



Cite this: RSC Adv., 2019, 9, 37119

Fluorine ion induced phase evolution of tin-based perovskite thin films: structure and properties†

Junsheng Wu, ^{ab} Fang Fang, ^a Zhuo Zhao, ^{ab} Tong Li, ^{ab} Rizwan Ullah, ^c Zhe Lv, ^b Yanwen Zhou^{*a} and David Sawtell^{*d}

To study the effect of fluorine ions on the phase transformation of a tin-based perovskite, $\text{CsSnI}_{3-x}(\text{F})_x$ films were deposited by using thermal vacuum evaporation from a mixed powder of SnI_2 , SnF_2 and CsI , followed by rapid vacuum annealing. The color evolution, structure, and properties of $\text{CsSnI}_{3-x}\text{F}_x$ films aged in air were observed and analyzed. The results showed that the colors of the films changed from black to yellow, and finally presented as black again over time; the unstable $\text{B}-\gamma\text{-CsSnI}_{3-x}\text{F}_x$ phase transformed into the $\text{Y}-\text{CsSnI}_{3-x}\text{F}_x$ phase, which is then recombined into the $\text{Cs}_2\text{SnI}_{6-x}\text{F}_x$ phase with the generation of SnO_2 in air. Fluorine dopant inhibited the oxidation process. The postponement of the phase transformation is due to the stronger bonds between F and Sn than that between I and Sn. The color changing process of the $\text{CsSnI}_{3-x}\text{F}_x$ films slowed that the hole concentrations increased and the resistivities decreased with the increase of the F dopant ratio. With the addition of SnF_2 , light harvesting within the visible light region was significantly enhanced. Comparison of the optical and electrical properties of the fresh annealed $\text{CsSnI}_{3-x}\text{F}_x$ films showed that the band gaps of the aged films widened, the hole concentrations kept the same order, the hole mobilities reduced and therefore, the resistivities increased. The double layer $\text{Cs}_2\text{SnI}_{6-x}\text{F}_x$ phase also showed 'p' type semi-conductor properties, which might be due to the incomplete transition of Sn^{2+} to Sn^{4+} , i.e. Sn^{2+} provides holes as the acceptor.

Received 15th September 2019
Accepted 30th October 2019

DOI: 10.1039/c9ra07415e
rsc.li/rsc-advances

1. Introduction

Within the past decade, inorganic-organic hybrid perovskite solar cells have drawn researchers' attention for their enhanced efficiencies due to their superior long diffusion length, with a conversion efficiency of approximately 3.8% in 2009¹ to greater than 23.32% today.² Since the early reports of lead (Pb)³ and tin (Sn)⁴ based perovskite solar cells, the efficiencies of Pb based perovskite solar cells have increased to nearly 22%, which exceeds those of poly-silicon solar cells.⁵ Although most studies are focused on MAPbX_3 perovskite materials, tin-based perovskite has attracted more and more attention due to the low toxicity of tin.⁶ So far, the literature^{7,8} has reported that both Sn doped and the entirely Sn based perovskite solar cells (ASnX_3) are not as efficient as Pb-based perovskite solar cells due to two primary reasons: (1) tin is easily oxidized from Sn^{2+} to Sn^{4+} oxidation states whilst exposed to air.^{9,10} The diffusion length of

the photogenic carriers is limited because of too many p-type carriers produced by the self-doping effect within the tin-based perovskite materials; (2) SnI_2 and MAI (CH_3NH_3^+) reacts quickly and rapidly crystallizes during spin-coating process^{9,11} and therefore, it is difficult to control the crystallinity of perovskite, lead to poor coverage and uniformity of the film. To combat this stable and defect free ASnX_3 films are required.

Inorganic materials generally have higher stability than organic materials. For this reason, the studies on inorganic perovskite materials have been proposed. As the standard ASnX_3 perovskite, the schematic crystal structure of $\text{B}-\gamma\text{-CsSnI}_3$, $\text{Y}-\text{CsSnI}_3$ and Cs_2SnI_6 are shown in Fig. 1, in which the $\text{B}-\gamma\text{-CsSnI}_3$ phase is black with a three dimensional perovskite structure,¹² the $\text{Y}-\text{CsSnI}_3$ phase is yellow with a one-dimensional double-chain structure,¹³ and Cs_2SnI_6 is black with double layer perovskite structure. The unstable $\text{B}-\gamma\text{-CsSnI}_3$ phase promptly transforms to the $\text{Y}-\text{CsSnI}_3$ phase, then forms into the Cs_2SnI_6 phase in air, accompanied with the formation of SnO_2 . The compound of Cs_2SnI_6 exhibits its stability in damp air due to the stable Sn^{4+} state. The 'vacancy ordered' double layer perovskite structure is formed by the missing half of the Sn atoms located in the center of octahedral, and then reconstitutes to the discontinuous regular octahedral structure (SnI_6).²⁻¹⁴

^aSchool of Chemical Engineering, University of Science and Technology Liaoning, 114051 Liaoning, Anshan, China

^bInstitute of Surface Engineering, University of Science and Technology Liaoning, Anshan, 114051 Liaoning, China. E-mail: zhouyanwen@ustl.edu.cn

^cDepartment of Physics, Beijing Normal University, 100875 Beijing, China

^dSurface Engineering Group, Manchester Metropolitan University, Manchester M1 5GD, England, UK. E-mail: d.sawtell@mmu.ac.uk

† Electronic supplementary information (ESI) available. See DOI: [10.1039/c9ra07415e](https://doi.org/10.1039/c9ra07415e)



