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# Thermal stability, electrochemical and structural characterization of hydrothermally synthesised cobalt ferrite (CoFe<sub>2</sub>O<sub>4</sub>)

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Monophasic nano-crystalline CoFe<sub>2</sub>O<sub>4</sub> (CFO) nanoparticles of high purity have been synthesised through a low temperature hydrothermal route, which does not involve hazardous chemicals, or conditions. The easy, green procedure involves a hydrothermal treatment at 135 °C of an aqueous suspension of the oxalate salts of the precursors. No further purification or annealing procedure was necessary to obtain the crystalline nano-structured oxide. The nanoparticles were characterized structurally and chemically by powder X-ray diffraction (PXRD), Inductively Coupled Plasma Spectrometry (ICP-MS) and Scanning Electron Microscopy (SEM), thus confirming the successful synthesis of the CoFe<sub>2</sub>O<sub>4</sub> particles with the expected crystal phase and stoichiometry and an almost complete inverse spinel structure. From the nanoparticles pellets were pressed to investigate the electronic conduction properties using electrochemical impedance spectroscopy (EIS). At low temperatures, the conductivity measurements reveal a semiconducting behavior originating from hopping between Co sites and a total conductivity dominated by the grain boundary contribution. At higher temperatures ( $T > 400$  °C) a metallic–insulator transition occurs, which is attributed to additional hopping of electrons between the Fe sites.

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## Introduction

Due to their magnetic, electronic, and dielectric properties, transition metal spinel ferrites are of great importance for modern electronic industry and have been studied intensively in the last century.<sup>1–4</sup> The material properties can be tailored over a wide range by changing the oxide composition and its structural features. The high tunability of the properties arises from doping of the spinel structure with the general formula  $[M_{(1-\delta)}^{(II)}Fe_{\delta}^{(III)}]_{Td}[M_{\delta}^{(II)}Fe_{(2-\delta)}^{(III)}]_{Oh}O_4$ , where the  $M^{(II)}$  and  $Fe^{(III)}$  cations are coordinated by oxide ions in 8 out of 64 tetrahedral ( $T_d$ ) or 16 out of 32 octahedral ( $O_h$ ) sites per unit cell. A normal spinel ferrite (e.g. bulk  $Zn^{(II)}Fe_2^{(III)}O_4$ ) has all the  $M^{(II)}$  in  $T_d$  coordination sites and all the  $Fe^{(III)}$  in  $O_h$  sites, while an inverse spinel (e.g. bulk  $Ni^{(II)}Fe_2^{(III)}O_4$ ) has all the  $M^{(II)}$  ions sitting in  $O_h$  sites and half  $Fe^{(III)}$  sitting in  $O_h$  and half in  $T_d$  sites. Bulk CoFe<sub>2</sub>O<sub>4</sub> (CFO) in particular is a ferrimagnetic semiconducting material with an inverse spinel structure with Co<sup>(II)</sup> sitting on the octahedral sites.<sup>5,6</sup> Amongst others CFO exhibits a high

coercivity, Curie temperature, and chemical stability,<sup>7–11</sup> making CFO an interesting material for many applications, such as recording devices, microwave absorption electro-magnetic wave shielding devices, sensors, spintronic applications, and as electrodes in Li ion batteries.<sup>1,12–17</sup> In the last years, ferrite nanoparticles especially have attracted much attention as their properties are not only affected by the chemical composition, but also by the size of the crystallites.<sup>6,18–23</sup> For example, in CFO nanoparticles an increase of the coercivity was reported for decreasing particle size.<sup>6,22</sup> Not only the magnetic properties, but also the electrical properties vary with size, shape and the preparation condition of the particles.<sup>6,8,24,25</sup> The reason is that nano-structured oxide ferrites rarely exhibit a pure inverse or direct spinel structure.<sup>26,27</sup> Most of the prepared ferrites have a partial inversion degree  $\delta$  in its crystal structure, which depends on the preparation procedure and on the further treatment of the oxide, i.e. if  $\delta$  is equal to 0 the spinel is direct and the  $M^{(II)}$  are sitting only on  $T_d$  sites, while the  $Fe^{(III)}$  occupy only the  $O_h$  sites. Some of us have recently shown that the synthesis route can remarkably affect the inversion degree.<sup>28</sup> To synthesise ferrite nanoparticles several methods have been developed during the last years. Solid-states methods, such as mechanical milling<sup>29,30</sup> are advantageous in scalability and in absence of purification steps, especially for industrial scale up. Nevertheless, the limited contact area between the reactant causes undesired sub-products and slow reaction rates are needed to be compensated with increased temperature of about 800 °C.<sup>31–34</sup> Amongst wet-chemical methods, including sol-gel,

