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Triperyleno[3,3,3]propellane triimides: achieving a new generation of quasi-$D_{3h}$ symmetric nanostructures in organic electronics†

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Rigid three-dimensional (3D) polycyclic aromatic hydrocarbons (PAHs), in particular 3D nanographenes, have garnered interest due to their potential use in semiconductor applications and as models to study through-bond and through-space electronic interactions. Herein we report the development of a novel 3D-symmetric rylene imide building block, triperyleno[3,3,3]propellane triimides (6), that possesses three perylene monoimide subunits fused on a propellane. This building block shows several promising characteristics, including high solubility, large π-surfaces, electron-accepting capabilities, and a variety of reactive sites. Further, the building block is compatible with different reactions to readily yield quasi-$D_{3h}$ symmetric nanostructures (9, 11, and 13) of varied chemistries. For the 3D nanostructures we observed red-shift absorption maxima and amplification of the absorption coefficients when compared to the individual subunits, indicating intramolecular electronic coupling among the subunits. In addition, the microplates of 9 exhibit comparable mobilities in different directions in the range of $10^{-3}$ cm$^2$ V$^{-1}$ s$^{-1}$, despite the rather limited intermolecular overlap of the π-conjugated moieties. These findings demonstrate that these quasi-$D_{3h}$ symmetric rylene imides have potential as 3D nanostructures for a range of materials applications, including in organic electronic devices.

Introduction

Due to the prospects of distinctive electronic and optical properties, considerable attention has been given to the synthesis of polycyclic aromatic hydrocarbons (PAHs). A–D Most PAHs can be classified into one- (1D) and two-dimensional (2D) PAHs, such as 1D linear polyacenes and 2D nanographenes, with both of them showing robust and large charge-carrier mobilities in transistors. 2–3 In contrast, three-dimensional (3D) PAHs remain an underexplored molecular construct due to numerous synthetic challenges, in particular for rigid 3D-symmetric PAHs, 8–11 even though they have been suggested as materials systems for organic electronics, molecular rotors, host–guest chemistries, liquid crystals, and supramolecular assemblies. 12–13

The classic 3D organic building block triptycene was first synthesized in 1942 by Bartlett et al. (Fig. 1). 16 Triptycene and its derivatives, for instance, have been widely incorporated into polymer backbones to disrupt chain packing and increase the free volume, qualities that are beneficial for improving gas diffusion, sensory response, thermal stability, Young’s modulus, and luminescence quantum efficiency. 17–20 On the other hand, there has been interest in the synthesis of triptycene-extended PAHs, wherein additional five- and six-membered aromatic rings are fused to the triptycene core. 21–28

Fig. 1 Representation of the structural evolution of 3D-symmetric building blocks for rigid 3D PAHs, from triptycene (left) to trinaphtho[3,3,3]propellane (middle) to triperyleno[3,3,3]propellane triimides (right).
This interest is partly due to the π-extended triptycenes possessing large internal molecular free volume (IMFV), which can provide an approach to prepare nanoporous materials or two-dimensional (2D) polymers. Importantly, such molecules can serve as key components in molecular electronic materials, as they can possess extended π-conjugation in three directions in a rigid structure, which will influence the condensed-phase aggregation and alter the stability, solubility, electronic properties, and device performance. For example, Wang et al. recently demonstrated a π-extended triptycene fused with perylene diimides (PDI) as a non-fullerene acceptor for organic solar cells. Further, Nuckolls et al. reported a promising class of 3D nanographenes based on triptycene, which were used as electron-extracting layers, which helped to decrease contact resistances in perovskite solar cells and led to exceptional materials that exhibit remarkably stable electron-accepting properties to the resulting 3D materials. In particular, 3D-symmetric molecules provide an ideal model to investigate through-bond coupling effects at the imide nitrogen atoms, the cation. Thus, we subjected these molecules to ensure adequate solubility for further synthetic modification. Thus, we subjected these to typical PMI reactions to functionalize the bay positions. In analogy with the standard PMI chemistry, bromination of 6 in chloroform gave hexabrominated 7 in 45% yield. To explore the scope of the Pd-catalyzed cross-coupling strategy, reactions of 7 with different substrates (4-tert-butylphenylboronic acid, 2-trimethyltin-5-tribisopropylsilylthiophene, and n-butyl-acetylene) were studied (Scheme 2). We found that the desired product could undergo 6-fold cyclization smoothly to produce the desired product in 75% yield and 11 in 75% yield under standard photo-induced cyclization conditions. During our initial synthetic effort for compound 13, we observed no cyclization products under the typical alkyne cyclization conditions with either PtCl2 or some Brønsted acids. Finally, the cyclization product 13 was obtained by treatment of 12 with 1,8-diazabicyclo[5.4.0]undec-7-ene (DBU) in 25% yield. It should be emphasized that the t-butyl, n-butyl, and triisopropylsilyl (TIPS) groups are necessary to ensure adequate solubility for efficient cyclization. The feasibility of the synthetic strategy suggests that the bromine-containing building block 7 is amenable to a wide range of large and rigid 3D nanostructures with tunable geometries through PMI/PDI chemistry. These rigid 3D-symmetric compounds are soluble in common solvents and could be characterized by spectroscopic analysis, including HMBC and 1D NOE.

