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1. Introduction

Two-dimensional semiconductor transition metal dichalcogenides (TMDCs) have sparked immense efforts for their applications in next-generation optoelectronic and electronic devices, such as light-emitting diodes (LEDs), field-effect transistors (FETs), solar cells and photodetectors, due to their unique optical and electrical properties.¹⁻⁹ The excitons in the TMDCs can be formed via the recombination of photo-excited electron-hole pairs based on the attractive coulombic interactions.¹⁰ The formed excitons can radiatively decay to the ground state accompanied by the emission of light.11 However, it remains practically difficult to achieve high-efficiency photoluminescence, as the radiative recombination of excitons is significantly suppressed by the dominated assembly of negative trions in n-doped monolayer TMDCs.¹²⁻¹⁵ Several strategies were attempted to introduce p-type doping to modulate the electronic property of monolayer TMDCs, thus resulting in the enhancement of radiative recombination.16-19

Surface group-modified MXene nano-flake doping of monolayer tungsten disulfides†

Exciton/trion-involved optoelectronic properties have attracted exponential amount of attention for various applications ranging from optoelectronics, valleytronics to electronics. Herein, we report a new chemical (MXene) doping strategy to modulate the negative trion and neutral exciton for achieving high photoluminescence yield of atomically thin transition metal dichalcogenides, enabled by the regulation of carrier densities to promote electron-bound trion-to-exciton transition *via* charge transfer from TMDCs to MXene. As a proof of concept, the MXene nano-flake-doped tungsten disulfide is demonstrated to obtain an enhanced PL efficiency of up to ~five folds, which obviously exceeds the reported efficiency upon electrical and/or plasma doping strategies. The PL enhancement degree can also be modulated by tuning the corresponding surface functional groups of MXene nano-flakes, reflecting that the electron-withdrawing functional groups play a vital role in this charge transfer process. These findings offer promising clues to control the optoelectronic properties of TMDCs and expand the scope of the application of MXene nano-flakes, suggesting a possibility to construct a new heterostructure junction based on MXenes and TMDCs.

Such strategies, including substitutional doping during growth,¹⁶ back gate or top liquid gate,²⁰ vertical heterostructure and/or gas molecule adsorption,^{21,22} and band structure engineering,²³ had been applied for manipulating the carrier concentrations in monolayer TMDCs. Nevertheless, the requirement for complicated and precise technological processes is the main hindrance to obtain the desired optoelectronic properties of monolayer TMDCs, which becomes a great limitation to the fundamental study of monolayer TMDCs. These daunting challenges stimulate the exploration of an alternative method to manipulate the optoelectronic properties of TMDCs. Chemical doping via the drop casting method has offered a foray into a promising and feasible method to modify the charge carrier densities of two-dimensional (2D) semiconductors.19,24 Till date, only few known chemicals have been utilized to control the carrier density in monolayer TMDCs to regulate their optoelectronic properties.17,22,25,26

Recently, a new type of 2D transition metal carbide (MXene) has been produced and synthesized from the MAX phase.²⁷ The MAX phases are commonly represented by the chemical formula of $M_{n+1}AX_n$ (Fig. 1a), where A represents the main group element, X is C and/or N, M is an early transition metal, and n = 1, 2 or 3.²⁸ Owing to the wet etching of the MAX phase in a fluoride-containing solution to remove A element and then delaminate the layers, the surface of the resulting MXene is always terminated by a mixture of functional groups of -F, -O and/or -OH with the corresponding chemical formula of $M_{n+1}-X_nT_x$ (Fig. 1b), where T represents the terminated functional group, and x is the atomic percentage of terminated functional