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co-precipitation or thermal composition, hydrothermal synthesis is an important low temperature alternative as structurally uniform particles can be synthesised at temperatures below 150 °C.<sup>8,31,34,35</sup> However, due to the strong dependence of the electrical and magnetic properties of ferrites on their structural properties and degree of inversion a detailed characterization of the nanoparticles synthesized by hydrothermal methods is essential to optimize the materials' properties for future applications.

Here, we present the preparation and characterization of nanostructured monophasic CFO using a wet chemistry route based on water as solvent and mild reaction conditions, and affording the formation of the ferrites in extremely short treatment time (*i.e.* 1–3 hours).<sup>36,37</sup> The synthetic route employed is designed to take into account most of the green chemistry principles<sup>38</sup> such as atom economy, a low temperature or energy efficient process, and the use of non-hazardous solvents and chemicals to have an environmentally friendly alternative to the solid-state synthesis for the production of mixed oxides. Monophasic nanostructured CFO with a crystalline domain size of 15–20 nm were prepared in high yield and high atomic efficiency using water as solvent and a mild thermal treatment (135 °C). The CFO nanoparticles were treated at higher temperatures to perform dielectric characterization. To investigate the influence of the heat treatment and the resulting change in microstructure on the electrical transport properties, two different temperatures of 400 °C and 950 °C were chosen. Structural investigations reveal that nanoparticles with a nearly complete inverse spinel structure are obtained. The analysis of the electrical transport properties using electrochemical impedance spectroscopy shows a semiconducting behavior due to hole hopping between Co sites at temperatures below 400 °C, while a semiconductor to metal transition occurs at higher temperature due to additional electron hopping between Fe sites of the lattice.

## Results and discussion

### High-resolution scanning electron microscopy

The microstructure of the pellets prepared was analyzed using High-resolution scanning electron microscopy (HR-SEM) after the annealing process. As shown in Fig. 1, both pellets annealed

at 400 °C (CFO@400) and 950 °C (CFO@950) reveal a complex microstructure, whereas the pellet annealed at 950 °C is denser than the one annealed at 400 °C. Furthermore, beside single nanoparticles with diameters of about 20 nm also large grains with diameters up to 1 μm can be observed for the sample annealed at 400 °C. With increasing annealing temperature, the average diameter of the single grains increases resulting in different crystalline domains of several μm for the pellet annealed at 950 °C.

### XRD analysis

In order to determine the crystal structure of the obtained material, X-ray powder diffraction measurements were performed, which are shown in Fig. 2. The labeled reflections denote the reflections of a cubic Bravais lattice with a  $Fd\bar{3}m$  space group, the  $\text{CoFe}_2\text{O}_4$  reference pattern ICSD 98553 (ref. 39) is shown for comparison. All patterns show the characteristic reflections of cubic CFO. For the as-prepared sample (CFO@RT) as well as the pellet annealed at 400 °C (CFO@400) the diffraction pattern shows an additional, weak reflections at 49.7°, which probably arises due to a negligible, secondary phase as by-product. This additional reflection is not present in the pattern of the pellet annealed at 950 °C (CFO@950) confirming in this case a phase pure material. Compared to the CFO@RT pattern, the reflections of the CFO@400 sample did not change, while the CFO@950 pellet shows extremely sharp reflections due to the size increase of the crystalline domains during the annealing process. Using the Scherrer equation<sup>40</sup> on different reflections ( $hkl$ ),

$$d_{(hkl)} = \frac{K\lambda}{\beta_{(hkl)} \cos \theta_{(hkl)}} \quad (1)$$

being  $d_{(hkl)}$  the supposedly spherical crystallite dimension,  $\beta_{(hkl)}$  the full width at half maximum,  $\theta_{(hkl)}$  the half of the reflection diffraction angle. The particles sizes for the CFO@RT, CFO@400 and CFO@950 sample were estimated to be  $16.7 \pm 4$  nm,  $17.6 \pm 4$  nm and  $46.4 \pm 8$  nm, respectively. Thus, the

