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Continuous focusing, fractionation and extraction of anionic analytes in a microfluidic chip†

Vasileios A. Papadimitriou, * Loes I. Segerink  and Jan C. T. Eijkel

Electrokinetic focusing and separation methods, specifically ion concentration polarization focusing (ICPF), provide a very powerful and easy to use analytical tool for several scientific fields. Nevertheless, the concentrated and separated analytes are effectively trapped inside the chip in picoliter volumes. In this article we propose an ICPF device that allows continuous and selective extraction of the focused analytes. A theoretical background is presented to understand the dynamics of the system and a 1D model was developed that describes the general behavior of the system. We demonstrate the selective extraction of three fluorescent model anionic analytes and we report selective extraction of the analytes at a 300-fold increased concentration.

Introduction

Electrokinetic separation methods are powerful tools that have found application in a wide range of fields, from biological/chemical research to food industry and forensics. Since the early days of microfluidics many electrokinetic separation methods have been translated to lab-on-chip systems. A special class of electrokinetic techniques combines separation and focusing such as isoelectric focusing (IEF), electric field gradient focusing (EFGF) and isotachopheresis (ITP). These techniques require an application-specific electrolyte preparation which can be a tedious process and significantly reduce their usability for point-of-care systems. For example, ITP requires a leading and trailing electrolyte with electrophoretic mobilities higher and lower than the target analyte. In recent years ion concentration polarization focusing (ICPF) was introduced, which is capable of concentrating and separating analytes without the use of specific electrolytes. ICPF focuses analytes in the electric field gradient created by the phenomenon of ion concentration polarization (ICP).

ICPF was introduced by Wang *et al.*¹ more than a decade ago and it was capable of concentration factors of proteins and peptides in the order of millions. Since its introduction vast amount of research has been performed with this technique, where Quist *et al.*² demonstrated that apart from concentration also separation occurred.^{3–7,34} Though the technique offers a very powerful and simple analytical tool, a

disadvantage is that the focused and separated analytes are “trapped” inside the chip in picoliter volumes. Further analysis, for example by mass spectrometry, thus is impossible. Attempts to extract the preconcentrated analytes have been reported, but they require the use of Quake valves⁸ or magnetically actuated valves.⁹ Continuous extraction of the concentrated biomolecules and cells in non-selective fashion has been demonstrated by Kwak *et al.*¹⁰ In this article we propose an ICPF device modified to selectively and continuously extract the fractionated and concentrated analytes from their complex background, allowing further downstream analysis.

Theory

Concentration polarization and establishment of a depletion zone

In order to understand the dynamics of the device we will first briefly describe the process of ICPF. For a more extensive theoretical description of ICPF the reader is referred to the recent report by Ouyang *et al.*¹¹ In addition the feasibility of extraction of concentrated (cationic) analytes *via* ICPF was reported *via* numerical simulations of similar systems.^{12,13} ICPF is based on differential ionic migration in an electric field gradient. The electrical field gradient is created by applying a steady potential difference over a microchannel which contains a region in the axial direction with a salt concentration gradient. As the current density must be equal in the entire system, an electric field gradient results, which can be formulated as

$$E(x) = J \cdot \rho(x) \quad (1)$$

with J [$A\ m^{-2}$] the current density, ρ [$\Omega\ m$] the electrical resistivity of the solution and x the axial coordinate. In order to

BIOS Lab on a Chip group, MESA+ Institute for Nanotechnology, Max Planck Centre for Complex Fluid Dynamics and Technical Medical Centre, University of Twente, The Netherlands. E-mail: v.papadimitriou@utwente.nl

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form the concentration gradient in ICPF, the phenomenon of ICP is used.¹⁴ A typical configuration to generate an ICP is by connecting the microchannel considered to one or two other microchannels *via* a “cation-permselective zone” (Fig. 1a).⁵ The term “cation-permselective zone” indicates an area where predominantly cations are able to enter and pass through, and the anion transport is greatly reduced. This cation-permselective zone, typically a nanochannel⁹ or a cation-permselective membrane such as Nafion,⁴ creates a flux gradient of the anions and cations in the direction of the electric field. This flux gradient removes cations and anions from the anodic side of the barrier resulting in a zone with low concentration of all ionic species, called the depletion zone.

Fig. 1a shows the separation channel connected to two buffer channels *via* a Nafion membrane, which was patterned by capillarity¹⁵ in a manner previously reported.⁵ The buffer channels are grounded while two different potentials are applied at the ends of the separation channel (V^u at upstream and V^d at downstream reservoir with $V^u > V^d$). This configuration will create a depletion zone in the separation channel and simultaneously apply an axial electrical field. At the interface between the depletion zone and the bulk solution an E-field gradient is generated, where anionic analytes will be focused (see below).

The novelty of our device is, that to perform continuous extraction, we added two extraction channels to this configuration, intersecting the separation channel perpendicularly (Fig. 1b). The focused analytes are extracted by applying a small negative pressure (P_{ext}) in these extraction channels. The system variables relevant for the extraction process are the analyte position, analyte extraction rate and analyte concentration rate. The two reservoir potentials (V^u and V^d) and the extraction pressure (P_{ext}) all influence these variables and their effects are coupled. To describe this complex system an approximative one-dimensional (1D) model was created. It consists of two coupled parts: a model for the electrical currents (Fig. 2, blue) and a model for the hydrodynamic fluxes in the system (Fig. 2, green). The model is limited in its scope as it does not provide a description of the formation of the different analyte zones in the E-field gradient zone. However, it serves as a model to understand the coupling of the electrical and hydrodynamic phenomena of the device.

Electric model (Fig. 2)

As mentioned above, the cation-permselective Nafion membrane creates a local gradient of the cation and anion fluxes, creating the depletion zone in the separation channel. For simplicity of the model we assume that the current is carried exclusively by the background electrolyte (not by the analytes) and that the cation and anion transport numbers in the bulk solution are equal. We also assume that the membrane is perfectly cation selective. We denote the salt concentration in the depletion zone as C_{dz} . Under these

