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## A double helix of opposite charges to form channels with unique CO<sub>2</sub> selectivity and dynamics†

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Porous molecular materials represent a new front in the endeavor to achieve high-performance sorptive properties and gas transport. Self-assembly of polyfunctional molecules containing multiple charges, namely, tetrahedral tetra-sulfonate anions and bifunctional linear cations, resulted in a permanently porous crystalline material exhibiting tailored sub-nanometer channels with double helices of electrostatic charges that governed the association and transport of CO<sub>2</sub> molecules. The charged channels were consolidated by robust hydrogen bonds. Guest recognition by electrostatic interactions remind us of the role played by the dipolar helical channels in regulatory biological membranes. The systematic electrostatic sites provided the perfectly fitting *loci* of complementary charges in the channels that proved to be extremely selective with respect to N<sub>2</sub> (*S* = 690), a benchmark in the field of porous molecular materials. The unique screwing dynamics of CO<sub>2</sub> travelling along the ultramicropores with a step-wise reorientation mechanism was driven by specific host–guest interactions encountered along the helical track. The unusual dynamics with a single-file transport rate of more than 10<sup>6</sup> steps per second and an energy barrier for the jump to the next site as low as 2.9 kcal mol<sup>−1</sup> was revealed unconventionally by complementing *in situ* <sup>13</sup>C NMR anisotropic line-shape analysis with DFT modelling of CO<sub>2</sub> diffusing in the crystal channels. The peculiar sorption performances and the extraordinary thermal stability up to 450 °C, combined with the ease of preparation and regeneration, highlight the perspective of applying these materials for selective removal of CO<sub>2</sub> from other gases.

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## Introduction

The manufacture of materials with permanent porosity is one of the major challenges in current research. This goal can be achieved through various strategies that promote the formation of covalent and metal–organic frameworks as well as molecular architectures.<sup>1–8</sup> They are diversified by their greater or lesser tendency to form stable structures and their degree of crystalline periodicity. Highly crystalline porous molecular materials can be constructed easily by spontaneous assembly of precursor molecules, provided that their shapes and functionalities are adequately designed.<sup>9–11</sup> Indeed, fully organic porous materials can be assembled by means of soft and reversible interactions enabling the highest modularity, processability and large

single-crystal formation. Their excellent performances in gas capture and interaction with chemical and physical stimuli have been explored in recent years.<sup>12–29</sup>

Organic functions of opposite charges on complementary molecules promote intense attractive interactions, and reciprocal recognition such as positively charged ammonium groups with carboxylic and sulfonic anions; thus, electrostatic interactions can be exploited to sustain robust molecular frameworks.<sup>30–35</sup> Additionally, electrostatic charges can be exposed to the pores, making it possible to expand research towards unexplored developments.

In our approach, multiple-charged functional groups protruding in the space produce directional interactions, favourable to the design of porous molecular architectures, forming programmable charged channels. In general, the shape and symmetry of the building blocks to construct porous frameworks must be taken into account: specifically, highly symmetric and rigid tetrahedral struts (*S*<sub>4</sub> point group) constitute a primary element, which sustained permanently porous materials and yielded very high surface areas.<sup>36–44</sup>

Biological systems are promising sources for the formation of sub-nanometer channels for the transport of molecules and

