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Isothermal relaxation kinetics for the reduction and oxidation of SrFeO₃ based perovskites†

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The perovskite oxide SrFeO₃ has favourable redox properties for oxygen exchange applications, including oxygen separation and oxygen production chemical looping cycles. For such applications, lower temperature operation can improve the energy demand and feasibility of the process, but can also lead to kinetic limitations. Here we investigate the oxidation and reduction reaction kinetics of SrFeO₃ in the temperature range 450–750 K. Isothermal relaxation techniques are used to observe the reaction rates across this temperature range, using a thermogravimetric analysis system. Experimental data are analysed according to an isoconversional method and fit with a simple power law model to extract activation energies. The apparent activation energy of oxidation and reduction was found to be 92 ± 16 and 144 ± 17 kJ mol⁻¹ respectively. Comparison of oxidation and reduction kinetics together with considerations of particle size indicate that the oxidation reaction rate may be limited by diffusion in the bulk, while the reduction reaction rate is limited by the surface reaction. Furthermore, we also investigated the mixed perovskite Sr_{0.93}Ca_{0.07}Fe_{0.9}Co_{0.1}O₃, which exhibited a 4-fold increase in the oxidation rate.

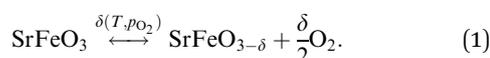
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1 Introduction

The perovskite oxide SrFeO₃ exhibits an oxygen non-stoichiometry which depends on the temperature and partial pressure of oxygen,^{1–3}



The oxygen non-stoichiometry δ , decreases with temperature,³ but has been found to persist at relatively low temperatures, $T < 673$ K.⁴ For this reason SrFeO₃ based perovskites have been extensively investigated as oxygen storage materials, which can be utilized in oxygen separation and production processes.^{5–8} These processes operate *via* a redox cycle, where oxygen is absorbed in an oxidation step and released again in a reduction step. They can use a temperature swing,⁹ and/or a pressure swing,^{10–12} to switch between oxidation and reduction of the oxide. In the literature these oxygen separation processes are also often referred to as thermochemical redox cycles,^{13,14} high

temperature pressure swing adsorption processes,¹⁰ or chemical looping air separation.¹⁵

For these oxygen separation applications, operating the redox cycle at lower temperatures can be very beneficial. Lower temperature implies less constraints on the reactor design and heat exchanger materials, a longer lifetime for the redox material, a potentially much lower energy demand, and therefore higher efficiency. However, lower temperatures will also hinder the kinetic activity of the redox reaction. Isothermal relaxation experiments on reduced SrFeO₃ powder have shown rapid oxidation kinetics for temperatures above 673 K,¹⁶ where thermodynamic equilibrium was reached in approximately 1–2 min at 723 K.^{6,16} However, at 523 K 300 hours was insufficient time for a dense disk of the reduced perovskite SrFeO_{2.5} to reach equilibrium.⁴ These results suggest that oxidation reaction kinetics can become a critical limitation for temperatures below 723 K.

Previous kinetic studies have mostly focused on intrinsic oxygen transport properties at temperatures above 723 K, with a targeted application of SOFC cathodes.¹⁷ Electrical conductivity relaxation techniques were used to determine surface exchange and diffusion coefficients for SrFeO₃ at temperatures in the range 973–1273 K.^{18,19} Relaxation kinetics of SrFeO₃ have also been studied using thermogravimetric analysis,²⁰ and oxygen uptake in a lab scale fixed bed reactor,²¹ but with both studies focused on temperatures above 723 K.

In this work we focus on a lower temperature range, 450–750 K, and use thermogravimetric analysis (TGA) with changes in

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temperature and abrupt changes in partial pressure to achieve the non-equilibrium conditions required for relaxation kinetics. We consider both the oxidation and reduction reactions of SrFeO₃, and a mixed variant Sr_{0.93}Ca_{0.07}Fe_{0.9}Co_{0.1}O₃ (SCFCO) from a previous study.²² This mixed perovskite has cobalt added in place of iron, to modify the redox behaviour, and calcium in place of strontium to mitigate the effect of the lower ionic radius of Co⁴⁺ with respect to Fe⁴⁺, and maintain the same Goldschmidt tolerance factor as SrFeO₃.²² This mixed variant of SrFeO₃ shows that the kinetics can be improved by the substitution of cobalt and calcium into the lattice.

2 Experimental

2.1 Sample preparations and characterization

SrFeO₃ powder samples used in the thermogravimetric analysis (TGA) were prepared *via* a solid state route. Fe₃O₄ and SrCO₃ were stoichiometrically mixed and annealed at 1373 K for 20 hours in an alumina crucible, after which the SrFeO₃ is formed in large sintered granules. These granules were then pestled into a powder. In order to allow for a larger mass to be placed in the TGA, the SrFeO₃ powder was compressed into an 8 mm diameter pellet, with a load of 4 metric tons, and then sintered at 1573 K for 20 hours, to improve structural stability. The resulting pellet weighed close to 500 mg, and mass and volume measurements indicated it had approximately 10–15% void space. Following the oxidation kinetics analysis, the pellet was used to study the reduction kinetics, but the first attempts at measuring the reduction, resulted in the pellet fracturing during the cycle. To avoid the possibility that a change in morphology might alter the kinetics, the pellet was ground back into a powder, and this was analysed instead. It was also this powder that the characterization data in Table 1 for the pellet is given.

The SCFCO sample was synthesized *via* the Pechini method,²² with detailed descriptions of this route given in earlier publications.^{23,24} After auto-combustion, the resulting fine oxide powder was ground and filled into alumina crucibles, which were then heated in a high temperature box furnace (Carbolite RHF 14/35) in order to produce the perovskite phase. The samples were treated twice at 800 °C for 10 hours followed by a treatment at 1300 °C for 20 hours. All samples were heated in air, and no intermediate powdering or mixing was carried out.

The perovskite phase formation of the samples were verified by powder X-ray diffraction (PXRD) using a STOE STADI P diffractometer equipped with a Dectris MYTHEN 1 K detector in transmission mode using a Ge monochromator for Cu-K_α radiation. The powders were scanned from 10–90°, with a step size of 0.01°, with a Rietveld refinement performed using FullProf suite.²⁵ The results are summarized in Table 1 (see ESI† for raw data and refinement plots), where it can be seen that the pellet showed a pure cubic perovskite structure, with a lattice parameter *a* which was slightly larger than that of SrFeO_{2.96}.²⁶ This expansion of the lattice can be explained by the fact that we do not expect a larger oxygen non-stoichiometry in our sample, which would expand the lattice due to more Fe³⁺ which has a larger ionic radius compared to Fe⁴⁺. The SCFCO sample was also primarily a cubic perovskite structure, but with two secondary phases, a brownmillerite and a Ruddlesden–Popper phase. The SrFeO₃ annealed at 1100 °C was split between two tetragonal phases one of which was the Ruddlesden–Popper structure.