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group.²⁹ MXenes have attracted exponential attention in diverse fields including supercapacitance,³⁰ electromagnetic interference (EMI) shielding,³¹ nonlinear photonics,^{32,33} Li-ion and other types of battery fabrication,³⁴ photo-detection,³⁵ energy storage,³⁶ photothermal therapy³⁷ as well as water purification³⁸ because of their unique chemical and physical properties. In addition, the superior hydrophilicity of MXenes allows it to be easily dispersed in solution, making it a promising candidate for solution processing. Notably, according to the recent theoretical studies, MXenes were theoretically indicated for controlling the electronic density of TMDCs via tuning the surface functional group; nevertheless, this idea has not yet been experimentally explored.³⁹ In this study, as a proof of concept, we deposited the MXene solution onto 1-layer (1L) tungsten disulfide (WS₂) via the drop-casting method to tune exciton/trion-involved optoelectronic properties of 1L WS₂ (Fig. 1c and d). MXene-decorated WS₂ successfully achieved a \sim 5-fold PL enhancement than the pristine WS₂, which could have originated from the effective charge transfer from WS₂ to MXenes induced by a strong electron-withdrawing effect of the attached surface functional group. We also reveal that this electron-withdrawing ability of MXenes can be well modulated by tuning its corresponding surface functional group via a postetch annealing method. This study not only extends the application area of MXenes, but also develops a simple and green strategy to achieve bright photoluminescence with high quantum yields in monolayer WS₂.

2. Experimental section

2.1. Synthesis of $Ti_3C_2T_x$ MXenes

MAX phase Ti_3AlC_2 with a purity of 98% was purchased from Beijing Forsman Pte Ltd. F- $Ti_3C_2T_x$ MXenes were made by

2.3. Mxene doping

MXenes were dispersed in an isopropanol (IPA) solution at a concentration of $\sim 10 \text{ mg mL}^{-1}$. The dopant of MXenes was deposited *via* a drop-casting method. The volume of doping was



etching MAX-phased Ti₃AlC₂ in 48% HF at 60 °C for over 24 h. After etching, the samples were washed and centrifuged (3500 rpm) for 20 min with DI water for 5 times to obtain pH \ge 5 in an aqueous phase. The resulting solution was vacuum-filtered using a micropore PTFE 0.2 µm filter to obtain the MXene powder. The delamination of MXene flakes was performed by hand-shaking of the sample for 5 min, followed by centrifugation at 3500 rpm for 1 h.³⁵ The resulting concentrated aqueous colloidal solution of Ti₃C₂T_x MXene sheets was used for the doping of 1 L WS₂. Due to the etching process, Ti₃C₂ MXene contains fluorine elements; hence, the resulting materials will be referred to as F-Ti₃C₂T_x in this study. N-doped Ti₃C₂T_x was prepared using the post-annealing method. F-Ti₃C₂T_x was placed in a tube CVD in a NH₃/Ar atmosphere at 400 °C for 4.5 h to obtain N-Ti₃C₂T_x.

2.2. Characterization

The Raman spectra, PL spectra and mapping were conducted using a WITec alpha 300 R system at excitation wavelengths of 488 and 532 nm, respectively. In order to avoid heating and optical doping effect, the power of the laser was kept as low as 10 μ W for the room temperature PL measurement. The AFM measurement was performed using a cyphers. The crystal structure, sample size, morphologies, and surface element were determined *via* X-ray diffraction (XRD), transmission electron microscopy (TEM) and X-ray photoelectron spectroscopy (XPS).