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ions, wherein the electrostatic distribution of charges on the channel walls and steric factors play a strategic role for building useful functions.<sup>45</sup> By analogy, we have engineered a permanently porous 3D molecular architecture containing sub-nm channels formed by a double helix with two intertwined and alternatively charged ribbons, characteristic of transmembrane channels. It was generated by the assembly of tetrahedral and tetrafunctional 4,4',4'',4'''-methane tetrabenzenesulfonate (TBS) anions and a bifunctional cation 1,4-diamidinumbenzene (DAB) and was found to be highly stable. The porous structure was capable of reversibly absorbing relevant gases such as CO<sub>2</sub> and CH<sub>4</sub> with an exceptional CO<sub>2</sub>/N<sub>2</sub> selectivity of 690 owing to the tight fitting and the complementary electrostatic pattern of the sub-nm channels. The direct detection of CO<sub>2</sub> pervading the porous molecular framework was provided by multinuclear 2D MAS NMR spectra, depicting the exceptional CO<sub>2</sub>-matrix association in the properly fitting pockets. Furthermore, the anisotropic <sup>13</sup>C NMR variable-temperature lineshapes unveiled the unique coordinated rotation and translation dynamics of CO<sub>2</sub> while it jumps from one site to the next, following the helical charge patterning of the channel. Intriguingly, CO<sub>2</sub> molecules receive the orientational encoding from the nanochannel sites at each diffusion turn of the screwing motion. Thus, by this unconventional means, we could establish the transport rate of a single gas molecule that is as fast as a million steps per second, corresponding to guest diffusion rates observed in

biological channels. DFT calculations and Molecular Dynamics (MD) supported the NMR determination of the preferred arrangement and the energy landscape perceived by CO<sub>2</sub> exploring the channels.

## Results and discussion

### Fabrication of charged sub-nanometer channels

The organotetrakisulfonate TBS self-assembled with organo-diamidinium DAB to form needlelike single crystals suitable for X-ray analysis (ESI†). After guest removal by thermal treatment under vacuum at 150 °C, the 3D architecture is maintained, yielding a new crystalline porous organosulfonate-amidinium salt (CPOS-5) (Fig. 1). It crystallizes in the tetragonal *I*4<sub>1</sub>*a* space group, which contains a quarter of TBS (a phenyl-SO<sub>3</sub><sup>−</sup> group) and half of DAB in the asymmetric unit (molecular formula equal to TBS(DAB)<sub>2</sub>). A 4-fold rotoinversion axis crosses the central carbon (C1), which orients four sulphonate groups along the vertices of a tetrahedron. Hence, organosulfonate anions and organoamidinium cations, driven by hydrogen bonds,

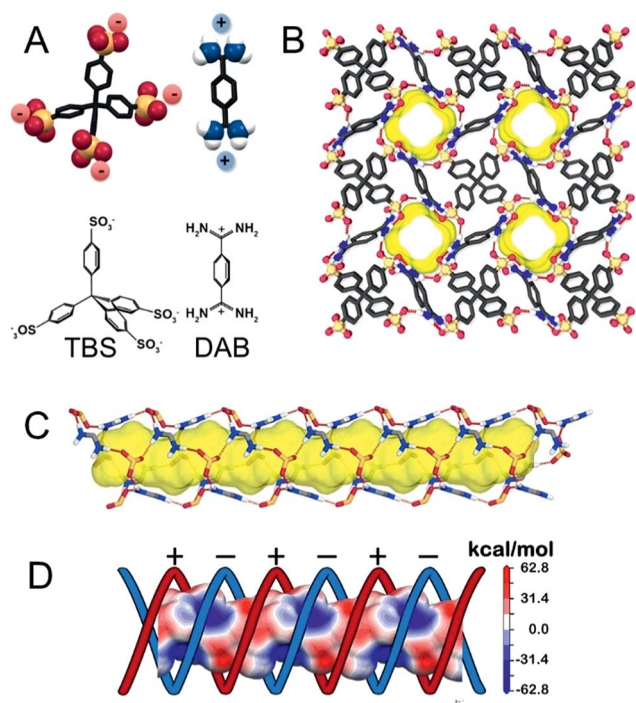


Fig. 1 (A) Schematic representation of anionic TBS and cationic DAB synthons of the porous organic framework. (B) View of the porous framework along the *c*-axis; yellow represents the empty channels (a rolling sphere of 1.4 Å was used). The channel running parallel to the *c*-axis: (C) the synthons are connected through a H-bonding network in a helical fashion; (D) electrostatic map showing the positively and negatively charged helical ribbons.