The particle size distribution of the powders was measured using a Horiba LA-950 laser scattering device. The powders were suspended in deionized water with a small amount of surfactant used to stop particles from getting trapped on the surface of the liquid. A scanning electron microscope was used to get qualitative visual confirmation of the particle size distribution measurements. The data is shown in Table 1, where it can be seen that the samples sintered at higher temperature have larger particles. The pellet shows the largest particles and broadest distribution. This can also be qualitatively seen in the SEM images in the ESI.† It should be noted here that in order to measure the pellets particle size distribution the pellet was first crushed using a mortar and pestle which may have reduced the average particle size.

2.2 Thermogravimetric analysis

For TGA experiments, a thermobalance system by Netzsch (STA 409 CD) was used. This system was used to investigate both equilibrium non-stoichiometry and isothermal relaxation kinetics. Both powder and pellet samples were placed on a platinum foil wrapped alumina sample holder, and mounted on an alumina pin. The oxygen partial pressure was controlled by using different mixtures of pure oxygen, argon, and synthetic air. The changes in sample mass Δm , were then recorded while

Table 1 Sample characterization data from XRD analysis and particle size distribution measurements. The main perovskite phase determined from XRD is given with the percentage mass in this phase (and side phases), and the lattice parameters. From the size distribution measurements we have the mean particle diameter by mass of material \bar{d} , the standard deviation in particle diameter σ_d , and $d_{10} - d_{90}$ which is the diameters bounding the lower 10% and upper 90% of the cumulative mass/volume

Sample	<i>T</i> [K]	XRD main structure	Wt% perovskite	Lattice parameters [Å]	\bar{d} [μm]	σ_d [μm]	$d_{10} - d_{90}$ [μm]
SFO powder	1373	<i>I4mmm</i> – tetragonal	50% ^a	<i>a</i> = 10.8970(7), <i>c</i> = 7.7170(5)	5.3	1.82	3.2–7.7
SFO pellet	1573	<i>Pm3m</i> – cubic	100%	<i>a</i> = 3.8695(1)	10.85	9.22	4.29–17.2
SCFCO	1573	<i>Pm3m</i> – cubic	80% ^b	<i>a</i> = 3.8543(1)	7.8	2.79	4.5–11.5

^a This SrFeO₃ powder sample showed two distinct perovskite structures, with approximately 50% in each phase, with both phases known to exhibit oxygen non-stoichiometry.^{27,28} ^b The SCFCO sample showed two additional secondary phases, a brownmillerite and a Ruddlesden–Popper structure both of which also exhibit oxygen non-stoichiometry.^{29,30}



the temperature and partial pressures were varied. Since SrFeO₃ displays non-stoichiometry even at ambient conditions, the changes in mass were always measured relative to an equilibrium reference point with a fixed temperature and oxygen partial pressure.

The relative change in non-stoichiometry $\Delta\delta$ in SrFeO_{3- δ} , can then be calculated from the mass changes,

$$\Delta\delta = \frac{\Delta m M_{\text{sample}}}{m M_{\text{O}}}, \quad (2)$$

where m is the sample mass, M_{O} is the molar mass of molecular oxygen and M_{sample} is the molar mass of the sample at the reference point. The reference point was $T = 803$ K and $p_{\text{O}_2} = 1$ bar, with two previous studies on the SrFeO_{3- δ} perovskite system indicating that this reference point corresponds to SrFeO_{2.77},^{4,6} giving $M_{\text{sample}} = 187.78$ g mol⁻¹. For SCFCO, the oxygen non-stoichiometry at the reference point was assumed to be the same as SrFeO₃. In order to extract equilibrium values of the non-stoichiometry, the partial pressure and temperature were varied step-wise (see Fig. S3 in the ESI†) and the data in plateau regions where Δm , T and p_{O_2} were approximately constant with respect to time were taken to be in equilibrium.

In order to measure isothermal relaxation kinetics the samples must be brought to a state where the non-stoichiometry is not at the equilibrium value $\delta \neq \delta_{\text{eq}}$. For oxidation kinetics, the samples were initially reduced under Ar 5.0 ($\approx 10^{-4}$ – 10^{-5} O₂) flow of 100 ml min⁻¹ at 800 °C. The samples were then cooled under the same Ar flow to the desired temperature, where re-oxidation is limited due to the lack of available oxygen. When the temperature has stabilised at the desired set point the gas flow is then switched to pure oxygen. The perovskite then begins to oxidise towards equilibrium and the resulting increase in mass was measured against time, giving an isothermal relaxation profile. Finally, the sample was brought to 400 °C in synthetic air after each kinetic run, as a reference point. Further experiments were carried out where the oxidising gas was a mixture with an oxygen percentage in the range 1–100% O₂, in order to investigate the effect of the oxygen partial pressure on the kinetics. For the SCFCO perovskite the reduction step in argon was performed at 600 °C instead of 800 °C, as this material was seen to undergo a phase transition at higher temperatures in a previous study.²²

For the reduction reaction the non-equilibrium state required for relaxation kinetics was achieved only by switching the gases. The samples were brought to the desired temperature set point in a pure oxygen flow of 200 ml min⁻¹, giving $p_{\text{O}_2} \approx 1$ bar. When the temperature had stabilised the gas flow was then switched to 190 ml min⁻¹ Ar and 10 ml min⁻¹ synthetic air (20% O₂) giving $p_{\text{O}_2} \approx 0.01$ bar. The more reducing atmosphere will cause the sample to lose oxygen, proceeding towards equilibrium with the resulting decrease in mass measured against time. Finally, the samples were brought to 450 °C in pure oxygen after each kinetic run, as the reference point.

A potentially limiting aspect of these isothermal relaxation measurements is the gas switching time, which can be affected by the mass flow controllers and diffusion along the gas flow. An analysis of this is given in the ESI,† which was validated

using a mass-spectrometer downstream of the system. The time required to go from an initial change in p_{O_2} , to greater than 90% switched (*i.e.* 0.9 bar O₂ for oxidation and 0.1 bar O₂ for reduction), was on the order of 20 seconds for oxidation and 7 seconds for the reduction reaction. Mass changes measured in this region are included in plots and in determining the conversion extent, but were not used to analyse the kinetic data.

Examples of the raw TGA data for equilibrium analysis, reduction and oxidation isothermal kinetics are all available in the ESI.†

3 Results

3.1 Equilibrium measurements

In previous studies equilibrium measurements have typically focused on high temperatures 773–1473 K.^{1,3} There is some limited data at lower temperatures,^{4,6} but in the present study we aim to have a broader set of equilibrium data at lower temperatures, so that we know more accurately what the final equilibrium non-stoichiometry should be when the reactions are complete (δ_{∞} in eqn (3)).