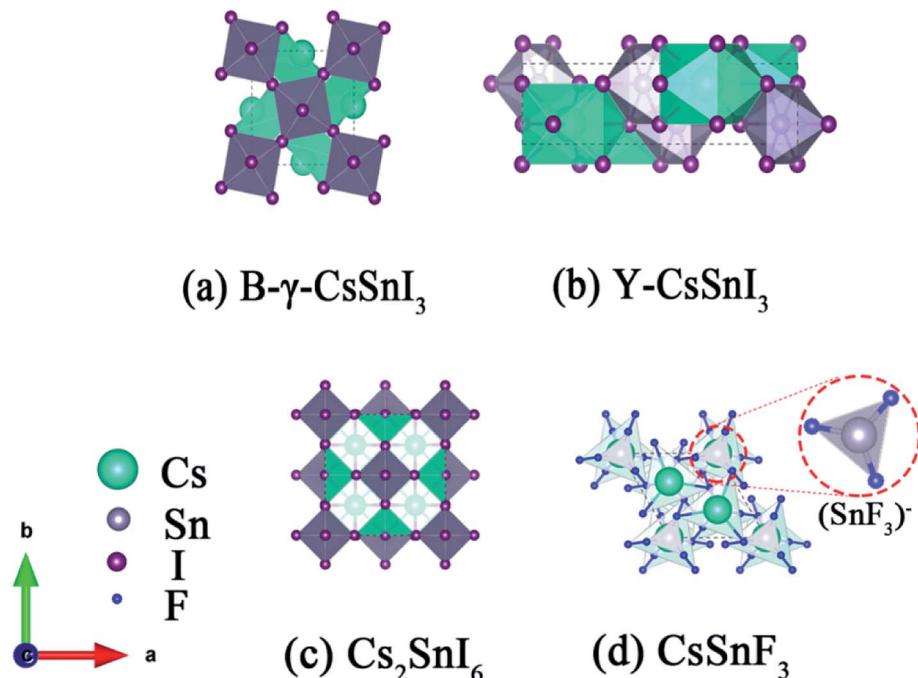


Fig. 1 Polyhedral model of (a) B- γ -CsSnI₃, (b) Y-CsSnI₃, (c) Cs₂SnI₆ and (d) CsSnF₃. SnI₆ octahedra are in gray, Cs is in green, Sn is in gray, the I anions are purple, and F is blue.

ASnX₃ perovskite structure (A = metal or NH₄ and X = halogen family) generally consists of (SnX₃)⁻ and M⁺, in which Sn²⁺ is unstable. Vilminot originally discussed evolution of the ionic-conductivity of the MSnF₃ phases (M = Na, K, Rb, Cs, NH₄, Tl) in 1985.¹⁵ For example, CsSnF₃ exhibits a crystal structure consisting of isolated (SnF₃)⁻ anionic polyhedral,¹⁶ shown in Fig. 1d. Each Sn²⁺ cation in CsSnF₃ is bonded to three fluorine atoms in a distorted triangular pyramidal coordination environment, *i.e.* the lone pair of 5s electrons of tin(II) and three pairs of electrons, shared unequally with the fluorine ions, occupied on the most stable configuration sp³ hybridization orbits.¹⁷

Researchers and experiments have shown that the tolerance factors of the stable perovskite materials should be between 0.78–1.05,¹⁸ in which the tolerance factor, $t = (R_A + R_X)/\sqrt{2}(R_B + R_X)$. In the formula, R_A is the radius of monovalent cation, R_B , the radius of divalent metal cations and R_X , the radius of the halide ion. If there are multiple ions in the A and/or B positions, the average radius should be taken. The perovskite structure tends to be stable when t is close to 1. By calculation, the ' t ' of CsSnI_{3-x}F_x is within 0.8734–0.9587, and t gets larger as the F ratio increases. That means that the perovskite structure of CsSnF₃ is more stable than that of CsSnI₃, and F doped in CsSnI₃ may delay the phase transformation of CsSnI₃.

The crystal parameters of the CsSnI(F)₃ phases are shown in Table 1. The Sn–I bond length is stretched whilst the B- γ -CsSnI₃ phase transformed into Y-CsSnI₃, and then shortened along with the initial formation of the stable Cs₂SnI₆ during the process of tin oxidation. There are many factors affecting bond length and bond energy, such as atomic radius, distance

between nuclei, repulsive force between lone pair electrons, feedback bond and so on. In ionic compound, ion radius is the main contributor to bond length. Therefore, the Sn–F bond length in CsSnF₃ is less than the bond length of Sn–I in CsSnI₃, the difference between them is 1.1321 Å due to the much smaller ionic radii of F (1.33 Å) compared to I (2.20 Å). Generally, the order of stability of the halide complexes of tin(II) is F > Cl > Br > I,¹⁷ tin(II) preferentially bonds with fluoride ion. Hence, SnF₂ was introduced as the doping source to improve the stability of tin-based perovskite materials. Meanwhile, as an antioxidant, fluoride ion inhibits the oxidation process of Sn²⁺,^{10,19} the fluorine also reduces the densities of the materials, which improves their performance when used as a photovoltaic material.²⁰

Even though the main method to prepare CsSnI₃ films is still the one-step solution based process, it is difficult to produce dense, pinhole free film due to the rapid crystallization of tin-based perovskite.^{9,20–22} The films prepared by this method are very sensitive to film formation conditions, such as annealing temperature,^{23,24} solution concentration,^{25,26} precursor solution composition^{27,28} and solvent selection.^{29–31} Due to the evaporation of the solvent and the volatilization of the materials during the process of annealing, the crystal tends to easily aggregate and shrink, and the morphology of the films are mainly cluster-like and needle-like, so it is easy to cause the devices' efficiencies to be uneven. The thermal vacuum evaporation technique is an effective approach to prepare high coverage homogeneous thin films and has been widely used in lead-based perovskites.^{32,33} Therefore, tin-based perovskite films should be possible to be prepared by this technique.^{10,34} Here, the mixture