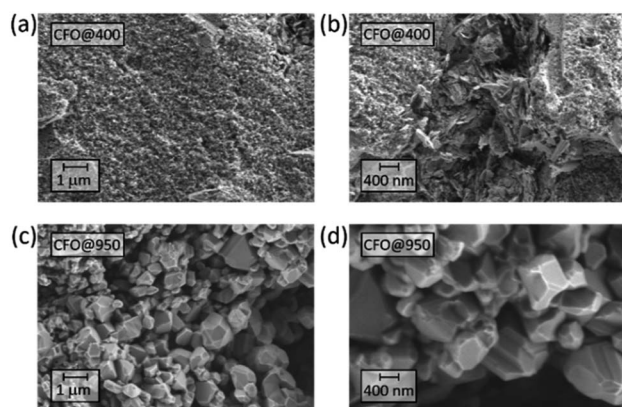


Fig. 1 HR-SEM of CFO@400 (a and b) and CFO@950 (c and d).

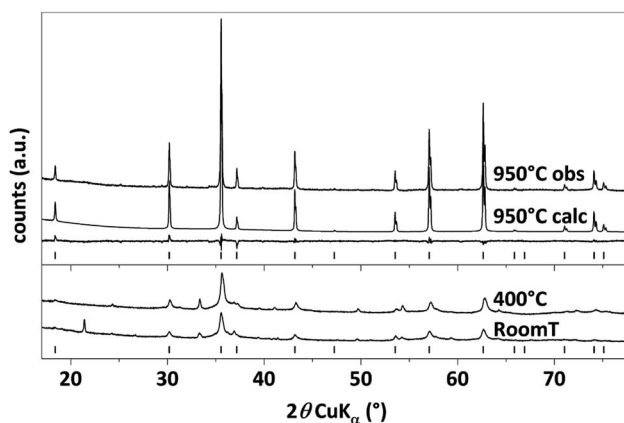


Fig. 2 PXRD of CFO@RT, CFO@400 and CFO@950. Calculated pattern from the CFO@950 Rietveld refinement and its difference from the observed data. The reference pattern ICSD 98553 is also shown for comparison.



growth of the crystallites occurs between 400 and 950 °C. It is worth noting that the Scherrer equation is less reliable for crystallites bigger than 100 nm because at these dimensions the widening of the reflections is caused by instrumental limitations.

To refine the atomic positions in the lattice and, therefore, to investigate the occurrence of a correlation between structural features and electrical behaviour, Rietveld refinement was performed on the CFO@950 pattern. A Chebyshev function<sup>41</sup> with ten coefficients was used to fit the pattern baseline. Only the  $T_d$  and  $O_h$  fractional position and Uiso parameter for both ions  $\text{Co}^{2+}$  and  $\text{Fe}^{3+}$  were refined applying the identity restriction to the same site type and the sum to complementary multiplicity restriction to the sites in the unit cell, while the oxygen position and fraction (Frac) were not refined.

The resulting fit scores a goodness-of-fit (GOF) equal to 1.51 and can be considered as reliable.<sup>42</sup> As shown in Table 1, the occupancy value for the metallic ions indicates a partial inversion  $\delta$  of 0.91 indicating an almost completely inverse spinel. On average, 46% of the  $O_h$  sites of the crystalline domain are occupied by  $\text{Co}^{2+}$  and 54% are occupied by  $\text{Fe}^{3+}$ , *i.e.* there is a statistical alternation of  $\text{Co}^{2+}$  and  $\text{Fe}^{3+}$  in adjacent  $O_h$  positions.

### ICP-MS

To validate the PXRD analyses the actual chemical composition of the mixed oxide was determined using inductively coupled plasma mass spectrometry (ICP-MS).

As shown in Table 2, the experimentally determined stoichiometry of the oxide coincides with the theoretical one within the experimental error range, thus confirming the reliability of the synthesis route.