SrFeO₃ is non-stoichiometric in oxygen, where the concentration of oxygen vacancies depends on the temperature and partial pressure (see eqn (1)).^{31,32} This oxide can exhibit a number of intermediate phases, which have been extensively reported on in the literature.^{4,33–35} Here we focus on the determination of the equilibrium non-stoichiometry at a range of temperature and partial pressures below 900 K.

The relative change in non-stoichiometry $\Delta\delta$ at equilibrium, was determined as described in the experimental details, and is plotted in Fig. 1. The data shown for SrFeO₃ was obtained using the ground pellet which was annealed at 1573 K, as this showed the purest phase in the XRD analysis. The data is given in terms of the relative non-stoichiometry $\Delta\delta$ so that SrFeO₃ and SCFCO can be directly compared, where $\Delta\delta = 0$ corresponds to $T = 803$ K and $p_{\text{O}_2} = 1$ bar. They have extremely similar changes in non-stoichiometry with temperature and oxygen partial pressure. This is not particularly surprising as they have the same phase and the species expected to change oxidation state

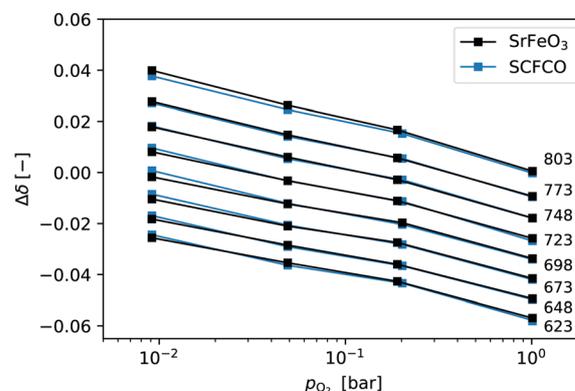


Fig. 1 Relative non-stoichiometry $\Delta\delta$ plotted against the oxygen partial pressure for temperatures in the range 623–803 K. Data at a given temperature are connected by lines and labeled on the right.



in the oxidation and reduction reactions is predominately Fe^{4+} - Fe^{3+} in both cases.

3.2 Oxidation kinetics

The oxidation kinetics were measured at isothermal conditions in a pure oxygen atmosphere in the temperature range 450–750 K. The mass changes as a function of time were used to calculate $\Delta\delta(t)$, using eqn (2). Some examples of the raw TGA data are given in the ESI.† In order to analyse such data it is convenient to define an extent of conversion X , given by,

$$X(t) = \frac{\Delta\delta_0 - \Delta\delta(t)}{\Delta\delta_0 - \Delta\delta_\infty}, \quad (3)$$

where $\Delta\delta_0$ is the non-stoichiometry at $t = 0$ and $\Delta\delta_\infty$ is the equilibrium value which the reaction is converging upon. We then have $X \in (0,1)$, where 0 corresponds to no reaction having taken place and 1 means we have reached chemical equilibrium. For oxidation reactions at temperatures $T < 600$ K, the samples do not reach equilibrium, and so we need to know $\Delta\delta_\infty$ at these temperatures. As there is no equilibrium data available in this region we simply use a linear fit of the data in Fig. 1 with $p_{\text{O}_2} = 1$ bar (see ESI,† for graph showing the linear extrapolation). The data certainly has a linear dependence in the temperature range measured, and the extrapolation should be relatively accurate provided we do not approach the fully oxidised state, SrFeO_3 with $\delta = 0$. This is indeed the case here where the smallest absolute δ value used from the extrapolation was $\delta = 0.12$, so $\text{SrFeO}_{2.88}$. We used the same linear extrapolation method for SCFCO.

A quick check of the effect of temperature on the oxidation reaction can be obtained by looking at the half life of the reactions $t_{1/2}$, defined by,

$$X(t_{1/2}) = 0.5. \quad (4)$$

Fig. 2 shows the variation of the half life with temperature, where it can be seen that for all materials the reactions proceed faster at higher temperature as expected. For SrFeO_3 , 100 and 200 mg of the powder, and a 500 mg pellet were investigated,

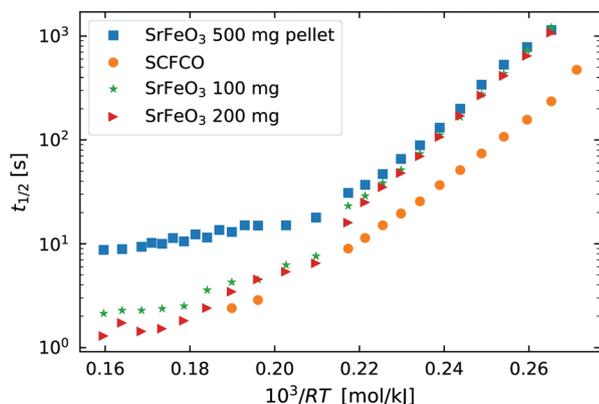


Fig. 2 The half life $t_{1/2}$ of the oxidation reaction $t_{1/2}$ vs. $\frac{10^3}{RT}$ for re-oxidation of the SrFeO_3 pellet, 200 mg of the SCFCO powder, and the SrFeO_3 powder sintered at 1373 K (100 and 200 mg).

where at temperatures below 550 K ($10^3/RT > 0.22$), they all have similar half lives. This is a good indication that in this region the reaction is controlled by the materials intrinsic kinetic properties and not by external effects such as mass transfer. Indeed the SrFeO_3 powder annealed at 1373 K does show slightly shorter half lives in this region, which we would expect due to the smaller particle size.

Above 550 K ($10^3/RT < 0.22$) the trends in $t_{1/2}$ change, particularly for the pellet where the half life no longer has a strong dependence on temperature. In this region the half life is on the order of 10 seconds for the 500 mg pellet. The total mass change to reach the $X = 0.5$ is on the order of 5 mg or 1.56×10^{-4} mol of O_2 , and the oxygen supply rate of 100 ml min^{-1} corresponds to $7.34 \times 10^{-5} \text{ mol s}^{-1}$ of O_2 . So the minimum $t_{1/2}$ by mass transfer limitations is just over 2 seconds. Given that the gas switching in the TGA requires approximately 20 seconds (see gas switching section in the ESI†), we conclude that the kinetics above 550 K are mass transfer limited.

