental composition of (c) $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ and (d) $\text{Sr}_{0.2}\text{Ca}_{0.8}\text{FeO}_{3-\delta}$.

Fig. 2 Oxygen storage capacity as a function of composition in the temperature range 500–900 °C and differential mass as a function of temperature for (a and c) $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ and (b and d) $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$.

The variation in the OSC with the composition of the perovskites was more evident in the temperature range 500–700 °C, showing that the substitution of A or B-site cations promotes the release of O_2 at low temperatures. Fe B-site substitution by Co by 20% increased the observed OSC from 0.8 wt% (when $x = 0$) to 1.25 wt% at 500 °C. Increasing the amount of Co further resulted in an even higher OSC, reaching a maximum of the OSC = 1.42 wt% for $x = 0.6$ at 500 °C. In the case of Sr A-site substitution by Ca, the highest OSC of 2.15 wt% was observed for $x = 0.2$ at 500 °C. The substitution of Sr by Ca did not yield an increased OSC at high temperatures (700–900 °C), likely due to the high $p_{\text{O}_2, \text{eq}}$ in this temperature range.

The dependence of the composition of the perovskite on the amount of O_2 released as a function of temperature was assessed further by temperature programmed reduction (TPR).

For samples of composition $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ ($0 \leq x \leq 0.8$), a mass loss due to the release of O_2 occurred at temperatures as low as 330 °C (Fig. 2(c)). Importantly, an increased substitution of Fe by Co resulted in an increased rate of mass loss (with a maximum rate for $x = 0.6$), *i.e.* an increased reducibility of the metal oxide at low temperature. To summarise, B-site substitution of Fe by Co represents a possibility to release O_2 at relatively low temperatures compared to the original material. In the case of Ca (A-site) substituted perovskites ($\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$, $0 \leq x \leq 0.25$), reported in Fig. 2(d), an increasing quantity of Ca in the perovskite increased the rate of reduction in N_2 compared to the original $\text{SrFeO}_{3-\delta}$.

Similar to the case of Co B-site substitution, the range of temperatures in which the materials released the entire quantity of O_2 decreased from 330–750 °C (SrFeO_3) to 330–550 °C.

The cyclic redox stability of the OCs $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ and $\text{SrFe}_{0.2}\text{Co}_{0.8}\text{O}_{3-\delta}$, *i.e.* the two most promising materials of the Ca- or Co-substituted materials, respectively, was studied using

Fig. 3 Cyclic redox stability of $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ over 30 cycles at 500 °C and SEM micrographs of $\text{Ca}_{0.2}\text{Sr}_{0.8}\text{FeO}_{3-\delta}$ before and after 30 redox cycles at 500 °C.

both a TGA and a fixed bed reactor. Fig. 3 plots the oxygen capacity of $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ over 30 cycles (TGA at 500 °C). There was no measurable loss in the performance of $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ and the cyclic OSC was ~2.1 wt% in all cycles but the first two, in which the OSC was ~2.0 wt%. Similarly, the morphology of the material remained unaltered during the cyclic test, as observed by SEM (Fig. 3).

XRD patterns collected before and after the cycling experiment did not show any substantial difference (Fig. S3(a), ESI[†]), in particular no phase separation was observed. Rietveld full pattern refinement confirmed an unchanged cell parameter (Fig. S3(b), ESI[†]). The cyclic redox stability of $\text{SrFe}_{0.2}\text{Co}_{0.8}\text{O}_{3-\delta}$ was also investigated in the TGA at 500 °C over 30 redox cycles (Fig. S4, ESI[†]). For $\text{SrFe}_{0.2}\text{Co}_{0.8}\text{O}_{3-\delta}$ we observed an activation period over the first 10 cycles during which the amount of O_2 absorbed by the OC increased with cycle number. This activation period was followed by a stable reactivity with no loss in performance with cycle number. Similar to $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$, when comparing the XRD patterns of $\text{SrFe}_{0.2}\text{Co}_{0.8}\text{O}_{3-\delta}$ before and after the cycling experiment no apparent change in the structure of the OC was observed (Fig. S5(a), ESI[†]). The Rietveld refinement of the full XRD pattern of the OC collected after the cyclic experiment confirmed its unchanged crystal structure (Fig. S5(b), ESI[†]). This suggests that the observed activation of the OC was not related to a chemical phenomenon, but possibly a physical phenomenon, such as an improved mass transport within the material. Indeed, a slight increase in the pore volume and porosity of the OC $\text{SrFe}_{0.2}\text{Co}_{0.8}\text{O}_{3-\delta}$ (but not $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$) was observed upon cycling (Table S2, ESI[†]), which may explain the observed activation period for this material.