**X-ray crystallography**

Single crystals of 9 suitable for X-ray diffraction analysis were grown by slow diffusion of methanol into toluene solution at room temperature. 9 crystallizes in the P2₁/c space group with
four molecules per unit cell. As shown in Fig. 2A and B, the molecule adopts a quasi-$D_{3h}$ symmetric conformation with a dihedral angle of 116° between the individual subunits. The central carbon–carbon bond (denoted as C1–C2) length of the bridging ethylene unit in 9 is quite long, 1.62 Å, a feature due to the large PDI sizes and bulky alkyl chains, which combine to provide considerable strain to the central carbon–carbon bond. The surrounding carbon–carbon bond lengths are 1.51–1.53 Å, which are close to that of typical sp$^2$–sp$^2$ carbon–carbon bonds. 9 has a maximal distance of 2.1 nm between the ends of two subunits and a width of 1.2 nm and a length of 1.6 nm across one subunit.

The molecules in the crystal pack in 2D layers. As shown in Fig. 2C, in an individual layer, two molecules are parallel to each other, with a distance of 3.1 nm between their symmetry axes, with the third molecule situated in the centre of this cavity and its symmetry axis perpendicular to that of neighbor molecule. Hence, there are no large voids in the layers. This crystal packing is stabilized via multiple short C–H⋯π contacts (in the range of 2.72–2.87 Å) among the alkyl chains on one molecule and the PDI monomers on adjacent molecules; no intermolecular π⋯π interactions are observed due to the alkyl chains, which block the overlap between the π-faces of the subunits. This packing contrasts with triptycene derivatives, which usually crystallize into layer structures that are driven by intermolecular π⋯π interactions between the individual acene subunits and produce large voids between the individual subunits.24,29
Electronic structure and optical and redox properties

To gain insight into the electronic and optical properties of these molecules, the electronic structure was evaluated via density functional theory (DFT) calculations at the OT-ωB97X-D/6-31g(d,p) level of theory (further details are given in the ESI). Calculations were carried out for the four full molecules (6, 9, 11, and 13) and for each of the four monomeric units. The frontier HOMO and LUMO of 6, 9, 11 and 13 are assorted linear combinations of the HOMO and LUMO (respectively) of the individual monomers. As shown in Fig. 3, the HOMO is delocalized across all three units in 6, 9, 11, and 13. For 6 and 9, the HOMO-1 and HOMO-2 are energetically degenerate, are approximately 0.3 eV energetically stabilized when compared to the HOMO, and are delocalized on only two monomer units (Table S2†). The pattern differs in 11 and 13, where the degeneracy of the HOMO-1 and HOMO-2 is broken, with the HOMO-1 also delocalized over all three units (as with the HOMO), while the HOMO-2 spans only two monomers; here, the HOMO is approximately 0.2 eV higher in energy than the HOMO-1. All four molecules share the same pattern for the frontier LUMO, LUMO+1, and LUMO+2, where the LUMO and LUMO+1 are degenerate and delocalized over two of the three monomers while the LUMO+2 is delocalized over all three monomers.