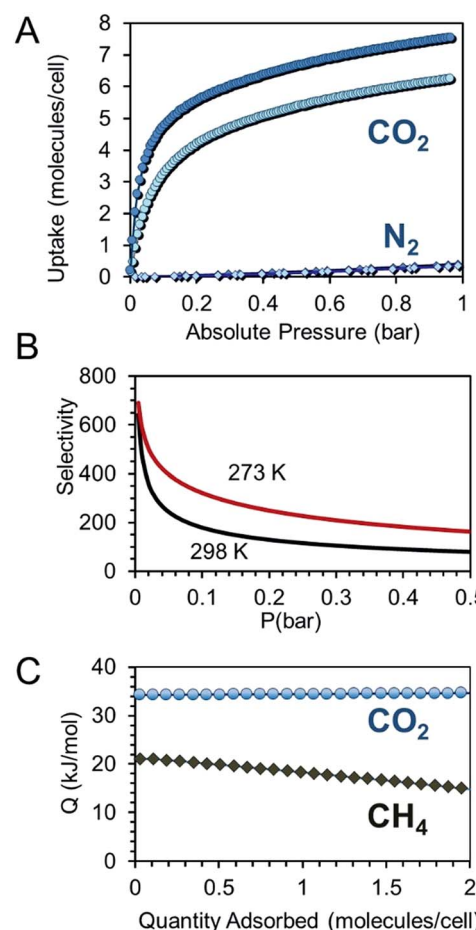


Fig. 2 (A) CO<sub>2</sub> and N<sub>2</sub> isotherms of CPOS-5 collected at 273 K (blue circles and diamonds, respectively) and 298 K (light-blue circles and diamonds, respectively). (B) CO<sub>2</sub>/N<sub>2</sub> selectivity values of CPOS-5 versus pressure calculated starting from a 15/85 mixture and applying IAST calculation. (C) Isothermic heats of adsorption (*Q*<sub>st</sub>) versus CO<sub>2</sub> and CH<sub>4</sub> loading.

assemble in a double helix of electrically charged ribbons (Fig. 1C). Both negatively and positively charged moieties draw alternatively charged helices that lead to the formation of 1D-channels (Fig. 1D), running parallel to the *c*-axis. The empty channels with a cross-section of  $5.3 \times 6.8 \text{ \AA}^2$  generate a porosity of  $743.03 \text{ \AA}^3$  per cell corresponding to 14.6% of the cell volume (radius of the rolling sphere =  $1.2 \text{ \AA}$ ). The  $^{13}\text{C}$  MAS NMR spectrum confirmed the absence of guests in the cavities (ESI†). The porous compound exhibits high stability to water and temperature (up to  $400^\circ\text{C}$ ) owing to the connection of four organosulfonate anions to DAB monomers through charge-assisted hydrogen bonds in all directions.

### Gas capture and $\text{CO}_2$ selectivity

$\text{CO}_2$  adsorption isotherms were collected on CPOS-5 to demonstrate its permanent porosity (Fig. 2). Interestingly, at  $273 \text{ K}$  the  $\text{CO}_2$  adsorption isotherm displayed a Langmuir profile, which at  $1 \text{ bar}$  reached the value of  $48 \text{ cm}^3 \text{ g}^{-1}$ , representing 8  $\text{CO}_2$  molecules per unit cell (Fig. 2A). Notably,  $\text{CO}_2$  uptake was already extremely high at low pressure and exhibited a loading as high as  $59 \text{ mg g}^{-1}$  ( $273 \text{ K}$ ) at  $0.1 \text{ bar}$ , suggesting excellent affinity of the quadrupolar gas molecules with charge-decorated pore

surfaces. At  $0.1 \text{ bar}$  and at room temperature, typical conditions of  $\text{CO}_2$  pressure in flue gas from power plants, CPOS-5 adsorbs  $39 \text{ mg g}^{-1}$ , outperforming many well-known MOFs with greater surface areas such as HKUST-1, UMCM-150( $\text{N}_2$ ), MIL-47, IRMOF-3 and ZIF-8,<sup>46</sup> porous hydrogen-bonds and supramolecular organic frameworks such as SOF-1,<sup>23</sup> SOF-7,<sup>47</sup> HOF-3 (ref. 27) and CB6 (ref. 17) and its absorption value is comparable to that of the best performing HOF-5.<sup>24</sup>