In the fixed bed experiments, the number of redox cycles was increased to 100. Fig. 4, plots the p_{O_2} for the OC $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$, as a function of time over 100 redox cycles performed at 500 °C. For comparison, the p_{O_2} of the blank experiment (reactor filled with SiC of the same volume as in the experiment with OC) is given by empty circles. $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$, shows no significant difference in the measured p_{O_2} between cycle 1 (red) and cycle 100 (black), confirming the stable cyclic performance of the material, observed in the previous TGA experiments. The area between the p_{O_2} profiles obtained from the OC and the blank experiment using SiC allows to calculate

Fig. 4 pO_2 measured at the outlet of the fixed bed reactor for cycle 1 (red) and cycle 100 (black) with the respective evolution of $\Delta\delta$ (dashed lines) for $Sr_{0.8}Ca_{0.2}FeO_{3-\delta}$ at 500 °C.

the total amount of O_2 released by the OC over 40 min and was determined as $\sim 640 \mu\text{mol g}_{\text{OC}}^{-1}$, which is $\sim 97\%$ of the total O_2 that could be released based on the TGA-determined OSC ($\text{OSC}_m = 2.15 \text{ wt}\%$). Fig. 4 also plots the evolution of the oxygen non-stoichiometry $\Delta\delta$ (dashed lines) for cycle numbers 1 and 100; no apparent difference between the two cycle numbers is observed.

The same type of cyclic experiment was performed also for $SrFe_{0.2}Co_{0.8}O_{3-\delta}$ (Fig. S6, ESI†). In agreement with what has been observed in the TGA, the pO_2 profiles for cycle 1 (green) and cycle 100 (black) were different. The amount of O_2 released by $SrFe_{0.2}Co_{0.8}O_{3-\delta}$ increased from cycle 1 to 10 from $\sim 67\%$ to $\sim 96\%$ of the total O_2 that could be released based on the TGA-determined OSC ($\text{OSC}_m = 1.2 \text{ wt}\%$ or $\sim 370 \mu\text{mol } O_2 \text{ g}_{\text{OC}}^{-1}$). The activation phase (first 10 cycles) of the O_2 release was followed by a stable reactivity for the remaining 90 cycles. Fig. S7 (ESI†) compares the pO_2 profiles of the 100th cycle for $SrFe_{0.2}Co_{0.8}O_{3-\delta}$ and $Sr_{0.8}Ca_{0.2}FeO_{3-\delta}$. In both materials, the O_2 release was rapid initially and slowed down as the maximal value of $\Delta\delta$ was approached. In contrast to monometallic oxides that undergo a phase transition at constant pO_2 (such as the redox pair $\text{CuO}/\text{Cu}_2\text{O}$), the flow of O_2 released by the perovskite decreases with time because the $pO_{2,\text{eq}}$ depends on δ . This effect has consequences for the design of chemical looping applications, requiring e.g. semi-batch reactors or operation in counter-current flow if a steady stream of O_2 is to be obtained.

Thermodynamic assessment

A thermodynamic analysis was performed on $SrFe_{1-x}Co_xO_{3-\delta}$ and $Sr_{1-x}Ca_xFeO_{3-\delta}$ to assess the suitability of the investigated oxides for CLAS. Fig. S8 (ESI†) shows an overview of the thermogravimetric measurements (weight change as a function of time, temperature and pO_2) acquired for $SrFe_{1-x}Co_xO_{3-\delta}$, with x ranging from 0 to 0.8. It was observed that upon heating from ambient temperature to 800 °C the materials with a higher degree of substitution of Fe by Co (Fig. S8, ESI†) released a lower amount of O_2 than unsubstituted $SrFeO_{3-\delta}$. This suggests that the $pO_{2,\text{eq}}$ of the substituted samples decreased with increasing x . In addition, the amount of O_2 released was fully recovered (i.e. the material could be fully reoxidised) at each