UV-vis absorption spectra of 6, 9, 11, and 13 in chloroform are shown in Fig. 4. In the high-energy (short wavelength) region, the absorption spectra of these compounds, except 6, show well-defined vibronic fine structures. In the low-energy (longer wavelength) region, the absorption spectra of the bay-extended compounds (9, 11 and 13) show a clear blue-shift relative to the parent 6, and there is a red-shift in the absorption in the order of 13 > 9 > 11. This trend is consistent with those previously observed in bay-extended PDI55 and PMI.44,59 It should be noted that compound 6 exhibits the characteristic absorption of typical perylene imide with maxima at 554 nm (2.24 eV), which is slightly red-shifted by 46 nm (0.20 eV)

Fig. 3 Pictorial representations of selected frontier molecular orbitals of 6, 9, 11, and 13 as determined at the OT-ωB97X-D/6-31g(d,p) level of theory. Pictorial representations of the monomeric units are shown in the ESI.†
compared to that of PMI (508 nm; 2.44 eV), indicating that there is some electronic communication among the PMI subunits. Remarkably, at the longest wavelength absorption (554 nm), the maximum absorption coefficient is 150 844 M⁻¹ cm⁻¹, which is significantly larger than the sum of the three PMI units (3 × 33 305 M⁻¹ cm⁻¹ = 99 915 M⁻¹ cm⁻¹) (Fig. S1 in the ESI†), a phenomenon that has been observed in triptycene-based 3D graphene nanostructures.28 A similar red-shift and amplification of the absorption coefficient (122 700 M⁻¹ cm⁻¹) at the maximum absorption wavelength (433 nm) of 13 are also found in contrast to the coronene monomide subunit (418 nm, 33 020 M⁻¹ cm⁻¹).29 Compounds 6 and 9 show strong emission bands at 588 nm and 514 nm, with fluorescence quantum yields of 85% and 89%, which are higher than those of 11 (505 nm, 22%) and 13 (418 nm, 32%).

Time-dependent DFT (TDDFT) calculations at the OToB97X-D/6-31g(d,p) level of theory were performed to dig deeper into the nature of the molecular excited-state characteristics. For the full molecules, the excited-states contain significant contributions from several one-electron transitions (Table S2†), and state degeneracy makes the description of the excitations somewhat complex. Further, for 6, 11 and 13, the low-lying excitations with appreciable oscillator strengths are composed of two degenerate excitations, while for 9 this degeneracy is broken. As reported by Mullen and co-workers, the experimental spectrum of 13 does show a low-energy, low-oscillator strength transition that can be described as having PMI → coronene character.45,44,59 For 6 and 13, natural transition orbital (NTO) analyses show that the first excitations with large intensity are combinations of two transitions where the holes and electrons are delocalized on two monomers, while the second excitation is composed of a combination of two monomer-localized transitions. In 9, the pattern is reversed, though the lowest-energy transition has the holes and electrons mainly on one monomer. For 11, the two degenerate transitions are composed of two mostly monomer-localized hole–electron pairs, though in each there is some extension onto the second monomer. Overall, there is generally a good agreement among the UV-vis absorption spectra and computed TDDFT excitations.

Adiabatic ionization potentials and electron affinities (AIP and AEA, respectively) and spin densities were determined by DFT. Both the AIP and AEA decrease when moving from the monomer to the full 3D molecule, with the AEA decreasing more; we note that the decrease of each value is very similar across all four molecules (Table S2†). The spin densities, which reveal the distribution of the excess charge in the oxidized or reduced states, of the monomers display some density on the terminal sp³ carbons in both the anion and cation (Fig. S9†). In the full molecules, the radical-cation is fully delocalized – a feature consistent with the fact that the HOMO for each molecule was delocalized over each monomer unit. The radical-anion, on the other hand, is only partially delocalized. We recall that the LUMO and LUMO+1 are degenerate orbitals that are both delocalized across only two subunits; we deduce that the nature of the radical-anion spin density arises from a linear combination of this orbital character.