The isosteric heat of adsorption  $Q_{\text{st}}$  calculated using the van't Hoff equation accounted for  $34.5 \text{ kJ mol}^{-1}$ , which is in agreement with the binding energy determined by DFT calculations  $\Delta E = 33.8 \text{ kJ mol}^{-1}$  (ESI†). The  $Q_{\text{st}}$  value is comparable to those of some high-performance porous crystalline materials such as SOF-1,<sup>23</sup> MIL-53,<sup>48</sup> HKUST-1 (ref. 49) and NaX zeolites,<sup>50</sup> which contain channels decorated with charged moieties.

The prominent difference in the  $\text{CO}_2$  uptake with respect to  $\text{N}_2$  under the same pressure and temperature conditions prompted us to calculate the  $\text{CO}_2/\text{N}_2$  selectivity using Ideal Adsorbed Solution Theory (IAST) applied to single-component isotherms starting from a  $15 : 85$  ( $\text{CO}_2 : \text{N}_2$ ) mixture in order to simulate the flue gas conditions (Fig. 2B). The  $\text{CO}_2/\text{N}_2$  selectivity at  $273 \text{ K}$  exhibited values in the  $320\text{--}690$  range at low



Fig. 3 CPOS-5 loaded with  $^{13}\text{C}$ -enriched  $\text{CO}_2$  (0.75 bar at RT): (A)  $^{13}\text{C}$  MAS NMR spectrum; (B) 2D  $^1\text{H}$ – $^{13}\text{C}$  HETCOR MAS NMR spectra collected at room temperature and at two distinct contact times (ct). The host– $\text{CO}_2$  cross peaks are highlighted in red. Chemical structures of TBS and DAB with labels (below). (C) Electrostatic potential projected on  $\text{CO}_2$  molecules by the DAB and TBS host moieties. The  $\text{CO}_2$  molecules were localized by DFT calculations.





pressures (<0.1 bar) and 110 at 1 bar. The CO<sub>2</sub>/N<sub>2</sub> selectivity at room temperature was 540, as calculated using the ratio of Henry's constants (ESI†). This is one of the highest selectivities shown to date in porous molecular materials.<sup>19,20,24,47</sup>

Furthermore, the isosteric heat of adsorption calculated for methane is 21.0 kJ mol<sup>-1</sup>, which is remarkably high compared to those of typical 3D frameworks with open metal sites or higher surface areas (Fig. 2C).<sup>9</sup> It is comparable to SOF-1 ( $\Delta H = 20.8$  kJ mol<sup>-1</sup>) and one of the highest reported values in microporous polymers and crystals.<sup>23,51–53</sup>

### Screwing dynamics and diffusion rate of CO<sub>2</sub> in the channels

The remarkable CO<sub>2</sub>-matrix association allowed direct observation of the gas exploring the nanochannels by NMR spectroscopy. The <sup>13</sup>C MAS NMR spectrum showed two dominant signals at  $\delta_c = 125.5$  and 125.8 ppm for <sup>13</sup>C-enriched CO<sub>2</sub>: the former identifying free CO<sub>2</sub> and the latter CO<sub>2</sub> housed in the confined space of the crystal (Fig. 3A). The downfield signal is

generated by the magnetic susceptibility effect experienced by CO<sub>2</sub> molecules embedded in the polar environment of the crystal. In the cross-polarization (CP) spectrum, the  $\delta_c = 125.8$  ppm signal is largely dominating, for the fast magnetization transfer even at room temperature, from the matrix hydrogens to the CO<sub>2</sub> carbons lying a short distance away (<1 nm),<sup>54</sup> proving that the gas is exceptionally well-retained (ESI†). This is a remarkable result given the well-known difficulty of transferring nuclear polarization in conventional systems from a solid towards a gas, because of the short gas residence time, especially at room temperature.