temperature. Interestingly, the sample with the composition $SrFe_{0.2}Co_{0.8}O_{3-\delta}$ showed a very different behaviour compared to all the other samples in the temperature range 500–700 °C. The sample weight of this material did not change in the pO_2 range 9×10^{-4} – 1×10^{-1} bar, which implies that it is not required to employ an atmosphere of extremely low pO_2 to extract a significant amount of O_2 from the OC. For example, when switching the gas atmosphere from air ($pO_2 = 2.09 \times 10^{-1}$ bar) to N_2 ($pO_2 = 7.09 \times 10^{-5}$ bar), the OSC of $SrFe_{0.2}Co_{0.8}O_{3-\delta}$ was 1.21 wt% at 500 °C, while the OSC reduced only marginally to 1.1 wt% when operating between air and 5% O_2 ($pO_2 = 5 \times 10^{-2}$ bar). For comparison, the OSC of unsubstituted $SrFeO_{3-\delta}$ at 500 °C decreased from 0.8 wt% (when switching the gas atmosphere between air and N_2) to 0.1 wt% (when switching between air and 5% O_2). From a practical standpoint, $SrFe_{0.2}Co_{0.8}O_{3-\delta}$ represents a good candidate for CLAS due to its ability to release and take up O_2 at relatively low temperatures (i.e. < 600 °C) even in “oxygen-rich” environments (for example for pO_2 between 0.1 and 0.2 bar). In contrast, other compositions and the unsubstituted $SrFeO_{3-\delta}$ require very low – and practically difficult to realise – pO_2 to release a large quantity of O_2 , which would lower the economic attractiveness of CLAS.

The initial oxygen non-stoichiometry δ_0 of all samples was determined by iodometric titration at ambient temperature and is reported in Table S2 (ESI†). The substitution of the B-site cation Fe by Co (up to 60% substitution) yielded only a slight change in the initial δ_0 ($\delta_0 \approx 0.22$), while in the case of 80% substitution, $\delta_0 \approx 0.32$. On the other hand, substitution of A-site Sr by Ca strongly increased the available amount of O_2 (i.e. $\delta_0 \approx 0.15$) compared to the original $SrFeO_{3-\delta}$ ($\delta_0 \approx 0.22$). The increased amount of available redox-active lattice oxygen compared to the original $SrFeO_{3-\delta}$ when partially substituting the A- or B-site of the perovskite is reflected by the overall higher OSC of the substitute material, as shown in Fig. 2.

Fig. 5(a) shows exemplarily the evolution of the oxygen non-stoichiometry (symbols) of $SrFe_{1-x}Co_xO_{3-\delta}$ ($0 \leq x \leq 0.8$) as a function of pO_2 at 500 °C (the dotted lines represent qualitative guidelines). The initial value of δ of each sample at 500 °C was higher than that measured at ambient temperature due to the increase of $pO_{2,\text{eq}}$ with temperature; i.e. there was a release of O_2 upon heating the sample. Generally, all Co-substituted

Fig. 5 (a) Oxygen non-stoichiometry as a function of pO_2 for $SrFe_{1-x}Co_xO_{3-\delta}$ at 500 °C (b) Arrhenius plot of $SrFe_{0.8}Co_{0.2}O_{3-\delta}$.

samples showed a similar behaviour, except for the case of $x = 0.8$, for which the oxygen non-stoichiometry decreased sharply at relatively high oxygen partial pressures (*i.e.* $pO_2 = 1 \times 10^{-1}$ bar). Importantly, the value of $3 - \delta$ was found to be ~ 2.5 at low pO_2 (*i.e.* $pO_2 = 7.09 \times 10^{-5}$ bar, pure N_2) for all materials at 500 °C. This suggests that the perovskite structure was completely reduced to its defect-order variant brownmillerite ($ABO_{2.5}$). In the absence of a strongly reducing gas (*e.g.* H_2) that would thermodynamically enable a further reduction, the O_2 released from the perovskite structure corresponds to the maximal amount of lattice oxygen that can be released for this class of materials. Overall, the trend of the oxygen non-stoichiometry as a function of pO_2 recorded at 600 °C and 700 °C (Fig. S9(a) and (b), ESI†) was qualitatively similar to the behaviour at 500 °C. Increasing the temperature further to 800 °C (Fig. S9(c), ESI†), no significant difference in the evolution of $3 - \delta$ with pO_2 could be observed between each sample. Moreover, at 800 °C, the measured values of $3 - \delta$ were always greater than 2.5 (*i.e.* ~ 2.54) for the lowest oxygen partial pressure studied (*i.e.* $pO_2 = 7.09 \times 10^{-5}$ bar), suggesting that a phase change to brownmillerite did not occur at this temperature. *In situ* XRD measurements confirmed that at 800 °C the release of O_2 occurred without any phase transition (this is discussed in detail in the following section).⁴⁰