Table 1 Crystallographic refinement details for the tin based perovskite polymorphs

Sample	Type	Bond length (Å)	Lattice constant (Å per °)			Ref.	
B- γ -CsSnI ₃	Cs-I	4.09495	<i>a</i>	8.6885	$\alpha = \beta = \gamma$	90	38
			<i>b</i>	12.3775			
Y-CsSnI ₃	Sn-I	3.1685	<i>c</i>	8.6384			
	Cs-I	4.0680	<i>a</i>	10.350	$\alpha = \beta = \gamma$	90	38
Cs ₂ SnI ₆	Sn-I	3.2475	<i>b</i>	4.7632			
			<i>c</i>	17.684			
CsSnF ₃	Cs-F	3.06629	<i>a</i> = <i>b</i>	7.18763	$\alpha = \beta$	90	43
	Sn-F	2.11543	<i>c</i>	16.08594	γ	120	
CsSnI _{3-x} F _x (non-optimized)	Cs-I	4.0215	<i>a</i>	8.688	$\alpha = \beta = \gamma$	90	
	Sn-I	3.1143	<i>b</i>	12.378			
	Cs-F	3.3794	<i>c</i>	8.6430			
	Sn-F	3.2934					

powder of SnI₂ and CsI were evaporated onto the glass slide substrates by thermal vacuum evaporation method to form fully covered, dense, pinhole free CsSnI₃ perovskite film. Also, by adding SnF₂ powder into the mixture, the CsSnI_{3-x}F_x films were prepared as well. The evolution of the color, structure and properties of the annealed freshly and annealed aged CsSnI_{3-x}F_x films over time was observed, measured and analyzed to explore the effect of F doping.

2. Experimental details

2.1. Sample preparation

Cesium iodide (99.9% CsI), tin(II) fluoride (99.99% SnF₂) and tin iodide (99.9% SnI₂) were produced by Ying Kou You Xuan Trade Co., Ltd. China. Acetone (99.7%) and ethanol (99.7%) were purchased from Sinopharm Chemical Reagent Co., Ltd. China. Deionized water was filtered in the laboratory. All the reagents were of analytical grade and used as received.

The process of sample preparation was shown in Fig. S1.† Weighted SnI₂, SnF₂ and CsI powder was mixed in a mortar and placed in a tungsten boat of size 50 × 15 × 2 mm, refer to Table 2. The AC power (50 Hz frequency) was applied to the tungsten boat through two connected electrodes, the glass slides, CAT. no. 7101 with a size of 25.4 × 76.2 mm, were ultrasonically cleaned in acetone for 900 s, diluted by deionized water and ethanol before being loaded into the DM 450C vacuum chamber. The glass slide was held above the tungsten boat at a separation of 150 mm. The chamber was pumped down to 2 × 10⁻³ Pa and the SnI₂, SnF₂ and CsI was evaporated at a voltage

of 70 V and current of 140 A. The evaporation process referred to Tong.³⁵ The films were finally annealed at 210 °C for 4 min in argon gas by the Rapid Thermal Processor of RTP-500V.

2.2. Measurement techniques

The thicknesses of the films were measured by using a KLA-Tencor Alpha-step D-100 type profilometer on a step created on the films by masking the glass substrates. The phases of the CsSnI_{3-x}F_x films were measured by X'Pert powder X-ray diffractometer (XRD) in glancing angle scanning mode at 0.5° incident angle with Cu K α X-ray from 10° to 80°, and analyzed by High-Score software.³⁶ The electrical and optical properties of the CsSnI_{3-x}F_x films were measured by a HALL 8800 Hall Effect Measurement device and a CARY 5000 UV-Vis-NIR spectrometer over the wavelength range of 300 to 1000 nm, respectively. AFM tests were performed in ambient conditions at room temperature with scanning probe microscopy (Bruker, Multimode 8 with controller V). A Pt/Ir-coated tip on a Si cantilever (tip radius of 20 nm, force constant of 2.8 N m⁻¹ and a resonant frequency of 75 kHz) was used to characterize the topography of the films. The typical tip-scanning velocity was 2 μ m s⁻¹.

3. Results & discussion

3.1. Morphological structure

To demonstrate the macro-evolution process of the fluorine doped cesium tin iodine (CsSnI_{3-x}F_x) films, photographs of the colors of the freshly annealed and aged CsSnI_{3-x}F_x films exposed to ambient air were shown in Fig. 2a. With increasing

Table 2 The chemical composition of CsSnI_{3-x}F_x thin films

Sample	CsI (mol)	SnI ₂ (mol)	SnF ₂ (mol)	Proportion
CsSnI ₃	0.0025	0.0025	0	100 : 100 : 0
CsSnI _{2.88} F _{0.11}	0.0025	0.0024	0.000094	100 : 96.25 : 3.75
CsSnI _{2.78} F _{0.22}	0.0025	0.0023	0.000190	100 : 92.5 : 7.5
CsSnI _{2.67} F _{0.33}	0.0025	0.0022	0.000280	100 : 88.75 : 11.25