The surprising efficiency of gas-solid CP was exploited for the success of 2D <sup>1</sup>H–<sup>13</sup>C CP NMR, which allowed for both *in situ* spectroscopic detection of CO<sub>2</sub> and the identification of the specific interaction sites (Fig. 3B).<sup>33,55,56</sup> Even at a contact time as short as 0.5 ms and room temperature, CO<sub>2</sub> carbon nuclei receive magnetization from the amidinium hydrogens (at  $\delta_H = 7.4$  ppm), showing the short distances of the gas from these



**Fig. 4** (A) Variable temperature <sup>13</sup>C NMR spectra of the CPOS-5/<sup>13</sup>CO<sub>2</sub> sample under static conditions. The chemical shift anisotropic lineshape of solid CO<sub>2</sub>, simulated starting from the main tensor components, is reported below for comparison. (B) Inclination angle ( $\theta$ ) of the CO<sub>2</sub> main-axis with respect to the reorientation-axis z about which the 90° ( $\phi$  angle) jump occurs. (C) Screwing mechanism for CO<sub>2</sub> rototranslational dynamics in the channel: the 90° rotation angle of CO<sub>2</sub> projection onto the *ab*-plane and along the channel *c*-axis for CO<sub>2</sub> reorientation from one site to the next. (D) Overlay of the CO<sub>2</sub> positions along the channel as calculated by the DFT transition state method. (E) Helical trajectory of one CO<sub>2</sub> oxygen atom (red spheres) about the channel axis (blue bar). (F) Arrhenius plot of reorientational rates ( $k$ ) versus the inverse of temperature. (G) Calculated energy profile for the rototranslation of CO<sub>2</sub> along the channel.



species lining the channel walls. The short internuclear distances were confirmed by *ab initio* DFT calculations, which pointed out the CO<sub>2</sub> location and a distance of 3.24 Å between the N–H hydrogen and CO<sub>2</sub> carbon. The close proximity of the CO<sub>2</sub> carbon with the N–H group is imposed by the interaction of one of the CO<sub>2</sub> oxygens with the amidinium cation. Also the aromatic C–H group of DAB undergoes favorable interactions with CO<sub>2</sub> oxygens as demonstrated by the highly positive electrostatic potentials (Fig. 3C: the red regions on the oxygen atom of CO<sub>2</sub>). Actually, at a longer contact time of 1 ms the 2D NMR spectrum highlights the correlation between the aromatic C–H of DAB and CO<sub>2</sub> carbon (distance of 4.18 Å). With regard to the organosulfonate anions, DFT calculations show CO<sub>2</sub> molecules clamped in between the oxygens of two TBS sulfonate groups (Fig. 3C: the blue regions on the carbon atom of CO<sub>2</sub>), keeping the C–H hydrogens of TBS far apart from the CO<sub>2</sub> carbon (distance of 5.75 Å), consistent with the absence of NMR cross-correlations. From the combined NMR and modelling approach, we can conclude that the charge distribution of the host complements the charges on the CO<sub>2</sub> molecule and determines the stabilization of the gas in the channels.

Remarkably, multiple aspects of the dynamics of the CO<sub>2</sub> molecules in the channels came to light from variable temperature <sup>13</sup>C NMR spectra recorded in the static mode (Fig. 4A). Actually, the spectra depicted axial chemical shift anisotropy (CSA) restricted to a fraction of the total spectral width of solid CO<sub>2</sub> in the bulk.<sup>57,58</sup> From the NMR lineshape simulation we were able to determine the mechanism of motion, the dynamics and the inclination of CO<sub>2</sub> molecules in the channels. The lineshape at room temperature describes a rapid 4-site re-orientation of CO<sub>2</sub> molecules by a 90° jump mechanism (Fig. 4B). Nonetheless, the CO<sub>2</sub> main axis maintains an inclination angle equal to  $\theta = 56.7^\circ$  with respect to the rotation axis. By lowering the temperature to 215 K, the inclination increases to 58.1° due to the increased loading and the larger contribution of CO<sub>2</sub>–CO<sub>2</sub> interactions (ESI†). Notably, by the spectral analysis at variable temperature, it was possible to measure the 90° jump rates (*k*) that are higher than 10<sup>6</sup> Hz at room temperature and still over 10<sup>5</sup> Hz at 215 K.

