

Cite this: *J. Mater. Chem. C*, 2018, 6, 8923

# Weak thermal quenching of the luminescence in the $\text{Ca}_3\text{Sc}_2\text{Si}_3\text{O}_{12}:\text{Ce}^{3+}$ garnet phosphor†

Suchinder K. Sharma,<sup>‡</sup> Yuan-Chih Lin,<sup>‡</sup> Irene Carrasco,<sup>§</sup> Tobias Tingberg,<sup>c</sup> Marco Bettinelli<sup>b</sup> and Maths Karlsson<sup>\*,a</sup>

We report results of the luminescence properties of the three garnet type phosphors  $\text{Ce}^{3+}$ -doped  $\text{Ca}_3\text{Sc}_2\text{Si}_3\text{O}_{12}$  (CSSO: $\text{Ce}^{3+}$ ),  $\text{Sr}_3\text{Y}_2\text{Ge}_3\text{O}_{12}$  (SYG: $\text{Ce}^{3+}$ ) and  $\text{Y}_3\text{Al}_5\text{O}_{12}$  (YAG: $\text{Ce}^{3+}$ ), investigated using optical spectroscopy techniques and vacuum referred binding energy (VRBE) diagram analysis. By monitoring the temperature dependence of the luminescence decay time we establish an excellent, intrinsic, thermal stability of luminescence in CSSO: $\text{Ce}^{3+}$ , with a nearly constant decay time ( $\approx 60$  ns) up to, at least,  $T = 860$  K. In comparison, SYG: $\text{Ce}^{3+}$  and YAG: $\text{Ce}^{3+}$  exhibit a significant reduction of the luminescence decay time upon heating, starting at around  $T = 280$  K and  $T = 550$  K, respectively, suggesting a lower internal thermal stability of luminescence in these two garnet phosphors. These findings are supported by the energy separation between the  $\text{Ce}^{3+} 5d_1$  level and the conduction band (CB) of the respective hosts, which are found at 1.36 eV (CSSO: $\text{Ce}^{3+}$ ), 0.45 eV (SYG: $\text{Ce}^{3+}$ ), and 1.17 eV (YAG: $\text{Ce}^{3+}$ ), respectively, as predicted by their VRBE diagrams. The performance of CSSO: $\text{Ce}^{3+}$  was evaluated by applying the phosphor on a blue InGaN LED. The system shows a luminous efficacy of optical radiation of  $243 \text{ lm W}^{-1}$  and a linear response with increasing applied voltage, suggesting it is a highly promising phosphor for future technological applications, particularly at high temperature operating environments.

Received 13th June 2018,  
Accepted 12th July 2018

DOI: 10.1039/c8tc02907e

rsc.li/materials-c

## 1 Introduction

Phosphor-converted white-light emitting diodes (pc-WLEDs) based on blue InGaN LEDs offer many advantages compared to traditional light sources due to their high efficiency, tuneable color, and long lifetimes.<sup>1–6</sup> The phosphors consist of a crystal-line host lattice containing a small amount of luminescent ions, which convert the blue light to longer wavelengths *via* electronic transitions. Garnet crystals, of general formula  $\text{A}_3\text{B}_2\text{C}_3\text{O}_{12}$ , where B and C may be the same or different atoms and as characterized by a 160-atom body-centered cubic unit cell with a lattice constant  $L$  [Fig. 1(a)], are widely considered to be one of the most important families of hosts for luminescent ions, especially for trivalent lanthanide ions such as  $\text{Ce}^{3+}$ .

The A, B, and C atoms are (8-fold) dodecahedrally, (6-fold) octahedrally, and (4-fold) tetrahedrally coordinated to O atoms that are shared between neighboring cation–oxygen polyhedra.<sup>7</sup> The choice of cations provides a means to tune the local coordination environment of the luminescent ions and thus important materials properties of relevance for luminescence, such as the centroid shift ( $\epsilon_c$ ) and crystal field splitting ( $\Delta$  and  $\Delta E_{5d_{1-2}}$ ) of the luminescent ions;<sup>7–12</sup> see Fig. 1(b) for a schematic illustration of the effect on the electronic energy levels of a free  $\text{Ce}^{3+}$  ion when inserted in a host.

The best known garnet is by far yttrium aluminium garnet ( $\text{Y}_3\text{Al}_5\text{O}_{12}$ , YAG) that, when doped with  $\text{Ce}^{3+}$  (YAG: $\text{Ce}^{3+}$ ), emits light in the green-yellow range due to inter-configurational  $5d-4f$  transitions and constitutes an important phosphor for use in pc-WLEDs.<sup>13</sup> The YAG structure ( $L = 12.01 \text{ \AA}$ ) is built up of  $\text{AlO}_6$  octahedra and  $\text{AlO}_4$  tetrahedra, and  $\text{YO}_8$  dodecahedra, and the  $\text{Ce}^{3+}$  substitutes for  $\text{Y}^{3+}$  due to matching ionic charge and radii (*cf.*  $1.143 \text{ \AA}$  for  $\text{Ce}^{3+}$  and  $1.019 \text{ \AA}$  for  $\text{Y}^{3+}$ , in 8-fold coordination).<sup>14–16</sup> Two novel, green-emitting, analogous garnet systems to YAG: $\text{Ce}^{3+}$ , which can shed further light on the luminescent properties of garnet phosphors, are  $\text{Ce}^{3+}$ -doped  $\text{Ca}_3\text{Sc}_2\text{Si}_3\text{O}_{12}$  (CSSO: $\text{Ce}^{3+}$ ) and  $\text{Ce}^{3+}$ -doped  $\text{Sr}_3\text{Y}_2\text{Ge}_3\text{O}_{12}$  (SYG: $\text{Ce}^{3+}$ ). CSSO and SYG are isostructural with  $L = 12.25 \text{ \AA}$  and  $L = 13.08 \text{ \AA}$ , respectively.<sup>17,18</sup> On the basis of considerations of the ionic radii,  $\text{Ce}^{3+}$  is expected to substitute for  $\text{Ca}^{2+}/\text{Sr}^{2+}$  (*cf.*  $1.143 \text{ \AA}$  for  $\text{Ce}^{3+}$  and  $1.12 \text{ \AA}$  for  $\text{Ca}^{2+}$  and  $1.26 \text{ \AA}$  for  $\text{Sr}^{2+}$ , in 8-fold coordination),<sup>15</sup>

