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## High-performance ultra-violet phototransistors based on CVT-grown high quality SnS<sub>2</sub> flakes†

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van der Waals layered two-dimensional (2D) metal dichalcogenides, such as SnS<sub>2</sub>, have garnered great interest owing to their new physics in the ultrathin limit, and become potential candidates for the next-generation electronics and/or optoelectronics fields. Herein, we report high-performance UV photodetectors established on high quality SnS<sub>2</sub> flakes and address the relatively lower photodetection capability of the thinner flakes *via* a compatible gate-controlling strategy. SnS<sub>2</sub> flakes with different thicknesses were mechanically exfoliated from CVT-grown high-quality 2H-SnS<sub>2</sub> single crystals. The photodetectors fabricated using SnS<sub>2</sub> flakes reveal a desired response performance ( $R_\lambda \approx 112 \text{ A W}^{-1}$ , EQE  $\approx 3.7 \times 10^4\%$ , and  $D^* \approx 1.18 \times 10^{11} \text{ Jones}$ ) under UV light with a very low power density (0.2 mW cm<sup>-2</sup> @ 365 nm). Specifically, SnS<sub>2</sub> flakes present a positive thickness-dependent photodetection behavior caused by the enhanced light absorption capacity of thicker samples. Fortunately, the responsivity of thin SnS<sub>2</sub> flakes (e.g.  $\sim 15 \text{ nm}$ ) could be indeed enhanced to  $\sim 140 \text{ A W}^{-1}$  under a gate bias of +20 V, reaching the performance level of thicker samples without gate bias (e.g.  $\sim 144 \text{ A W}^{-1}$  for a  $\sim 60 \text{ nm}$  flake). Our results offer an efficient way to choose 2D crystals with controllable thicknesses as optimal candidates for desirable optoelectronic devices.

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## Introduction

Two-dimensional (2D) materials, such as graphene since its first isolation from bulk crystals, have created opportunities to access novel features at reduced dimensionality.<sup>1–3</sup> Inspired by this, there has been growing interest to find new 2D structures, because fundamentally and technologically interesting properties (charge carriers and photons are confined in a 2D plane, *etc.*) have been identified in the ultrathin limit, but most attention so far has been paid to layered transition metal dichalcogenides (TMDs).<sup>4–7</sup> Furthermore, graphene shows a zero-band gap which traditionally has limited its application in optoelectronic devices.<sup>8</sup> In addition, layered semiconducting TMDs with a general description of MX<sub>2</sub> (a metal atom, M = Mo, W, Sn, *etc.* is sandwiched between two adjacent chalcogen layers X = S, Se, Te, *etc.*) have exhibited extraordinary properties in many fields, emerging as suitable candidates for next-generation electronic or photoelectronic systems beyond the current silicon era.<sup>9–13</sup>

The surprising and appealing characteristics found in the low-dimensional limit stimulate the exploration of other van der Waals layered materials.<sup>14–16</sup> Nowadays, various layered semiconductors can be prepared *via* top-down (e.g. mechanical exfoliation) or bottom-up (e.g. chemical vapour deposition, CVD) methods.<sup>17,18</sup> Nevertheless, in view of sustainable development, earth-abundant 2D materials are more essential for widespread use in modern devices. Tin disulfide (SnS<sub>2</sub>, belonging to the IV-VIA group) comprises earth-abundant constituents (Sn and S), which could vastly support the industrial requirements.<sup>19</sup> Bulk SnS<sub>2</sub> crystals (usually representing a bandgap of 2.1–2.31 eV) have long been explored for possible applications in photovoltaics and photoelectrochemistry.<sup>20,21</sup> In addition, recent evidence has already pointed out the fascinating progress in applying 2D SnS<sub>2</sub> structures in lithium ion batteries,<sup>22</sup> field-effect devices<sup>23</sup> and photodetectors,<sup>24</sup> benefiting from the simple exfoliation from bulk crystals and the controlled bottom-up synthesis.<sup>25,26</sup> In particular, due to a considerable absorption coefficient (10<sup>6</sup> cm<sup>-1</sup>),<sup>27</sup> light sensors using 2D SnS<sub>2</sub> structures (most are obtained by CVD) exhibit superior UV-vis sensing performance such as high responsivity and good stability.<sup>28–30</sup> Note that, however, the thickness dependent responsivity of 2D material (including SnS<sub>2</sub> flakes) based photodetectors (*i.e.* the relatively lower photodetection capability in thinner flakes) may make them inappropriate for the upcoming smart systems with high integration density because the chip manufacturing process in development

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requires thinner (monolayer of few-layer) semiconductor channels.<sup>31–33</sup> Therefore, some prospective strategies should be proposed to address this problem and make these prototypes appealing for practical applications. Unfortunately, this issue has not been well considered in previous optoelectronic applications.