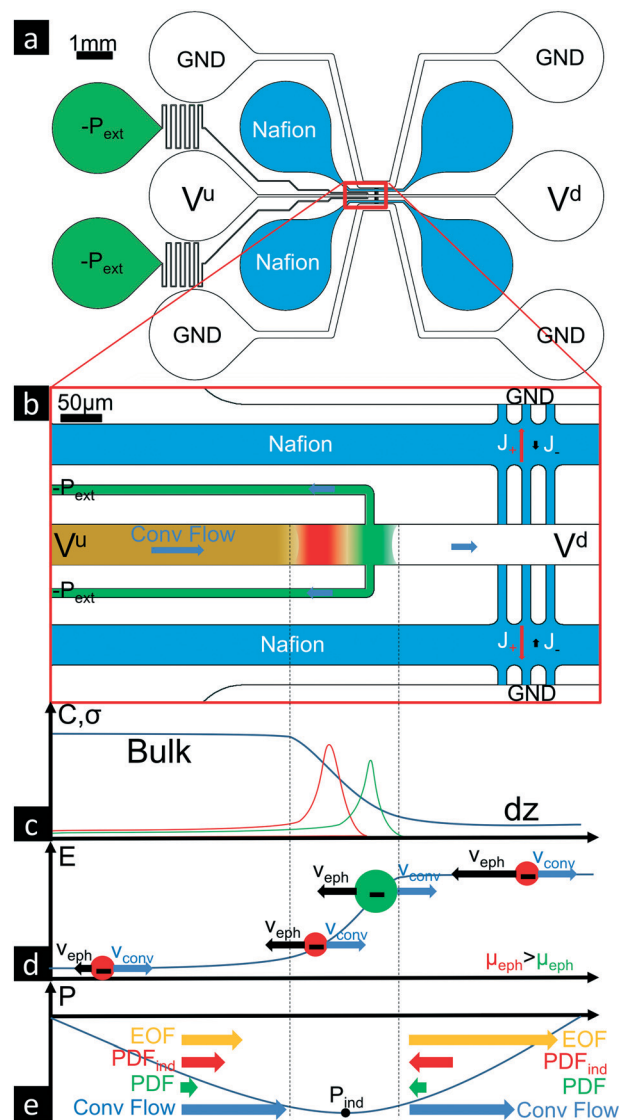


Fig. 1 Device outline (a) and concept of our ICPF device for extraction (b). (a) A potential is applied in the upstream (V^u) and downstream (V^d) reservoir with $V^u > V^d$. A low negative pressure is applied in the extraction channels marked with green. The side channels are filled with a buffer and they are grounded. A depletion zone (shown in white (b)) is formed due to ICP when an E-field is applied across an exclusion zone (Nafion) (b and c). At the interface between dz and the bulk (shown in orange) the concentration gradient (c) creates an E-field gradient zone when a second E-field is applied across it. All ions move towards the dz with a constant net flow (v_{conv}) but due to their difference in electrophoretic mobility (μ_{eph}) they experience different electrophoretic velocity (v_{eph}) in the E-field gradient zone resulting in separation. The analytes will start focusing at the position in the E-field that electrophoretic (v_{eph}) and bulk flow velocity (v_{conv}) are equal and opposite (d). Due to non-uniformities of the electro-osmotic flow (EOF) an induced pressure (P_{ind}) is created at the focusing location (e). 1D area averaged flow velocities. The focused analytes can be placed at the extraction channel and will be extracted under low negative pressure, which needs to be lower than the induced pressure ($P_{ext} < P_{ind}$) note: on figure (e) the EOF arrow accounts for both the primary and secondary EOF due to the extended space charge close to the Nafion membrane.¹¹



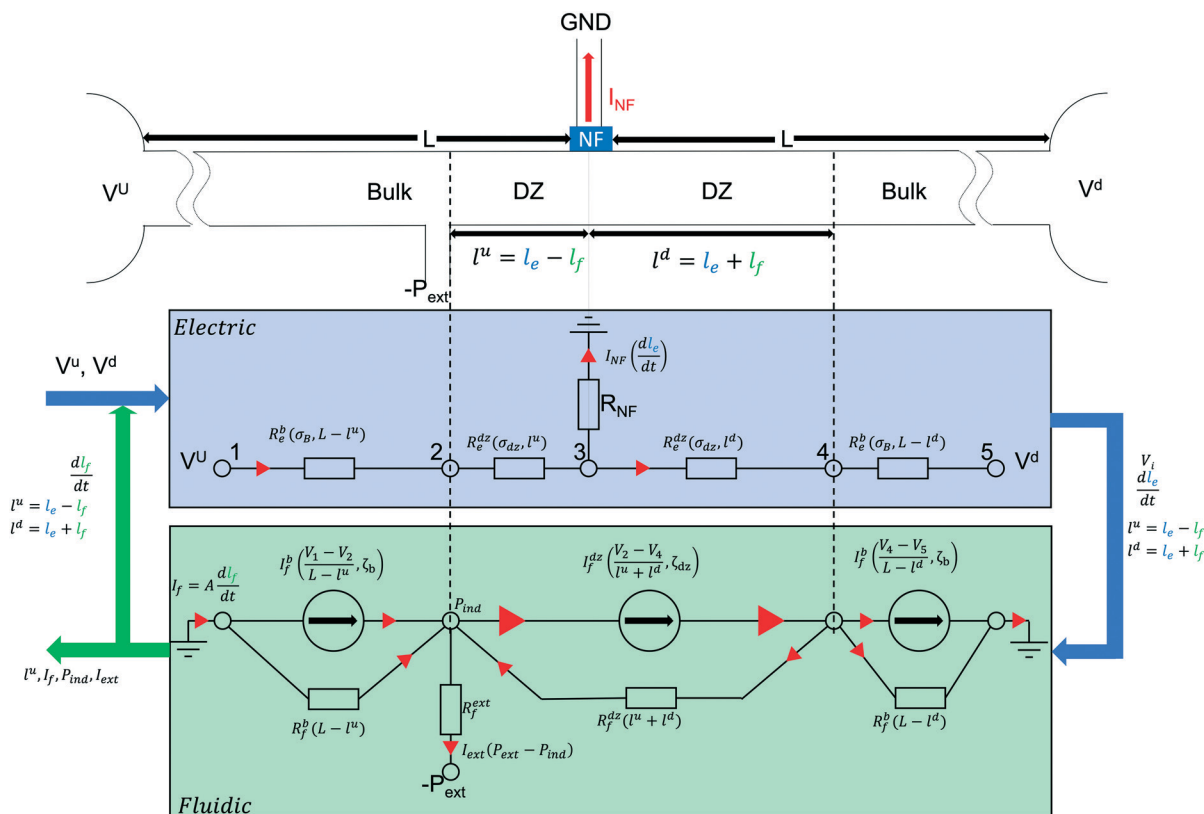


Fig. 2 Schematic of the 1D model for our ICPF extraction device. Each resistance and current is a function of the parameters stated in the brackets next to it. The electric model (blue) describes the electric response of the system and the growth of the depletion zone

($\frac{dl^c}{dt}$) due to ICP. The inputs for the electric model are the actuation potentials and some system parameters such as dimensions, conductivities and concentrations. The outputs, which are $\frac{dl^c}{dt}$ and all the potentials (in each node V_i) between bulk and depletion zone

interfaces, are fed to the fluidic model (green). The fluidic model (green) with the use of the potentials and zone sizes (provided by the electric model) calculates the E-fields and EOFs (represented as current sources) in each zone. The unequal current sources induce a secondary current (representing induced pressure-driven flow) through the hydraulic resistors (R_f^i). The total current ($\frac{dl^f}{dt}$) (i.e. bulk flow = EOF +

PDF + PDF_{ind}) is calculated and updates the zone sizes which is fed back to electric circuit. The process is repeated until convergence and the converged solution for focusing position (l^u), bulk flow in the separation channel (concentration rate) and extraction flow for each time

step is exported.

conditions, the growth rate of the depletion zone length in the channel, dl_e/dt [m s^{-1}], depends on the current running through the Nafion membrane as

$$\frac{dl_e}{dt} = \frac{I}{2(C_b - C_{dz})FA_c} \quad (2)$$

where I [A] is the electric current through the membrane, F [C mol^{-1}] is the Faraday constant, C_b [mol m^{-3}] is the bulk ion concentration and A_c [m^2] is the cross-sectional area of the separation channel. The growth rate of the depletion zone thus scales linearly with the electric current through the Nafion membrane.

The conductivity in the depletion zone is generally measured to be several orders of magnitude lower than in the bulk,⁴ and hence the total resistance of the system increases when the depletion zone grows and occupies an increasing part of the channel. Since constant potentials are applied to the reservoirs, the current through the channel and the Nafion membrane decreases as the depletion zone grows, resulting in a reduced growth rate of this zone as time proceeds. The electric model shown in Fig. 2 was created to calculate at any moment in time the size of each zone (bulk, depletion zone) both upstream and downstream from the Nafion membrane, as well as the various currents and the potentials at each interface between the zones.