The respective thermogravimetric experiment for $Sr_{1-x}Ca_xFeO_{3-\delta}$, in which the A-site Sr was substituted by Ca, is reported in Fig. S10 (ESI†) while the measured $3 - \delta$ values as a function of pO_2 for $Sr_{1-x}Ca_xFeO_{3-\delta}$ ($x = 0, 0.1, 0.2, 0.25$) are reported in Fig. S11(a)–(d) (ESI†). Generally, the evolution of $3 - \delta$ with pO_2 for the A-site substituted materials was found to be similar to the B-site substituted materials. All the A-site substituted perovskites gave $3 - \delta = 2.5$ at the lowest pO_2 studied (*i.e.* 7.09×10^{-5} bar, pure N_2) at 500 °C, 600 °C and 700 °C, suggesting the transition to brownmillerite at this oxygen partial pressure. Increasing the temperature to 800 °C, a value of $3 - \delta = 2.5$ was observed only in the case of $Sr_{0.8}Ca_{0.2}FeO_{3-\delta}$ at the lowest pO_2 studied (*i.e.* 7.09×10^{-5} bar, pure N_2), while the measured values of $3 - \delta$ for the other samples were ~ 2.57 at this pO_2 , suggesting that a phase change to brownmillerite did not occur at this temperature.

To extract the partial molar enthalpy and entropy of vacancy formation from the TGA measurements of the substituted perovskites, the negative natural logarithm of pO_2 was plotted as a function of the inverse absolute temperature (Arrhenius plot) for constant $\Delta\delta$. Fig. 5(b)) exemplarily shows the Arrhenius plot for $SrCo_{0.8}Fe_{0.2}O_{3-\delta}$. The Arrhenius plot of all the other Co-substituted samples can be found in the ESI† (Fig. S12(a)–(d)). As we increase the temperature range compared to related works (albeit on different materials)^{35,41–43} we observe two distinct regimes, both of which can be approximated by a linear relationship. Similar observations were made also for Ca-substituted materials. For all samples investigated in this work, the transition temperature between the two regimes (linked to a structural change of the bulk perovskite phase) was ~ 770 °C. Similarly, for $SrCo_{0.8}Fe_{0.2}O_{3-\delta}$ Liu *et al.*⁴⁴ identified also 770 °C as the temperature above which the O_2 release occurs without a

Fig. 6 Evolution of the partial molar enthalpy (Δh°_O) and entropy (Δs°_O) of formation of vacancies as a function of the change in the oxygen non-stoichiometry $\Delta\delta$ in (a) and (c) the low temperature ($T < 770$ °C) and (b) and (d) the high temperature regime ($T > 770$ °C), respectively, for $SrFe_{1-x}Co_xO_{3-\delta}$. Gibbs free energy change of vacancy formation (Δg°_O) as a function of temperature in the (e) low temperature ($T < 770$ °C) and (f) high temperature regime ($T > 770$ °C), respectively. The Gibbs free energy change in (e) and (f) was plotted for $\Delta\delta = 0.025$ with the error bounds representing two standard deviations due to the errors associated with Δh°_O and Δs°_O .

phase transition from perovskite to brownmillerite, known as order-disorder transition.⁴⁰ From the Arrhenius plot, the thermodynamic properties of the OCs (*i.e.* the standard partial molar enthalpy and entropy of vacancy formation) were determined as a function of the change in the oxygen non-stoichiometry $\Delta\delta$. The partial molar enthalpy of formation as a function of $\Delta\delta$ for $SrFe_{1-x}Co_xO_{3-\delta}$ ($0 \leq x \leq 0.8$) is shown in Fig. 6(a) and (b), for the low ($T < 770$ °C) and high (770–900 °C) temperature range, respectively. Generally, the samples in which Fe was substituted by Co showed smaller partial molar enthalpy of formation compared to the unsubstituted $SrFeO_{3-\delta}$, when $T < 770$ °C. Additionally, in the same range (*i.e.* $T < 770$ °C), an increase in the oxygen non-stoichiometry yielded an increase in Δh°_O for $SrFe_{1-x}Co_xO_{3-\delta}$ with $x = 0, 0.2, 0.4, 0.6, 0.8$, consistent with what has been found in thermodynamic studies performed on perovskites such as $La_{1-x}(Ca,Sr)_xMn_{1-y}Al_yO_3$.⁴¹ The unsubstituted sample $SrFeO_{3-\delta}$ gave the highest partial molar enthalpy of formation of ~ 40 kJ mol^{−1}, when $\Delta\delta = 0.09$. The smaller values of the partial molar enthalpy of formation of the Co-substituted samples compared to the original $SrFeO_{3-\delta}$