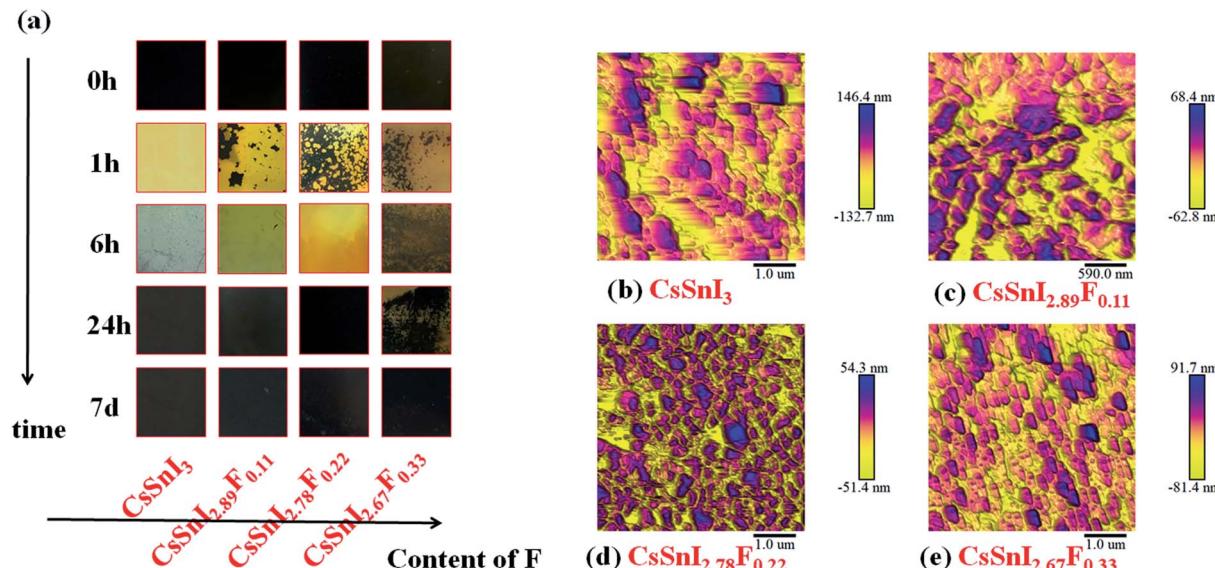


Fig. 2 (a) Macroscopic images of the evolution in the cesium tin iodine doped fluorine ($\text{CsSnI}_{3-x}\text{F}_x$) films. (b–e) AFM images of the cesium tin iodine doped fluorine ($\text{CsSnI}_{3-x}\text{F}_x$) with different fluorine content, (b–e) correspond to CsSnI_3 , $\text{CsSnI}_{2.89}\text{F}_{0.11}$, $\text{CsSnI}_{2.78}\text{F}_{0.22}$, $\text{CsSnI}_{2.67}\text{F}_{0.33}$, respectively.

time, the colors of the films changed from black to yellow, and finally presented as black again. These colors served as a good indication of the oxidation progress of $\text{B-}\gamma/\text{Y-CsSnI}_3$ to Cs_2SnI_6 in air.³⁷ It is evident from a comparison of the photographs in Fig. 2a that doping F slowed down the transformation process of the perovskite phases, since the color of F doped films were predominantly yellow after six hours exposure in air, whereas that of the CsSnI_3 film without fluorine dopant had almost completely blackened. The morphologies of the corresponding $\text{CsSnI}_{3-x}\text{F}_x$ films were examined by AFM, respectively (see Fig. 2b–e). Strikingly, no pinhole appeared in entire scope in the $\text{CsSnI}_{3-x}\text{F}_x$ films. Furthermore, the grain sizes became finer as the amount of doped fluorine increased.

3.2. Phase structure

Corroborating evidence for the process of phase transition in $\text{CsSnI}_{3-x}\text{F}_x$ films exposed in air was provided by X-ray

diffraction (XRD), shown in Fig. 3a and b, which were consistent with the variations in morphology. Firstly, for the annealed $\text{CsSnI}_{3-x}\text{F}_x$ films (see Fig. 3a), the peaks at 25.11° , 29.08° and 41.45° were attributed to $\text{B-}\gamma\text{-CsSnI}_3$ phase (022), (220) and (224) planes (Ref. Code: 01-043-1162), respectively. Meanwhile, the well-matched peaks of 27.61° and 39.44° , assigned to CsI phase (110) and (200) planes (Ref. Code: 01-077-2185), were observed. However, the annealed XRD patterns also showed Y-CsSnI_3 phase (111) and (121) planes (Ref. Code: 01-071-1898). The $\text{B-}\gamma\text{-CsSnI}_3$ peaks of the annealed $\text{CsSnI}_{3-x}\text{F}_x$ films became more distinct with the increase of F dopant ratios whilst the Y-CsSnI_3 peaks became weaker. Although stoichiometric ratios were strictly controlled, the sharp and high intensity peaks of CsI also appeared in the XRD patterns, this can be attributed to incomplete reaction with SnI_2 , because SnI_2 (vaporization point 993.15 K in ambient air) is easier to evaporate and decompose (unstable Sn^{2+}) than CsI .

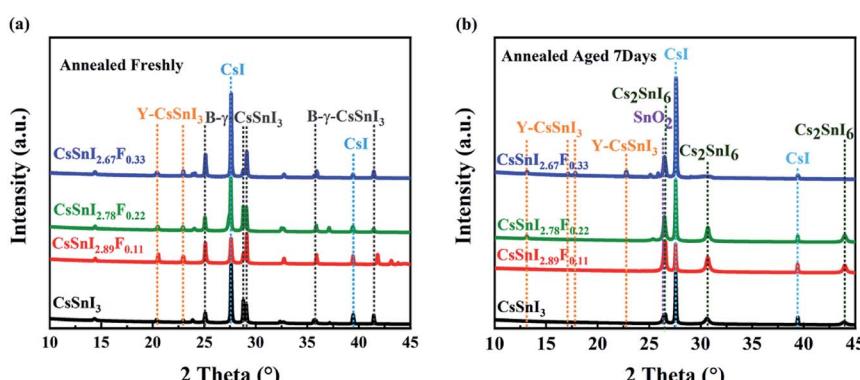


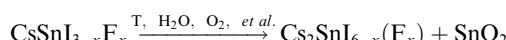
Fig. 3 Evolution of XRD patterns of $\text{CsSnI}_{3-x}\text{F}_x$ films with time. (a) XRD patterns of the $\text{CsSnI}_{3-x}\text{F}_x$ films annealed freshly, (b) XRD patterns of the annealed $\text{CsSnI}_{3-x}\text{F}_x$ films exposed in air for 7 days.