In this paper, we have achieved high-performance UV photodetectors established on  $\text{SnS}_2$  flakes and employed a gate-tunable strategy to address the challenge posed by the positive thickness-dependent light sensing behaviour of 2D materials.  $\text{SnS}_2$  flakes with different thicknesses (that can be thinned to  $\sim 7$  nm in the work) were mechanically exfoliated from CVT-grown high-quality 2H- $\text{SnS}_2$  single crystals. We describe several characterization methods to identify the phase and microstructure of these multilayer  $\text{SnS}_2$  samples. The photodetectors fabricated by using individual  $\text{SnS}_2$  flakes reveal a desired response performance ( $R_\lambda \approx 112 \text{ A W}^{-1}$ , EQE  $\approx 3.7 \times 10^4\%$ , and  $D^* \approx 1.18 \times 10^{11} \text{ Jones}$ ) under UV light with a very low power density ( $0.2 \text{ mW cm}^{-2}$  @ 365 nm). Specifically, the  $\text{SnS}_2$  flakes present a relatively low photodetection ability in the thinner flakes. But it is found that the responsivity of thin flakes (e.g.  $\sim 15$  nm) could be enhanced to  $\sim 140 \text{ A W}^{-1}$  aided by a positively gated bias (+20 V), and it is comparable to the performance of a non-gated  $\sim 60$  nm flake. This indeed provides a potential way to address the positive thickness-dependent detection capability mainly caused by the enhanced light absorption capacity of thicker samples, and our findings imply that such earth-abundant and environmentally friendly tin-based chalcogenides are desirable for sustainable “green” optoelectronics applications.

## Experimental section

### Preparation of multilayer $\text{SnS}_2$ flakes

The starting materials here are  $\text{SnS}_2$  bulk crystals, grown *via* a chemical vapor transport (CVT) route with pure iodine as the transport agent. The pre-mixed powders of Sn (99.99%, Aladdin) and S (99.99%, Alfa Aesar) at a stoichiometric ratio of 1 : 2 with additional iodine (99.8%, Aladdin, 5 mg  $\text{cm}^{-3}$ ) were vacuum sealed ( $>10^{-4}$  Torr) in a quartz tube. The quartz tube was then placed in a two-zone furnace. The reactant zone was slowly heated up ( $\sim 10$  h) to 800 °C while the other end was set to 750 °C. The growth process was maintained for  $\sim 10$  h, followed by a naturally cooling process down to room temperature. Afterwards, strip-like products were obtained. Thin  $\text{SnS}_2$  flakes were mechanically exfoliated from the as-synthesized crystals aided by adhesive tape and then dry-transferred onto a freshly cleaned  $\text{SiO}_2/\text{Si}$  substrate (with a dielectric layer,  $\sim 300$  nm thick  $\text{SiO}_2$ ). Thin samples were roughly identified using an optical microscope (Olympus, CX41) in combination with a charge-coupled device.

### Materials characterization

The phases of the crystals were characterized by X-ray diffraction (XRD) on a Rigaku Miniflex 600x powder diffractometer. XPS measurements were performed using a Thermo ESCALAB 250XI X-ray photoelectron spectrometer. The thicknesses of the

exfoliated  $\text{SnS}_2$  samples were measured by atomic force microscopy (AFM, Agilent 5500). A home-built Raman spectroscope/microscope (iHR320, Horiba) was utilized to acquire Raman spectra and spatially resolved Raman maps with an incident laser of 532 nm while employing a 405 nm light for photoluminescence (PL) measurements. A transmission electron microscope (TEM, JEM-2100F) was employed to evaluate the morphologies and crystal structure of thin  $\text{SnS}_2$  flakes. The UV-vis spectrum was measured using a spectrophotometer (MPC-3100, Shimadzu).