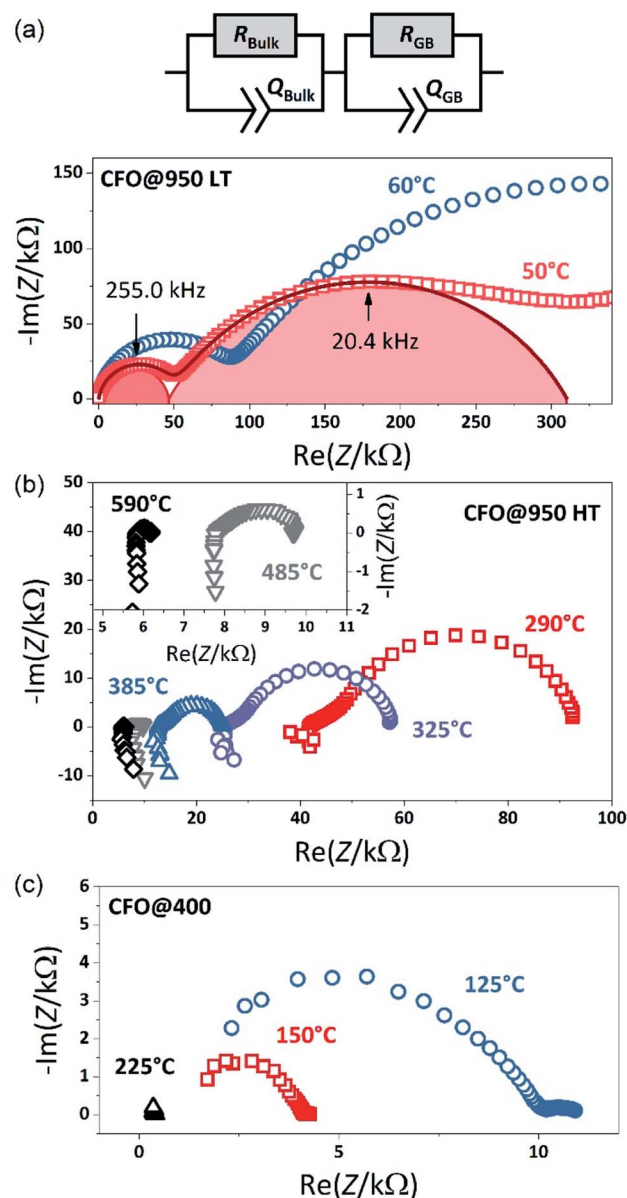
### Impedance measurements

To investigate the electrical properties of the CFO pellets, temperature-dependent impedance measurements were performed. Fig. 3(a) shows the impedance spectra in the Nyquist plots of the CFO@950 pellet, which exhibits two distinct semicircles in the temperature range between −45 °C and 50 °C. To analyze the impedance data an equivalent circuit with two  $RQ$  elements connected in series was used. Each  $RQ$  element consists of a resistance  $R$  and a constant phase elements (CPE)  $Q$  representing one of the non-ideal semicircles. According to the brick layer model of microcrystalline ceramic materials,<sup>43–47</sup> the semicircle in the high frequency range is interpreted as the

**Table 2** Theoretical and experimental stoichiometry of the hydro-thermally synthesised CFO

Expected stoichiometry	Experimental stoichiometry
$\text{CoFe}_2\text{O}_4$	$\text{Co}_{1.001}\text{Fe}_{2.008}\text{O}_{(4+\delta)}$

transport process inside the crystallites, *i.e.* the bulk properties, while the semicircle at lower frequencies represents the transport process across the grain boundaries. At higher temperature



**Fig. 3** (a) Nyquist plot of the impedance spectrum of CFO@950 measured in the temperature range between −45 and 70 °C (just 60 and 50 °C shown) and the equivalent circuit used to describe the data. Nyquist plot of the impedance spectrum of (b) CFO@950 measured between 290 and 590 °C, and (c) of CFO@400 measured between 125 and 225 °C.

**Table 1** Metal ions occupancies of CFO@950 obtained from Rietveld refinement of the PXRD diffraction pattern

Ion	Site	Frac	Uiso	GoF
$\text{Co(II)}_{T_d}$	$T_d$	0.09	$2 \times 10^{-4}$	1.51
$\text{Fe(III)}_{T_d}$	$T_d$	0.91	$2 \times 10^{-4}$	
$\text{Co(II)}_{O_h}$	$O_h$	0.46	$2.1 \times 10^{-3}$	
$\text{Fe(III)}_{O_h}$	$O_h$	0.54	$2.1 \times 10^{-3}$	
O	32	1.00	$1.2 \times 10^{-2}$	