(vaporization point 1553 K in ambient air) in the process of thermal evaporation. The variations of the XRD patterns show that fluorine doping obstructed the transformation from B- γ -CsSnI₃ to Y-CsSnI₃ perovskite phases.

After aging seven days in air, the films returned to a light black color. This phenomenon is consistent with the trend of Qiu's experiment.³⁷ The new diffraction peaks (see Fig. 3b) at 26.44°, 30.61°, and 43.91° were attributed to the (222), (004) and (044) orientations of Cs₂SnI₆ phase (Ref. Code: 00-051-0466). However, the aged XRD patterns showed (120), (121), (130) and (210) planes of the Y-CsSnI₃ phase. Furthermore, several diffraction peaks at 27.57°, 39.41°, and 26.58° assigned to CsI and SnO₂ (Ref. Code: 00-077-0448) were observed as well. Again, the Y-CsSnI₃ phase was much obvious with the increase of F dopant ratios, which meant F doping delayed the phase transformation, and Y phase existed in the highly F doped films even after aging for seven days.

Based on these findings *via* the above XRD data, partial CsI, SnI₂, and SnO₂ peaks were observed during the evolution process of the CsSnI_{3-x}F_x films, the chemical reactions involved may be as follows:



Therefore, it is clear that this process was due to oxidation of Sn in the compounds.⁶ Sn²⁺ ions are quite sensitive to external oxygen, especially in humid environments, which can be oxidized to more stable Sn⁴⁺ analogues. This oxidation process may fundamentally destroy the charge neutrality of the CsSnI₃ perovskite structure and lead to phase transition. The Sn-F bonds form when fluorine replaces iodide in the lattice. The Sn-F bond length is shorter than that of Sn-I, resulting in a more stable crystal structure and weakening the role of oxygen, refer to Fig. S2.† In short, doping with fluorine slowed down the oxidation progress of the Sn²⁺ in CsSnI₃ into Sn⁴⁺. The B- γ -CsSnI_{3-x}(F)_x phase eventually transferred into Cs₂SnI_{6-x}(F)_x and was accompanied by the formation of SnO₂ over time, but the formation processes are delayed by the additional F dopant. The remnant Y-CsSnI₃ phase proved not only to delay the phase transition, but also the existence of Sn²⁺. The hybrid of Sn⁴⁺ with Sn²⁺ resulted in the Cs₂SnI_{6-x}(F)_x films showing their p type property.

3.3. Electrical properties of CsSnI_{3-x}F_x films

The electrical properties and thickness of the CsSnI_{3-x}F_x films at room temperature were measured by a HALL 8800

Hall Effect Measurement device and KLA-Tencor Alpha-step D-100 type profilometer respectively. The results were shown in Table 3 for the freshly annealed and Table 4 for those aged seven days after annealing. In this case, the thicknesses of the films were used to calculate the resistivity of the films. Both the annealed and aged CsSnI_{3-x}F_x films were p type semiconductors with the holes provided by the Sn vacancies. Intrinsic defects such as Sn vacancies in the ternary Cs-Sn-I system gave rise to p-type conductivity²⁰ and DFT calculations have shown that the formation energy of V_{Sn} defects was the lowest among all defects.³⁸ In Table 3, the freshly annealed CsSnI₃ film exhibited the carrier densities of $\sim 10^{14}$ cm⁻³, and the carrier density of CsSnI_{3-x}F_x films increased to $\sim 10^{16}$ cm⁻³ with the increase of SnF₂ content. It's also worth noting that the resistivities of CsSnI_{3-x}F_x films were 1-2 orders of magnitude smaller than that of undoped film. The reduction of the Sn vacancy concentrations can be attributed to the strong bonding energy between F and Sn.

After aging in air for seven days, the electrical properties of the CsSnI_{3-x}F_x films kept the same order. As described in Section 3.2 Phase structure, the main composition of the CsSnI_{3-x}F_x films was Cs₂SnI₆. This is accompanied by the generation of a mass of Sn vacancies. The carrier densities of the aged CsSnI_{3-x}F_x films were almost at the same level in comparison to those of the annealed CsSnI_{3-x}F_x films. The resistivities of CsSnI_{3-x}F_x films were still relatively small. After aging for seven days, the phase transformation occurred, the oxidation of tin was completed, and the double layer 'vacancy ordered' phase Cs₂SnI_{6-x}F_x formed. As the results, the Sn vacancies of the aged films were high, and their carrier density increased sharply. However, the rate of increase of the charge carriers of the aged films decreased with the increase of SnF₂ content. The phases were stable due to F doping, the creation processes of the Sn vacancies were delayed and therefore, the changes of the electrical properties were slower. This fully illustrates SnF₂ hindering the oxidation process of Sn²⁺, which blocked the formation of Cs₂SnI₆.