Fluidic model (Fig. 2)

The fluxes of the analyte ions i , J_i [mol m⁻² s⁻¹] are given by the Nernst–Planck equation. We formulate the fluxes as a function of axial position x as:

$$J_i(x) = J_{i,\text{diff}}(x) + J_{i,\text{conv}}(x) + J_{i,\text{eph}}(x) \quad (3)$$

Here $J_{i,\text{diff}}$, $J_{i,\text{conv}}$ and $J_{i,\text{eph}}$ are the flux contributions of diffusion, convection and electrophoresis of analyte ions i , respectively. In our initial analysis we neglect diffusion for reasons of simplicity ($J_{i,\text{diff}}(x) = 0$).

We will now discuss the separate contributions to the flux of analyte i in eqn (3) in more detail.

The contribution by convection is

$$J_{i,\text{conv}}(x) = v_{\text{conv}}C_i(x) \quad (4)$$

Here C_i [mol m⁻³] is the concentration of the species i and v_{conv} [m s⁻¹] the linear convective liquid velocity assuming constant channel cross section. It has been demonstrated that the area-averaged bulk flow velocity in a channel can be obtained by a linear combination of pressure-driven (v_{PDF} [m s⁻¹]) and electroosmotic flow (v_{EOF} [m s⁻¹]).¹⁶ In our channel a pressure distribution is created by the suction that is applied at the extraction channels, as well as by the axial variation of EOF stemming from the conductivity gradient, that by the requirement of mass conservation locally generates a restoring pressure-driven flow. Below we will briefly treat both contributions.

Electroosmotic flow. The EOF magnitude varies along the channel length, as it is determined by the local electric field, which results from the potential difference between upstream (V^u) and downstream (V^d) reservoirs and the local resistivity (eqn (1)). At any location along the channel the magnitude of the EOF follows from the Helmholtz–Smoluchowski equation:

$$v_{\text{EOF}}(x) = \frac{\varepsilon\varepsilon_0\zeta}{\eta} E(x) \quad (5)$$

where η [kg m⁻¹ s⁻¹] is the viscosity of the liquid, ζ [V] is the zeta potential (assumed constant in the model) and ε , ε_0 [F m⁻¹] are the relative and vacuum permittivity. In our system EOF generally is the dominant convective velocity so the net convective velocity vector follows the EOF direction. Since we have a negative zeta potential and $V^u > V^d$ the bulk solution flows from the upstream to the downstream reservoir. In the depletion zone of the separation channel, the high E-field results in a high EOF velocity while in the remainder of the channel the low E-field results in a low EOF velocity. EOF is also significantly affected by the zeta potential (ζ) which is directly influenced by the pH. Mogi *et al.*¹⁷ reports a significant change of the pH close the depletion zone. Here we will neglect this effect.

EOF-induced pressure driven flow. Since aqueous electrolytes are non-compressible and because of mass conserva-

tion, a pressure-driven flow is induced towards the interface between the bulk and the depletion zone, where a negative pressure is generated. An EOF-induced pressure is for example also reported by Herr *et al.*¹⁶ in the case of channels with a non-uniform zeta potential.

In addition, close to the membrane a space charge region exists due to ICP.^{18–22} The space charge region has a large net positive charge, so the action of the electrical field in the bulk is to move it to the cathodic reservoir (same direction as the main EOF). Because this EOF (usually named electroosmosis of the second kind¹¹) is so much larger than anywhere else in the channel, we get the PDF back-flow vortices. A small part of the extra ‘push’ on the solution is translated to a PDF towards the cathode. Summarizing, the space charge region at the Nafion acts as a small and strong pump.

The device is operated in the over-limiting current regime¹⁹ where these EOF induced vortices have a strong contribution to the local 2D flow velocity profile. The 2D flow velocity profile that will result from the combination of all PDF and EOF contributions is a rich scientific topic,^{4,11,23,24,30} which for simplicity we here reduce to an area-averaged 1D flow velocity (Fig. 1e).

Adding both contributions to the bulk flow velocity ($v_{\text{conv}} = v_{\text{EOF}}(x) + v_{\text{PDF}}(x)$), we can write the convective flux as:

$$J_{i,\text{conv}}(x) = v_{\text{EOF}}(x)C_i(x) + v_{\text{PDF}}(x)C_i(x) \quad (6)$$

Extraction pressure driven flow. The extraction channels are electrically floating; hence they have no contribution of EOF. We extract the analytes by applying an external pressure at the extraction channel intersections with the separation channel (P_{ext}), which pressure must be lower than the local EOF-induced negative pressure ($P_{\text{ext}} < P_{\text{ind}}$). It must be noted that without an applied extraction pressure, or when the applied pressure is lower than the induced pressure (P_{ind}), the flow is directed from the extraction channels towards the separation channel. In our system this effect can however be neglected, as the thin and long extraction channels have high hydrodynamic resistance and the induced and applied external pressures result in minute volume flows, approximately two orders of magnitude smaller than the EOF volume flow.

Analyte electrophoresis and focusing. Anions migrate in the opposite direction of the convective flow with an electrophoretic velocity ($v_{i,\text{eph}}$) which scales linearly with the local electric field ($E(x)$ [V m⁻¹]).

$$v_{i,\text{eph}}(x) = \mu_{i,\text{eph}}E(x) \quad (7)$$

Here $\mu_{i,\text{eph}}$ [m² s⁻¹ V⁻¹] is the electrophoretic mobility. At the interface between the depletion zone and the bulk a region exists with a concentration and thus E-field gradient. The further the analytes move up the electric field in this zone, the higher their electrophoretic velocity. Within a certain mobility window, each ionic species will have a specific point in the gradient where its electrophoretic velocity is equal and opposite to the convective flow velocity, at which



point it will be focused (Fig. 1d).² The analytes will thus be separated and create focused bands, sorting themselves out based on their electrophoretic mobilities with the lowest mobility ion closest to the depletion zone, where the E-field is highest.

Analyte extraction. By tuning the focusing location of a specific analyte by the applied potentials to the intersection point of separation channel and extraction channels, and by applying a low negative pressure at the extraction channels, we will be able to selectively extract the focused analytes whose mobilities fall within a specific electrophoretic mobility window. This mobility window will depend on the electric field gradient and the width of the extraction channel.

Fluidic model. The fluidic model (Fig. 2) derives the convective bulk flow in the main channel. An electric equivalent of the fluidic system is used where the volumetric flows, pressures and hydraulic resistances are represented by currents, potentials and electric resistances, respectively. Similar approaches have been reported before,^{25–29} but in our system the local magnitude of the EOF (generated by the local E-field) is provided by the coupled electric model. The EOF of the second kind can be accounted in the model *via* the use of high zeta potential in the upstream depletion zone. The EOF in each section of the channel (bulk, depletion zone) is modelled as a current source with a current magnitude that is calculated from the linear velocity (eqn (4)) multiplied by the cross-sectional area. The length of each zone, and the potential differences applied across it are provided by the coupled electric model. The zone lengths are denoted as l^u and l^d for the depletion zone, $L - l^u$ and $L - l^d$ for the bulk solution located upstream and downstream, where L is the distance from reservoir to the Nafion membrane (Fig. 2). In addition to the EOF, the model derives the induced hydrodynamic pressure by introducing feedback hydraulic resistances based on the length of each zone and the channel cross-sectional geometry. The feedback current that is induced through the resistors to equalize the difference between the current sources, results in a feedback voltage over the hydraulic resistances, which represents the induced hydrodynamic pressure. Finally, the extraction pressure and PDF are simulated *via* a potential source at the end of the extraction channel (extraction resistance, R_f^{ext}). The resulting current through the separation and extraction channel corresponds to the sum of the convective flows (EOF, PDF) described by eqn (5).