explain their increased reducibility at low temperatures ($T < 770$ °C). The partial molar entropy of vacancy formation was generally constant (~ 7 J mol⁻¹ K⁻¹) with increasing $\Delta\delta$ when at $T < 770$ °C (Fig. 6(c)). In the high temperature range (770–900 °C), the partial molar enthalpy of vacancy formation decreased with increasing $\Delta\delta$ for all the Co-substituted samples, a result also reported for lanthanum manganites in the context of solar thermochemical water splitting.^{42,43} Overall, in the temperature range 770–900 °C, higher values of $\Delta h_{\text{O}}^{\circ}$ compared to the unsubstituted SrFeO_{3- δ} were found for each $\Delta\delta$, translating into a decreased reducibility of the substituted samples in this temperature range. In the same temperature range, an increase in the oxygen non-stoichiometry yielded an increase in the partial molar entropy of vacancy formation (Fig. 6(d)) showed generally increasing values with increasing $\Delta\delta$. Interestingly, the Co-substituted samples show higher $\Delta s_{\text{O}}^{\circ}$ compared to the unsubstituted sample.

The Arrhenius plots and the thermodynamic properties (partial molar entropy and enthalpy of formation of vacancies) obtained for Sr_{1-x}Ca_xFeO_{3- δ} are given in Fig. S13 and S14 (ESI[†]). In the low temperature range (*i.e.* $T < 770$ °C), the partial molar enthalpy of formation was relatively constant (~ 20 kJ mol⁻¹) with increasing $\Delta\delta$ for all substituted samples. Dou *et al.*³⁸ found similar values (*i.e.* ~ 25 kJ mol⁻¹) for Sr_{0.8}Ca_{0.2}FeO_{3- δ} . Similarly, to the case of substitution of Fe by Co, smaller values of $\Delta h_{\text{O}}^{\circ}$ compared to unsubstituted SrFeO_{3- δ} confirmed an increased reducibility induced by the substitution of Sr by Ca. In the temperature range 700–900 °C, higher values of $\Delta h_{\text{O}}^{\circ}$ compared to unsubstituted SrFeO_{3- δ} were found for each $\Delta\delta$. Table S3 in ESI[†] summarises the extrapolated partial molar enthalpy and entropy of vacancy formation for all samples.

Fig. 6(e) and (f) exemplary show the resulting partial molar Gibbs free energy change of vacancy formation $\Delta g_{\text{O}}^{\circ} = \Delta h_{\text{O}}^{\circ} - T\Delta s_{\text{O}}^{\circ}$ when $\Delta\delta = 0.025$. It was found that $\Delta g_{\text{O}}^{\circ}$ was always smaller for all substituted samples compared than the unsubstituted SrFeO_{3- δ} when $T < 770$ °C. Note that this is consistent with the increased reducibility of all substituted samples observed in this temperature range compared to the unsubstituted SrFeO_{3- δ} . In the high temperature range (*i.e.* 770–900 °C), $\Delta g_{\text{O}}^{\circ}$ of the substituted samples was always smaller than SrFeO_{3- δ} .

In situ XRD measurements

In situ XRD was performed for all the substituted perovskites (and the reference, unsubstituted perovskite) to probe the evolution of the crystal structure as a function of temperature. The measurements were performed isothermally at different temperatures in pure N₂ ($p_{\text{O}_2} = 7.09 \times 10^{-5}$ bar). Fig. 7(a) shows an exemplary result of the evolution of the structure of SrFe_{0.8}Co_{0.2}O_{3- δ} with scan number under a flow of pure N₂ ($p_{\text{O}_2} = 7.09 \times 10^{-5}$ bar) at 600 °C. With scan number (*i.e.* with the progress of O₂ release) the main peak of the perovskite phase (110) shifted towards lower diffraction angles (indicating an increase in the perovskite unit cell parameter) due to the

Fig. 7 *In situ* XRD showing the evolution of the structure of SrFe_{0.8}Co_{0.2}O_{3- δ} when exposed to N₂ (a) at $T = 600$ °C and (b) $T = 900$ °C. The scan time of one diffractogram over the 2θ range of $31\text{--}35^\circ$ was ~ 75 s.

release of O₂. When all the O₂ was released from the perovskite, a phase transition to brownmillerite occurred at this temperature (*i.e.* 600 °C). The same experiment, performed at 900 °C (Fig. 7(b)), shows that only the position of main perovskite peak (110) changed as O₂ was being released, but no separate reflections corresponding to brownmillerite were observed (in agreement with TG experiments). When reducing SrFe_{0.8}Co_{0.2}O_{3- δ} under N₂ at different temperatures, the *in situ* XRD measurements showed that for temperatures ≤ 750 °C the release of O₂ occurred through a phase transition from perovskite to brownmillerite, while for temperatures > 750 °C only the perovskite peak was observed, implying that no brownmillerite phase was formed.