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3. Result and discussion

3.1. MXene nano-flake preparation and characterization

The MXene $(F-Ti_3C_2T_x)$ nano-flakes can be conventionally prepared via a sequence of etching and delamination processes (see Experimental section). The typical morphology of MXene nano-flakes was characterized via transmission electron microscopy (TEM). As shown in Fig. 2a, the thin $F-Ti_3C_2T_x$ nano-flakes are dispersed as a stack of few layers. To further characterize the quality of the exfoliated $F-Ti_3C_2T_x$ nano-flakes, X-ray diffraction (XRD) and Raman spectroscopy were also performed. The XRD pattern (Fig. 2b, bottom, black line) of Ti_3AlC_2 MAX phase shows the typical (002), (004), (101), (103), (104), (105), and (107) XRD peaks at 2θ degree of 9.78°, 19.47°, $34.10^\circ,\,36.95^\circ,\,39.30^\circ,\,42.07^\circ,$ and $48.88^\circ,$ respectively, which is in accordance with previous reports.40 The disappearance of a strong (104) peak at $2\theta = 39.30^{\circ}$ demonstrates the successful preparation of $F-Ti_3C_2T_x$ nano-flakes through the removal of Al atom (Fig. 2b, top, red line). Furthermore, the as-prepared F- $Ti_3C_2T_x$ nano-flakes showed a lower 2θ degree (9.14°) and a broader peak compared to that of Ti₃AlC₂, which suggests a change in the *c*-lattice parameter and decreased the structure order in $F-Ti_3C_2T_x$ nano-flakes, reconfirming the successful delamination.41 The Raman spectrum shown in Fig. 2c exhibits several vibration modes with peaks at 262.91, 403.46, 520.41 and 597.47 cm⁻¹, which could be assigned to nonstoichiometric

titanium carbides.42 Additional peaks at 1347.58 and 1573.93 cm⁻¹ could have originated from the D and G bands of carbon layers in the as-prepared F-Ti₃C₂T_x nano-flakes.⁴² X-ray photoelectron spectroscopy (XPS) investigations were conducted to characterize the surface chemical compositions of the asprepared $F-Ti_3C_2T_r$ nano-flakes. The full-spectrum XPS (Fig. 2d) reveals several elements including titanium, carbon, oxygen and fluoride, demonstrating the presence of -O, -F and -OH terminations on the surface of F-Ti₃C₂T_x nano-flakes.^{41,43} The high-resolution XPS spectra of F-Ti₃C₂T_r nano-flakes in F regions can be fitted using two peaks centred at 683.99 and 682.5 eV, indicating the existence of both Ti-F and Al-F, respectively. The O 1s core level can be deconvoluted into four components at 533.80, 532.91, 531.84 and 530.73 eV, which can be attributed to H₂O, Ti-OH, C-Ti(OH)_r and C-Ti-O_r, respectively. The XPS spectrum of Ti 2p (Fig. S1[†]) verifies the presence of Ti-C, Ti-O and Ti-F, respectively. Moreover, the highresolution C 1s XPS spectrum shows three fitting peaks, corresponding to C-Ti, C-C, and C-O, respectively (Fig. S2[†]).

3.2. Preparation and characterization of tungsten disulfide

To demonstrate the exciton modulation of 1L WS₂ *via* MXene deposition, we prepared a monolayer WS₂. The standard micromechanical method was used to exfoliate 1L-WS₂ from bulk crystals *via* a Scotch tape and transferred onto a silicon (Si) wafer with a 280 nm thick thermal oxidation (SiO₂) coated layer.⁴⁴ The exfoliated WS₂ flakes were determined by an optical microscope image, Raman spectrum, atomic force microscopy



Fig. 2 (a) TEM images, (b) XRD patterns and (c) Raman spectrum of $F-Ti_3C_2T_x$ nano-flakes. (d–f) XPS spectra of $F-Ti_3C_2T_x$ nano-flakes in (d) full spectrum region, (e) F 1s region and (f) O 1s region. Shaded colored areas in (e) and (f) represent curve-fitting results. The scale bar in (a) is 50 nm.

