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## Microphase separation of highly amphiphilic, low $N$ polymers by photoinduced copper-mediated polymerization, achieving sub-2 nm domains at half-pitch†

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The lower limit of domain size resolution using microphase separation of short poly(acrylic acid) homopolymers equipped with a short fluorinated tail, posing as an antagonist 'A block' in pseudo AB block copolymers has been investigated. An alkyl halide initiator with a fluorocarbon chain was utilized as a first 'A block' in the synthesis of low molecular weight polymers (1400–4300 g mol<sup>-1</sup>) using photoinduced Cu(II)-mediated polymerization allowing for very narrow dispersity. Poly(*tert*-butyl acrylate) was synthesized and subsequently deprotected to give very low degrees of polymerization ( $N$ ), amphiphilic polymers with low dispersity ( $D = 1.06$ – $1.13$ ). By exploiting the high driving force for demixing and the well-defined 'block' sizes, we are able to control the nanostructure in terms of domain size (down to 3.4 nm full-pitch) and morphology. This work demonstrates the simple and highly controlled synthesis of polymers to push the boundaries of the smallest achievable domain sizes obtained from polymer self-assembly.

## Introduction

The remarkable properties and tunability of block copolymers (BCPs) offer themselves to a wide range of applications, such as nanoporous membranes for drug delivery,<sup>1</sup> nanolithography and microelectronics.<sup>2–4</sup> These applications have driven the need for smaller domain sizes by feasible synthesis, rendering BCPs a plausible candidate. The thermodynamic driving force of BCP self-assembly (phase separation) is directly related to a balance between the incompatibility of blocks (the Flory–Huggins interaction parameter,  $\chi$ ) and the total degree of

polymerization of the polymer ( $N$ ), therefore, reaching the limits of ultrasmall domain spacings requires high  $\chi$ -low  $N$  BCPs.<sup>5</sup> One of the most widely studied block copolymers for microphase separation in the bulk (solid state) is polystyrene-*b*-poly(methyl methacrylate),<sup>6–9</sup> although the smallest achievable lamellar domain spacing is limited to 17.5 nm full-pitch, as a result of relatively low  $\chi$ .<sup>10</sup> Self-consistent mean-field theory (SCFT) predicts that the order–disorder transition (ODT) of a perfectly symmetrical diblock copolymer occurs when  $\chi N > 10.5$ , and materials are typically disordered below this value.<sup>11</sup> Dispersity ( $D$ ) is a further important factor affecting bulk microphase separation, in which an increase in  $D$  has been shown to change morphology, increase domain spacing and decrease overall degree of order.<sup>12–14</sup> There have been several recent reports of sub-10 nm domain sizes.<sup>4,15–18</sup> In order to push the lower size limit even further, SCFT suggests that the combination of a fluorinated block with a block comprising of highly polar repeat units could lead to domain spacings as small as 2 nm.<sup>19</sup>

The nature of the C–F bond presents an extreme in resistance to external stimuli, low flammability and high hydrophobicity to polymers with high fluorine content.<sup>20</sup> Fluoropolymers can give a wide variety of materials with potential structures and morphologies from semi-crystalline to fully amorphous, and applications across coatings and electronics.<sup>21</sup> The distinctive properties of fluorine contributes to a high degree of incompatibility of blocks in many BCP combinations, with the potential to achieve nanoscale domain spacings in microphase separation.<sup>16</sup> Such polymers have been previously synthesized by reversible-deactivation radical polymerization (RDRP) techniques, which is an optimum choice due to the ability to control molecular weight, dispersity and architecture.<sup>22,23</sup> Photoinduced Cu(II)-mediated controlled radical polymerization has proven to be a versatile approach to BCP synthesis, allowing the use of various activated monomers in both aqueous and organic solvents, with reduced side reactions,<sup>24</sup> and without the need for rigorous deoxygenation.<sup>25</sup> Photoinduced Cu(II)-polymerisation in con-

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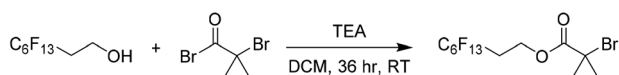
junction with aliphatic tertiary amines<sup>26,27</sup> provides spatiotemporal control, very low dispersity and high end-group fidelity *via* simple implementation.