Bulk phase transitions during reduction (and oxidation) are often an undesired characteristic in chemical looping applications (*i.e.* CLAS) as they may induce abrupt volume changes in the materials affecting negatively their mechanical stability. For large scale processes where, for example, fluidised beds may be

Fig. 8 *In situ* XRD patterns of Sr_{0.8}Ca_{0.2}FeO_{3- δ} when exposed consecutively to (i) pure N₂, (ii) $p_{\text{O}_2} = 7.9 \times 10^{-3}$ bar, (iii) $p_{\text{O}_2} = 8.6 \times 10^{-3}$ bar and (iv) $p_{\text{O}_2} = 2.09 \times 10^{-1}$ bar (a) at $T = 600$ °C and (b) $T = 900$ °C; (c) and (d) represent the collected diffractograms of figure (a) and (b), respectively. The scan time of one diffractogram over the 2θ range of $31\text{--}35^\circ$ was ~ 75 s.

used, a poor mechanical stability can result in attrition and loss of material. The change in the molar volume upon phase transition from the perovskite structure to brownmillerite is <1.5% and hence mechanical stresses within the particles may be less of a concern.

Turning back to $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ ($x = 0, 0.4, 0.6, 0.8$) and $\text{Sr}_{1-x}\text{Ca}_x\text{FeO}_{3-\delta}$ ($x = 0.1, 0.25$), Fig. S14 (EIS), the formation of a brownmillerite phase when releasing O_2 (in N_2) was observed for all samples at $T < 750^\circ\text{C}$, while the perovskite structure remained stable for $T > 750^\circ\text{C}$. Only in the case of 20% substitution of Sr by Ca a phase transition from perovskite to brownmillerite was observed also at $T > 750^\circ\text{C}$ (Fig. 8(a) and (b)), consistent with the measured $3 - \delta = 2.5$ for this sample (Fig. S10(d), ESI[†]). For both substituted and unsubstituted perovskites, the regeneration of the perovskite structure was possible only in air ($p\text{O}_2 = 2.09 \times 10^{-1}$ bar).

4. Conclusions

We have studied the influence of Ca A- or Co B-site substitution of SrFeO_3 on the materials' chemical-physical properties at temperatures ranging from 400 to 900 $^\circ\text{C}$, relevant for chemical looping applications. We have shown that phase pure materials can be obtained when limiting the substitution of A- or B-site cations to the range of $0 \leq x \leq 0.8$ or $0 \leq x \leq 0.25$ for Co B-site or Ca A-site substitution, respectively. Thermogravimetric analysis revealed that the chemical-physical properties (quantity of oxygen released, desorption/sorption rate and reducibility) of the materials are altered by the partial substitution of their A- or B-site cations. In particular, it was found that the substitution of Fe by Co increased the reducibility of the metal oxide in the low temperature regime (*i.e.* $T = 400\text{--}600^\circ\text{C}$) when compared to the unsubstituted $\text{SrFeO}_{3-\delta}$. Our quantitative thermodynamic assessment confirmed a decreased partial molar enthalpy of formation of $\text{SrFe}_{1-x}\text{Co}_x\text{O}_{3-\delta}$ in this temperature range. Among the Co-substituted samples, $\text{SrFe}_{0.2}\text{Co}_{0.8}\text{FeO}_{3-\delta}$ showed a high OSC ~ 1.1 wt% when operating between air and 5% O_2 at 500 $^\circ\text{C}$. The ability to effectively release and take up O_2 at relatively low temperatures (*i.e.* $< 600^\circ\text{C}$) without the need for gas atmospheres of extremely low $p\text{O}_2$ makes this material a good candidate for CLAS.

Sr A-site substitution by Ca improved the OSC and among all the materials investigated, $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ exhibited the highest OSC (2.15 wt%) at 500 $^\circ\text{C}$ when operating between air and N_2 . The redox cyclic stability of $\text{SrFe}_{0.2}\text{Co}_{0.8}\text{FeO}_{3-\delta}$ and $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ was investigated over up to 100 redox cycles in a fixed bed reactor and no apparent alteration in the crystalline structure nor loss in activity was observed.

When probing the evolution of the oxygen non-stoichiometry as a function of temperature and $p\text{O}_2$, two regimes were observed for the substituted perovskites. The transition temperature between the two regimes is $\sim 770^\circ\text{C}$. For $T < 770^\circ\text{C}$ the O_2 release in N_2 ($p\text{O}_2 = 7.09 \times 10^{-5}$ bar) was accompanied by a phase transition from a perovskite to a brownmillerite phase, while the perovskite structure remained stable for

$T > 770^\circ\text{C}$, as confirmed by *in situ* XRD measurements. The only exception was $\text{Sr}_{0.8}\text{Ca}_{0.2}\text{FeO}_{3-\delta}$ that showed a phase transition during the reduction in N_2 over the entire temperature range studied (500–900 $^\circ\text{C}$).