Finally, Fig. 2 also shows the half life for the oxidation of SCFCO powder, which can be seen to have considerably faster kinetics than SrFeO_3 . For example at 473 K, SCFCO has a half life of just 107 s, while the SrFeO_3 powder had a half life of 426 s. This factor of four increase in the rate of reaction can not be explained by the physical morphology of samples, as the SCFCO sample had a similar particle size distribution as the SrFeO_3 pellet (see Table 1). We can therefore assume, that the mixed perovskite has improved intrinsic oxygen transport kinetics, which could be better oxygen lattice diffusion or surface exchange properties.

3.2.1 Kinetic analysis. In summary Fig. 2 tells us that the SrFeO_3 pellet is mass transfer limited at temperatures above 550 K, and that SCFCO has improved oxidation kinetics over SrFeO_3 . For further analysis we consider the SrFeO_3 pellet, as this sample showed the highest purity (see the ESI,† Fig. S3) and the best signal to noise ratio in the TGA analysis. For comparison we also apply the same analysis to SCFCO.

The mechanism by which the oxidation takes place can be broken into a number of steps, (1) oxygen gas adsorption and dissociation, (2) charge transfer and incorporation of the adsorbed oxygen into the oxide, and (3) diffusion into the bulk.^{36–38} In the methods used here we do not have the possibility to probe the individual reaction steps. For simplicity it is therefore easier to consider the adsorption, dissociation and incorporation of oxygen into the bulk, as the surface reaction. We then have a surface reaction followed by a diffusion process, which for spherical particles can be analytically modelled.³⁹ However, here we have a distribution of particles, which rules out the use of such a simple analytical approach. Instead, we first apply the isoconversional method from the ICTAC Kinetics Committee recommendations for performing kinetic computations on thermal analysis data.⁴⁰

To implement this approach, we assume that the rate of reaction can be parametrized in terms of the temperature T , conversion extent X , and the oxygen gas concentration c_{O_2} ,

$$\frac{dX}{dt} = k(T)f(X)c_{\text{O}_2}^n \quad [\text{s}^{-1}] \quad (5)$$



where $k(T)$ is the rate constant, $f(X)$ is a function of the conversion extent, and n is the order of the oxygen concentration dependence. We further assume that the rate constant takes an Arrhenius form,

$$k(T) = k_0 \exp\left(\frac{-E_a}{RT}\right) [\text{s}^{-1} (\text{mol m}^{-3})^{-n}]. \quad (6)$$

where E_a is the activation energy and k_0 is the pre exponential factor. Note that in order to have the correct units on the right hand side of eqn (5), the pre exponential factor k_0 is given units of $[\text{s}^{-1} (\text{mol m}^{-3})^{-n}]$. The activation energy in this case is often referred to as an apparent activation energy, as it shows the effect of temperature on the overall kinetic process, but not the intermediate steps involved in the reaction. This can correspond to a physical activation barrier for the rate limiting step, but it is not necessarily the case.

It is worth noting here that the ICTAC Kinetics Committee recommendations suggest a generic function of partial pressure $h(p_{\text{O}_2})$ instead of concentration.⁴⁰ However, in the classical literature on this topic concentration is preferred.⁴¹ This is an important distinction as the concentration for a fixed partial pressure depends on the temperature, where according to the ideal gas law we have,

$$c_{\text{O}_2} = \frac{n}{V} = \frac{p_{\text{O}_2}}{RT} [\text{mol m}^{-3}]. \quad (7)$$

In order to determine the value of n , SrFeO₃ relaxation kinetics were measured at oxygen partial pressures pressures 1, 0.2, 0.05 and 0.01 bar, and at temperatures of 523 K and 573 K. This data range was selected as at lower temperatures and lower partial pressures the reaction proceeds too slowly to accurately measure the rate, and at higher temperatures the pure oxygen case was mass transfer limited. To analyse the data we note that the initial slope of the mass change is proportional to the rate of change of the conversion extent $\left(\frac{dm}{dt}\right)_{t=0} \propto \left(\frac{dX}{dt}\right)_{t=0}$. Therefore a log-log plot of $\left(\frac{dm}{dt}\right)_{t=0}$ vs. c_{O_2} should yield n as the slope. Fig. 3 shows such a plot, where the slope of the lines were found to be $n = 0.68$ and 0.71 . We therefore took the value of n to be 0.695 in the rest of the analysis. We use the same value for all samples as the oxygen concentration dependence should not change much between samples. A sensitivity analysis also showed that varying n between 0.4 and 1 had very little effect ($<2\%$ change) on the apparent activation energies.

With the value of n we can now apply the isoconversional method. For a fixed value of X , eqn (5) can be written as,

$$c_{\text{O}_2}^{-n} \left(\frac{dX}{dt}\right)_{X=\text{const.}} = k(T)f(X)_{X=\text{const.}} \quad (8)$$

and taking the logarithm of both sides and using eqn (6) gives

$$\log\left(c_{\text{O}_2}^{-n} \left(\frac{dX}{dt}\right)_{X=\text{const.}}\right) = \frac{-E_a}{RT} + \log(k_0 f(X)_{X=\text{const.}}) \quad (9)$$

We can see that a log plot of $c_{\text{O}_2}^{-n} \left(\frac{dX}{dt}\right)_{X=\text{const.}}$ vs. $\frac{10^3}{RT}$ will have the activation energy as the slope and $\log(k_0 f(X))$ as the intercept.

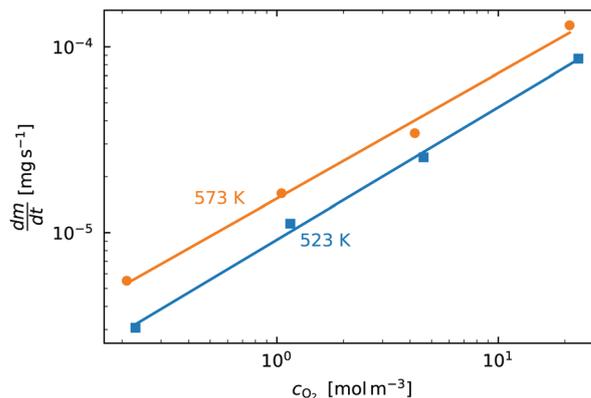


Fig. 3 The initial slope $\left(\frac{dm}{dt}\right)_{t=0}$ plotted against the oxygen concentration for c_{O_2} for data taken recorded at temperatures 523 K and 573 K.

Since we currently assume that we do not know anything about $f(X)$, we only extract the activation energy.

Fig. 4(a) shows the conversion extent vs. time for the SrFeO₃ pellet in the temperature range where the kinetics were not limited by mass transfer. In order to analyse the data we need to take the time derivative of X . Taking the numerical derivative of the original time series data is problematic, as noise in the original the mass data, taken at short intervals of 1 second, is amplified. Instead we take a subset of the experimental X , with larger discrete steps of 0.05 in the range 0 to 1, with these points marked by square symbols in Fig. 4(a). The numerical derivative was then calculated using second order accurate central differences, and first order accurate one-side difference at the boundary points. This resulted in a smoother derivative of the experimental rate which was suitable for further analysis.