Cyclic voltammetry (CV) and differential pulse voltammetry (DPV) studies were performed in dichloromethane with 0.1 M tetrabutylammonium hexafluorophosphate (TBAPF₆) as the supporting electrolyte, and the onset oxidation-reduction potentials were determined relative to Fc/Fc⁺ (4.8 eV) (Fig. 5). In contrast to the CV of PMI exhibiting two reduction waves and one oxidation wave, the CV of 6 reveals a splitting into four reduction waves and three oxidation waves. The observed splitting is due to the coulombic interactions from the electrons in different PMI subunits, which indicates electronic communication between different PMI subunits – a result consistent with evaluation of the DFT-derived MO for the neutral molecules and spin densities of the radical-cation and radical-anion states. The first half-wave reduction potential of 6 is −1.42 V, which is slightly more positive than that (−1.46 V) of PMI,
suggesting weak conjugation between the individual PMI subunits,\(^9\) again consistent with the radical-anion spin densities and AEA. The CV of 9 also shows four reduction waves and three oxidation waves, and 11 exhibits four reduction waves and two oxidation waves. Interestingly, 13 shows only two reduction waves and two oxidation waves. Moreover, we observed that the oxidation waves of these compounds nearly overlap, which is further evidence that electron communication between the individual subunits is relatively weak. According to their onset potentials, electron affinities/ionization potentials (EA/IP) of 6, 9, 11, and 13 are estimated to be \(-3.46/-5.56, -3.29/-5.70, -3.31/-5.74, \) and \(-3.13/-5.89 \text{ eV}, \) respectively (relative to \( \text{Fc/Fc}^+ \), 4.8 eV).

Semiconductor properties

As a proof-of-concept, we evaluated the intrinsic charge-carrier mobility of 9 in the form of single-crystalline plates, which were prepared by slow diffusion of isopropyl alcohol into a concentrated toluene solution. Fig. 6A shows the SEM images of 9 that form rhombic plates. The SAED pattern of 9 could be indexed based on the single crystal data, which indicate that the (100) plane of the crystal is parallel to the surface of the substrate (Fig. 6B). The single crystalline field-effect transistors were fabricated by an “organic nanowire mask” technique with four electrodes to explore the transport properties along the different directions of the plates (Fig. 6C).\(^{33}\) Fig. 6D shows the transfer characteristics of microplates of 9 tested under ambient conditions, which exhibit hole charge-carrier transport behavior. The conducting channels of 1 ↔ 2, 2 ↔ 3, 3 ↔ 4, and 4 ↔ 1 exhibit very similar mobilities of 4.92 × 10^{-3}, 4.55 × 10^{-3}, 4.95 × 10^{-3}, and 4.51 × 10^{-3} cm² V⁻¹ s⁻¹, respectively (Table S5). Although these mobilities are by no means large, 10⁻³ cm² V⁻¹ s⁻¹ is in fact quite good considering the limited intermolecular overlap of the \( \pi \)-conjugated units.\(^3\) Importantly, the distinct structure with a quasi-\( D_{3h} \) symmetric geometry could serve as a model to investigate the anisotropic charge-carrier transport in different lattice planes.

Conclusions

In summary, we described the straightforward synthesis of 3D-symmetric rylene imide, triperyleno[3,3,3]propellane triimides, that shows attractive properties, such as high solubility, large \( \pi \)-surfaces, electron accepting ability, and many reactive sites. Though reaction with different substrates, we demonstrated that this rylene imide is compatible with different reactions, such as Pd-catalyzed cross-coupling reactions and cyclo-dehydrogenation reactions, which can be applied to synthesize a variety of complex 3D nanostructures with different molecular geometries. Here, a series of large, rigid 3D-symmetric nanostructures were successfully synthesized by functionalization of the bay positions of the building block. Compared to their subunits, these molecules show red-shifted and amplified absorptions. There is notable electronic communication among the monomer units, which influences the nature of the excited states and redox characteristics. The combination of facile synthetic accessibility, high chemical stability, molecular diversification, and tunable electronic properties makes the 3D rylene imide and its derivatives promising materials for use in organic electronic devices. Further studies to explore the relationship between the molecular geometry and the intramolecular electronic communication as well as for synthesis of even larger, rigid 3D PAHs for organic devices are now in progress.

Conflicts of interest

There are no conflicts to declare.

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Notes and references