(AFM) and PL measurement to determine their exact layer number. The typical optical microscope image of 1L-WS₂ flakes on SiO₂/Si is shown in Fig. S3a.[†] The thickness of exfoliated flakes is revealed by the AFM height profile based on a noncontact mode, and measured along the white dot line crossing the flakes. The thickness of the 1-layer WS₂ is about 0.8 nm, as shown in Fig. S3b and c,† in which the slight drop in the boundaries may be induced by the absorbed residues and/or molecules. The Raman spectrum (Fig. S4[†]) of exfoliated flake excited with a 488 nm laser exhibits two strong peaks at about 355 and 417 cm⁻¹, which can be defined as in-plane E_{2g}^1 and outof-plane A_{1g} vibration modes, respectively. The difference (Δ) between A_{1g} and E_{2g}^1 peaks is ~62.3 cm⁻¹, ($\Delta = A_{1g}$ – E_{2g}^{1}), confirming the monolayer feature of the exfoliated WS₂ flake.45 In addition, the PL mapping measurement was adapted to verify the monolayer characteristics of the exfoliated flakes. The bright emission of the monolayer WS₂ (Fig. S5a⁺) is observed in the dot area in Fig. S3a.[†] However, on the account of the transition from direct band gap monolayer to indirect band gap multi-layer, the PL intensity of other area in Fig. S3a⁺ cannot be observed.17 It should be noted that the intensities of the PL mapping are slightly uneven, which may be caused by the non-uniformity natural n-doping effect in exfoliated 1L WS₂.

The weaker PL emission area should have a higher n-doping level compared to the higher PL intensity area, which would lead to the increased emission from low-intensity negative trions formed *via* the combination of excess electrons and neutral excitons (Fig. S5b†) in the extracted PL spectrum.⁴⁶ Moreover, the spectrum exhibits a peak at around 1.97 eV for monolayer WS₂ (Fig. S5b†), which is similar to that reported previously. As the strong spin–orbit interaction induced energy difference of A and B excitons in the emission spectrum is as large as ~0.4 eV, the emission from B exciton was not detected under a 532 nm laser excitation.⁴⁷

3.3. Exciton/trion modulation via MXene doping

To test the superiority of F–Ti₃C₂T_x nano-flakes in controlling the carrier density of WS₂, F–Ti₃C₂T_x nano-flakes were deposited onto 1L WS₂ by drop-casting. A direct evidence to indicate the successful electron modulation in WS₂ is the enhanced PL emission. Pristine WS₂ shows a weak PL emission with a peak centred at ~1.97 eV (Fig. 3a). The low emission intensity and relatively small peak energy could have originated from the heavy n-doping feature in pristine 1L WS₂.⁴⁶ After the coupling of F–Ti₃C₂T_x and 1L WS₂, the PL intensity of 1L WS₂/F–Ti₃C₂T_x has remarkably increased, which is almost ~5.0 folds more



Fig. 3 (a) PL spectra of 1L WS₂ and 1L WS₂/F-Ti₃C₂T_x measured at room temperature. The doping of F-Ti₃C₂T_x nano-flakes clearly modified the PL intensity and peak positions. (b) Fitted PL spectra of 1L WS₂ and 1L WS₂/F-Ti₃C₂T_x. Note that the peaks in PL spectra were fitted using the Lorentzian fitting by assuming two peaks of the neutral exciton (X) and trion (X⁻) peaks. (c) The integrated PL intensity of the neutral excitons (I_X), trions (I_{X^-}), and the sum (I_{total}) of I_X and I_{X^-} of 1L WS₂ and 1L WS₂/F-Ti₃C₂T_x, respectively. (d) Raman spectra of 1L WS₂ and 1L WS₂/F-Ti₃C₂T_x. A slightly blue shift by ~2.0 cm⁻¹ was observed in the A_{1q} mode of 1L WS₂/F-Ti₃C₂T_x compared to that of pristine 1L WS₂.