Fig. 4(b) shows the plot of the isoconversional method used to extract the activation energies. All of the lines have comparable slopes and thus also similar activation energies. Taking the average of the slopes as the activation energy and the standard deviation as the error gives $E_a = 92 \pm 17$ [kJ mol⁻¹]. Yoo *et al.* determined an apparent activation energy for diffusion of $E_a = 110 \pm 10$ [kJ mol⁻¹], via electrical conductivity relaxation experiments in the temperature range 1073–1273 K. The similarity in these values which have overlapping errors, could indicate that the oxidation reaction is diffusion controlled in the temperature region considered. This is supported by the data in Fig. 2 where the pellet is seen to have a half life on the order of 1.5 to 2 times that of the finer SrFeO₃ powder. The SEM images (ESI,† Fig. S5) and particle size distribution measurements (Table 1) indicate that the SrFeO₃ powder from the SrFeO₃ pellet has twice the average particle size compared to the finer SrFeO₃ powder. If the reaction were surface controlled we might expect a much larger difference in kinetics, as they would depend directly on the surface area. However, given that the SrFeO₃ powder showed different phases we should be cautious in directly comparing these samples.

3.2.2 Kinetic model. Given that the activation energy does not show any strong dependence on the conversion extent X in Fig. 4(b), it should be possible to choose an appropriate $f(X)$ to



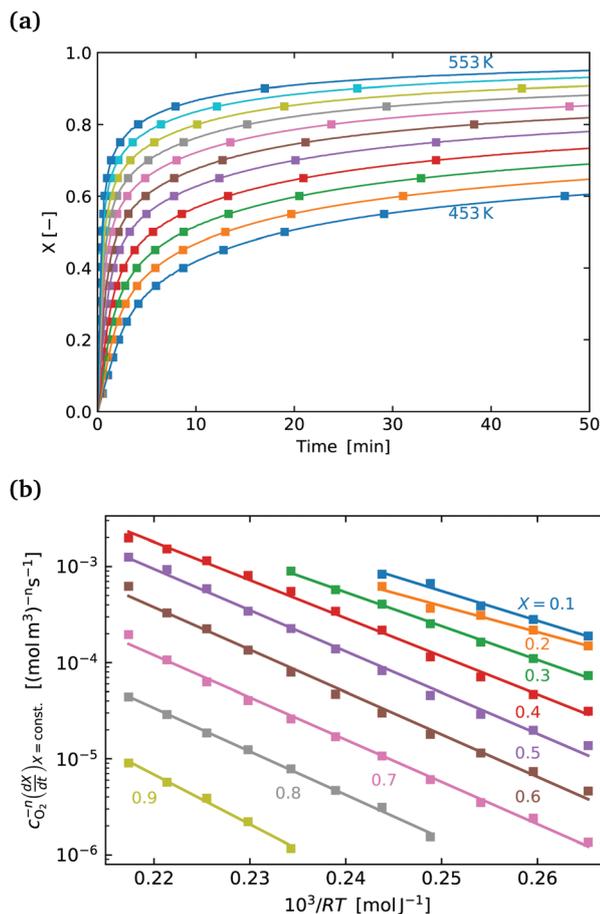


Fig. 4 (a) Conversion extent X vs. time for temperatures in the range 453–553 K in steps of 10 K. (b) $c_{O_2}^{-n} \left(\frac{dX}{dt}\right)_{X=\text{const.}}$ plotted on a log scale vs. $\frac{10^3}{RT}$ where the slope of each line is the activation $-E_a$ in kJ mol^{-1} . The labels indicate the fixed value of X for each line.

have a complete kinetic model for the system. For gas solid reaction systems with a uniform particles size that show a decreasing rate of reaction with time, one would usually employ a straightforward power law,

$$f(X) = (1 - X)^m \quad (10)$$

where the value of m depends on the reaction mechanism and rate limiting step of the process.⁴¹ In the case that we have a particle size distribution, smaller particles will be fully reacted faster than the larger particles, which complicates the modelling process. Kimura *et al.* have analytically shown that for a log-normal particle distribution, that the power law was still valid,⁴² and verified the result experimentally.^{43,44} Since we have approximate log-normal particle distributions here we apply this simple power law, so that our complete rate equations is given by,

$$\frac{dX}{dt} = \left(\frac{dX}{dt}\right)_{t=0} (1 - X)^m \quad (11)$$

where the initial rate is given by,

$$\left(\frac{dX}{dt}\right)_{t=0} = c_{O_2}^n k_0 \exp\left(\frac{E_a}{RT}\right) \quad (12)$$

taking the logarithm of both sides of eqn (11) gives,

$$\log\left(\frac{dX}{dt}\right) = m \log(1 - X) + \log\left(\left(\frac{dX}{dt}\right)_{t=0}\right) \quad (13)$$

so that a log-log plot of $\frac{dX}{dt}$ vs. $1 - X$ should be linear with a slope of m and intercept $\log\left(\left(\frac{dX}{dt}\right)_{t=0}\right)$.

Fig. 5(a) shows the experimental data obtained for the SrFeO₃ pellet and a fit made according to eqn (11). For small values of X the data diverges from the model, which is more significant at high temperatures. This behaviour is due to the fact that we do not have a perfect switch to an oxidising atmosphere. It takes approximately 20 seconds to go from low oxygen concentration to 90% oxygen atmosphere around the pellet. The effect is more evident at higher temperatures where the kinetics are faster and thus more of the data is affected by this switching region. When fitting the data we therefore only considered the regions outside this switching time, where the

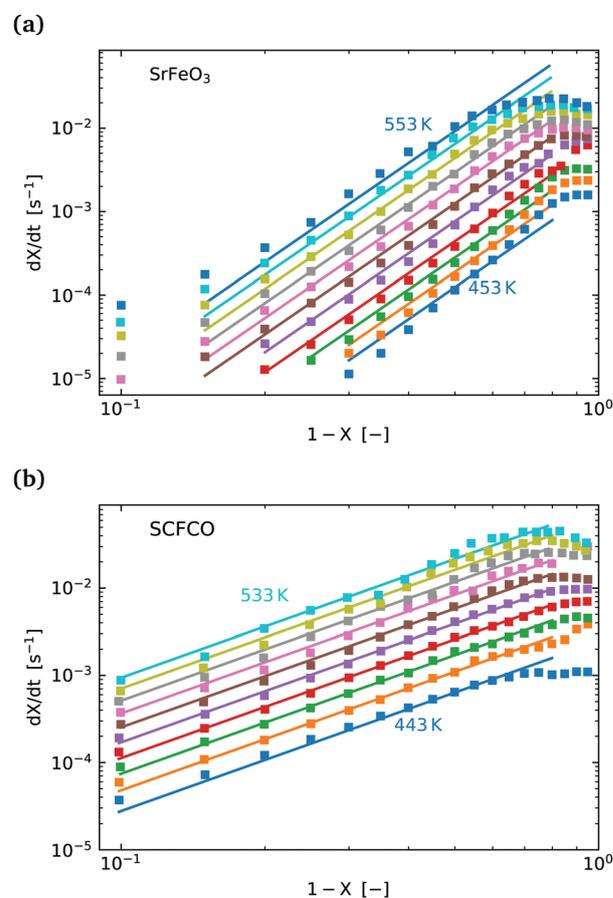


Fig. 5 (a) $\frac{dX}{dt}$ vs. $1 - X$ with a log-log scale. The markers show the experimental data from Fig. 4. The lines show a fit of eqn (11) to the section which shows a linear dependence. (b) The same plot and model fit for data obtained for SCFCO.