The basic principle of the model is shown in Fig. 2. An in-depth analysis is presented in the ESI† along with a description of the effect of the down-stream and up-stream potential on the size of the depletion zone (*i.e.* focusing location), on the concentration rate and on the maximum induced negative pressure. In the ESI† the model results are compared and explained based on previous theoretical and experimental works.^{16–21} The fitting parameters of the model are the zeta potential in the bulk and in the depletion zone, and the salt concentration in the depletion zone.^{32,33} For the latter we used the parameter “depletion factor” which describes how

many times the average concentration in the depletion zone is lower than the bulk. An estimate for the depletion factor was taken from literature.²⁴

Separation and extraction resolution

To determine the separation resolution of the system, we need to distinguish between two operational modes, namely peak mode and plateau mode, as can be found also in other focusing techniques such as ITP.^{2,35} A schematic of the two modes is shown in Fig. 3.

Operation in peak mode. In this mode (Fig. 3a) the analytes are in very low concentration compared to the bulk electrolyte so that we can assume that they do not contribute to the conductivity and the local E-field in the E-field gradient zone. This is the most common mode for typical applications. In our experiments for example the transport number of the analyte (Bodipy) is approximately 2×10^{-6} making its contribution to the electric field negligible. In the peak mode

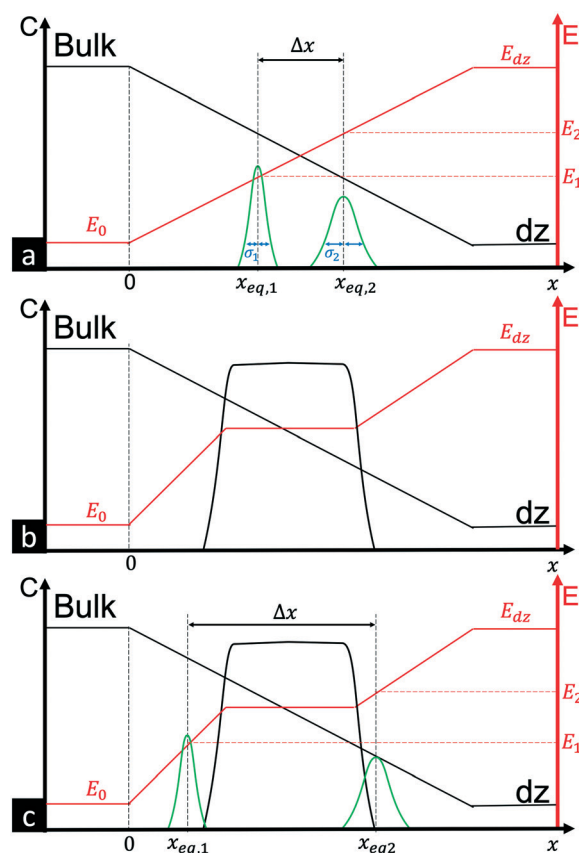


Fig. 3 a) Peak mode accumulation in the E-field gradient zone. The analytes are in low concentration, so they do not affect the local electric field (green concentration peaks not in scale). The analytes are focused at the E-field magnitude (E_1 , E_2) where they acquire a zero net velocity b) plateau mode. The analyte reaches the maximum concentration that satisfies the Kohlrausch regulation function (KRF) and forms a plateau at its concentration and the local electric field. c) Using a plateau mode analyte as an electrophoretic spacer. The plateau mode analyte pushes the peak mode analytes further apart. The peak mode analytes focus at the same E_1 , E_2 values as in (a).



the concentrating analytes form a Gaussian concentration profile. Since the early days of separation sciences^{36,37} the shape of the concentrated peak has been investigated. In our analysis we follow the approach from J. C. Giddings³⁸ where a focusing influence (in our case an electric field gradient) works against a defocusing influence (*i.e.* diffusion). We also assume a constant electric field gradient $\left(\frac{d^2E}{dx^2} = 0\right)$. This approach results in the concentration profile:

$$C_i(x) = C_{i,0} \exp \left(- \frac{\left(x - \frac{v_{\text{conv}} - E_0 \mu_i}{\frac{dE}{dx}} \mu_i \right)^2}{2 \frac{V_T}{z_i \frac{dE}{dx}}} \right) \quad (8)$$

where z_i is the valence of the analyte i , C_0 [mol m⁻³] is the maximum concentration of the focused analyte, x [m] is the direction along the separation channel with $x = 0$ denoting the location between bulk and depletion zone where the E-field gradient starts and E_0 the electric field in the bulk solution. Furthermore V_T [V] is the thermal potential ($V_T = k_b T/e$ with T [K] the temperature and k_b [J K⁻¹] Boltzmann's constant). Eqn (8) describes a Gaussian profile with mean

$$x_{\text{eq},i} = \frac{v_{\text{conv}} - E_0 \mu_{i,\text{eph}}}{\frac{dE}{dx}} \mu_{i,\text{eph}} \quad (9)$$

and variance σ^2 [m²]

$$\sigma_i^2 = \frac{V_T}{z_i \frac{dE}{dx}} \quad (10)$$

The derivation of eqn (8) and (9) is presented in the ESI.† The focusing location of species i is given by eqn (9). In peak mode we then obtain the resolution R_s

$$R_s = \frac{\Delta x}{2(\sigma_1 + \sigma_2)} \quad (11)$$

as

$$R_s = \frac{v_{\text{conv}}}{2\sqrt{\frac{dE}{dx}} V_T} \cdot \frac{\left(\frac{1}{\mu_1} - \frac{1}{\mu_2} \right)}{\left(\frac{1}{\sqrt{z_1}} + \frac{1}{\sqrt{z_2}} \right)} \quad (12)$$

In eqn (12) the first term is device- and actuation potential-dependent while the second term depends on ana-

lyte properties. As shown in eqn (9) the focusing location is inversely proportional to the mobility. Analytes with low electrophoretic mobility will thus focus closer to the depletion zone where the E-field is high. It follows from eqn (12) that the resolution of such slower analytes, for constant $\Delta\mu = \mu_1 - \mu_2$, will be better than for the faster analytes that focus at a location with lower E-field.

If the analytes are in peak mode, the mobilities of the analytes that are extracted from their focusing position by an extraction channel of determined width lie in a mobility window that can be calculated by equating Δx to the extraction channel width. If the mobility difference between two analytes in peak mode is smaller than that mobility window for a specific extraction channel width, then the extraction of a single analyte is impossible, and adjacent analyte(s) will be extracted along at varying concentration.

Finally, only analytes within a specific mobility window will focus in the device, which window is given by the maximum and minimum electric field and the convective flow velocity. As the analytes require an electrophoretic velocity equal and opposite to the convective flow hence this mobility window (assuming that the analyte i is peak mode and does not contribute to the local conductivity and electric field) is given by

$$\frac{v_{\text{conv}}}{E_0} > \mu_i > \frac{v_{\text{conv}}}{E_{\text{dz}}} \quad (13)$$

Operation in plateau mode. If the concentration of an analyte approaches the background electrolyte concentration, it will significantly contribute to the conductivity and its effect on the electric field cannot be neglected. In this case the analyte will concentrate no further and instead of a Gaussian peak a plateau will be formed at the maximum concentration (Fig. 3b), which will widen in time. At the location of the plateau the E-field will be constant ($dE/dx = 0$). The maximum concentration for a specific analyte can be calculated from the Kohlrausch regulation function (KRF [mol V s m⁻⁵])³⁹

$$\text{KRF} = \sum_i \frac{z_i C_i}{\mu_i} \quad (14)$$

A conservation law can be derived that requires that the KRF value remains constant across all zones in a capillary (*i.e.* the bulk zone and any plateau formed from an analyte) on application of an electrical field. The E-field gradient is not affected on either side of the plateau hence other less abundant analytes will still be focusing there in peak mode. A known substance can be added to the solution with an electrophoretic mobility intermediate to the two analytes to push them apart (Fig. 3c). Such a compound is usually named electrophoretic spacer which has been previously reported in ICPF.^{3,40}



Experimental

In order to test the proposed system and procedure, polydimethyl siloxane (PDMS, DOWSIL™ 184 silicone elastomer kit)(1 : 10 crosslinker to polymer ratio) chips were fabricated. A silicon wafer was spin coated with 35 μm of MicroChem® SU-82050 negative photoresist and patterned with photolithography mask of the pattern of Fig. 1a. The supplier's instructions for exposure and development were followed to create the SU-8 soft lithography mould. The mould was used for standard soft lithography of PDMS chips⁴¹ which were bonded after O₂ plasma treatment to microscopy glass slides.