Since the carrier (hole or electron) concentration of semiconductor depends on the inherent defect concentration, the control of carrier concentration is a necessary condition to optimize the performance of solar cells. It works better as the hole-transport material (HTM) when hole concentration and conductivity are high, but it may work better as the light absorber material (LAM) when hole concentration and conductivity are moderate.³⁹ From the data of electrical properties, by tuning the F doping amount, hole concentration and

Table 3 The electrical properties and thickness of CsSnI_{3-x}F_x thin films annealed freshly

Sample	R_s ($\Omega \square^{-1}$)	ρ ($\Omega \text{ cm}$)	N/P (cm^{-3})	μ ($\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1}$)	Thickness (nm)	Type
CsSnI ₃	2.14×10^5	22.71	6.12×10^{14}	449.31	1060	P
CsSnI _{2.88} F _{0.11}	1.93×10^4	4.65	5.84×10^{15}	229.55	2410	P
CsSnI _{2.78} F _{0.22}	4.99×10^3	1.44	6.01×10^{16}	226.67	2880	P
CsSnI _{2.67} F _{0.33}	1.49×10^3	0.394	7.29×10^{16}	217.34	2650	P



Table 4 The electrical properties and thickness of $\text{CsSnI}_{3-x}\text{F}_x$ thin films annealed aged 7 days in air

Sample	R_s ($\Omega \square^{-1}$)	ρ ($\Omega \text{ cm}$)	N/P (cm^{-3})	μ ($\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1}$)	Thickness (nm)	Type
CsSnI_3	3.62×10^5	38.75	1.42×10^{15}	114.96	1060	P
$\text{CsSnI}_{2.88}\text{F}_{0.11}$	4.07×10^4	9.79	8.55×10^{15}	74.55	2410	P
$\text{CsSnI}_{2.78}\text{F}_{0.22}$	1.44×10^4	4.16	6.21×10^{15}	241.46	2880	P
$\text{CsSnI}_{2.67}\text{F}_{0.33}$	3.54×10^5	93.75	1.45×10^{16}	4.59	2650	P

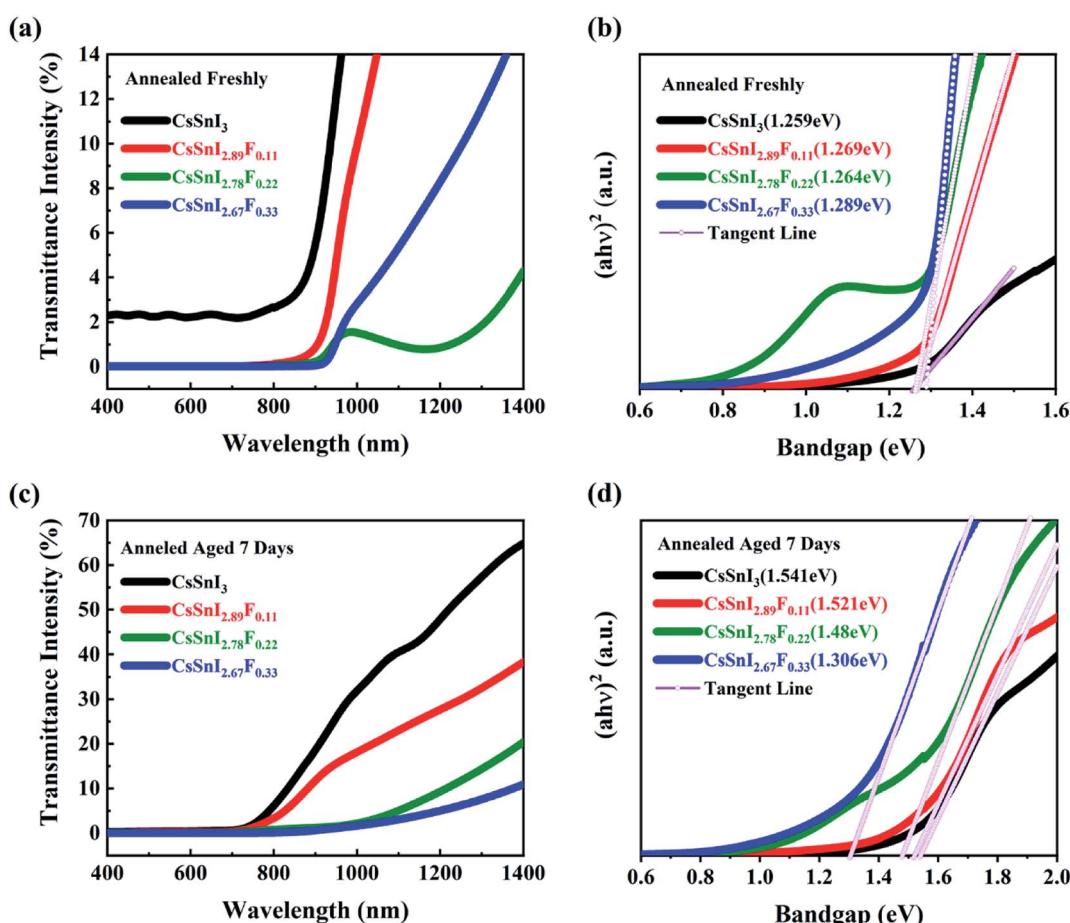
conductivity can be controlled, so that the film can be selectively arranged as HTM and LAM.

3.4. Optical properties

The absorptive spectra of the annealed and aged $\text{CsSnI}_{3-x}\text{F}_x$ films within the range of 400–1400 nm were presented in Fig. 4a and c, respectively. Within the UV range, the absorptivity of the films was almost the same. The absorption edges of the annealed and aged $\text{CsSnI}_{3-x}\text{F}_x$ films exhibited obvious differences. With the addition of SnF_2 , the light harvesting in the visible light region of the solar spectrum of the $\text{CsSnI}_{3-x}\text{F}_x$ films was significantly enhanced. These differences indicate the difference of the band gap between samples. The spectra of the absorption *via* optical band gaps of the annealed and aged films

were shown in Fig. 4b and d, which were calculated by the formula of $\alpha h\theta = A(h\theta - E_g)^{1/2}$. The Tauc plot⁴⁰ is used to evaluate the optical energy gap. The bandgaps of all the $\text{CsSnI}_{3-x}\text{F}_x$ films are between 1.25 eV and 1.3 eV as shown in Fig. 4b which was consistent with those previously reported.⁴¹