points appear to have a linear dependence on X . To fit the model each temperature data set was initially fit linearly to determine m , the average of these individual fit values was taken and used to fit the model to the entire dataset.

After the entire data set is fit using the same mean slope for each temperature, the intercepts from Fig. 5, can be used to determine the activation energy E_a and the pre-exponential factor k_0 , according to eqn (12). This analysis is shown graphically in the ESI,† where an Arrhenius plot can be seen to yield a clear linear dependence.

This fitting procedure and the iso-conversional analysis was repeated for SCFCO and the SrFeO₃ powder, with the fit parameters for all materials given in Table 2, and the experimental data and model shown in Fig. 5 for the pellet and SCFCO (see ESI,† for the full set of analysis plots). The SrFeO₃ powder sample had very similar activation energies compared to the pellet, but with a larger standard deviation in activation energies determined *via* the isoconversional method, as can be seen by the large error (Table 2). Comparison of the k_0 values for the SrFeO₃ pellet and finer powder shows a factor of two increase for the finer powder, which could indicate diffusion control, where we might expect $k_0 \propto \bar{d}$ (for a surface reaction $k_0 \propto \bar{d}^2$). However, given that both materials are actually different phases (see Table 1), this point should be treated with some caution.

SCFCO on the other hand has a considerably lower activation energy compared to SrFeO₃, which explains the faster kinetics. The $1 - X$ power law dependence m is also lower than that of SrFeO₃, and the model can be seen to have a better fit of the experimental data, particularly for larger X .

Considering all of the oxidation kinetics results, we conclude that the kinetics may be diffusion controlled and that the SCFCO sample has a lower activation energy of diffusion. Indeed, Wang *et al.* also reported that the material Ba_{0.5}Sr_{0.5}Co_{0.4}Fe_{0.6}O₃ has much improved oxygen transport properties compared to SrFeO₃.¹⁷ This could be attributed in both cases to the addition of cobalt into the lattice, which can introduce more oxygen vacancies due to the fact that cobalt likely remains in the Co³⁺.²² However, it important to note that such simplified conclusion are rather speculative. In reality the oxygen transport in the lattice can have many dependencies, including the temperature, oxygen non-stoichiometry, electric conductivity, and lattice strain/distortions. It is therefore very difficult to draw specific physical conclusions from the apparent activation energies.

Table 2 Oxidation kinetic parameters including, activation energies determined by the isoconversional method (iso- X), and the model fit parameters for eqn (11) and (12), for the SrFeO₃ pellet, the 200 mg of SrFeO₃ powder sintered at lower temperature (see Table 1), and the SCFCO sample

Sample	E_a iso- X [kJ mol ⁻¹]	E_a model [kJ mol ⁻¹]	m [—]	k_0 [s ⁻¹ (mol m ⁻³) ^{-n}]
SFO pel.	92 ± 17	95.3	3.94 ± 0.5	1.652 × 10 ⁷
SFO pow.	99 ± 28	97.9	3.98 ± 0.77	3.571 × 10 ⁷
SCFCO	79 ± 6.5	79.5	1.94 ± 0.11	6.293 × 10 ⁵

3.3 Reduction kinetics

The reduction kinetics were measured at isothermal conditions by switching the atmosphere from pure oxygen to a 1% oxygen in inert gas mixture, in the temperature range 573–723 K at intervals of 25 K. Below 573 K the kinetics were seen to be very slow and above 723 K the half life was approaching the same values where we saw a switch to mass transfer controlled kinetics for the oxidation. We also note here that for reduction kinetic analysis only the ground pellet and the SCFCO sample were analysed.

To monitor the kinetics the mass change as a function of time was again used to calculate $\Delta\delta(t)$, using eqn (2). Some examples of the raw TGA data are given in the ESI,† (Fig. S8). We again use the extent of conversion X , given by eqn (3), where the value of $\Delta\delta_\infty$ was taken to be the equilibrium values seen in Fig. 1(a). For temperatures outside the range a linear extrapolation of the data was used.

As a starting point we again look at the effect of temperature on the half life of the reactions $t_{1/2}$, shown in Fig. 6. The first thing worth noting is that the reduction reaction is considerably slower than the oxidation reaction for a given temperature.

At $\frac{10^3}{RT} \approx 0.2$ ($T = 603$ K) the half life for the oxidation of the SrFeO₃ pellet was around 15 seconds, whereas for the reduction reaction it was around 6000 seconds. This certainly indicates that the rate of the reduction reaction is not controlled by the same physical regime as that of the oxidation reaction. In-fact since diffusion plays the same role in both oxidation and reduction, it must be that the rate of the reduction reaction is controlled by the surface reaction kinetics. Otherwise we would expect similar or faster reaction times compared to the oxidation reaction.

Again the SCFCO sample shows faster reaction kinetics, but the difference is not as pronounced as the oxidation case. Here the SrFeO₃ has a $t_{1/2}$ which is typically double that of SCFCO, compared to a factor of four difference for oxidation.

3.3.1 Kinetic analysis and modelling. The mechanism by which the reduction reaction takes place can be broken into two parts, (1) a surface reaction where oxygen gas molecules form on the surface and then dissociate, followed by (2) the diffusion of oxygen vacancies into the bulk. In this case we

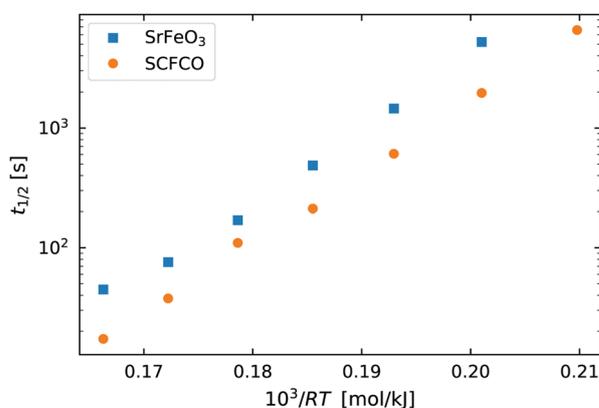


Fig. 6 The half life of the reduction reaction $t_{1/2}$ vs. $\frac{10^3}{RT}$ for the ground SrFeO₃ pellet annealed at 1573 K, and the SCFCO powder.