### Device fabrication and measurements

The fabrication of photodetectors based on individual  $\text{SnS}_2$  flakes relies on a standard UV lithography (URE-2000/25) process followed by thermal evaporation (JSD 300) of desired electrode metals (Cr/Au,  $\sim 10/60$  nm). The photoresponse measurements were executed on a probe station (ZFT-50T) equipped with two sourcemeters (Model 2450, Keithley). Light sources (THORLABS) of different wavelengths with tunable power were applied in the photodetection tests.

## Results and discussion

To get  $\text{SnS}_2$  samples with varied thicknesses, we employed tape-assisted mechanical isolation from raw crystals synthesized *via* a chemical vapour transport (CVT) approach. Previous research suggests that  $\text{SnS}_2$  occurs as different polytypes associated with different interlayer stacking of S-Sn-S layers; a low-temperature ( $<800$  °C) synthesis process tends to produce the 2H-polytype.<sup>34,35</sup> Fig. 1a illustrates the 3D structure model of 2H- $\text{SnS}_2$ ; the S-Sn-S atomic planes (the distance between adjacent planes is  $\sim 0.6$  nm) with covalent bonding are held together by weak van der Waals force.<sup>36</sup> In this work,  $\text{SnS}_2$  crystals were prepared *via* a CVT process in a two-zone tube furnace, as schematically shown in Fig. 1b (more information has been provided in the Experimental section). During the CVT growth, gas transport is realized by temperature gradient and thus impact on the final products (Fig. S1†). In this work, a 50 °C temperature gradient is in favour of growing high quality  $\text{SnS}_2$  single crystals, in agreement with a previous report,<sup>37</sup> and large sized crystals (given in the inset of Fig. 1c, lateral size  $> 3$  mm) with a clean surface could be found at the lower temperate zone. The phase structure of the crystals was identified through the powder X-ray diffraction (XRD) technology, as depicted in Fig. 1c. The products exhibit a hexagonal structure (JCPDS PDF number 23-0677,  $a = b = 0.36$  nm,  $c = 0.59$  nm). Strong diffraction reflections emerging around  $\sim 15.0^\circ$ ,  $\sim 30.3^\circ$  and  $\sim 46.1^\circ$  are indexed to the (0001), (0002) and (0003) planes, respectively. The predominance of the (000N) ( $N = 0, 1, 2, \text{etc.}$ ) peaks suggests that the Z-direction with the (0001) plane as the basal plane is the preferential orientation for these growing  $\text{SnS}_2$  crystals.<sup>38</sup> The valence states of the  $\text{SnS}_2$  crystals were characterized by X-ray photoelectron spectroscopy (XPS). The spectra have been calibrated by employing absorbed C (1s) as the reference. Expected Sn and S elements from these crystals were detected in the XPS survey (Fig. 1d). Two peaks at 486.5 and 495.1 eV (Fig. 1e)

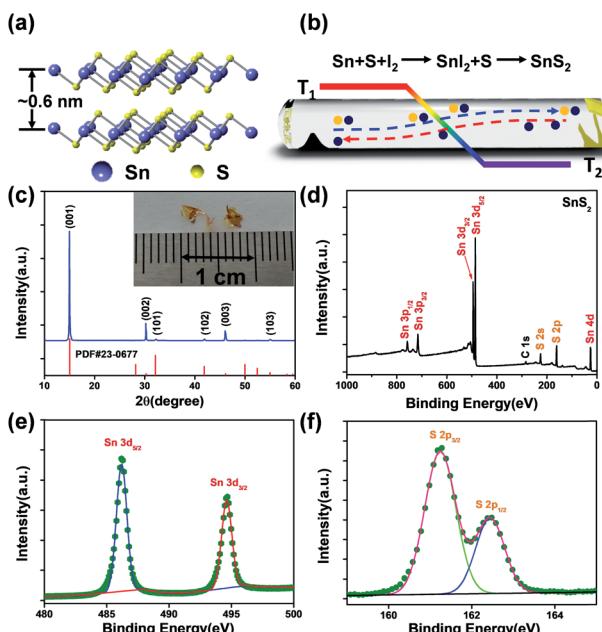


Fig. 1 Synthesis of SnS<sub>2</sub> crystals. (a) Crystal structure (side view) of layered SnS<sub>2</sub>. (b) Schematic diagram showing the CVT process. (c) XRD patterns of the as-synthesized products, inset: the photograph of SnS<sub>2</sub> single crystals. (d) XPS survey of SnS<sub>2</sub> crystals, (e) Sn3d XPS and (f) high resolution S 2p XPS of SnS<sub>2</sub> single crystals.

originate from Sn3d<sub>5/2</sub> and Sn3d<sub>3/2</sub> of Sn<sup>4+</sup>.<sup>39</sup> Fig. 1f illustrates the high-resolution core XPS spectrum of S 2p (~162.6 eV), an indication of the existence of S2p<sub>3/2</sub> and S2p<sub>1/2</sub> orbitals.<sup>40</sup> These results evidence the formation of layered SnS<sub>2</sub> crystals with high quality.