than that of pristine 1L WS₂ (Fig. 3c). Meanwhile, the peak position is also blue shifted to \sim 2.02 eV. To clearly figure out the variation in PL properties caused by F-Ti₃C₂T_x nano-flake doping, the Lorentzian fitting was performed. The fitted results of 1L WS₂/F-Ti₃C₂T_x and 1L WS₂ are shown in the top and bottom panels of Fig. 3b, respectively. The peaks of $1L WS_2$ $F-Ti_3C_2T_x$ and 1L WS₂ can be well fitted by assuming two peaks of neutral exciton (X) and negative trion (X⁻) peaks.²⁵ According to the fitted results, the dominant peak is assigned to the negative trion with a spectral weight (I_{X^-}/I_{total}) of up to ~0.82 (Fig. 3b and c). It should be noted that the neutral exciton peak becomes predominant and the weight of the negative trion of 1L WS₂/F-Ti₃C₂T_x is greatly decreased to \sim 0.18. The variation in the PL intensity and peak position could be ascribed to the transition from negative trions to neutral excitons after F- $Ti_3C_2T_x$ nano-flake doping. The exciton transition could be well explained by the p-type doping, which can significantly decrease the electron density of 1L WS2.24 According to the previous report, the A1g vibration mode of 1L WS2 Raman spectrum is very susceptible to electron coupling. The decrease in the electron density could cause blue-shifting of the A1g vibration mode.17 Therefore, in order to confirm the change in the electron density in 1L WS₂/F-Ti₃C₂T_x, the Raman spectra of 1L WS₂/ $F-Ti_3C_2T_x$ was measured. As shown in Fig. 3d, the Raman spectrum of 1L WS₂/F-Ti₃C₂T_x shows a blue-shifted A_{1g} vibration mode of $\sim 2 \text{ cm}^{-1}$, indicating the decrease in the electron density of 1L WS₂/F-Ti₃C₂T_x. Consequently, the variation in PL and Raman properties could originate from the withdrawal of electrons from 1L WS₂ to $F-Ti_3C_2T_x$ nano-flakes owing to the surface electron-withdrawing group and physically and/or

chemically adsorbed $\rm H_2O$ of $\rm F\text{--}Ti_3C_2T_{\it x}$ nano-flakes (Fig. 1b and d). 17,45

3.4. Effect of surface functional group of MXenes on PL enhancement

In order to figure out the vital role of the surface electronwithdrawing group and adsorbed H₂O of MXene nano-flakes on the enhancement of the PL emission of 1L WS₂, a control experiment was performed using an as-synthesized nitrogen (N)-substituted MXene $(N-Ti_3C_2T_x)$ as a dopant. $N-Ti_3C_2T_x$ was constructed by annealing HF-etched $Ti_3C_2T_x$ in ammonia at 700 °C and characterized by XPS and Raman spectra.40 The fullspectrum XPS reveals several elements, including titanium, carbon, oxygen and nitrogen, demonstrating the presence of -O, -OH and -N terminations on the surface of N-Ti₃C₂T_x nanoflakes (Fig. 4a). It should be noted that the strong electronwithdrawing F functional group was successfully eliminated at high temperatures, which would significantly decrease the electron-withdrawing ability of N-Ti₃C₂T_x nano-flakes.⁴⁸ The high-resolution XPS spectra of the N-Ti₃C₂T_x nano-flakes in N regions can be fitted using two peaks centred at 401.45 and 399.25 eV (Fig. 4b), indicating the existence of both C-N and Ti-N, respectively. The O 1s core level can be deconvoluted into two components at 532.21 and 530.73 eV, which could be attributed to $C-Ti(OH)_x$, and $C-Ti-O_x$, respectively (Fig. 4c). This demonstrates the removal of physically and/or chemically adsorbed H_2O at high temperatures, suggesting a further decrease in the electron-withdrawing ability of N-Ti₃C₂T_x nano-flakes. The Raman spectrum of N-Ti₃C₂T_x nano-flakes is similar to that of F-Ti₃C₂T_x nano-flakes owing to the preserved Ti₃C₂ structure



Fig. 4 (a-c) X-ray photoelectron (XPS) spectra of $F-Ti_3C_2T_x$ nano-flakes in (a) full-spectrum region, (b) N 1s region and (c) O 1s region. Shaded colored areas in (b) and (c) represent curve-fitting results. (d) Raman spectrum of $N-Ti_3C_2T_x$ nano-flakes. (e) PL spectra of 1L WS₂ and 1L WS₂/N-Ti₃C₂T_x measured at room temperature. (f) The integrated total (I_{total}) PL intensity of 1L WS₂ and 1L WS₂/F-Ti₃C₂T_x.