Lastly, our results show that in spite of the high OSC of the perovskite materials studied, the challenge to obtain a constant flow of O_2 released remains as the O_2 release varies with the oxygen non-stoichiometry δ . This intrinsic feature of non-stoichiometric materials would require engineering solutions (*e.g.* moving bed reactor design in a counter-current flow configuration) to control the variation in the oxygen non-stoichiometry and, therefore, ensure a constant flow of O_2 at the outlet of the reactor if demanded by the chemical looping process.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- 1 M. Bui, C. S. Adjiman, A. Bardow, E. J. Anthony, A. Boston, S. Brown, P. S. Fennell, S. Fuss, A. Galindo, L. A. Hackett, J. P. Hallett, H. J. Herzog, G. Jackson, J. Kemper, S. Krevor, G. C. Maitland, M. Matuszewski, I. S. Metcalfe, C. Petit, G. Puxty, J. Reimer, D. M. Reiner, E. S. Rubin, S. A. Scott, N. Shah, B. Smit, J. P. M. Trusler, P. Webley, J. Wilcox and N. Mac Dowell, *Energy Environ. Sci.*, 2018, **11**(5), 1062–1176.
- 2 S. Fuss, W. F. Lamb, M. W. Callaghan, J. Hilaire, F. Creutzig, T. Amann, T. Beringer, W. De Oliveira Garcia, J. Hartmann, T. Khanna, G. Luderer, G. F. Nemet, J. Rogelj, P. Smith, J. L. Vicente, J. Wilcox, M. Del Mar Zamora Dominguez and J. C. Minx, *Environ. Res. Lett.*, 2018, **13**(6), 063002.
- 3 J. C. Minx, W. F. Lamb, M. W. Callaghan, S. Fuss, J. Hilaire, F. Creutzig, T. Amann, T. Beringer, W. De Oliveira Garcia, J. Hartmann, T. Khanna, D. Lenzi, G. Luderer, G. F. Nemet, J. Rogelj, P. Smith, J. L. Vicente, J. Wilcox and M. Del Mar Zamora Dominguez, *Environ. Res. Lett.*, 2018, **13**(6), 063001.
- 4 B. Mckee, Solutions for the 21st century, zero emissions technologies for fossil fuels, Technology status report, Energy, DOI:10.1016/B978-0-12-802748-6.01001-9.
- 5 H. I. Mathekg, B. O. Oboirien and B. C. North, *Int. J. Energy Res.*, 2016, **40**(7), 878–902.