To understand in-depth the NMR determinations of the gas diffusing in the tight-fitting channels, molecular dynamics combined with *ab initio* calculations provided a model for the energy profile and interaction sites perceived by CO<sub>2</sub> travelling in the channels. Each site accommodates a CO<sub>2</sub> molecule at a time and imparts to the CO<sub>2</sub> molecule an average inclination angle  $\theta$  of 57°–58°, in perfect agreement ( $\pm 0.7^\circ$ ) with the NMR data (ESI†). Regarding the fast 90° re-orientation, it cannot occur within the same site without a very high energy cost, for steric and electrostatic reasons. In contrast, the next site along the channel (symmetry related by a four-fold screw axis, 4<sub>1</sub>) can host CO<sub>2</sub> after its 90° turn (Fig. 4C). Thus, CO<sub>2</sub> reorientation, observed by NMR, implies rotation correlated with sliding motion to the next site with the rate of 10<sup>6</sup> steps per second (Fig. 4D). Indeed, the unique feature of this diffusion phenomenon is the concerted roto-translation mechanism (*i.e.* simultaneous rotation and translation along the *c* axis) occurring with extremely rapid dynamics of the confined CO<sub>2</sub>, which follows a helical trajectory by a screwing path along a single file (Fig. 4E). From the NMR jump

rates (*k*) vs. temperature, an Arrhenius plot can be derived, establishing a small activation energy of 2.9 kcal mol<sup>−1</sup> (Fig. 4F), in excellent agreement with the energy barrier of 2.1 kcal mol<sup>−1</sup> for a diffusion step of CO<sub>2</sub> molecules, during the helical transport, as calculated by DFT analysis (Fig. 4G).

## Conclusions

Multifunctional tetra-anionic and di-cationic molecules combine together to form frameworks with the property of being permanently porous. The crystalline framework was designed by the use of intrinsically rigid molecules as constructive elements and ionic interactions between sulfonate anions and amidinium cations. The ionic patterning of the crystalline channels allows the establishment of high interactions of 35 kJ mol<sup>−1</sup>, ideal for CO<sub>2</sub> capture/release cycles. This high energy value stably retains the quadrupolar molecules of CO<sub>2</sub>. Such CO<sub>2</sub> sequestration was so efficient that the CO<sub>2</sub>/N<sub>2</sub> selectivity values at low pressures (<0.1 bar) range from 320 to 690.

The interaction strength of the CO<sub>2</sub> with the host was recognized by the striking capability of CO<sub>2</sub> to receive nuclear polarization from the framework at room temperature and to develop the anisotropic patterns in the static <sup>13</sup>C NMR spectra. A unique screwing mechanism of CO<sub>2</sub>, which follows the helical track of the charges on the channel walls, was observed. The rate of a single step (>10<sup>6</sup> steps per second) in the diffusion trajectory along the tunnel could be unconventionally established by NMR. The orientational and dynamical features are enforced by the computational models of gas–framework interactions. The contribution of the charges patterning the channel walls is of paramount importance for directing the association of a gas, such as CO<sub>2</sub>, with the host. These pieces of evidence encourage the use of multiple charge molecules of opposite polarity to spontaneously self-organize into functional materials endowed with high thermal stability and exploit the rich modularity of polyanions and polycations to design structures tailored to a specific gas for capture and sequestration.

Moreover, the fabrication of artificial charged channels in crystalline materials with a high degree of order and uniformity as well as a controllable structure promotes the understanding of the role of electrostatic interactions in the transport properties of transmembrane channels of biological systems.<sup>45,59</sup> Indeed, the CO<sub>2</sub> diffusion in helical narrow channels, such as in aquaporins, plays a strategic role in the selection and concentration of gases in plant cells.<sup>59</sup> Our newly fabricated crystalline channels share helical and sub-nanometric structural features, and transport properties in a single file through the polar environment, thus achieving high selectivity.

## Conflicts of interest

There are no conflicts to declare.

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