(Fig. 4d). It is worth noting that the D and G bands almost disappear, indicating the disordered carbon sheets of N-Ti₃C₂T_x nano-flakes induced by the N-intercalation. To verify the evolution of the PL emission of 1L WS₂ *via* N-Ti₃C₂T_x nano-flake doping, the PL spectrum of 1L WS₂/N-Ti₃C₂T_x was measured. As shown in Fig. 4e, although 1L WS₂/N-Ti₃C₂T_x also shows an improved PL emission with an exciton transition from negative trion to neutral exciton, the PL enhancement is much lower (~2.41 folds) compared to that of 1L WS₂/N-Ti₃C₂T_x (Fig. 4f). The decreased PL enhancement of 1L WS₂/N-Ti₃C₂T_x could be attributed to the decreased charge transfer from 1L WS₂ to N-Ti₃C₂T_x owing to the elimination of a strong electron-withdrawing functional group of F and adsorbed H₂O.^{17,45}

4. Conclusions

In summary, we succeeded in achieving the p-type doping of 1L-WS₂ via surface functional group-modified MXene nano-flakes. Upon doping with MXene nano-flakes, the blue-shifted A1g vibration mode and enhanced PL efficiency of neutral excitons can be obviously observed. These findings can serve as a direct proof of the regulated optoelectronic properties of 1L WS₂, demonstrating that the p-doping effect is dominant in surface functional group-modified MXene nano-flake-doped 1L-WS₂. The modulation of the optoelectronic properties of 1L WS₂ by drop-coating with MXene nano-flakes could be ascribed to the effective charge transfer from 1L WS₂ to MXenes induced by strong electron-withdrawing characteristics inherited from the surface fluorine functional group and/or adsorbed H₂O presented on the MXene nano-flakes. In addition, the electronwithdrawing ability of MXenes can be facilely controlled by tuning the type of surface functional groups. These findings suggest that 2D MXenes are promising materials for regulating the optical and electronic properties of 1L-WS₂, offering a new avenue to construct new heterostructure junctions, such as high-performance photodetectors, ambipolar transport field effect transistors, and flexible solar energy harvesting devices, and broadening the horizon for harnessing the unique properties of MXenes.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- M. Amani, D. H. Lien, D. Kiriya, J. Xiao, A. Azcatl, J. Noh, S. R. Madhvapathy, R. Addou, S. KC, M. Dubey, K. Cho, R. M. Wallace, S. C. Lee, J. H. He, J. W. Ager, X. Zhang, E. Yablonovitch and A. Javey, *Science*, 2015, **350**, 1065–1068.
- 2 W. Choi, N. Choudhary, G. H. Han, J. Park, D. Akinwande and Y. H. Lee, *Mater. Today*, 2017, **20**, 116–130.