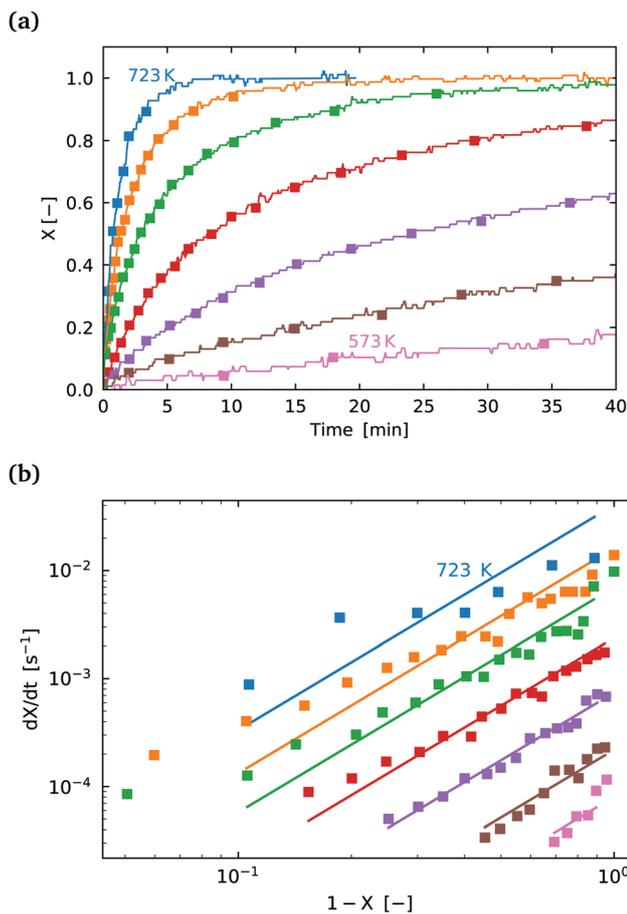


Fig. 7 (a) Conversion extent X vs. time for temperatures in the range 573–723 K in steps of 25 K for SrFeO₃ reduction. (b) $\frac{dX}{dt}$ vs. $1 - X$ with a log–log scale. The markers show the experimental data from (a) above. The lines show a fit of eqn (11) to the section which shows a linear dependence.

assume that the rate only depends on the temperature T , and the conversion extent X ,

$$\frac{dX}{dt} = k(T)f(X) \text{ [s}^{-1}\text{]} \quad (14)$$

where it is assumed that the rate constant $k(T)$ has an Arrhenius dependence according to eqn (6). The reduction does not require oxygen gas to proceed, so that we have omitted any terms dependent on the oxygen gas concentration. This is a simplification, since as we approach equilibrium the oxidation will start to compete with the reduction, and slow the reaction rate. This slowing of the reaction rate as we approach equilibrium is assumed to be part of $f(X)$ term. We again apply the isoconversional method to determine the activation energy. We also apply the same model given in eqn (11), where the initial rate loses the oxygen gas concentration term,

$$\left(\frac{dX}{dt}\right)_{t=0} = k_0 \exp\left(\frac{-E_a}{RT}\right). \quad (15)$$

Fig. 7(a) shows the conversion extent vs. time for the SrFeO₃ sample. The data has a lot more noise than that measured for

Table 3 A table showing the reduction kinetics activation energy determined by the isoconversional method (iso- X), and the model fit parameters for eqn (11) and (15), for the ground SrFeO₃ pellet, and the SCFCO sample. See the ESI for a full set of graphs showing the analysis

Sample	E_a iso- X [kJ mol ⁻¹]	E_a model [kJ mol ⁻¹]	m [–]	k_0 [s ⁻¹]
SFO pellet	144 ± 16	145.5	2.1 ± 0.9	1.27 × 10 ⁹
SCFCO	142 ± 5	146.2	2.3 ± 0.8	3.45 × 10 ⁹

oxidation, which is due to the fact that the overall mass changes were much smaller. For reduction kinetics 220 mg of powder sample were placed in the TGA and the switching of partial pressures resulted in a total mass change of approximately 1.3 mg. In comparison, the oxidation of the 500 mg sample gave a total mass change of more than 10 mg, and a much better signal to noise ratio. The noise in the data means that the analysis has a larger uncertainty, as can be seen by the error in m in Table 3.

The isoconversional analysis and model fit of both samples show similar apparent activation energies, which are much larger than those seen for the oxidation reaction. This offers further evidence that the reaction control cannot be the same as is seen for the oxidation reaction. Fig. 7(b) shows the experimental data fit with the $(1 - X)^m$ model, with the fit parameters given in Table 3.

We conclude that the kinetic regime for the reduction reaction is different from the oxidation reaction, with a much larger energy barrier, and by process of elimination it must be surface reaction controlled. From a simplified energetic perspective, one might expect the activation energy of this reaction to at least be larger than the change in enthalpy for the reaction. From the literature the change in enthalpy per mole of atomic oxygen is in the range 70–90 [kJ mol⁻¹].⁶ Indeed this is the case with values of activation energy close to 150 [kJ mol⁻¹] measured here. Most of this change in enthalpy would also be attributed to the reaction at the surface, supporting the result that the reduction reaction is limited by the surface reaction.

4 Conclusions

The kinetic data collected in this work should be of high value for designing applications which make use of SrFeO₃ as an oxygen exchange material. In particular the low temperature oxidation and reduction kinetics are very important for oxygen separation applications. The results show that the substitution of some cobalt in place of iron, while also adding calcium to maintain the cubic perovskite structure in Sr_{0.93}Ca_{0.07}Fe_{0.9}Co_{0.1}O₃, greatly improved the oxidation kinetics. The presence of some secondary phases complicates the physical interpretation of this result, but this improvement in kinetics would allow cycling materials at lower temperatures. This can improve the energy balance and practicality of chemical looping oxygen separation processes. Analysis of the kinetic data indicated that the reduction reaction is limited by the surface reaction kinetics, and with a lower level of certainty the oxidation reaction was in the diffusion controlled regime.



Conflicts of interest

There are no conflicts of interest to declare.