The optical spectra of $\text{CsSnI}_{3-x}\text{F}_x$ films under exposure to ambient air for seven days are shown in Fig. 4c and d. The absorption edges of all the $\text{CsSnI}_{3-x}\text{F}_x$ films exhibited blue shift. All the bandgaps for the F doped films got smaller, and was close to 1.48 eV derived in Qiu's report.³⁷ This is a result of SnF_2 eliminating the formation of weak unidentified reflections due to the Y- CsSnI_3 phase.⁴² After the $\text{CsSnI}_{3-x}\text{F}_x$ films under exposure to ambient air for seven days, the Y- CsSnI_3 phase only was shown in the highly F doped films, which the fully

Fig. 4 Optical spectra of $\text{CsSnI}_{3-x}\text{F}_x$ films: (a and b) annealed freshly; (c and d) under exposure to ambient air for 7 days.

illustrates the role of SnF_2 . The transmittance of the $\text{Cs}_2\text{SnI}(\text{F})_6$ films within visible wavelength were strongly affected by doping F ions due to the decomposition of $\text{CsSnI}_{3-x}\text{F}_x$. It was further shown that SnF_2 blocked the formation of Cs_2SnI_6 by the changes of bandgap. Note that the sunlight absorption in the red and near-infrared regions has been a challenge for solar cells, and the excellent optical properties of $\text{CsSnI}_{3-x}\text{F}_x$ films made up for this defect.

4. Conclusion

CsSnI_3 transformed into stable double layer perovskite Cs_2SnI_6 phase in air was investigated for potential solar cell applications. By conducting a series of experiments, it was determined that the phase transformation in air can be slowed down by doping F ions into CsSnI_3 . The mechanism of phase transition delay is that Sn^{2+} preferentially bonds with F^- , and thus the process of Sn^{2+} in CsSnI_3 to be oxidized into Sn^{4+} was slowed down. The $\text{B}-\gamma\text{-CsSnI}(\text{F})_3$ phase eventually formed $\text{Cs}_2\text{SnI}(\text{F})_6$ accompanied by the formation of the SnO_2 phase over time, but the transition processes are delayed by the additional F dopant.

Furthermore, SnF_2 improved the carrier concentration and conductivity of the film, in which the resistivities of $\text{CsSnI}_{3-x}\text{F}_x$ films were less than $10 \Omega \text{ cm}$ except $\text{CsSnI}_{2.67}\text{F}_{0.33}$. The $\text{CsSnI}_{3-x}\text{F}_x$ films also absorbed more visible and infrared light with increased F doping than the pure CsSnI_3 film. It was further shown that SnF_2 blocked the formation of Cs_2SnI_6 by comparing the changes of bandgaps. The usage of $\text{CsSnI}_{3-x}\text{F}_x$ films as HTM or LAM in solar cells is envisioned to be a promising method of improving efficiency.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

The work at University of Science and Technology Liaoning is supported by National Natural Science Foundation of China (No. 51672119, 51774180, 51972155) and Royal Society International Exchanges 2017(IEC\NSFC\170535). Acknowledges go to the support from the Project of Young Teachers in the Higher Education Institutions of USTL (No. 2017QN13) and the Innovation Team Project of USTL (No. 2017TD01). Thanks go to the colleagues at Manchester Metropolitan University for their co-operation in this study.