Reversible addition–fragmentation polymerization (RAFT) has also been frequently used to synthesize high  $\chi$  BCPs.<sup>5,28</sup> Atom-transfer radical polymerization (ATRP) with Cu(I) has been used previously, but in combination with anionic polymerization.<sup>4,16</sup> A feature of ATRP is that the ester bond of the alkyl halide may be susceptible to post-polymerization modification conditions and promote unwanted hydrolysis. Herein, we hypothesise the combination of highly hydrophobic fluorine and highly hydrophilic acid units will result in high incompatibility between blocks, and we exploit photo-induced Cu-mediated polymerization for the synthesis of low molecular weight poly(acrylic acid)-fluoro block co-oligomers (BCOs) with very low dispersity ( $D = 1.06$ – $1.13$ ), allowing for very low  $N$ , which results in, to the best of our knowledge, the smallest domain sizes from polymer aggregation to date.

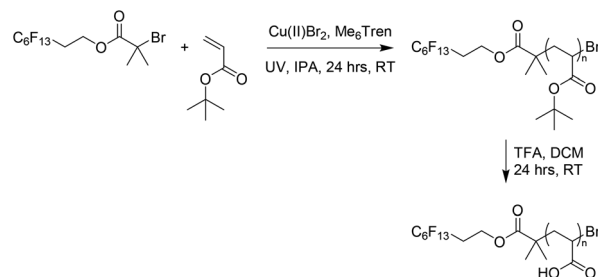
## Results and discussion

An alternative approach to the synthesis and subsequent self-assembly of BCOs is reported herein. The organic synthesis of an alkyl halide initiator with a PTFE-like side chain serves as a truly discrete first oligomeric ‘block’ (with effectively three repeat units) for a series of polymers with incredibly low  $N$ . Esterification of a fluorinated alcohol resulted in an alkyl halide initiator, PFOBiB (perfluorooctyl bromoisobutyrate), with a fluorocarbon moiety of 13 fluorine atoms (Scheme 1). This polymerization allows for very low dispersity ( $D < 1.10$ ), very high initiator efficiency and very low degree of polymerization (DP) materials such that it is possible to distinguish between polymers with small incremental differences such as  $DP = 5, 10, 15$ , etc. This is not the case when dispersity is much broader or initiator efficiency is not close to 100% as with much of the controlled radical polymerization chemistry reported, yet is key to our strategy.

Our rationale was that PFOBiB would itself be a block (F13) with discrete dispersity, and possess a high  $\chi$  interaction parameter with a polar homopolymer, which maintains a low  $N$ . PFOBiB was characterized by  $^1H$ ,  $^{13}C$  and  $^{19}F$  NMR spectroscopy (Fig. S1–3†), the excessive splitting observed in the  $^{13}C$  spectrum is a result of carbon–fluorine coupling. 1, 2 and 3-bond  $J$  values range from 160–300, 15–50 and 5–20 Hz, respectively,<sup>29</sup> producing clear splitting of a qt and tt for terminal fluorine environments. The Cu(II)-photo mediated polymerizations of methyl acrylate in DMSO were carried out to test efficacy as an effective initiator. Successful polymerizations showed excellent end group fidelity and very low dispersity



**Scheme 1** Synthesis of ATRP initiator PFOBiB.



**Scheme 2** Synthesis of  $F_{13}$ -PAA $_n$  polymers.

polymers at low  $N$  (Table S1†). Similar reactions were carried out alongside these, utilizing ethyl  $\alpha$ -bromoisobutyrate (EBiB) initiator as a control, revealing that the presence of fluorine in the initiator had little effect on the initiator efficiency.