Thin flakes with varied thicknesses (Fig. 2a) were obtained through mechanical exfoliation from large crystals exceeding 10 μm (limited by the AFM scanning range) accompanied by a clean surface and pristine state. The height of multilayer SnS<sub>2</sub> flakes often presents lateral dimension changes shown in Fig. 2b measured at the edge between the thicker and thinner part indicating a thickness of ~7 nm (corresponding to about ten layers). Raman spectroscopy was employed to quantify and map SnS<sub>2</sub> flakes. Previous reports suggest the existence of different crystal polytypes for layered SnS<sub>2</sub> crystals, 4H- and 2H-phases, respectively.<sup>41</sup> The most intense Raman peak for the 4H polytype emerges at 313.5 cm<sup>-1</sup>, relative to the joint contribution of A<sub>1</sub> and E phonon modes (this peak is very close to the A<sub>1g</sub> mode of 2H-SnS<sub>2</sub> at 315 cm<sup>-1</sup>), while the E-mode generates a doublet at 200 and 214 cm<sup>-1</sup>, respectively.<sup>19</sup> In parallel, in 2H-SnS<sub>2</sub> crystals, the E<sub>g</sub> mode generates a single, intense line around 205 cm<sup>-1</sup>, allowing a facile distinction against the 4H polytype.<sup>28</sup> The observed Raman active modes in Fig. 2c, *i.e.* two main peaks at 205 (E<sub>g</sub>, in plane) and 315 cm<sup>-1</sup> (A<sub>1g</sub>, out of plane) occur in the thick flakes (>200 nm, the inset of Fig. 2c), hence, providing an unambiguous fingerprinting of 2H polytypic crystals in this work. Here, the E<sub>g</sub> modes weaken and become unobservable with the reduction of flake thickness (down to the nanometer level), which could be presumably attributed to the reduction of in-plane scattering centers in the ultrathin SnS<sub>2</sub>

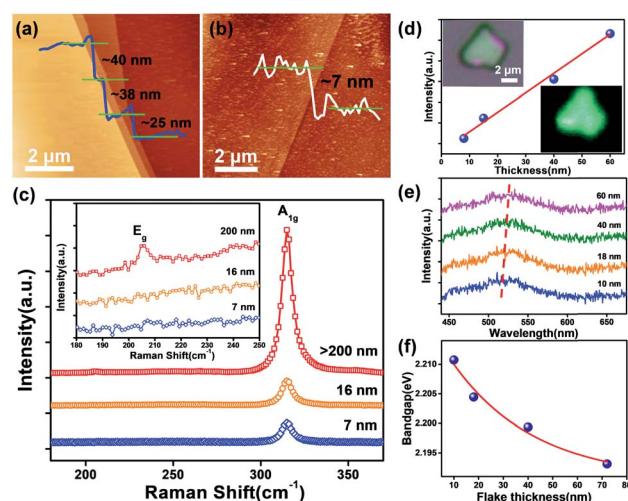
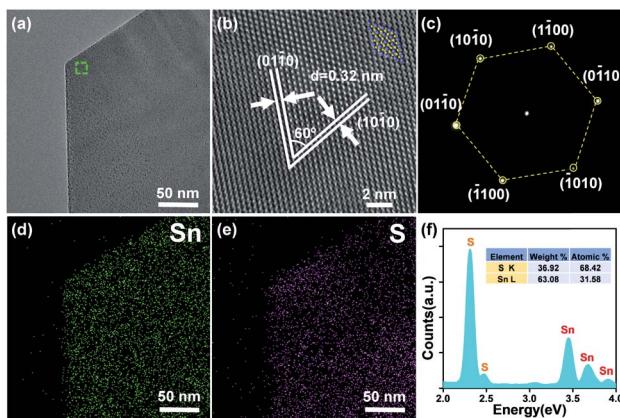