References

- 1 A. Kojima, K. Teshima, Y. Shirai and T. Miyasaka, *J. Am. Chem. Soc.*, 2009, **131**, 6050–6051.
- 2 Q. Jiang, Y. Zhao, X. Zhang, X. Yang, Y. Chen, Z. Chu, Q. Ye, X. Li, Z. Yin and J. You, *Nat. Photonics*, 2019, **13**, 460–466.
- 3 J. H. Im, H. S. Kim and N. G. Park, *APL Mater.*, 2014, **2**, 591–713.
- 4 I. Chung, B. Lee, J. He, R. P. Chang and M. G. Kanatzidis, *Nature*, 2012, **485**, 486–489.
- 5 N. R. E. Laboratory, Best Research-Cell Efficiencies chart, http://www.nrel.gov/ncpv/images/efficiency_chart.jpg.
- 6 C. C. Stoumpos, C. D. Malliakas and M. G. Kanatzidis, *Inorg. Chem.*, 2013, **52**, 9019–9038.
- 7 P. P. Boix, S. Agarwala, T. M. Koh, N. Mathews and S. G. Mhaisalkar, *J. Phys. Chem. Lett.*, 2015, **6**, 898–907.
- 8 P. Pyykko, *Chem. Rev.*, 1988, **88**, 563–594.
- 9 N. K. Noel, S. D. Stranks, A. Abate, C. Wehrenfennig, S. Guarnera, A. A. Haghhighirad, A. Sadhanala, G. E. Eperon, S. K. Pathak and M. B. Johnston, *Energy Environ. Sci.*, 2014, **7**, 3061–3068.
- 10 D. Moghe, L. Wang, C. J. Traverse, A. Redoute, M. Sponseller, P. R. Brown, V. Bulović and R. R. Lunt, *Nano Energy*, 2016, **28**, 469–474.
- 11 L. Glasser, *Inorg. Chem.*, 2012, **51**, 2420.
- 12 K. Yamada, S. Funabiki, H. Horimoto, T. Matsui, T. Okuda and S. Ichiba, *Chem. Lett.*, 1991, **20**, 801–804.
- 13 P. Mauersberger and F. Huber, *Acta Crystallogr., Sect. B: Struct. Crystallogr. Cryst. Chem.*, 1980, **36**, 683–684.
- 14 A. E. Maughan, A. M. Ganose, M. M. Bordelon, E. M. Miller, D. O. Scanlon and J. R. Neilson, *J. Am. Chem. Soc.*, 2016, **138**, 8453–8464.
- 15 S. Vilminot, W. Granier, A. Soufiane, L. Cot, J. M. Letoffe and P. Claudy, *Rev. Chim. Miner.*, 1985, **22**, 125–133.
- 16 T. Thao Tran and P. Shiv Halasyamani, *J. Solid State Chem.*, 2014, **210**, 213–218.
- 17 W. B. Schaap, J. A. Davis and W. H. Nebergall, *J. Am. Chem. Soc.*, 1954, **76**, 5226–5229.
- 18 C. A. Randall, A. S. Bhalla, T. R. Shrout and L. E. Cross, *J. Mater. Res.*, 1990, **5**, 829–834.
- 19 T. M. Koh, T. Krishnamoorthy, N. Yantara, C. Shi, W. L. Leong, P. P. Boix, A. C. Grimsdale, S. G. Mhaisalkar and N. Mathews, *J. Mater. Chem. A*, 2015, **3**, 14996–15000.
- 20 M. H. Kumar, S. Dharani, W. L. Leong, P. P. Boix, R. R. Prabhakar, T. Baikie, C. Shi, H. Ding, R. Ramesh and M. Asta, *Adv. Mater.*, 2014, **26**, 7122–7127.
- 21 F. Hao, C. C. Stoumpos, P. Guo, N. Zhou, T. J. Marks, R. P. H. Chang and M. G. Kanatzidis, *J. Am. Chem. Soc.*, 2015, **137**, 11445–11452.
- 22 S. J. Lee, S. S. Shin, Y. C. Kim, D. Kim, T. K. Ahn, J. H. Noh, J. Seo and S. I. Seok, *J. Am. Chem. Soc.*, 2016, **138**, 3974–3977.
- 23 A. Dualeh, N. Tétreault, T. Moehl, P. Gao, M. K. Nazeeruddin and M. Grätzel, *Adv. Funct. Mater.*, 2014, **24**, 3250–3258.
- 24 G. E. Eperon, V. M. Burlakov, P. Docampo, A. Goriely and H. J. Snaith, *Adv. Funct. Mater.*, 2014, **24**, 151–157.
- 25 T. Leijtens, B. Lauber, G. E. Eperon, S. D. Stranks and H. J. Snaith, *J. Phys. Chem. Lett.*, 2014, **5**, 1096–1102.
- 26 W. Qi, Y. Shao, Q. Dong, Z. Xiao, Y. Yuan and J. Huang, *Energy Environ. Sci.*, 2014, **7**, 2359–2365.
- 27 L. Po-Wei, L. Chien-Yi, C. Chu-Chen, Z. Fan, S. T. Williams, X. Xu-Kai, L. Jiangjen and J. K.-Y. Alex, *Adv. Mater.*, 2014, **26**, 3748–3754.
- 28 Y. Zhao and K. Zhu, *J. Phys. Chem. C*, 2014, 118.
- 29 D. Shen, X. Yu, X. Cai, M. Peng, Y. Ma, X. Su, L. Xiao and D. Zou, *J. Mater. Chem. A*, 2014, **2**, 20454–20461.
- 30 S. Paek, N. Cho, H. Choi, H. Jeong, S. L. Jin, J. Y. Hwang, J. K. Lee and J. Ko, *J. Phys. Chem. C*, 2014, **118**, 25899–25905.



31 K. Hak-Beom, C. Hyosung, J. Jaeki, K. Seongbeom, W. Bright, S. Seyeong and K. J. Young, *Nanoscale*, 2014, **6**, 6679–6683.

32 M. Liu, M. B. Johnston and H. J. Snaith, *Nature*, 2013, **501**, 395.

33 Q. Chen, H. Zhou, Z. Hong, S. Luo, H. S. Duan, H. H. Wang, Y. Liu, G. Li and Y. Yang, *J. Am. Chem. Soc.*, 2014, **136**, 622–625.

34 B. Wu, Y. Zhou, G. Xing, Q. Xu, H. F. Garces, A. Solanki, T. W. Goh, N. P. Padture and T. C. Sum, *Adv. Funct. Mater.*, 2017, **27**, 1604818.

35 X. R. Tong, Z. Zhao, J. S. Wu, Y. W. Zhou, P. Kelly and P. Gao, *Nanosci. Nanotechnol. Lett.*, 2018, **10**, 329–336.

36 T. Degen, M. Sadki, E. Bron, U. Konig and G. Nenert, *Powder Diffr.*, 2014, **29**, S13–S18.

37 X. Qiu, B. Cao, Y. Shuai, X. Chen, Z. Qiu, Y. Jiang, Y. Qian, H. Wang, H. Zeng and L. Jian, *Sol. Energy Mater. Sol. Cells*, 2017, **159**, 227–234.

38 I. Chung, J.-H. Song, J. Im, J. Androulakis, C. D. Malliakas, H. Li, A. J. Freeman, J. T. Kenney and M. G. Kanatzidis, *J. Am. Chem. Soc.*, 2012, **134**, 8579–8587.

39 Z. Xiao, Y. Zhou, H. Hosono and T. Kamiya, *Phys. Chem. Chem. Phys.*, 2015, **17**, 18900–18903.

40 J. D. Neufville and H. Rockstad, 1974.

41 L. Byunghong, C. C. Stoumpos, Z. Nanjia, H. Feng, M. Christos, Y. Chen-Yu, T. J. Marks, M. G. Kanatzidis and R. P. H. Chang, *J. Am. Chem. Soc.*, 2014, **136**, 15379–15385.

42 Y. Zhou, H. F. Garces, B. S. Senturk, A. L. Ortiz and N. P. Padture, *Mater. Lett.*, 2013, **110**, 127–129.

43 S. Körbel, M. A. L. Marques and S. Botti, *J. Mater. Chem. C*, 2016, **4**, 3157–3167.