- 6 C. Fu and T. Gundersen, *Energy*, 2012, **44**(1), 60–68.
- 7 F. G. Kerry, *Industrial Gas Handbook: Gas Separation and Purification*, CRC Press, Boca Raton, 2007.
- 8 H.-W. Häring, *Ind. Gases Process.*, Wiley-VCH Verlag GmbH & Co. KGaA, 2007, pp. 9–109.
- 9 W. F. Castle, *Int. J. Refrig.*, 2002, **25**, 158–172.
- 10 B. Moghtaderi, *Energy Fuels*, 2010, **24**, 190–198.
- 11 K. Shah, B. Moghtaderi, J. Zanganeh and T. Wall, *Fuel*, 2013, **107**, 356–370.
- 12 P. N. Dyer, R. E. Richards, S. L. Russek and D. M. Taylor, *Solid State Ionics*, 2000, **134**(1–2), 21–33.
- 13 I. Pfaff and A. Kather, *Energy Procedia*, 2009, **1**(1), 495–502.
- 14 S. Haider, A. Lindbräthen, J. A. Lie and M. B. Hägg, *Sep. Purif. Technol.*, 2018, **205**, 251–262.
- 15 W. B. Jensen, *J. Chem. Educ.*, 2009, **86**, 1266–1267.
- 16 H. Song, K. Shah, E. Doroodchi and B. Moghtaderi, *Energy Fuels*, 2014, **28**, 163–172.
- 17 C. Zhou, K. Shah and B. Moghtaderi, *Energy Fuels*, 2015, **29**(4), 2074–2088.
- 18 K. Shah, B. Moghtaderi and T. Wall, *Energy Fuels*, 2012, **26**(4), 2038–2045.
- 19 R. H. Görke, W. Hu, M. T. Dunstan, J. S. Dennis and S. A. Scott, *Appl. Energy*, 2018, **212**, 478–488.
- 20 K. Wang, Q. Yu and Q. Qin, *J. Therm. Anal. Calorim.*, 2013, **112**, 747–753.
- 21 K. Wang, Q. Yu, Q. Qin and W. Duan, *Chem. Eng. Technol.*, 2014, **37**, 1500–1506.
- 22 G. Azimi, H. Leion, T. Mattisson and A. Lyngfelt Proc. 10th Int. Conf. Greenhouse Gas Technology (GHGT-10), Amsterdam, The Netherlands, 2010.
- 23 K. J. Whitty, J. S. Lighty and T. Mattisson, *Handb. Chem. Looping Technol.*, 2018, 93–122.
- 24 F. Donat, W. Hu, S. A. Scott and J. S. Dennis, *Ind. Eng. Chem. Res.*, 2015, **54**(26), 6713–6723.
- 25 X. Zhu, K. Li, L. Neal and F. Li, *ACS Catal.*, 2018, **8**, 8213–8236.
- 26 B. Bulfin, J. Vieten, D. E. Starr, A. Azarpira, C. Zachäus, M. Hävecker, K. Skorupska, M. Schmücker, M. Roeb and C. Sattler, *J. Mater. Chem. A*, 2017, **5**, 7912–7919.
- 27 J. Vieten, B. Bulfin, P. Huck, M. Horton, D. Guban, L. Zhu, Y. Lu, K. A. Persson, M. Roeb and C. Sattler, *Energy Environ. Sci.*, 2019, **12**, 1369–1384.
- 28 N. Miura, H. Ikeda and A. Tsuchida, *Ind. Eng. Chem. Res.*, 2016, **55**, 3091–3096.
- 29 H. Ikeda, A. Tsuchida, J. Morita and N. Miura, *Ind. Eng. Chem. Res.*, 2016, **55**, 6501–6505.
- 30 M. Kubicek, A. H. Bork and J. L. M. Rupp, *J. Mater. Chem. A*, 2017, **5**(24), 11983–12000.
- 31 X. Zhu, Q. Imtiaz, F. Donat, C. R. Müller and F. Li, *Energy Environ. Sci.*, 2020, **13**, 772–804.
- 32 T. Motohashi, T. Ueda, Y. Masubuchi and S. Kikkawa, *J. Phys. Chem. C*, 2013, **117**, 12560–12566.
- 33 A. Mishra, T. Li, F. Li and E. E. Santiso, *Chem. Mater.*, 2019, **31**, 689–698.
- 34 E. I. Leonidova, I. A. Leonidov, M. V. Patrakeev and V. L. Kozhevnikov, *J. Solid State Electrochem.*, 2011, **15**(5), 1071–1075.
- 35 E. Bakken, T. Norby and S. Stølen, *Solid State Ionics*, 2015, **176**(1–2), 217–223.
- 36 C. Y. Lau, M. T. Dunstan, W. Hu, C. P. Grey and S. A. Scott, *Energy Environ. Sci.*, 2017, **10**(3), 818–831.
- 37 R. H. Görke, E. J. Marek, F. Donat and S. A. Scott, *Int. J. Greenhouse Gas Control*, 2020, **94**, 102891.
- 38 J. Dou, E. Krzystowczyk, X. Wang, T. Robbins, L. Ma, X. Liu and F. Li, *ChemSusChem*, 2020, **13**, 385–393.
- 39 E. J. Popczun, D. N. Tafen, S. Natesakhawat, C. M. Marin, T.-D. Nguyen-Phan, Y. Zhou, D. Alfonso and J. W. Lekse, *J. Mater. Chem. A*, 2020, **8**(5), 2602–2612.
- 40 Y. Tomura, I. Oikawa and H. Takamura, *Phys. Rev. Mater.*, 2019, **3**, 125802.
- 41 M. Ezbiri, M. Takacs, D. Theiler, R. Michalsky and A. Steinfeld, *J. Mater. Chem. A*, 2017, **5**, 4172–4182.
- 42 A. H. Bork, E. Povoden-Karadeniz, A. J. Carrillo and J. L. M. Rupp, *Acta Mater.*, 2019, **178**, 163–172.
- 43 M. Takacs, M. Hoes, M. Caduff, T. Cooper, J. R. Scheffe and A. Steinfeld, *Acta Mater.*, 2016, **103**, 700–710.
- 44 L. M. Liu, T. H. Lee, L. Qiu, Y. L. Yang and A. J. Jacobson, *SrCo_{0.8}Fe_{0.2}O_{3-δ}*, *Mater. Res. Bull.*, 1996, **31**, 29–35.