( $T > 50\text{ }^{\circ}\text{C}$ ) the onset of an additional transport process, *i.e.* a third semi-circle, is observed which is caused by the electric charge transfer at the Pt electrode as often observed for oxide materials.<sup>48,49</sup> However, this third semi-circle is not considered in the fitting to ensure continuity between the different temperature regimes by using the same equivalent circle. The corresponding fits to the data are shown in Fig. 3 as solid lines. From the constant phase elements  $Q_i$  the capacitance of the corresponding transport process can be calculated:<sup>49–51</sup>

$$C_i = R_i^{(1-n)} Q_i^{1/n} \quad (2)$$

For the semicircle in the high frequency range a capacitance  $C_1$  of about  $1.2\text{--}1.4 \times 10^{-11}\text{ F}$  is obtained, which is a typical value for the bulk material.<sup>43,45</sup> The second semicircle shows a capacitance of about  $5 \times 10^{-10}\text{ F}$ . The second semicircle therefore can be attributed to the response of the grain boundaries, which show values between  $C_2 = 10^{-10}$  to  $10^{-8}\text{ F}$  depending on the size of the grains in the particle,<sup>46,52,53</sup> *i.e.* with decreasing grain size the capacity shifts to lower values resulting in an overlap of the bulk and grain boundary semicircle in nanocrystalline samples.<sup>46,49,52</sup> Not only the grain size, but also the temperature can result in a merging of the two semicircles due to the temperature dependence of the resistance as shown in Fig. 3(b). With increasing temperature, the bulk and the grain boundary resistances decrease resulting in a dominant grain boundary contribution and, finally, only an ohmic behavior is observed for temperatures above  $600\text{ }^{\circ}\text{C}$  accompanied with inductive contributions from the measurement setup. Thus, in the temperature range between  $140\text{ }^{\circ}\text{C}$  and  $700\text{ }^{\circ}\text{C}$  only the total resistance was extracted from the impedance spectra.

Fig. 3(c) finally shows the impedance spectra obtained from the CFO@400 pellet measured  $50\text{ }^{\circ}\text{C}$  and  $130\text{ }^{\circ}\text{C}$ . Compared to the CFO@950 pellet, only the tail of one semicircle is observed, representing the grain and the grain boundary contribution, accompanied with a small semicircle at low frequencies again arising from the electrode. Thus, the total resistance of the sample is obtained at the intercept point of the first semicircle on the x-axis. The merging of the semicircles at high frequencies representing the grain and the grain boundary contribution imply that the crystallites are in the nanosize range<sup>46,49,54</sup> in agreement with the SEM and XRD observation.

In magnetite the electronic conductivity above the Verwey transition arises from polaron hopping of the electrons between on the octahedral coordinated  $\text{Fe}_B^{2+}$  and  $\text{Fe}_B^{3+}$  sites centres throughout the oxide.<sup>14,55,56</sup> Although  $\text{CoFe}_2\text{O}_4$  has a similar crystal structure as magnetite, replacing  $\text{Fe}_B^{2+}$  with  $\text{Co}_B^{2+}$  reduces the electrical conductivity as cobalt prefers a  $\text{Co}^{2+}$  valence state and has no weakly bound electrons, which hinders electron hopping between the octahedral sites.<sup>14</sup> Instead, hole hopping between  $\text{Co}_B^{3+}$  and  $\text{Co}_B^{2+}$  can occur. In general, the total conductivity in CFO is therefore given by:<sup>57</sup>

$$\sigma = \sigma_e + \sigma_h = n_e e \mu_e + n_h e \mu_h, \quad (3)$$

with  $\sigma_e$  and  $\sigma_h$  being the electron and hole conductivity,  $e$  is the elementary charge, and  $n_{e,h}$  and  $\mu_{e,h}$  are the corresponding carrier concentrations and mobilities, respectively. Which of both conductivities dominates the overall transport behavior of cobalt-ferrite strongly depends on the composition as well as on the stoichiometry. While  $\text{Co}_{1-x}\text{Fe}_{2+x}\text{O}_4$  exhibits n-type conductivity, where electron hopping between  $\text{Fe}_B^{2+}$  and  $\text{Fe}_B^{3+}$  sites dominates, CFO with iron excess ( $\text{Co}_{1+x}\text{Fe}_{2-x}\text{O}_4$ ) shows p-type conductivity due to hole hopping between  $\text{Co}_B^{2+}$  and  $\text{Co}_B^{3+}$ .<sup>57,58</sup>