The second block was synthesized by polymerization of *tert*-butyl acrylate (*t*-BA) with subsequent deprotection to give poly(acrylic acid) (PAA) (Scheme 2), which has been previously reported in BCP self-assembly.<sup>30–32</sup> Polymerization of a protected monomer was chosen as it better matches the solubility profile of PFOBiB. Furthermore, the direct polymerization of AA can be problematic.<sup>33</sup> Reactions were carried out in IPA rather than DMSO to prevent a biphasic system occurring due to the insolubility of *Pt*-BA in DMSO.<sup>34</sup> Polymers of varying  $N$  were synthesized ( $N = 10$  to  $30$ ) with 8  $CF_2/CH_2$  repeat units in the initiator to give a range of different PAA volume fractions ( $f_{PAA}$ ) due to the constant  $F_{13}$  block, thus covering a controlled range across the phase diagram to induce a change in nanostructure at higher PAA contents. These polymers were obtained in high conversion ( $>99\%$ ) and low dispersity (Table 1, Fig. S4†), which is inherently minimized by the lack of a need for sequential monomer addition.  $F_{13}$ -*Pt*-BA $_n$  polymers were subsequently deprotected using trifluoroacetic acid (TFA) to afford the  $F_{13}$ -PAA $_n$  white powders. The loss of the *t*-butyl peak in the  $^1H$  NMR spectra indicates a successful deprotection (Fig. S5†).  $F_{13}$ -PAA $_n$   $N$  values were calculated by comparing the methyl groups of the  $F_{13}$ -group to the polymer *t*-butyl peak from  $^1H$  NMR spectra of the equivalent  $F_{13}$ -*Pt*-BA $_n$ , as recognizable protons were lost after deprotection. Similarly,

**Table 1** Molecular weight characteristics for  $F_{13}$ -PAA $_n$  polymers

Sample	$N_{AA}^a$	$M_n$ , GPC-SEC <sup>b</sup>	$D_{GPC-SEC}^c$	$M_n$ , theo. <sup>d</sup> (g mol <sup>−1</sup> )	$M_n$ , MALDI (g mol <sup>−1</sup> )
$F_{13}$ -PAA <sub>4</sub>	4	1400	1.08	801	741
$F_{13}$ -PAA <sub>5</sub>	5	1450	1.06	873	N/A
$F_{13}$ -PAA <sub>6</sub>	6	1400	1.09	945	848
$F_{13}$ -PAA <sub>9</sub>	9	1900	1.11	1162	978
$F_{13}$ -PAA <sub>11</sub>	11	1950	1.10	1306	1317
$F_{13}$ -PAA <sub>15</sub>	15	2650	1.08	1594	N/A
$F_{13}$ -PAA <sub>18</sub>	18	3400	1.10	1810	1744
$F_{13}$ -PAA <sub>25</sub>	25	4300	1.13	2515	2097

<sup>a</sup> Degree of polymerization calculated from  $F_{13}$ -*Pt*-BA $_n$  ( $^1H$  NMR).

<sup>b</sup>  $M_n$ , GPC-SEC taken from GPC-SEC of  $F_{13}$ -*Pt*-BA $_n$  in THF. <sup>c</sup> Dispersity calculated from GPC-SEC of  $F_{13}$ -*Pt*-BA $_n$  in THF. <sup>d</sup>  $M_n$ , theo predicted from  $N_{NMR}$  values.



the dispersities were taken from  $F_{13}$ -Pt-BA<sub>n</sub> as  $F_{13}$ -PAA<sub>n</sub> failed to elute properly in both THF and DMF GPC-SEC solvents. This was attributed to the highly amphiphilic nature combined with very low molecular weight.  $F_{13}$ -Pt-BA<sub>n</sub> has a significantly higher mass therefore and drastically different solubility to  $F_{13}$ -PAA<sub>n</sub>, retention times could be obtained. The  $F_{13}$ -block has a more negative refractive index than the GPC solvents, which also contributed to abnormalities in the molecular weight distributions for  $F_{13}$ -PAA<sub>n</sub>.