Fig. 2 Characterization of SnS<sub>2</sub> flakes. (a) and (b) Typical AFM images of ultrathin SnS<sub>2</sub> flakes acquired *via* mechanical exfoliation. (c) Thickness dependent Raman spectra of SnS<sub>2</sub> flakes, inset: enlarged view of the characteristic peak at ~205 cm<sup>-1</sup>. (d) Change in the intensity of the A<sub>1g</sub> peak with layer thickness, inset: an optical image of a SnS<sub>2</sub> flake and the corresponding Raman map of the A<sub>1g</sub> mode. (e) PL spectrum of SnS<sub>2</sub> with different thicknesses. (f) The layer-dependent bandgap of few-layer SnS<sub>2</sub>.

flakes.<sup>42,43</sup> But the A<sub>1g</sub> (out of plane) mode illustrates a significant increase in the peak intensity with the increasing thickness (Fig. 2d), probably arising from the enhanced light absorption capacity of the thicker samples,<sup>44,45</sup> which may have an influence on the device performance. Raman mapping (Fig. 2c) of a SnS<sub>2</sub> flake (the optical image is depicted in Fig. 2d) using the characteristic line at 315 cm<sup>-1</sup> demonstrates its uniform polytype. 2H-SnS<sub>2</sub> flakes exhibit isotropy in the (0001) plane as confirmed by the polarized Raman characterization of the A<sub>1g</sub> vibration mode. Therefore, the origination dependent photo-responsivity should not be a problem for their application in photodetection.<sup>46</sup> The PL spectra of 2D SnS<sub>2</sub> flakes in Fig. 2e consist of a single feature, attributed to the nature of their indirect band-gap structure, with a value of ~2.20 eV similar to the reported data of 2H-SnS<sub>2</sub>.<sup>28</sup> Moreover, the peak shifts to lower energies with increasing thickness (Fig. 2f). During PL measurements, holes will combine with the photo-excited electrons *via* Coulomb interactions with the binding energy lying in the band-gap region; however, strong spatial confinement and a reduced screening effect in ultrathin samples could result in a significantly enhanced excitonic effect causing a blue shift in the bandgap.<sup>47</sup>

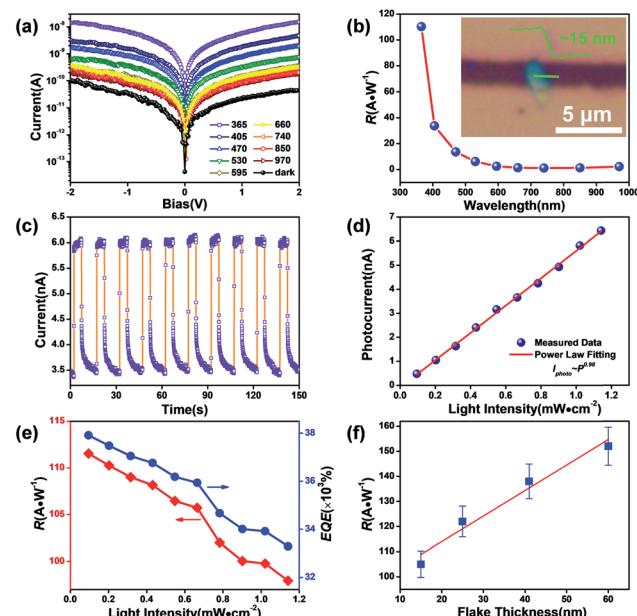
The microscopic structure and the compositions of the 2D SnS<sub>2</sub> crystals were determined using a transmission electron microscope (TEM) and an energy-dispersive X-ray detector (EDX). Fig. 3a shows a low magnification view (bright field) of a single flake with no observable terraces. A direct view of the lattice structure (Fig. 3b, reflected by the HRTEM image of the labeled region in Fig. 3a) suggests the hexagonal arrangement (atomic mode given on the top) of Sn and S atoms in 2H SnS<sub>2</sub> crystals. The plane distance of ~0.32 nm well matches the *d*-



**Fig. 3** TEM identification of  $\text{SnS}_2$  flakes. (a) Low-magnification TEM image of an ultrathin  $\text{SnS}_2$  flake. (b) High resolution TEM characterization with atomic mode on the top and (c) the corresponding SAED pattern of the  $\text{SnS}_2$  flake. EDX elemental mapping of (d) Sn and (e) S revealing uniform distributions. (f) EDX spectrum of the  $\text{SnS}_2$  flake; the inset illustrates the atomic ratio of Sn and S.