To analyse the conduction process in more detail, the activation energy was determined from the temperature dependence of the conductivity. Fig. 4 shows the conductivity of both pellets in an Arrhenius-type representation. Both pellets show in the temperature range below  $400\text{ }^{\circ}\text{C}$  a semiconducting behavior, *i.e.* an increase in conductivity with increasing temperature, which is typical for hopping transport in ferrites.<sup>17,58–61</sup> The temperature-dependence of the conductivity is then given by:

$$\sigma(T) = \frac{\sigma_0}{T} \exp\left(-\frac{E_A}{k_B T}\right) \quad (4)$$

where  $E_A$  denotes the activation energy of the hopping process. Fitting the experimental data of the CFO@950 sample to the expression  $\ln(\sigma T) = \ln(\sigma_0) - E_A/k_B T$  gave an activation energy of  $0.50 \pm 0.015\text{ eV}$  and  $0.64 \pm 0.02\text{ eV}$  for the bulk and the grain boundary contribution, respectively. The bulk conductivity is two magnitudes larger than the grain boundary conductivity, *i.e.* the grain boundaries act as barriers for the electronic transport. The total conductivity and the corresponding activation energy ( $E_A = 0.62 \pm 0.02\text{ eV}$ ) of the pellet are therefore dominated by the grain boundaries. As discussed above, the bulk and the grain boundary contribution could only be separated below  $100\text{ }^{\circ}\text{C}$ . Thus, at higher temperature only the activation energy of the total conductivity was determined with  $E_A = 0.63 \pm 0.02\text{ eV}$ , which is in excellent agreement with the low temperature value. The activation energy obtained is also comparable to values reported in literature for nanocrystalline

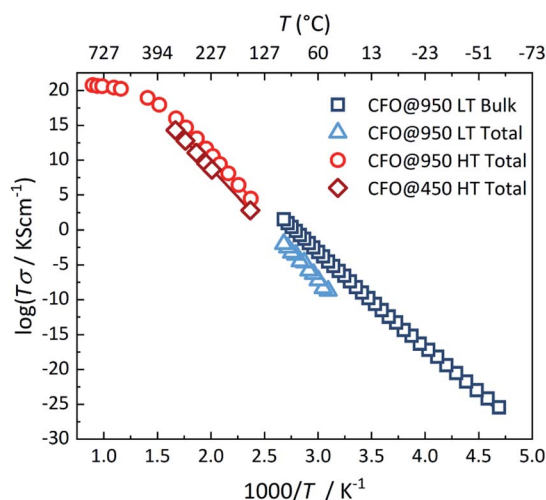


Fig. 4 Arrhenius plot of the conductivities of CFO@400 and CFO@950.





CoFe<sub>2</sub>O<sub>4</sub>, which are in the range between 0.55 to 0.65 eV.<sup>57,58,62,63</sup> Furthermore, it reveals that hole hopping between the Co sites is the dominant transport mechanism at low temperatures as the activation energy for electron hopping between Fe sites is in the range of 0.2 eV.<sup>57–59</sup> For the CFO@400 pellet the same activation energy was determined. However, the total conductivity is about two orders of magnitude smaller than for the CFO@950 sample which is due to the smaller grains of the sample and the resulting larger grain boundary contribution. Above 300 °C, the slope of the total conductivity of the CFO@950 sample changes and a much smaller activation energy of about  $0.17 \pm 0.014$  eV is observed. This behavior is attributed to a change in the transport mechanism. As discussed above holes as well as electrons contribute to the transport in cobalt-ferrites, whereas each can dominate the overall behavior depending on the composition and the stoichiometry of the ferrite. With increasing temperature, the cobalt ferrite is reduced, *i.e.* oxygen is released from the lattice, resulting in an increased density of electrons due to charge compensation. Thus, at high temperatures also electron hopping between the Fe<sub>B</sub><sup>2+</sup> and Fe<sub>B</sub><sup>3+</sup> site centres takes place. As the electron mobility is larger than the hole mobility, the overall transport behavior at high temperatures is dominated by the electronic conductivity resulting in the observed decrease of the activation energy in agreement with previous reports.<sup>62,63</sup>

## Experimental

### Chemicals

Tetraethylammonium hydroxide (20% w/w in water) (TENOH) and sodium hydroxide were purchased from Sigma Aldrich (Milan, Italy). Iron(III) nitrate nonahydrate was purchased from Merck (Darmstadt, Germany). Oxalic acid dihydrate (99.8%) and cobalt(II) chloride hexahydrate were purchased from Carlo Erba (Rodano, Milan, Italy).

