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# 1. Introduction

Hybrid organic-inorganic lead halide perovskites with outstanding photovoltaic performance have been identified as desirable light absorption materials for perovskite solar cells.1-3 The power conversion efficiency of perovskite solar cells has been improved from an intial 3.8%<sup>4</sup> to 23.7%.<sup>5</sup> In a typical ABX<sub>3</sub> perovskite structure, A is a monovalent cation such as methylammonium (CH<sub>3</sub>NH<sub>3</sub><sup>+</sup>, MA), formamidinium  $(CH(NH_2)_2^+, FA)$ , and  $Cs^+$ . B is a metal ion (usually Pb<sup>2+</sup>,  $Sn^{2+}$ , or  $Ge^{2+}$ ) and X is a halide anion (usually I<sup>-</sup>, Br<sup>-</sup>, or Cl<sup>-</sup>). Recently, composition engineering on ABX<sub>3</sub> perovskites has been carried out on the cation and halide. This has proved to be a powerful approach to tune the stability and band gap.<sup>6-11</sup> The most popular cations for composition engineering are Cs, MA, and FA.12-16 These works have demonstrated that mixed cations have advantages in improving the stabilities and optoelectronic properties of hybrid perovskites. Therefore, there is a strong desire to explore new stable perovskites with suitable photovoltaic performances.

# Ethylammonium as an alternative cation for efficient perovskite solar cells from first-principles calculations

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Mixed-cation lead halide perovskites have emerged as a new class of promising photovoltaic materials for perovskite solar cells. Formamidinium (FA), methylammonium (MA), and Cs cations are widely studied in the field of mixed-cation hybrid halide perovskites. In this work, we have investigated ethylammonium (CH<sub>3</sub>CH<sub>2</sub>NH<sub>3</sub>, EA) as an alternative cation to explore the stabilities and electronic properties of mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites. The results indicate that replacing MA with EA is a more effective way to improve the stabilities of the mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites except for MA<sub>0.75</sub>EA<sub>0.25</sub>PbI<sub>3</sub>. The band gap of MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> slightly increases with *x* from 0.25 to 1.00, which is quite different from the MA–FA mixed-cation perovskites. The results indicate that the *c* axis distortion of the Pb–1–Pb bond angles can play a greater role in tuning the band gap. Moreover, the mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites show comparable absorption abilities in the visible light region to the pure MAPbI<sub>3</sub> structure. We hope that our study will be greatly helpful for further experiments to find more efficient perovskite materials in the future.

In general, the mixed-cation composition engineering can be evaluated by the Goldschmidt tolerance factor (t), which is defined as follows:<sup>17</sup>

$$t = \frac{R_{\rm A} + R_{\rm X}}{\sqrt{2}(R_{\rm B} + R_{\rm X})} \tag{1}$$

where  $R_A$  and  $R_B$  are the ionic radii of A and B cations respectively, and  $R_X$  is the ionic radius of the anion. The results suggest that the value of the tolerance factor *t* could be between 0.81 and 1.11.<sup>18</sup> To maintain a cubic structure, the ideal *t* should be close to 1. When t > 1.11, the A-site cation is too large and usually precludes the formation of a perovskite. When t < 0.81, the A-site cation is too small, and this often leads to other structures.

The radius of the ethylammonium cation (CH<sub>3</sub>CH<sub>2</sub>NH<sub>3</sub>, EA) is about 2.3 Å,<sup>19</sup> which is greater than that of MA (1.8 Å) but smaller than that of FA (2.6 Å).<sup>20</sup> Based on the eqn (1), the tolerance factor of MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> ranges from 0.83 to 0.94, as shown in Fig. 1. The tolerance factors of the MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites are consistent with the relevant conclusions.<sup>18</sup> This indicates that incorporating EA into MAPbI<sub>3</sub> is a effective way to form stable mixed-cation MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites. However, the EAPbI<sub>3</sub> structure was synthesized in 2012.<sup>21</sup> Its structure was described as a 2H perovskite type with a larger band gap of 2.2 eV.<sup>21,22</sup> The incorporation of EA cations with larger effective radius to improve material stability has been investigated experimentally.<sup>19,23,24</sup> A small amount of EA cations ( $x \le 0.3$ ) have been successfully incorporated into MAPbI<sub>3</sub> to form the stable mixed MA<sub>1-x</sub>EA<sub>x</sub>-PbI<sub>3</sub> perovskites.<sup>25</sup> Recently, a new 3D perovskite with a chemical

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Fig. 1 Correlations between the tolerance factor and the structures of the  $MA_{1-x}EA_xPbI_3$  perovskites.

structure of  $(1,3\text{-}Pr(NH_3)_2)_{0.5}PbI_3$  was successfully synthesized based on a larger organic cation of propane-1,3-diammonium  $(1,3\text{-}Pr(NH_3)_2^{2+})$ .<sup>26</sup> The results showed that this new perovskite exhibited a tetragonal structure and enhanced stability.<sup>26</sup> These results indicate that EA as a alternative cation can form a stable 3D perovskite structure under proper experimental conditions. In this work, we have investigated the structural stabilities and electronic properties of MA-EA mixed-cation MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites, whose electronic properties are substantially different from the MA-FA mixed-cation perovskites.