spacing of the  $\{10 - 10\}$  planes for hexagonal phase  $\text{SnS}_2$ .<sup>48</sup> The selected area electron diffraction (SAED) pattern (obtained by applying incident electrons parallel to the  $c$ -axis) in Fig. 3c shows well-sequenced diffraction spots with a six-fold symmetry, indicating high-quality crystallinity of this crystal caused by a vertically stacking layer plane along the  $[0001]$  direction. The composition of the crystals was verified by adopting energy-dispersive X-ray spectroscopy (EDX). EDX mapping results in Fig. 3d and e suggest the homogeneous distribution of Sn and S elements in the flake. In parallel, the EDX spectrum (Fig. 3f) clearly reveals signals of Sn and S with an approximate atomic ratio of 2 : 1. The aforementioned demonstrations signify the high purity of the single-crystalline  $\text{SnS}_2$  flakes.

Photodetectors established on individual flakes were employed to systematically estimate the optoelectronic properties of the  $\text{SnS}_2$  crystals. The set-up diagram of a two-terminal light sensor is illustrated in Fig. S3a,<sup>†</sup> where the incident light is perpendicular to the  $\text{SnS}_2$  flake. Fig. S3b<sup>†</sup> and 4a present the current–voltage ( $I_{\text{DS}}-V_{\text{GS}}$ ) characteristics of the representative  $\text{SnS}_2$  light sensor in darkness and under illumination with varied wavelengths. The linear curves result from the near ohmic barrier between the Cr/Au electrodes and  $\text{SnS}_2$  channel. The plot of responsivity (Fig. 4b,  $R_{\lambda} = I_{\text{ph}}/PS$ ,  $I_{\text{ph}} = I_{\text{light}} - I_{\text{dark}}$ ,  $S$  refers to the activated area of  $\sim 5.17 \mu\text{m}^2$  and  $P$  is the power intensity of incident light,  $V_{\text{DS}} = 1 \text{ V}$ ) versus wavelength provides a quantitative assessment of the photoresponse ability of the  $\text{SnS}_2$  flake. The optical image of the  $\text{SnS}_2$  channel with a thickness of  $\sim 15 \text{ nm}$  and a length of  $\sim 3 \mu\text{m}$  is illustrated in the inset. The cut-off wavelength is about 550 nm, close to its deduced bandgap,  $E_{\text{indirect}} \approx 2.18 \text{ eV}$  from the UV-vis absorption spectrum in Fig. S4.<sup>†</sup><sup>29</sup> The light sensor exhibits high photoresponsivity in the ultra-violet range. The stability and reproducibility of the  $\text{SnS}_2$  based photodetector (Fig. 4c, measured at a bias of 1 V) are uncovered via tracking current under periodic



**Fig. 4** Photodetectors established on ultrathin  $\text{SnS}_2$  flakes. (a)  $I-V$  characteristics measured in darkness and under incident light of varied wavelengths (under comparable light intensity). (b) Spectral responsivity of the  $\text{SnS}_2$  flake based photodetector, inset: the representative two-terminal devices designed on a thin  $\text{SnS}_2$  flake. (c) Time resolved current of the light sensor measured at  $V_{\text{DS}} = 1 \text{ V}$  under 365 nm ( $0.2 \text{ mW cm}^{-2}$ ). (d) Power law fitting photocurrents versus light intensities. (e) Responsivity ( $R_{\lambda}$ ) and external quantum efficiency (EQE) plotted as a function of light intensity. (f) Thickness dependent responsivity of  $\text{SnS}_2$  flake based devices.