A droplet of Nafion® perfluorinated resin solution (20 wt% – Sigma-Aldrich) was introduced in the two reservoirs marked with blue in Fig. 1a. The Nafion solution filled and pinned in the channel *via* capillary forces. The Nafion solution was dried at 60 °C for 30 minutes to create the solid cation selective Nafion membrane. During the drying process Nafion experiences a significant shrinkage. Because of the shrinkage occasionally Nafion detaches from the PDMS and gaps are formed that allow undesired flow from the separation to the buffer channels. To eliminate this flow, the buffer channels were filled with 2% wt UltraPure™ Agarose (Invitrogen) in 1 \times phosphate buffered saline (PBS) (Sigma-Aldrich).

For the fluorescent microscopy an Olympus IX51 was used and images/videos were captured with FLIR Grasshopper® 3 color camera. As model cation analytes BODIPY™ 492/515 disulfonate (BDP) (Invitrogen), Alexa Fluor™ 647 carboxylic acid, tris(triethylammonium) (AF647) (Invitrogen) salt, and Cascade Blue™ hydrazide trisodium salt (CB) (Invitrogen) were chosen and diluted in 1 \times PBS or 0.1 \times PBS which served as a background electrolyte buffered at pH 7.4. A calibration curve between fluorescent analyte concentration and fluorescent intensity for each analyte was prepared for the specific microscope and camera settings used in the experiments. The concentration factors were then calculated using this curve.

The electric potentials were applied by two Keithley 2410 sourcemeter power supplies and the extraction pressure *via* a Fluigent MFCS-EZ.

Results and discussion

Analyte position control

As described in the model, the analyte focusing position is regulated by regulating the size of the depletion zone, which scales with the ratio of the upstream over downstream potential. An example of the analyte position control is shown in Fig. 4 where the potential control is used to locate the focused analyte band at the intersection with the two extraction channels. The device could subsequently successfully be used for prolonged extraction of a single fluorescent concentrated analyte. In our devices we achieved reproducible and stable extraction for up to 20 minutes with the limit being only posed by the small size of our reservoirs. Longer extractions

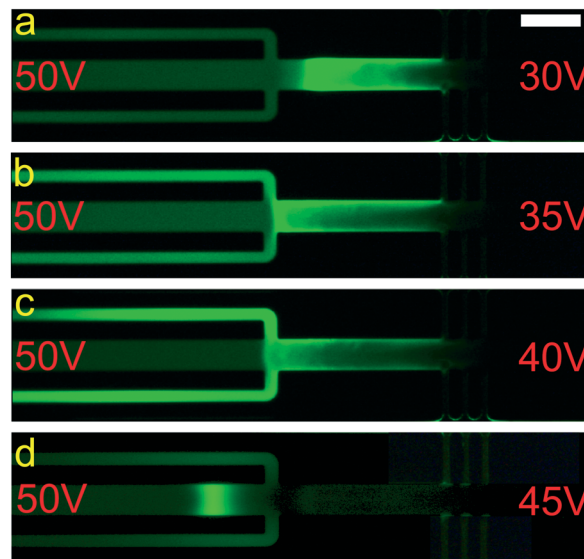


Fig. 4 Fluorescent microscopy images demonstrating focusing, band formation, control and extraction (1.5 μM BDP in 0.1 \times PBS). a) Focusing of the analyte away from the extraction channel. The concentration of the dye in the extraction channel is similar to the bulk. b) Movement of the focused analyte towards the extraction channel and partial release. c) Focusing and simultaneous extraction. d) Focusing at the other side of extraction channel. If a second slower analyte was present it could be extracted. Scale bar – 100 μm .

require larger reservoirs in order to avoid hydrostatic pressure changes and pH changes due to electrolysis.

In Fig. 5 the focusing position is plotted against the ratio of upstream (V^u) and downstream (V^d) potential. The upstream potential was kept steady at 110 V and the downstream potential was varied from 110 V to 10 V and back to 110 V with steps of 10 V over a period of 35 minutes. The

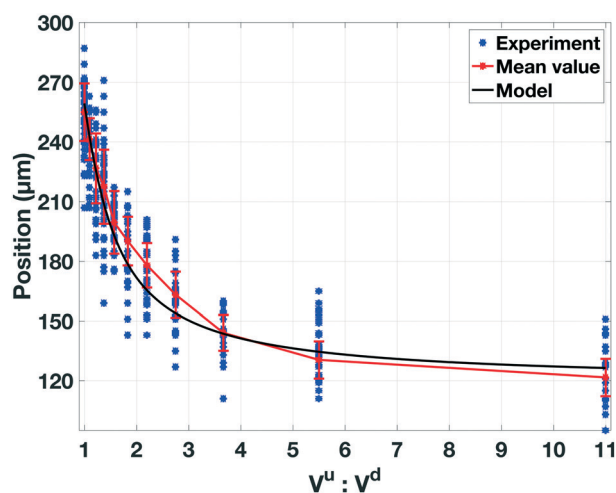


Fig. 5 Analyte position vs. upstream over downstream potential ratio. The blue points correspond to the distance of the analyte (BDP in 1 \times PBS) from the Nafion membrane. The red line connects the average value for each potential ratio and the red bars correspond to the standard deviation. The black line shows the behaviour predicted by our model.



blue points denote the measured position of the analyte over time, while the red line connects the average values for each potential ratio (red bars: standard deviation of the focusing location over time). The instabilities in the position mainly stem from electroconvective vortices and instabilities in the depletion zone.^{24,31} The higher the flow rate (higher $V^u:V^d$) the smaller the deviation from the mean focusing value (standard deviation of $9\ \mu\text{m}$ at $V^u:V^d$ of 11), while for lower flow rates the standard deviation increases ($19\ \mu\text{m}$ for $V^u:V^d$ of 1.22).

Analyte selection control

We also investigated the use of the analyte positioning control for the selection of focused analyte to be extracted. In Fig. 6 the selection between two negatively charged analytes, namely BDP and AF647 is shown. The electrophoretic mobilities of the two analytes as determined by on-chip electrophoresis were $\mu_{\text{BDP}} = 2.11 \times 10^{-8}\ \text{m}^2\ \text{V}^{-1}\ \text{s}^{-1}$ (a mobility of $1.76 \times 10^{-8}\ \text{m}^2\ \text{V}^{-1}\ \text{s}^{-1}$ is reported in the literature⁴²) and $\mu_{\text{AF647}} = 1.58 \times 10^{-8}\ \text{m}^2\ \text{V}^{-1}\ \text{s}^{-1}$. The slower analyte (*i.e.* the red AF647)

is expected to focus closer to the depletion zone at the right of the figure where the E-field is higher, while the green BDP will focus closer to the bulk solution on the left of the figure, at the lower E-field. By tuning the actuation potentials (V^u , V^d) we could select which of the focused analytes positionally overlaps with the extraction channel and will be extracted.