## 2. Computational details

All of the first-principles calculations were performed with the Vienna Ab initio Simulation Package (VASP) based on

DFT.<sup>27</sup> The projector-augmented wave method was used for the electron-ion interactions and the Perdew-Burke-Ernzerhof (PBE)28 generalized gradient approximation (GGA)29 was used to describe the exchange-correlation effect of the electrons. The cut-off energy for the plane wavefunctions was set to 550 eV. A Monkhorst–Pack 6  $\times$  6  $\times$  6  $\times$  6 k-point mesh over the first Brillouin zone was used for structural optimization and a denser  $8 \times 8 \times 8$  *k*-point mesh was used for the electronic and optical properties. The total energy was converged to  $10^{-5}$  eV. The atom coordinates were fully optimized until the residual forces were smaller than 0.01 eV  $Å^{-1}$ . It is well known that van der Waals interactions play an important role in hybrid organic-inorganic halide perovskite materials.30,31 Thus, the DFT-D3 functional, which includes dispersion interactions, was used for the structural optimization.32

It is well known that the PBE functional always underestimated the band gaps of semiconductors. In the previous DFT calculation, our calculated result shows that the band gap of MAPbI<sub>3</sub> is 1.50 eV with PBE functional,<sup>33</sup> which is quite close experimental value of 1.55 eV.<sup>34</sup> Thus, the PBE functional could give the reasonable band gaps for hybrid lead iodine perovskites. Besides, the calculated optical properties of cationdoped MAPbI<sub>3</sub> with PBE functional also can show the right results. Note that DFT-D3 method was adopted to account for van der Waals interactions in all calculations.

To evaluate the stabilities of the pure and mixed perovskites, we calculated the formation energies of the pure and EA-doped MAPbI<sub>3</sub> perovskites. The crystal structures of MAI was taken from the ref. 35 The crystal structure of EAI was obtained from monoclinic (C2/m) CH<sub>3</sub>CH<sub>2</sub>NH<sub>2</sub>I<sup>36</sup> with full relaxation of the structural parameters. The stable structure of hexagonal PbI<sub>2</sub> was used. The energies of PbI<sub>2</sub>, MAI, and



Fig. 2 Optimized stable geometries of the mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites.

EAI were calculated using the same computational parameters.

#### 3. Results and discussion

#### 3.1 Geometric structures

The simulated mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites are performed by gradually reducing the percentage of MA from the starting 14cm tetragonal phase. The simulated unit cell contains four MAPbI<sub>3</sub> units, so we can simulate three values of the MA-EA ratio, namely 3: 1 (x = 0.25), 2: 2 (x = 0.50), 1: 3 (x = 0.75), and the pure phase of x = 0 and 1. For x = 0.25, 0.50, and 0.75, all of the possible substitution sites are considered. The optimized structures of the mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites are shown in Fig. 2. For the structures with x = 0.50 and x = 0.75, all the possible substitution sites are approximately equivalent with the energy differences within 0.02 eV. However, the results show that different structures with x = 0.25 have different stabilities, with calculated maximum energy differences within 0.08 eV. The most stable structures of  $MA_{1-x}EA_xPbI_3$  with x = 0.25, 0.50, and 0.75 are chosen to investigate their structural stabilities, electronic and optical properties.

The optimized lattice constants of the mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites are listed in Table 1. In the tetragonal MAPbI<sub>3</sub> structure, the relaxed lattice constants a is 8.76 Å and c is 12.95 Å, which is good agreement with the experimental results.<sup>37</sup> The volume slightly increases with increasing EA percentage because of large-size cation doping.