- 3 S. Manzeli, D. Ovchinnikov, D. Pasquier, O. V. Yazyev and A. Kis, *Nat. Rev. Mater.*, 2017, **2**, 17033.
- 4 X. Yu, P. Yu, D. Wu, B. Singh, Q. Zeng, H. Lin, W. Zhou, J. Lin,
 K. Suenaga, Z. Liu and Q. J. Wang, *Nat. Commun.*, 2018, 9, 1545.
- 5 X. Duan, C. Wang, A. Pan, R. Yu and X. Duan, *Chem. Soc. Rev.*, 2015, **44**, 8859–8876.
- 6 C. L. Tan and H. Zhang, Chem. Soc. Rev., 2015, 44, 2713-2731.
- 7 J. Pu and T. Takenobu, Adv. Mater., 2018, 30, 1707627.
- 8 S. C. Dhanabalan, J. S. Ponraj, Z. Guo, S. Li, Q. Bao and H. Zhang, *Adv. Sci.*, 2017, 4, 1600305.
- 9 S. Wang, A. Robertson and J. H. Warner, *Chem. Soc. Rev.*, 2018, **47**, 6764–6794.
- 10 J. S. Ross, S. Wu, H. Yu, N. J. Ghimire, A. M. Jones, G. Aivazian, J. Yan, D. G. Mandrus, D. Xiao, W. Yao and X. Xu, *Nat. Commun.*, 2013, 4, 1474.
- 11 Y. Zhang, C. Lim, Z. Dai, G. Yu, J. W. Haus, H. Zhang and P. N. Prasad, *Phys. Rep.*, 2019, **795**, 1–51.
- 12 L. J. Li, E. C. O'Farrell, K. P. Loh, G. Eda, B. Ozyilmaz and N. A. Castro, *Nature*, 2016, **529**, 185–189.
- 13 J. Pei, J. Yang, X. Wang, F. Wang, S. Mokkapati, T. Lü, J. Zheng, Q. Qin, D. Neshev, H. H. Tan, C. Jagadish and Y. Lu, *ACS Nano*, 2017, **11**, 7468–7475.
- 14 J. D. Lin, C. Han, F. Wang, R. Wang, D. Xiang, S. Qin, X. Zhang, L. Wang, H. Zhang, A. T. S. Wee and W. Chen, *ACS Nano*, 2014, 8, 5323–5329.
- 15 J. Pei, J. Yang, T. Yildirim, H. Zhang and Y. Lu, *Adv. Mater.*, 2019, **31**, 1706945.
- 16 Z. Hu, Z. Wu, C. Han, J. He, Z. Ni and W. Chen, *Chem. Soc. Rev.*, 2018, 47, 3100–3128.
- 17 N. Peimyoo, W. Yang, J. Shang, X. Shen, Y. Wang and T. Yu, *ACS Nano*, 2014, **8**, 11320–11329.
- 18 X. Chen and A. R. McDonald, *Adv. Mater.*, 2016, **28**, 5738–5746.
- 19 S. Zhang, H. M. Hill, K. Moudgil, C. A. Richter, A. R. Hight Walker, S. Barlow, S. R. Marder, C. A. Hacker and S. J. Pookpanratana, *Adv. Mater.*, 2018, **30**, 1802991.
- 20 C. Cong, J. Shang, Y. Wang and T. Yu, *Adv. Opt. Mater.*, 2016, 6, 1700767.
- 21 Z. Wang, Z. Dong, Y. Gu, Y. Chang, L. Zhang, L. Li, W. Zhao,
 G. Eda, W. Zhang, G. Grinblat, S. A. Maier, J. K. W. Yang,
 C. Qiu and A. T. S. Wee, *Nat. Commun.*, 2016, 7, 11283.
- 22 Y. I. Jhon, Y. Kim, J. Park, J. H. Kim, T. Lee, M. Seo and Y. M. Jhon, *Adv. Funct. Mater.*, 2016, **26**, 7551–7559.
- 23 Y. Sun, Z. Zhou, Z. Huang, J. Wu, L. Zhou, Y. Cheng, J. Liu,
 C. Zhu, M. Yu, P. Yu, W. Zhu, Y. Liu, J. Zhou, B. Liu,
 H. Xie, Y. Cao, H. Li, X. Wang, K. Liu, X. Wang, J. Wang,
 L. Wang and W. Huang, *Adv. Mater.*, 2019, 31, 1806562.