illumination of  $0.2 \text{ mW cm}^{-2}$  @ 365 nm. Enlarged views of the rising and decay sides are illustrated in Fig. S3c,<sup>†</sup> from which the response time ( $\tau_{\text{rising}}$ ) and recovery time ( $\tau_{\text{decay}}$ ) are calculated to be 40 and 160 ms, respectively. Fig. 4d displays the dependence of photocurrent on light intensity, which can be well described by a power law,  $I_{\text{ph}} \sim P^{\beta}$ , where  $\beta$  is an exponent determined by trap states on the surface of the photo-sensitive media.<sup>49</sup> In general, a low power intensity would benefit the occupation of surface states by holes separated from photo-induced electron–hole pairs, followed by a rapid recombination with the negatively charged oxygen. But abundant electron–hole pairs will be generated under a higher light intensity. The reduction of the hole-trap states at the surface until the complete occupation of surface traps will contribute to faster recombination (in several picoseconds) between extra charges; subsequently, these fresh pairs will not contribute to the charge transfer process, but form non-radiative carrier-recombination centers, thus leading to a decline (Fig. 4e) of responsivity and quantum efficiency ( $\text{EQE} = hcR_{\lambda}/e\lambda$ ,  $h$  is Planck's constant).<sup>50,51</sup> The sub-linear behaviour in Fig. 4d with a fitting value  $\beta \sim 0.98$  (very close to 1) may be associated with a low concentration of traps or defects in these  $\text{SnS}_2$  flakes.<sup>52</sup> The parameters of our sensors are comparable or superior to those of documented 2D  $\text{SnS}_2$  and other 2D material-based UV sensors as summarized in Table 1, potentially offered by the high-quality, large specific surface area.<sup>53</sup> In addition, we find a positive thickness-



**Table 1** Comparison of the 2D UV photodetector performance with those reported by others.  $T$ : thickness, ME: mechanical exfoliation,  $D^*$ : detectivity

Material	Synthesis	Electrode	$T$ [nm]	$\lambda$ [nm]	$R_\lambda$ [A W $^{-1}$ ]	EQE [%]	$D^*$ [Jones]	$\tau$ [ms]	Ref.
WO <sub>3</sub>	CVD	Cr/Au	12	365	293	997	—	40/80	57
$\beta$ -Ga <sub>2</sub> O <sub>3</sub>	Oxidant	Cr/Au	6	254	3.3	$1.6 \times 10^3$	$4 \times 10^{12}$	30/60	58
Bi <sub>2</sub> Te <sub>3</sub>	ME	Pt	—	325	26.82	102	$1.29 \times 10^9$	280/1600	59
SnS <sub>2</sub>	CVD	Ti/Au	—	390	—	150	—	8/150	60
SnS <sub>2</sub>	CVD	Cr/Au	10	350	260	$9.3 \times 10^4$	$1.9 \times 10^{10}$	20/16	28
SnS <sub>2</sub>	CVD	Cr/Au	114	100–800	1.568	480.1	—	42/40	61
SnS <sub>2</sub>	ME	Cr/Au	15	365	112	$3.7 \times 10^4$	$1.18 \times 10^{11}$	40/160	This work

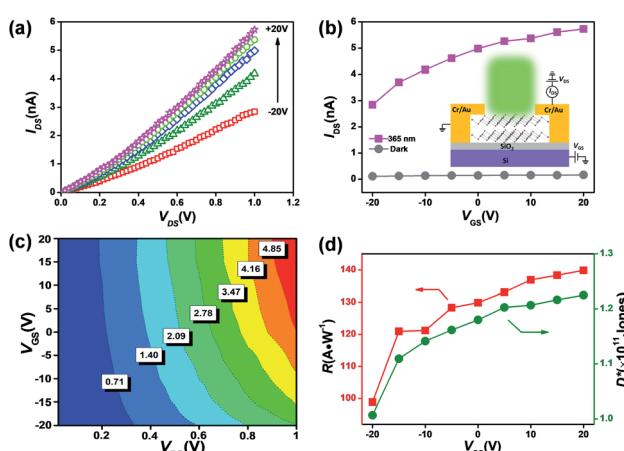
dependent photodetection capability of these 2D SnS<sub>2</sub> flakes with a similar active zone ( $R_\lambda$  up to  $144$  A W $^{-1}$  for a  $\sim 60$  nm thick SnS<sub>2</sub> flake, in Fig. 4f), which has also been documented in other research.<sup>54</sup> For a given wavelength range, the photocurrent of photosensitive flakes is proportional to the absorption if we consider a constant namely the internal quantum efficiency ( $\eta$ ), which can be expressed as  $I_{ph} = \alpha \times d \times \eta$ , where  $\alpha$  is the absorption and  $d$  is the thickness of the flake.<sup>55</sup> Therefore, the dependence of photoresponsivity on flake thickness may originate from enhanced absorption in thicker samples.<sup>56</sup>