As shown in experiment A of Fig. 6, in the extraction channels only the fluorescence of the targeted analyte (BDP) is present, while the fluorescence of the “interfering” analyte (AF647) remains below the limit of detection of our microscope (ESI† video). A similar result is shown for experiment B but now with AF647 as target analyte and BDP as “interfering” analyte. The lack of fluorescence of the interfering species in the extraction channel demonstrates that our method in this case could be used for selective extraction.

As shown in Fig. 6 there is a clear difference in the resolution between experiments A and B before the extraction. This difference is caused by the difference in potential difference applied in the two experiments (30 V in A and 5 V in B). As described by eqn (12) the resolution improves with a lower electric field gradient (*i.e.* lower potential difference) at the cost of the concentration rate. The poor resolution results in the substantial overlap between the concentrated analytes seen in Fig. 6A. A selective extraction of analyte in that case is still possible but requires a small width of the extraction channel relative to the width of the desired concentrated analyte band.

The extraction of fluorescent analytes is a straightforward, reproducible and robust process, since we can monitor our analyte of interest until its position overlaps with the extraction channel. Extraction of non-fluorescent analytes would be a more cumbersome process, since the electrophoretic mobility of the analyte of interest must be known with respect to the electrophoretic mobilities to the other species in the solution. A possible experimental solution is to use a fluorescent substance with a mobility close to analyte of interest as a marker, placing it adjacent to the extraction channel.

Higher concentration factors

The factor with which the extracted analyte can be concentrated is tuneable. A steady extracted concentration factor is achieved when the extracted analyte flux matches the arriving analyte flux through the main channel towards the extraction location. A typical operational scheme for this purpose is the following: i) the process is initiated with actuation potentials such that the focusing location of the analyte is not overlapping the extraction channel; ii) a sufficient waiting time (typically a concentration rate is 1 times the bulk concentration per second) ensures that the desired concentration factor is achieved; iii) the potentials are adjusted to move the focused analyte band to the extraction channel; iv) a negative pressure is applied at the extraction channel to create an extraction flux that approximately matches the incoming analyte flux. An example of an extraction attempting to follow this scheme is shown in Fig. 7. We obtained an average

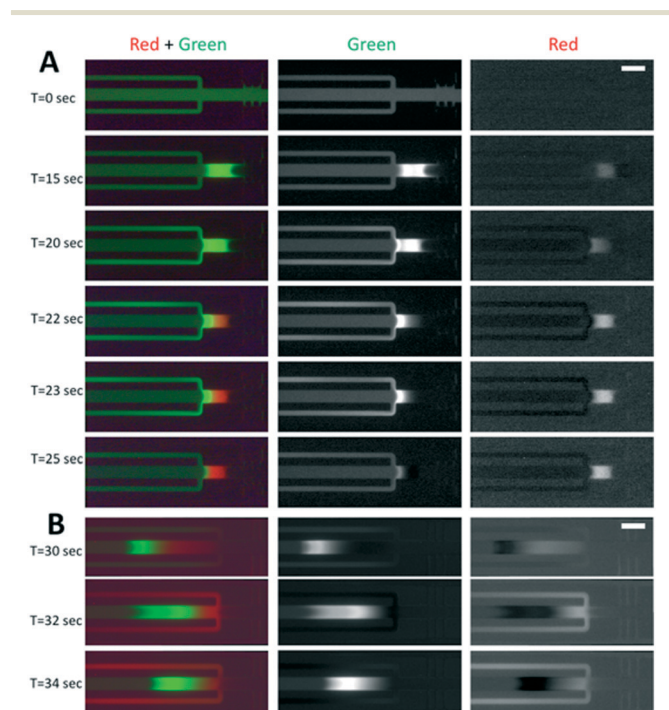


Fig. 6 Separation and extraction of single analytes from a mixture (BDP (green) and AF647 (red) in $0.1\times$ PBS). Red filter column indicates the AF647 concentration, green filter column the BDP and red + green the composite image. We focus and create bands of both analytes for approximately 15 s and then by changing the potentials we release BDP (A) or AF647 (B). The extraction channels remain dark in the red filter (A) and green filter (B) which indicates that only the targeted analyte is extracted. In both experiments an extraction pressure of $-3\ \text{mBar}$ was applied at $T = 20\ \text{s}$ (A) and $T = 30\ \text{s}$ (B). Before the application of the extraction pressure the extraction channels contain bulk liquid. Note: A and B are two separate experiments and the time values are not related. Actuation potentials during extraction (A): $V^u = 60\ \text{V}$ and $V^d = 30\ \text{V}$, (B) $V^u = 60\ \text{V}$ and $V^d = 55\ \text{V}$. Scale bars – $100\ \mu\text{m}$.



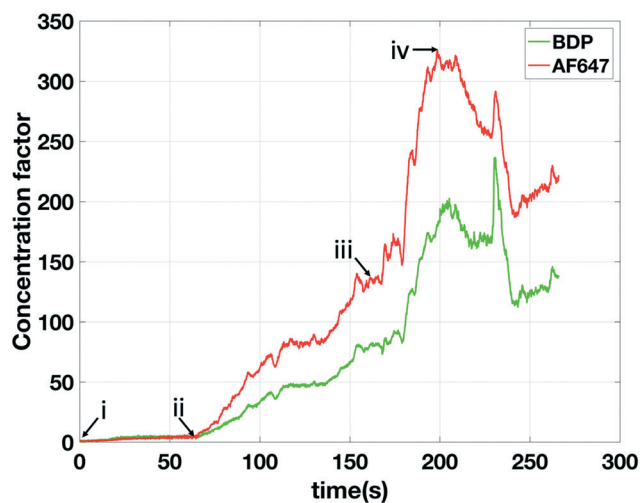


Fig. 7 Concentration factor as (concentration)/(bulk concentration) of analytes in the extraction channel versus time. i) At $t = 0$ s the device is turned on with actuation potentials of $V^u = 80$ V and $V^d = 20$ V and an extraction pressure of -15 mBar. The analytes are focusing at a location that does not overlap with the extraction channel hence the concentration at the extraction channel does not change. ii) At $t = 60$ s potentials are changed to $V^u = 80$ V and $V^d = 45$ V which moves the focused analytes to the extraction channel. The extraction pressure of -15 mBar creates an extraction flux that is lower than the incoming flux in the separation channel hence a constant increase in the concentration is seen. iii) At $t = 160$ s the actuation potentials remain constant and the pressure is reduced to -13 mBar hence an even lower extraction flux results in a higher concentration increase in the extraction channel. iv) At $t = 200$ s the actuation potentials remain constant and the pressure is increased back to -15 mBar. It was expected that the original concentration increase rate (slope between ii and iii) will be restored, but that only occurred after approximately 40 s. Electrolyte: $1\times$ PBS. Analytes: 100 nM BDP, 100 nM AF647.

concentration factor of about 200 in the continuously extracted analyte over several minutes. During the experiments (as also can be seen in Fig. 7) a slow response time of the pressure-driven flow was noticed varying between 5 to 40 seconds. We attribute this to instabilities of the pressure controller which was operated at its resolution limit. In addition, as can be seen in Fig. 7, the change of only 2 mBar in extraction pressure introduced at moment iii has a strong influence on the extraction rate. We conclude that a better controllability of the extraction process is needed. In future devices, which can be implemented *e.g.* by increasing the hydraulic resistance of the extraction channel.

Despite the high concentration factor of AF647 as seen in Fig. 7, in this experiment we obtained low selectivity since also BDP was extracted (at a concentration factor approximately 2/3 of AF647). Similar to Fig. 6A this was due to the poor separation resolution caused by the high potential difference.