- 24 C. R. Ryder, J. D. Wood, S. A. Wells and M. C. Hersam, *ACS Nano*, 2016, **10**, 3900–3917.
- 25 S. Mouri, Y. Miyauchi and K. Matsuda, *Nano Lett.*, 2013, **13**, 5944–5948.
- 26 Y. Tao, X. Yu, J. Li, H. Liang, Y. Zhang, W. Huang and Q. J. Wang, *Nanoscale*, 2018, **10**, 6294–6299.
- 27 M. Ghidiu, M. R. Lukatskaya, M. Zhao, Y. Gogotsi and M. W. Barsoum, *Nature*, 2014, **516**, 78–171.
- 28 M. Naguib, V. N. Mochalin, M. W. Barsoum and Y. Gogotsi, *Adv. Mater.*, 2014, **26**, 992–1005.
- 29 H. Lin, Y. Chen and J. Shi, Adv. Sci., 2018, 5, 1800518.
- 30 M. Hu, T. Hu, Z. Li, Y. Yang, R. Cheng, J. Yang, C. Cui and X. Wang, ACS Nano, 2018, 12, 3578–3586.
- 31 J. Liu, H. Zhang, R. Sun, Y. Liu, Z. Liu, A. Zhou and Z. Yu, *Adv. Mater.*, 2017, **29**, 1702367.
- 32 X. Jiang, S. Liu, W. Liang, S. Luo, Z. He, Y. Ge, H. Wang, R. Cao, F. Zhang, Q. Wen, J. Li, Q. Bao, D. Fan and H. Zhang, *Laser Photonics Rev.*, 2018, **12**, 1700229.
- 33 L. Wu, X. Jiang, J. Zhao, W. Liang, Z. Li, W. Huang, Z. Lin, Y. Wang, F. Zhang, S. Lu, Y. Xiang, S. Xu, J. Li and H. Zhang, *Laser Photonics Rev.*, 2018, **12**, 1800215.
- 34 M. Naguib, J. Come, B. Dyatkin, V. Presser, P. Taberna, P. Simon, M. W. Barsoum and Y. Gogotsi, *Electrochem. Commun.*, 2012, 16, 61–64.
- 35 S. Chertopalov and V. N. Mochalin, *ACS Nano*, 2018, **12**, 6109–6116.
- 36 B. Anasori, M. R. Lukatskaya and Y. Gogotsi, *Nat. Rev. Mater.*, 2017, 2, 16098.

- 37 C. Xing, S. Chen, X. Liang, Q. Liu, M. Qu, Q. Zou, J. Li, H. Tan, L. Liu, D. Fan and H. Zhang, ACS Appl. Mater. Interfaces, 2018, 10, 27631–27643.
- 38 L. Ding, Y. Wei, Y. Wang, H. Chen, J. Caro and H. Wang, Angew. Chem., Int. Ed., 2017, 56, 1825–1829.
- 39 Y. Liu, H. Xiao and W. A. Goddard, J. Am. Chem. Soc., 2016, 138, 15853–15856.
- 40 Y. Wen, T. E. Rufford, X. Chen, N. Li, M. Lyu, L. Dai and L. Wang, *Nano Energy*, 2017, 38, 368–376.
- 41 S. Tu, Q. Jiang, X. Zhang and H. N. Alshareef, *ACS Nano*, 2018, **12**, 3369–3377.
- 42 K. Krishnamoorthy, P. Pazhamalai, S. Sahoo and S. Kim, J. Mater. Chem. A, 2017, 5, 5726–5736.
- 43 S. J. Kim, H. Koh, C. E. Ren, O. Kwon, K. Maleski, S. Cho, B. Anasori, C. Kim, Y. Choi, J. Kim, Y. Gogotsi and H. Jung, ACS Nano, 2018, 12, 986–993.
- 44 X. Yu, S. Zhang, H. Zeng and Q. J. Wang, *Nano Energy*, 2016, 25, 34–41.
- 45 H. M. Oh, H. Jeong, G. H. Han, H. Kim, J. H. Kim, S. Y. Lee, S. Y. Jeong, S. Jeong, D. J. Park, K. K. Kim, Y. H. Lee and M. S. Jeong, ACS Nano, 2016, 10, 10446–10453.
- 46 P. K. Chow, R. B. Jacobs-Gedrim, J. Gao, T. Lu, B. Yu, H. Terrones and N. Koratkar, ACS Nano, 2015, 9, 1520–1527.
- 47 B. Zhu, H. Zeng, J. Dai and X. Cui, *Adv. Mater.*, 2014, 26, 5504–5507.
- 48 S. Lai, J. Jeon, S. K. Jang, J. Xu, Y. J. Choi, J. Park, E. Hwang and S. Lee, *Nanoscale*, 2015, 7, 19390–19396.