The positive thickness-dependent photodetection capability may be a challenge for the system-on-a-chip design where a thinner channel is required. Previous work employed the surface sensitization of SnS<sub>2</sub> nanosheets using a 2 nm thick HfO<sub>2</sub> nanolayer grown by atomic layer deposition (ALD).<sup>54</sup> However, this inevitably leads to a complex procedure or a higher production cost. A phototransistor, in which the increase of gate voltage gives rise to an increase in the channel current, may pave an alternative way to address this issue.<sup>62,63</sup> Here, the emphasis has been placed on the device (as schematically shown in the inset of Fig. 5b) performance operating

under illumination ( $0.2$  mW cm $^{-2}$  @ 365 nm). The light output characteristics (thickness  $\sim 15$  nm) are displayed in Fig. 5a. The output and transfer characteristics measured in darkness are illustrated in Fig. S5.<sup>†</sup> Both the SnS<sub>2</sub> channels reveal a typical n-type semiconducting behaviour. The correlation between the channel current and back-gate bias ( $I_{DS}$ - $V_{GS}$  curves) under illumination and dark conditions is plotted in Fig. 5b. An increase in the channel current can be found as  $V_{GS}$  increases, possibly due to the leading role of photo-generation in comparison with tunneling or thermionic currents.<sup>64</sup> The false-color plot in Fig. 5c provides direct evidence that higher photocurrent could be achieved under a high (positive) gate bias. Thereby, the indicators of light sensors relying on SnS<sub>2</sub> flakes, such as responsivity and detectivity (Fig. 5d), could be further improved (up to  $142$  A W $^{-1}$ ) through changing the gate voltage. Under a positive gate bias, the Fermi level in the (n-type) semiconductor will approach the conduction band and lead to a reduction of barrier height; as a result, more photo-excited charges could overcome the gate barrier and contribute to an increased carrier density.<sup>65,66</sup> Hence, this compatible manufacturing approach could address the positive thickness-dependent sensing capability caused by the enhanced absorption capacity in thicker samples.

## Conclusions

In summary, we report high-performance UV photodetectors established on SnS<sub>2</sub> flakes and address the relatively low photodetection capability in the thinner flakes *via* a compatible gate-tunable route. Multilayer SnS<sub>2</sub> flakes (that can be thinned to  $\sim 7$  nm in this work) were mechanically isolated from CVT-grown high-quality 2H-SnS<sub>2</sub> single crystals whose components are both inexpensive and earth-abundant. The phase and microstructure were unambiguously identified through several characterization techniques, including XRD, XPS, AFM, Raman and TEM. When exposed to UV illumination with a very low power density ( $0.2$  mW cm $^{-2}$  @ 365 nm), the light sensors using SnS<sub>2</sub> flakes exhibit high responsivity ( $112$  A W $^{-1}$ ), EQE ( $3.7 \times 10^4$  %) and detectivity ( $1.18 \times 10^{11}$  Jones), comparable or superior to those of reported SnS<sub>2</sub> and other 2D material-based UV photodetectors. Most importantly, SnS<sub>2</sub> flakes present a positive thickness-dependent photodetection behaviour, possibly attributed to the enhanced light absorption capacity of thicker samples. However, the responsivity of thinner flakes (e.g.  $\sim 15$



**Fig. 5** Phototransistors based on SnS<sub>2</sub> flakes. (a) Output characteristics of the SnS<sub>2</sub> ( $\sim 15$  nm) transistors operated under incident light of  $0.2$  mW cm $^{-2}$  @ 365 nm. (b)  $I_{DS}$ - $V_{GS}$  curves under 365 nm light and in darkness at  $V_{DS} = 1$  V, inset: the layout of the phototransistor. (c) False-color plot of the SnS<sub>2</sub> phototransistor exposed to incident light, where the color reflects the intensity of the measured photocurrent. (d) Photoresponsivity and detectivity measured at  $V_{DS} = 1$  V as a function of  $V_{GS}$ .



nm) can be further improved (up to 140 A W<sup>-1</sup>) under a gate bias of +20 V, comparable with the performance of a non-gated thick flake ( $\sim$ 144 A W<sup>-1</sup> for a  $\sim$ 60 nm flake). Hence, our results offer an efficient way to choose 2D materials with an optimal thickness, and such earth-abundant and environmentally friendly tin-based chalcogenides are highly desirable for sustainable “green” optoelectronics applications.

## Conflicts of interest

There are no conflicts to declare.

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