As described in the theory section the separation resolution can be improved using an electrophoretic spacer. In Fig. 8 an example of the use of a spacer is shown. A third fluorescent marker (CB) was added with a mobility higher than BDP and the added phosphate of the buffer, causing

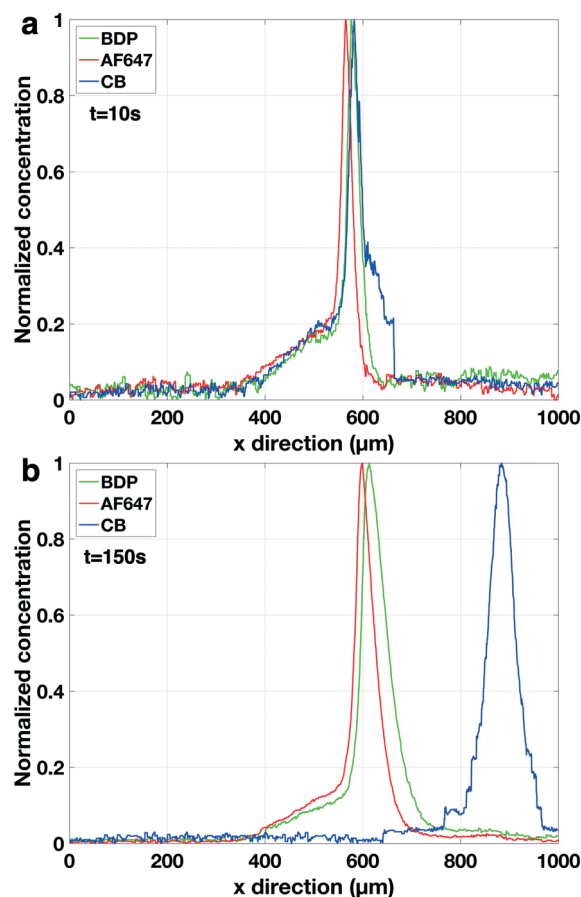


Fig. 8 Normalised concentration profile of three focused fluorescent analytes in peak mode. a) $t = 10$ s. The concentration of the electrophoretic spacer is still low comparable to the background electrolyte hence also the spacer is still in peak mode. b) $t = 150$ s. The concentration of the spacer has increased hence now it is in plateau mode pushing CB peak away from AF647 and BDP. The concentration was normalized based on the maximum fluorescent intensity for each fluorescent analyte independently. Electrolyte: $1\times$ PBS. Analytes: 100 nM BDP, 100 nM AF647. $V^u = 60$ V and $V^d = 35$ V.

the phosphate (H_2PO_4^-) to reach its plateau concentration and act as a spacer, while AF647 and BDP remained in peak mode and at their initial position.

Comparison to the theoretical model

In the experiments characterizing the analyte positioning (Fig. 4) a single fluorescent analyte was used and the focusing position (*i.e.* depletion zone length) was found to scale, as predicted by our model, with the ratio of upstream over downstream potential. We found experimentally that the baseline width (4σ) of the focused analyte varied between 57 μm (for $V^u = 60$ V, $V^d = 0$ V) and 141 μm (for $V^u = 60$ V, $V^d = 55$ V). Eqn (10) allows us to derive the local E-field gradients from these measured peak widths. As the model provides us with the electric field in the bulk section of the channel and in the depletion zone at the applied potentials, the length of the E-field gradient zone can now be estimated assuming a



constant electric field gradient. With this approach we derived an E-field gradient zone length of approximately 500 μm . Since the extraction channels have a width of 20 μm , this allows us safe overlap and a steady extraction of an analyte even in the case of a slight variation in the position of the focused analyte over time (standard deviation 19 μm in the focusing position as calculated from Fig. 5).

The σ of the Gaussian (green and red) peaks in the experiments at higher concentration factors (Fig. 8b) is approximately 26 μm indicating an electric field gradient of $1.849 \times 10^6 \text{ V m}^{-2}$ (based on eqn (10)). At this field gradient our model predicts a spacing between the peaks of BDP and AF647 of 23 μm comparing well to the experimentally observed spacing of 19 μm . Interestingly it was found that the σ of the blue peak (CB) increased from 19 μm (Fig. 8a) to 31 μm (Fig. 8b) during the experiment indicating that the electric field gradient is less steep further away from the depletion zone. Fig. 9 shows the maximum concentration factor of the three analytes over time in the separation channel (Fig. 9a) and in the extraction channel (9b). Our model predicts a concentration rate of 2.7 times the bulk concentration per second while an average concentration rate of 2.1 can be seen in Fig. 9a. Before the extraction pressure was applied ($t < 150 \text{ s}$) all three analytes focussed with approximately the same concentration rate. Once the extraction pressure was applied, AF647 was extracted together with a lower concentration of BDP, similarly to Fig. 7. However, no detectable amount of the CB was extracted as it was pushed away from the extraction channel by the electrophoretic spacer.

In Fig. 7 and 9b a strong variation can be seen in the maximum concentration of the extracted analytes. This can be attributed to hydrodynamic instabilities in the focusing location at high potentials (ESI† video) in ICPF, causing variations in the spatial overlap between the focused analyte band and the extraction channel.

Summarizing, the high electric field gradients that result in narrowly focused and highly concentrated bands, also drastically reduce the resolution of the method as can be seen from the denominator in eqn (12). The resolution of the system will thus be superior for lower potential differences between upstream and downstream reservoirs. This is further theoretically described in the ESI† and can also be seen in Fig. 6B. There thus exists a trade-off between high resolution and high concentration rate. This conundrum can be solved by increasing increase the value of the numerator in eqn (12) by a method independent of EOF. This can be done *e.g.* by additional PDF. The addition of a PDF will also shift the mobility window of analytes that will be focused in the device.

We can thus envision two potential future uses of this method. If the main interest is to obtain extracted analytes at enhanced concentration in continuous flow, the default operation (*i.e.* without the addition of spacers and/or added convective flow) of the device can be used, if needed in combination with pressure-driven flow. If the main interest is the purity of the extract, electrophoretic spacers can be used (if

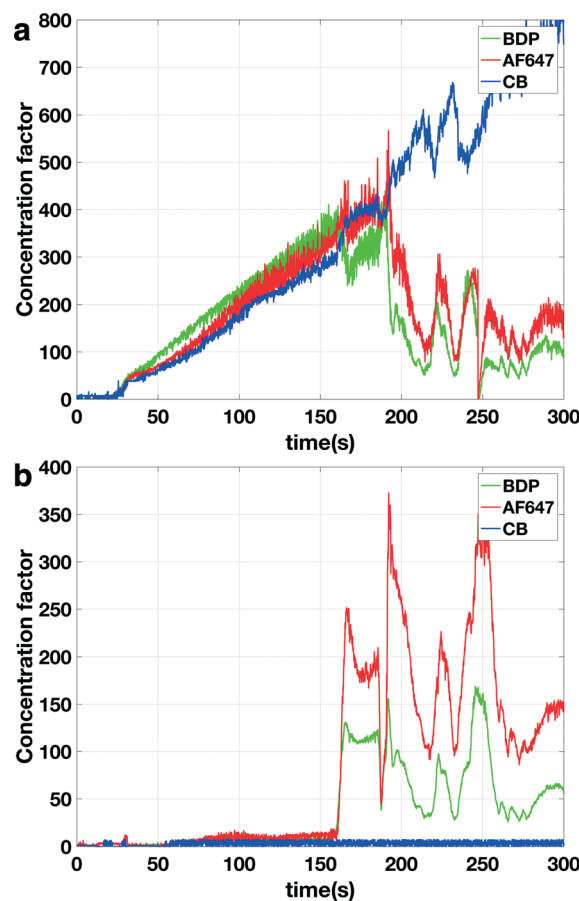


Fig. 9 Maximum concentration factor of three fluorescent analytes over time in the separation channel (a) and the extraction channel (b). a) Concentration factor of analytes in the separation channel. All analytes are concentrating with the same rate till $t = 150 \text{ s}$ when the extraction pressure of -17 mBar is applied. Once the extraction starts then the maximum concentration of AF647 and BDP drops in the separation channel while the concentration increase of CB is unaffected. b) Concentration factor of analytes in the extraction channel. Before the application of extraction pressure at $t = 150 \text{ s}$ no enrichment in the separation channel can be seen. Once the pressure is applied only AF647 and BDP are observed in the extraction channel.

the mobility of the analyte is known) or an external pressure driven flow can be added for improved separation resolution.