To further evaluate of the stabilities of the pure and mixed perovskites, we calculated the formation energies of the pure and EA-doped perovskites using the following equation:

$$\Delta H = E(\mathrm{MA}_{1-x}\mathrm{EA}_{x}\mathrm{PbI}_{3}) - [(1-x)E(\mathrm{MAI}) + xE(\mathrm{EAI}) + E(\mathrm{PbI}_{2})]$$
(2)

where  $E(MA_{1-x}EA_xPbI_3)$ , E(MAI), E(EAI), and  $E(PbI_2)$  refer to the formation energies of MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub>, MAI, EAI, and PbI<sub>2</sub> phases, respectively. The calculated results for the mixed  $MA_{1-x}EA_xPbI_3$  perovskites are shown in Fig. 3. For MAPbI<sub>3</sub>, our calculated formation energy is well consistent with the previous calculated values of -0.01 eV per f.u.38 and -0.02 eV per f.u.<sup>39</sup> All of the  $\Delta H$  values are negative for the MA<sub>1-x</sub>EA<sub>x</sub>-PbI<sub>3</sub> perovskites. This indicates that substituting MA with EA is energetically favourable. Interestingly, the stability of MA<sub>0.75</sub>EA<sub>0.25</sub>PbI<sub>3</sub> is much lower than that of MAPbI<sub>3</sub>. The similar results were also found for the mixed MA<sub>1-x</sub>Az<sub>x</sub>PbI<sub>3</sub>

Table 1 Calculated lattice parameters for the MA<sub>1-x</sub>EA<sub>x</sub>Pbl<sub>3</sub> perovskites

x	<i>a</i> /Å	b/Å	c/Å	$V/\text{\AA}^3$	
0.00	8.76	8.76	12.95	993.74	
0.25	8.92	8.70	12.97	1005.76	
0.50	9.05	8.67	12.98	1018.52	
0.75	9.00	8.81	13.02	1031.69	
1.00	8.97	8.94	13.04	1045.63	



perovskites.33 The stabilities of the mixed MA1-xEAxPbI3 perovskites are apparently improved when the EA percentage increases from 0.50 to 1.00.

#### 3.2 Electronic properties

-80.0

To obtain insight into the effect of EA-doped cation on the electronic properties, we performed band structure calculations for all of the structures. Fig. 4 shows the  $MA_{1-x}EA_xPbI_3$  band structures with x = 0.00, 0.25, 0.50, 0.75, 1.00. Notably, regardless of the EA percentage, the band structures of perovskites possess direct bandgaps. The band gap of MAPbI<sub>3</sub> at the PBE level is 1.50 eV, which agrees well with a previous theoretical result of 1.53 eV (ref. 40) and the experimental value of 1.55 eV.34 The band gaps are 1.57, 1.61, 1.64, 1.65 eV for  $MA_{1-x}EA_xPbI_3$  with x = 0.25, 0.50, 0.75, and 1.00, respectively. However, the band gap of EAPbI<sub>3</sub> is much lower than its experimental result (2.2 eV).<sup>21</sup> In fact, an orthorhombic structure in the space group of Pmmn was found for EAPbI3.21 EAPbI3 can form a stable 3D structure according to its radius. It indicates that the radius of EA may be measured. The band gaps of the mixed  $MA_{1-x}EA_xPbI_3$  perovskites exhibit a slightly increase with EA concentration, which is in good agreement with the experimental results.<sup>25</sup> However, the results show that the EA incorporation is significantly different from the FA incorporation although the two cations have similar radii. The (1,3- $Pr(NH_3)_2)_{0.5}PbI_3$  perovskite also showed a larger band gap (1.62) eV) than that of the MAPbI3 structure.26 These suitable band gaps of the mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites, from 1.50 to 1.65 eV, makes them highly promising for optoelectronic applications.

To understand the band structure modification, we calculated the projected density of states (PDOS) of the mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites. Fig. 5 shows the dominant PDOS near the band edgeds of  $MA_{1-x}EA_xPbI_3$  with x = 0.00, 0.25, 0.50, 0.75, and 1.00. The results indicate that the valence band maximum (VBM) of  $MA_{1-x}EA_xPbI_3$  originates mainly from the 5p orbital of I. The contribution of the conduction band minimum (CBM) is mainly composed of the 6p orbital of Pb.

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The organic cation does not contribute to the band edge states because the states of organic cation are far from the Fermi level.

Previous report has indicated that the electronic structures of hybrid perovskites are affected by structural distortion, and the main factor is the Pb–I–Pb bond angle.<sup>41</sup> The Pb–I–Pb bond angles of the  $MA_{1-x}EA_xPbI_3$  perovskites are shown in Table 2. The larger distortion of the Pb–I–Pb bond angles is from *ab* plane. The Pb–I–Pb bond angle from *ab* plane has little change except for EAPbI<sub>3</sub>.  $MA_{0.75}EA_{0.25}PbI_3$ gives rise to an increased band gap of 1.57 eV compared with the pure MAPbI<sub>3</sub> structure. This results indicate that the *c* axis distortion can play a greater role in tuning the band gap. It can be seen that the change of the Pb–I–Pb bond angles is responsible for the difference in band gaps of the  $MA_{1-x}$ - $EA_xPbI_3$  perovskites.