MALDI-ToF-MS further supported the characterization of  $F_{13}$ -PAA<sub>n</sub> (Fig. 1 and Fig. S6–12†). Profiles show an increasing mass unit of 72 g mol<sup>-1</sup>, consistent with the AA repeat unit. However, multiple end group distributions including H- and Br-termination suggests a loss of end group during hydrolysis. Br-Termination is not expected to be dominant as strongly acidic conditions are expected to reduce it. For all  $F_{13}$ -PAA<sub>n</sub> polymers, the main distribution showed consistency with vinyl termination, however, this is also consistent with cyclization of the terminal AA unit due to the presence of carboxylate anions.<sup>33</sup> It is likely that during hydrolysis, residual TFA continues to deprotonate AA units, thus promoting the cyclization of a lactone and elimination of HBr. The structure of the end group becomes important at low *N* as the loss of Br and concomitant rearrangement (Scheme 3) is substantial. The replacement of the carboxylic acid group with a rigid lactone ring would likely influence  $\chi$  at such low *N*, however, the effect of end group is beyond the scope of this study. MALDI-ToF-MS data were also used to calculate *M<sub>n</sub>* and *D* using the highest intensity peaks (Table 1), purely for confirmation of synthesis, as GPC-SEC (and <sup>1</sup>H NMR) was found unreliable for PAA and it is known that MALDI-ToF-MS can underestimate molecular weight. Despite being a less robust

method due to varying ionization efficiencies of different molecular weights, MALDI-ToF-MS has been shown to be valid for low dispersity and low *N* polymers.<sup>35</sup>

Glass transition temperatures (*T<sub>g</sub>*) obtained from differential scanning calorimetry (DSC) showed an increase in *T<sub>g</sub>* with *N*, which plateaus at ~110 °C (Fig. S13a and b†), therefore, microphase separation was achieved by thermal annealing (TA) the polymers in a PTFE cap at 120 °C overnight. As a result, variables associated with solvent vapour annealing (SVA) are avoided, such as evaporation rate, solvent concentration and solvent type/selectivity, which all have the potential to influence domain size and nanostructure.<sup>36</sup>

SAXS profiles show principal structure-factor peaks for all TA films (Fig. 2), which are not present in either the profiles for the PFOBiB liquid initiator (Fig. S14†) or the EBiB-PAA<sub>n</sub> non-fluorinated polymers (Fig. S15†). This indicates the formation of domains and corroborates the strong incompatibility present in  $F_{13}$ -PAA<sub>n</sub>. It should also be noted that SAXS studies of  $F_{13}$ -Pt-BA<sub>n</sub> polymers are not possible for direct comparison, as these are viscous liquids at room temperature and  $F_{13}$ -PAA<sub>n</sub> are solids.

First order peaks were converted into domain spacings (*d*) via eqn (1), whereby *q\** is the principal peak.

$$d = \frac{2\pi}{q^*} \quad (1)$$

Domain spacings were found to increase with increasing *N* (Fig. 3) where the longer PAA length forces the common  $F_{13}$

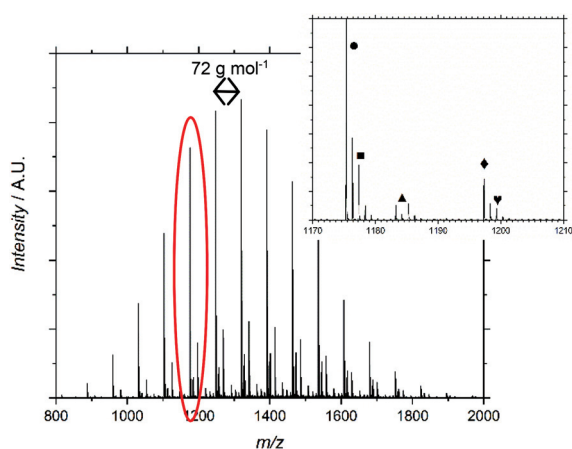
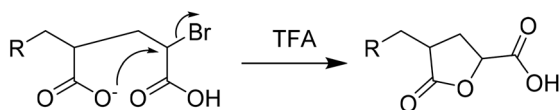


Fig. 1 MALDI-ToF-MS spectrum for  $F_{13}$ -PAA<sub>11</sub>. Inset shows which terminal group represents each distribution.