Conclusion

We demonstrated a device capable of selective extraction and concentration of anionic analytes out of a complex background solution. The maximum concentration factor of the extracted analytes found was above 300, as measured by an increased fluorescence intensity. In order to investigate the dynamics of the system and deepen our understanding of the phenomena at play, a simple two-part one-dimensional model (*i.e.* an electric and a fluidic) was developed that describes the input–output relations of the system (applied potentials, extraction pressure – focusing position, concentration rate, extraction rate) to a level that allowed prediction of



the system behaviour. Different future modes of operation are described based on the model and the experiments.

Conflicts of interest

There are no conflicts to declare.

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References

- 1 Y. C. Wang, A. L. Stevens and J. Han, *Anal. Chem.*, 2005, **77**(14), 4293–4299.
- 2 J. Quist, K. G. H. Janssen, P. Vulto, T. Hankemeier and H. J. Van Der Linden, *Anal. Chem.*, 2011, **83**, 7910–7915.
- 3 J. Quist, P. Vulto, H. Van Der Linden and T. Hankemeier, *Anal. Chem.*, 2012, **84**, 9065–9071.
- 4 S. J. Kim, L. D. Li and J. Han, *Langmuir*, 2009, **25**, 7759–7765.
- 5 V. Liu, Y.-A. Song and J. Han, *Lab Chip*, 2010, **10**, 1485.
- 6 J. Khandurina, S. C. Jacobson, L. C. Waters, R. S. Foote and J. M. Ramsey, *Anal. Chem.*, 1999, **71**, 1815–1819.
- 7 S. H. Ko, Y.-A. Song, S. J. Kim, M. Kim, J. Han and K. H. Kang, *Lab Chip*, 2012, **12**, 4472.
- 8 J. Choi, K. Huh, D. J. Moon, H. Lee, S. Y. Son, K. Kim, H. C. Kim, J. H. Chae, G. Y. Sung, H. Y. Kim, J. W. Hong and S. J. Kim, *RSC Adv.*, 2015, **5**, 66178–66184.
- 9 Y. Y. Chen, P. H. Chiu, C. H. Weng and R. J. Yang, *Biomicrofluidics*, 2016, **10**(1), 014119.
- 10 R. Kwak, S. J. Kim and J. Han, *Anal. Chem.*, 2011, **83**(19), 7348–7355.
- 11 W. Ouyang, X. Ye, Z. Li and J. Han, *Nanoscale*, 2018, **10**(32), 15119–15412.
- 12 L. Gong, W. Ouyang, Z. Li and J. Han, *J. Membr. Sci.*, 2018, **556**, 34–41.
- 13 L. Gong, Z. Li and J. Han, *Sep. Purif. Technol.*, 2019, **217**, 174–182.
- 14 Q. Pu, J. Yun, H. Temkin and S. Liu, *Nano Lett.*, 2004, **4**, 1099–1103.
- 15 V. A. Papadimitriou, L. I. Segerink, A. van den Berg and J. C. T. Eijkel, *Anal. Chim. Acta*, 2018, **1000**, 232–238.
- 16 A. E. Herr, J. I. Molho, J. G. Santiago, M. G. Mungal, T. W. Kenny and M. G. Garguilo, *Anal. Chem.*, 2000, **72**(5), 1053–1057.
- 17 K. Mogi, *Micromachines*, 2018, **9**(4), 167.
- 18 S. S. Dukhin, *Adv. Colloid Interface Sci.*, 1991, **35**, 173–196.
- 19 S. M. Rubinstein, G. Manukyan, A. Staicu, I. Rubinstein, B. Zaltzman, R. G. H. Lammertink, F. Mugele and M. Wessling, *Phys. Rev. Lett.*, 2008, **101**, 236101.
- 20 I. Rubinstein, B. Zaltzman and I. Lerman, *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.*, 2005, **72**, 011505.
- 21 I. Rubinstein and B. Zaltzman, *Math. Models Methods Appl. Sci.*, 2001, **11**(02), 263–300.
- 22 I. Rubinstein and B. Zaltzman, *Phys. Rev. E: Stat. Phys., Plasmas, Fluids, Relat. Interdiscip. Top.*, 2000, **62**, 2238.
- 23 F. C. Leinweber and U. Tallarek, *Langmuir*, 2004, **20**(26), 11637–11648.
- 24 S. J. Kim, Y. C. Wang, J. H. Lee, H. Jang and J. Han, *Phys. Rev. Lett.*, 2007, **99**, 044501.
- 25 A. Plecis and Y. Chen, *Anal. Chem.*, 2000, **72**(18), 4317–4321.
- 26 X. Xuan and D. Li, *J. Micromech. Microeng.*, 2003, **14**(2), 290.
- 27 A. Ajdari, *C. R. Phys.*, 2004, **5**(5), 539–546.
- 28 T. Glawdel, C. Elbuken, L. E. J. Lee and C. L. Ren, *Lab Chip*, 2009, **9**, 3243–3250.
- 29 K. W. Oh, K. Lee, B. Ahn and E. P. Furlani, *Lab Chip*, 2012, **12**(3), 515–545.
- 30 T. A. Zangle, A. Mani and J. G. Santiago, *Anal. Chem.*, 2010, **82**, 3114–3117.
- 31 J. C. De Valença, R. M. Wagterveld, R. G. H. Lammertink and P. A. Tsai, *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.*, 2015, **92**, 031003.
- 32 B. J. Kirby and E. F. Hasselbrink, *Electrophoresis*, 2004, **25**, 187–202.
- 33 B. J. Kirby and E. F. Hasselbrink, *Electrophoresis*, 2004, **25**, 203–213.
- 34 J. K. Sung and J. Han, *Anal. Chem.*, 2008, **80**, 3507–3511.
- 35 S. Rubin, O. Schwartz and M. Bercovici, *Phys. Fluids*, 2013, **26**, 012001.
- 36 D. A. MacInnes and L. G. Longworth, *Chem. Rev.*, 1932, **11**, 171–230.
- 37 A. Tiselius, *Trans. Faraday Soc.*, 1937, **33**, 524–531.
- 38 J. C. Giddings, *Unified Separation Science*, John Wiley & Sons, 1991.
- 39 R. Kohlrausch, *Ann. Phys.*, 1854, **167**, 179–214.
- 40 L. F. Cheow, A. Sarkar, S. Kolitz, D. Lauffenburger and J. Han, *Anal. Chem.*, 2014, **86**(15), 7455–7462.
- 41 D. B. Weibel, W. R. DiLuzio and G. M. Whitesides, *Nat. Rev. Microbiol.*, 2007, **5**, 209–218.
- 42 B. G. Splawn and F. E. Lytle, *Anal. Bioanal. Chem.*, 2002, **373**(7), 519–525.