To get more insight into the electronic properties, we next examine the Bader charge analysis for the pure and mixed perovskites. The results of Bader charge analysis are shown in Table 3. Charge transfer of organic cation and I ion gradually decrease with increasing x. However, charge transfer of Pb cation gradually increases with increasing x.

#### 3.3 Optical properties

To evaluate the impact of the EA incorporation on the optical performance of MA–EA mixed-cation perovskites, we have calculated the optical absorption coefficient of each structure. The absorption coefficient  $I(\omega)$  was given as below:<sup>42</sup>

$$I(\omega) = \sqrt{2}\omega \left[ \sqrt{\varepsilon_1(\omega)^2 + \varepsilon_2(\omega)^2} - \varepsilon_1(\omega) \right]^{1/2}$$
(3)

where  $\varepsilon_1(\omega)$  and  $\varepsilon_2(\omega)$  represent the real and imaginary parts of the dielectric function depending on the light frequency  $\omega$ . The absorption spectra of the perovskite compounds are shown in Fig. 6. The optical absorption of MAPbI<sub>3</sub> agrees well with the previous result.<sup>40</sup> It can be seen that the MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites have slightly weaker absorption than that of MAPbI<sub>3</sub> in the range of 300–400 nm, which are attributed to their band gaps. The MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites show the comparable absorption abilities in the visible light region compared with the pure MAPbI<sub>3</sub> ( $0.50 \le x \le 1.00$ ) perovskites can become the potential candidates for efficient perovskite solar cells.



Fig. 5 DOS structures of (a) MAPbl<sub>3</sub>, (b) MA<sub>0.75</sub>EA<sub>0.25</sub>Pbl<sub>3</sub>, (c) MA<sub>0.50</sub>EA<sub>0.50</sub>Pbl<sub>3</sub>, (d) MA<sub>0.25</sub>EA<sub>0.75</sub>Pbl<sub>3</sub>, and (e) EAPbl<sub>3</sub>.

Table 2 Pb–I–Pb bond angles and band gaps of the  $\mathsf{MA}_{1-x}\mathsf{EA}_x\mathsf{PbI}_3$  perovskites

	Pb–I–Pb bond		
x	In <i>ab</i> plane	Along <i>c</i> axis	$E_{\rm g}$ (eV)
0.00	150.0	175.6	1.50
0.25	149.9	174.2	1.57
0.50	149.6	171.2	1.61
0.75	151.2	171.0	1.64
1.00	155.8	171.5	1.65

5x10 <sup>5</sup>				EA-0. EA-0. EA-0. EA-0. EA-0.	00 25 50 75
5 4x10 <sup>5</sup>				—— EA-1.	00
∰ 3x10 <sup>5</sup> − 5					
ijd 2x10 <sup>5</sup> -					
1x10 <sup>5</sup> -					
0 <u>300</u>	400	500 Wavelen	600 gth (nm)	700	800

Table 3 Bader charge analysis of the  $\mathsf{MA}_{1-x}\mathsf{EA}_x\mathsf{PbI}_3$  perovskites (A represents the organic cation)

Α	Pb	Ι
+0.727	+0.941	-0.556
+0.719	+0.944	-0.554
+0.714	+0.946	-0.553
+0.710	+0.947	-0.552
+0.707	+0.949	-0.552
	A +0.727 +0.719 +0.714 +0.710 +0.707	A Pb +0.727 +0.941 +0.719 +0.944 +0.714 +0.946 +0.710 +0.947 +0.707 +0.949

Fig. 6 Calculated optical absorption spectra of the  $MA_{1-x}EA_xPbI_3$  perovskites.

# 4. Conclusions

6x10<sup>5</sup>

In summary, first-principles calculations based on DFT were carried out on structural stabilities, electronic and optical properties of the mixed MA<sub>1-x</sub>EA<sub>x</sub>PbI<sub>3</sub> perovskites. The results

indicate that incorporating EA into MAPbI<sub>3</sub> to form mixedcation  $MA_{1-x}EA_xPbI_3$  perovskite is a potentially effective way to improve its stability. The band gap of  $MA_{1-x}EA_xPbI_3$  slightly increases with *x* from 0.25 to 1.00, which is quite different from the MA–FA mixed-cation perovskites. The results indicate that the *c* axis distortion of the Pb–I–Pb bond angles can play a greater role in tuning the band gap. The mixed  $MA_{1-x}EA_xPbI_3$ perovskites show the comparable absorption abilities in the visible light region compared with the MAPbI<sub>3</sub> structure. Based on the above results, we hope that our study will be helpful for further experiments to find more efficient perovskite solar cells.

# Conflicts of interest

There are no conflicts to declare.

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