Scheme 3 Suggested end group cyclization of PAA.

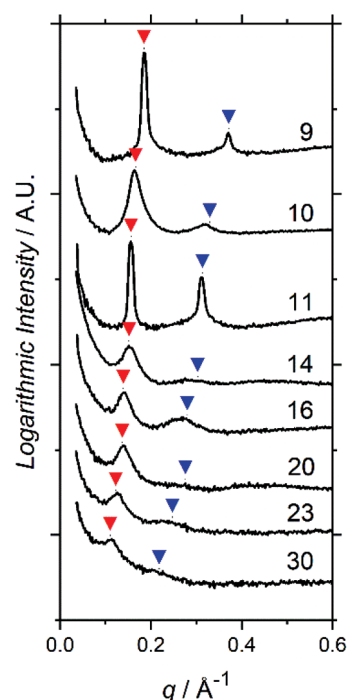
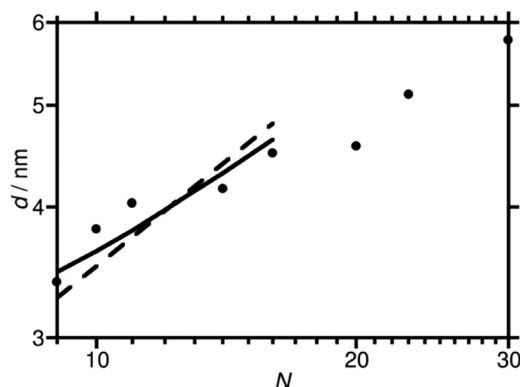


Fig. 2 SAXS profiles for  $F_{13}$ -PAA<sub>n</sub> polymers prepared by thermal annealing. Numbers indicate degree of polymerization (*N*). Red arrows indicate principal peaks (*q\**), blue arrows indicate theoretical position of second order peaks (2*q*).





**Fig. 3** Domain spacing ( $d$ ) plotted against degree of polymerization ( $N$ ). The dashed line shows a fit to the data for the lamellar phase using a formula for coil-coil polymers in the strong segregation regime. The full line shows the fit obtained using a formula for high  $\chi$  rod-coil polymers in the lamellar phase. The fit found with a formula for coil-coil polymers in the weak segregation regime is very close to that obtained with the high  $\chi$  rod-coil formula and is not plotted separately.

domains further apart. Domains ranged from 3.4–5.8 nm reported at full-pitch (which represents the average total distance between like domains), which to the best of our knowledge are the smallest reported for polymer self-assembly to date.<sup>15,16,37</sup> Presence of higher order peaks allows the morphology (nanostructures) to be ascertained. Second order peaks occurring at  $q/q^* = 2$  and  $q/q^* = \sqrt{3}$  represent lamellar (LAM) and hexagonally packed cylinders (HEX), respectively (Table 2). Strong Bragg peaks with narrow widths ( $N \leq 11$ ) imply a well-ordered material with good agreement with LAM nanostructure. Polymers where  $N \geq 14$  show weak intensity higher order peaks, indicating a weakly phase separated structure (W) with more ‘liquid-like’ order. Broadening of SAXS peaks arise for various reasons, the Williamson–Hall plot combines the Scherrer equation, which recognizes that line broadening increases as mean size of ordered domains decreases, with the Stokes and Wilson expression for strain broadening.<sup>38</sup> This explains that  $N = 10$  is likely kinetically trapped in a strained state, and proposes future *in situ* SAXS temperature studies to further investigate the ODT and  $\chi$  values of these

polymers.<sup>39</sup> At  $N \geq 20$ , the weak second order reflection appears to shift closer to  $q/q^* = \sqrt{3}$ , which would suggest a shift from lamellae morphology to hexagonally packed cylinders, although the SAXS features for these polymers are too weak to confirm such a transition.