### Sample preparation

According to the previously reported procedure<sup>37</sup> stoichiometric amount of the precursor metal salts and oxalic acid were dissolved in 10 mL deionised water in a 23 mL Teflon liner. The concentration of iron(III) and cobalt(II) were kept at  $\sim 0.1$  mol L<sup>−1</sup> and  $\sim 0.05$  mol L<sup>−1</sup>, respectively. 0.2 mL (0.14 mmol) of TENOH, acting as a peptising agent, were added to the solution. The precipitation of the precursor oxalate and hydroxides was induced by adding 1.5 mL of a 10 mol L<sup>−1</sup> NaOH aqueous solution achieving a pH of 12–13. The liner was then sealed and placed in a 23 mL Parr stainless steel acid digestion bomb, heated at 135 °C for 72 hours and left to cool down to room temperature. The obtained solid powders were isolated and purified by washing and centrifugation in three cycles, the first two at 5000 rotations per minute (RPM) for 5 minutes and the last one at 12 500 rpm for 12.5 minutes. Finally, the nanoparticles were dried at 80 °C. To perform electrochemical impedance spectroscopy measurements approximately 0.5 g of the sample powder was pressed to pellets with a diameter of 13 mm using a uniaxial press at 30 kN force for five minutes and subsequently an isostatic press at 550 bar for one hour. Two as

prepared pellets were sintered in air for 12 hours at 400° and 955 °C, respectively, in a Heraeus D-6450 oven.

### PXRD analyses

For the PXRD analysis a Panalytical Empyrean diffractometer was used with an angular accuracy of 0.001° and an angular resolution better than 0.01 employing the CuK<sub>α</sub> radiation. The X-ray patterns were collected with a  $\lambda$  of 1.5406 Å, a scan step of 0.013° (2 $\theta$ ) and 200 seconds acquisition per step time between 10 and 80°. The XRD profile was fitted using GSAS software.

### ICP-MS analyses

For the ICP-MS analyses, all reagents were of analytical grade and used as purchased: nitric acid (69% w/w Aristar for trace analysis, product key 450042N, batch K43187145 206). All solutions were prepared in milliQ ultrapure water obtained using a Millipore Plus System (Milan, Italy, resistivity 18.2 MOhm cm<sup>−1</sup>). The material weighted in using a five decimal digits analytical scale (g), digested in nitric acid (69% w/w) using the microwave digester CEM EXPLORER SPD 24/48 PLUS. Agilent's 7700 Series ICP-MS was tuned using a calibration standard (VI), 30 Analytes ULTRA SCIENTIFIC multi-standard ICP-MS IMS-120 tuning solution prepared in 3% w/w nitric acid was used as an internal standard and added to the sample solution *via* a T-junction. The tuning range was between 25 ppb and 1.5 ppm for Co and between 250 ppb and 14 ppm for Fe.

### EIS analyses

Electrical contact on the pellets were prepared by painting Ferro E1192 platinum paste on the two opposite circular faces of the pellets. The measurements were performed using a Parstat 2273 potentiostat coupled with a Nabertherm p320 furnace regulated with an Eurotherm 2416 thermostat for the measurements in the high temperature range (140 to 740 °C). In the low temperature range (−50 to 70 °C), a Novocool Cryosystem setup was used. The spectra were acquired between 100 mHz and 1 GHz with the Parstat 2273 and between 100 mHz and 2 GHz with the Novocool Cryosystem setup. The analysis of the data and the fitting procedure were performed using the Relaxys software.

### SEM imaging

The SEM images were acquired with a Zeiss Merlin microscope at a voltage between 13 and 10 kV and a current between 90 and 187 pA.

## Conclusions

Crystalline CoFe<sub>2</sub>O<sub>4</sub> nanoparticles were successfully synthesized using a low temperature green hydrothermal synthesis route. Structural analyses confirm a nearly complete inverse spinel structure and stoichiometry of the nanoparticles with sizes ranging between 15 and 20 nm. Impedance measurements on pellets as a function of temperature reveal that the total conductivity at temperatures below 400 °C is dominated by the



grain boundary contribution with an activation energy of about 0.63 eV, which arises from hole hopping between Co sites. At higher temperatures the activation energy strongly decreases to 0.17 eV due to an additional transport process of electron hopping between the Fe sites of the lattice.

## Conflicts of interest

There are no conflicts to declare.

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