Acknowledgements

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References

- C. Haavik, T. Atake and S. Stølen, *Phys. Chem. Chem. Phys.*, 2002, **4**, 1082–1087.
- A. Holt, T. Norby and R. Glenne, *Ionics*, 1999, **5**, 434–443.
- J. Yoo, C.-Y. Yoo, J.-H. Yu and A. J. Jacobson, *J. Am. Ceram. Soc.*, 2017, **100**, 2690–2699.
- Y. Takeda, K. Kanno, T. Takada, O. Yamamoto, M. Takano, N. Nakayama and Y. Bando, *J. Solid State Chem.*, 1986, **63**, 237–249.
- K. Krishnamurthy, D. Acharya and F. Fitch, Pilot-Scale Demonstration of a Novel, Low-Cost Oxygen Supply Process and its Integration with Oxy-Fuel Coal-Fired Boilers, Boc Group, Inc., the (a delaware corp) technical report, 2008.
- B. Bulfin, J. Lapp, S. Richter, D. Gubàn, J. Vieten, S. Brendelberger, M. Roeb and C. Sattler, *Chem. Eng. Sci.*, 2019, **203**, 68–75.
- J. Vieten, B. Bulfin, F. Call, M. Lange, M. Schmücker, A. Francke, M. Roeb and C. Sattler, *J. Mater. Chem. A*, 2016, **4**, 13652–13659.
- H. E. Bush, R. Datta and P. G. Loutzenhisser, *Sol. Energy*, 2019, **188**, 775–786.
- B. Bulfin, J. Vieten, C. Agrafiotis, M. Roeb and C. Sattler, *J. Mater. Chem. A*, 2017, **5**, 18951–18966.
- Z. Yang, Y. Lin and Y. Zeng, *Ind. Eng. Chem. Res.*, 2002, **41**, 2775–2784.
- H. Ikeda, S. Nikata, E. Hiraoka, A. Tsuchida and N. Miura, *Chem. Eng. Sci.*, 2016, **147**, 166–172.
- Z. Yang and Y. Lin, *Ind. Eng. Chem. Res.*, 2003, **42**, 4376–4381.
- M. Haänchen, A. Stiel, Z. Jovanovic and A. Steinfeld, *Ind. Eng. Chem. Res.*, 2012, **51**, 7013–7021.
- M. Ezbiri, K. M. Allen, M. E. Gálvez, R. Michalsky and A. Steinfeld, *ChemSusChem*, 2015, **8**, 1966–1971.
- B. Moghtaderi, *Energy Fuels*, 2009, **24**, 190–198.
- J. Vieten, B. Bulfin, D. E. Starr, A. Hariki, F. M. de Groot, A. Azarpira, C. Zachäus, M. Hävecker, K. Skorupska and N. Knoblauch, *et al.*, *Energy Technology*, 2019, **7**, 131–139.
- L. Wang, R. Merkle and J. Maier, *ECS Trans.*, 2009, **25**, 2497–2505.
- J. Yoo, A. Verma, S. Wang and A. J. Jacobson, *J. Electrochem. Soc.*, 2005, **152**, A497–A505.
- J. Yoo and A. J. Jacobson, *MRS Online Proc. Libr.*, 2002, **756**, EE3.10.
- M. Rekas, T. Bak, J. Nowotny, C. Sorrell, Y. Zhao, K. Foger and E. Vance, *J. Mater. Sci.: Mater. Electron.*, 2000, **11**, 691–696.
- S. Bychkov, I. Gainutdinov, S. Chizhik and A. Nemudry, *Solid State Ionics*, 2018, **320**, 297–304.
- J. Vieten, B. Bulfin, P. Huck, M. Horton, D. Guban, L. Zhu, Y. Lu, K. A. Persson, M. Roeb and C. Sattler, *Energy Environ. Sci.*, 2019, **12**, 1369–1384.
- J. Vieten, B. Bulfin, M. Roeb and C. Sattler, *Solid State Ionics*, 2018, **315**, 92–97.
- A. E. Danks, S. R. Hall and Z. Schnepf, *Mater. Horiz.*, 2016, **3**, 91–112.
- J. Rodriguez-Carvajal, *Nouveaux développements de FullProf Analyse de la microstructure et utilisation du recuit simulé pour la résolution de structures*, CEA/Saclay, France, 2001.
- J. Blasco, J. Stankiewicz and J. Garca, *J. Solid State Chem.*, 2006, **179**, 898–908.
- F. Deganello, L. Liotta, A. Longo, M. Casaletto and M. Scopelliti, *J. Solid State Chem.*, 2006, **179**, 3406–3419.
- S. Dann, M. Weller and D. Currie, *J. Solid State Chem.*, 1992, **97**, 179–185.
- F. Ramezanipour, J. E. Greedan, L. M. Cranswick, V. O. Garlea, J. Siewenie, G. King, A. Llobet and R. L. Donaberger, *J. Mater. Chem.*, 2012, **22**, 9522–9538.
- Y. Bréard, C. Michel, M. Hervieu, F. Studer, A. Maignan and B. Raveau, *Chem. Mater.*, 2002, **14**, 3128–3135.
- E. Bakken, S. Stølen, T. Norby, R. Glenne and M. Budd, *Solid State Ionics*, 2004, **167**, 367–377.
- C. Haavik, T. Atake, H. Kawaji and S. Stølen, *Phys. Chem. Chem. Phys.*, 2001, **3**, 3863–3870.
- B. Tofield, C. Greaves and B. Fender, *Mater. Res. Bull.*, 1975, **10**, 737–745.
- M. Schmidt and S. Campbell, *J. Solid State Chem.*, 2001, **156**, 292–304.
- P. Battle, T. Gibb and S. Nixon, *J. Solid State Chem.*, 1989, **79**, 75–85.
- P. Fielitz and G. Borchardt, *Phys. Chem. Chem. Phys.*, 2016, **18**, 22031–22038.
- R. D. Souza, *Phys. Chem. Chem. Phys.*, 2006, **8**, 890–897.
- N. A. Merino, B. P. Barbero, P. Eloy and L. E. Cadús, *Appl. Surf. Sci.*, 2006, **253**, 1489–1493.
- J. Crank, *The Mathematics of Diffusion*, Oxford University Press, Oxford, 1975.
- S. Vyazovkin, A. K. Burnham, J. M. Criado, L. A. Pérez-Maqueda, C. Popescu and N. Sbirrazzuoli, *Thermochim. Acta*, 2011, **520**, 1–19.
- O. Levenspiel, Inc., New York, 1972.
- S. Kimura, S. Tone and T. Otake, *J. Chem. Eng. Jpn.*, 1981, **14**, 491–493.
- S. kimura, Y. Takagi, S. Tone and T. Otake, *J. Chem. Eng. Jpn.*, 1983, **16**, 217–223.
- S. Kimura, *AIChE J.*, 1989, **35**, 339–342.