The extremely short length of the molecules leads to some differences from the standard behaviour of high  $\chi$  block copolymers. As a preliminary, rough illustration of this, we have performed some simple theoretical modelling with a view to gaining insight into these new materials to pave the way for future studies in this area. Accordingly, we compare fits to the plot of  $d$  against  $N$  for the lamellar phase obtained using three simple models, each based on different assumptions about the nature of the molecules. We use the values of  $N$  listed in Table 2, and have checked that using a definition of  $N$  that takes the volume of the repeat unit of one of the blocks as a reference volume (as is often done in fits to domain spacing data<sup>40</sup>) does not significantly change our results.

Firstly, we fit the plot of  $d$  against  $N$  for the lamellar phase with the standard strong segregation formula  $d = cN^{2/3}$ , where  $c$  is an adjustable parameter.<sup>41</sup> This formula is valid for high values of  $\chi N$ , and leads to a fit, shown with a dashed line in Fig. 3, that has a steeper slope than our data. Having noted this, there are two main ways in which the shortness of the molecules can be taken into account. The first is to suppose that the  $F_{13}$  block is so short that it should be treated as a rod. Support for this assumption is provided by simulations of PTFE,<sup>42</sup> which find that these polymers have a Kuhn length of 2.3 nm, longer than the molecular length of the  $F_{13}$  block itself. We therefore fit the data with a formula derived<sup>43</sup> for high  $\chi$  rod-coil block copolymers in the lamellar phase,  $d = cN^{2/3}f_{PAA}^{-1/3}$  (solid line in Fig. 3). This gives improved agreement. Alternatively, it can be supposed that the value of  $\chi N$  is low enough for the sample to be in the weak segregation regime,<sup>41</sup> where  $d = cN^{1/2}$ . The fit obtained using this formula is very similar to that obtained using the formula for rod-coil polymers (although it would give a straight line on the log-log plot above without the slight curvature of the rod-coil graph) and is not plotted separately.

In short, further investigations are needed to determine which of the final two models is more applicable here. Specifically, more data on the temperature- and  $N$ -dependence of the morphologies, together with a model with more microscopic detail, are needed. This more complete model would also take into account the possibility of some molecular alignment within the PAA blocks, which is suggested by the occurrence of high  $q$  peaks in the equivalent EBiB-PAA<sub>*n*</sub> polymers (*i.e.* non-fluorinated, Fig. S15†), and their absence in the liquid initiator precursors (EBiB or PFOBIB, Fig. S14†). Transmission electron microscopy (TEM) of  $F_{13}$ -PAA<sub>6</sub>, and  $F_{13}$ -PAA<sub>18</sub> showed morphologies which support the nanostructures revealed from SAXS. Many difficulties are associated with performing electron microscopy on these polymers. Charging effects prevented nanometer scale resolution as the film would burn, despite using low voltage (200 kV). The samples have an inherent lack of electron density, whereas polymers with aro-

**Table 2** Assembly properties of thermally annealed  $F_{13}$ -PAA<sub>*n*</sub> polymers

Sample	$N^a$	$d^b$ (nm)	$f_{PAA}$	Nanostructure
$F_{13}$ -PAA <sub>4</sub>	9	3.4	0.58	LAM
$F_{13}$ -PAA <sub>5</sub>	10	3.8	0.63	LAM
$F_{13}$ -PAA <sub>6</sub>	11	4.0	0.66	LAM
$F_{13}$ -PAA <sub>9</sub>	14	4.2	0.74	W/LAM
$F_{13}$ -PAA <sub>11</sub>	16	4.5	0.77	W/LAM
$F_{13}$ -PAA <sub>15</sub>	20	4.6	0.82	W/HEX
$F_{13}$ -PAA <sub>18</sub>	23	5.1	0.84	W/HEX
$F_{13}$ -PAA <sub>25</sub>	30	5.8	0.88	W/HEX

<sup>a</sup> Total degree of polymerization calculated from volume fraction of both blocks, see S17† for volume fraction calculation. <sup>b</sup> Reported at full-pitch. LAM = lamellae, HEX = hexagonally packed cylinders and W = weak, liquid-like order.





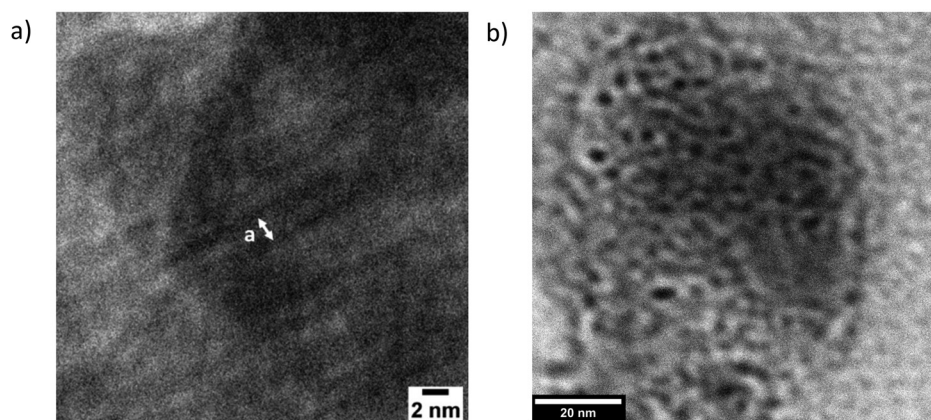


Fig. 4 TEM images of thermally annealed (a)  $F_{13}$ -PAA<sub>6</sub> and (b)  $F_{13}$ -PAA<sub>18</sub>.

matic moiety (e.g. poly(styrene)), overcome this problem.<sup>37,44</sup> Usually, samples would be stained with an electron rich substance to create contrast, however these are often of a larger particle size than the nanometer scale desired for  $F_{13}$ -PAA<sub>n</sub> polymers (i.e. Au nanoparticles ~20 nm). Despite these difficulties, TEM of TA  $F_{13}$ -PAA<sub>6</sub> (Fig. 4a) showed the linear pattern of LAM with a domain spacing calculated from TEM of  $d = \sim 2.6$  nm. The 'liquid-like' order of W/HEX can be seen in the TEM image of  $F_{13}$ -PAA<sub>18</sub> (Fig. 4b). The domain spacing was found to be  $\sim 5.4$  nm for  $F_{13}$ -PAA<sub>18</sub>, from an average of 12 measurements (Fig. S18†). The TEM-calculated domain spacings differ slightly from those obtained from SAXS, because unlike X-ray techniques, microscopy is not necessarily representative of the whole sample and it is not unusual for the surface structure to differ from the bulk, due to migration of groups to the hydrophobic air interface. It should also be noted that sizes could differ due to the different sample preparation methods for TEM and SAXS. SAXS films were prepared as thick films ( $\sim 1$  mm) on a PTFE surface, compared to drop-casted solutions ( $1 \text{ mg ml}^{-1}$ ) onto copper grids for TEM.

## Conclusions

In summary, we have synthesized polymers of low molecular weight with high block incompatibility and studied the self-assembling properties in the solid state. We have used thermal annealing to induce microphase separation, achieving a domain spacing as small as 3.4 nm reported full-pitch, which is, as far as we are aware, the lowest reported to date for this type of polymer assembly. SAXS studies indicated an increase in domain spacing with increasing PAA degree of polymerization and eventually a change in morphology (tentatively assigned to a shift from weakly ordered lamellae to weakly ordered hexagonally packed features) above PAA volume fractions of 0.82. In-depth study behind the theory and construction of the phase diagram of these novel polymers is currently in progress and will follow. TEM indicated the same nanostructures given by SAXS on representative samples, albeit

showing smaller domain spacings, particularly for our smallest polymer (2.6 nm *versus* 4.0 nm). This approach to high  $\chi$  polymers has proven to be a successful concept in polymer self-assembly for potential future use in the microelectronics industry's ever-growing need for smaller domain sizes.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

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