<sup>a</sup> Department of Chemistry and Chemical Engineering, Chalmers University of Technology, SE-412 96 Göteborg, Sweden. E-mail: maths.karlsson@chalmers.se; Tel: +46 31 772 6770

<sup>b</sup> Luminescent Materials Laboratory, Department of Biotechnology, University of Verona, and INSTM, UdR Verona, Strada Le Grazie 15, 371 34 Verona, Italy

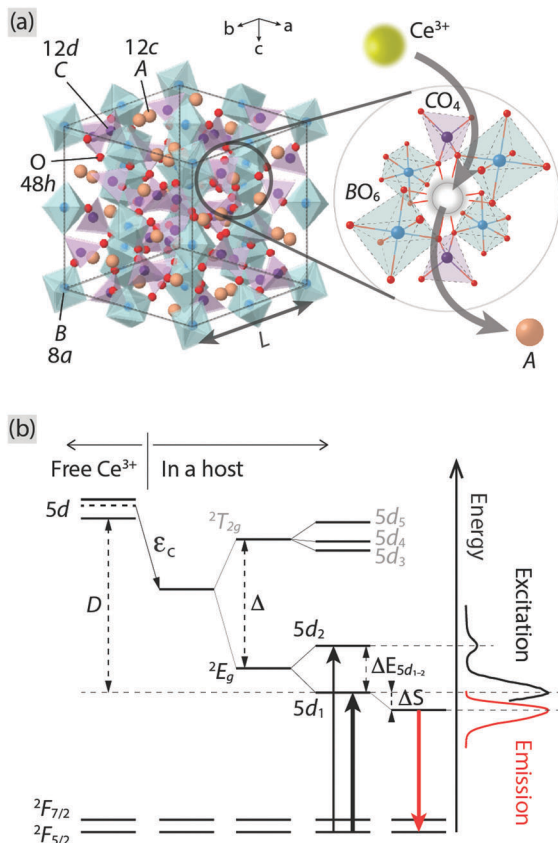
<sup>c</sup> Department of Microtechnology and Nanoscience, Chalmers University of Technology, 412 96 Göteborg, Sweden

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c8tc02907e

‡ Current address: Institute of Applied Physics, TU Bergakademie Freiberg, Leipziger Str. 23, 09596 Freiberg, Germany.

§ Current address: Advanced Technology Institute, Department of Electrical and Electronic Engineering, University of Surrey, Guildford GU2 7XH, UK.





**Fig. 1** (a) Illustration of an unit cell of a  $A_3B_2C_3O_{12}$  garnet crystal. The A, B, C, and O atoms occupy Wyckoff positions 12c, 8a, 12d, and 48h, respectively. The image to the right shows a close-up of the local coordination of an A atom, which is oriented in a way that the neighboring polyhedra can be clearly seen, showing the substitution of  $Ce^{3+}$  for an A ion. (b) A schematic energy level diagram for a free  $Ce^{3+}$  and when it is embedded in a host, in relation to the PL spectra of  $Ce^{3+}$ -doped  $Y_3Al_5O_{12}$  as an example (cf. Fig. 2).  $\epsilon_c$ : centroid shift,  $\Delta$ : crystal field splitting due to a cubal coordination,  $\Delta E_{5d_{1-2}}$ : crystal field splitting due to tetragonally distorted coordination,  $\Delta S$ : Stokes shift,  $D$ : red-shift of the lowest 5d level of  $Ce^{3+}$ ,  ${}^2T_{2g}$ : triply degenerate level of the upper 5d levels of  $Ce^{3+}$ ,  ${}^2E_g$ : doubly degenerate level of the lower 5d levels of  $Ce^{3+}$ , and  ${}^2F_{5/2}$  and  ${}^2F_{7/2}$ : the split 4f levels of  $Ce^{3+}$  due to the spin–orbit coupling.

but, due to the charge mismatch, two  $Ce^{3+}$  ions are expected to substitute for three  $Ca^{2+}/Sr^{2+}$  ions to maintain charge neutrality. Simultaneously, cation vacancies are formed as a charge compensating effect.

A critical challenge in the development of pc-WLEDs is to develop phosphors with high luminance and high thermal stability.<sup>19–21</sup> In the literature available on  $YAG:Ce^{3+}$ ,  $CSSO:Ce^{3+}$ , and  $SYG:Ce^{3+}$ , the thermal stability of luminescence of  $YAG:Ce^{3+}$  has been studied extensively.  $YAG:Ce^{3+}$  exhibits a 20% loss of the room temperature intensity at  $T = 500$  K and more than 50% at  $T = 700$  K, depending on  $Ce^{3+}$  dopant concentration and excitation wavelength.<sup>22,23</sup> The lowering of the emission intensity as a function of increasing temperature has been attributed to a combination of a change in the 4f–5d absorption strength, as well as (intrinsic) thermal quenching of luminescence, as inferred by a decrease in the luminescence decay time of the

$Ce^{3+}$  emission.<sup>22</sup> For  $CSSO:Ce^{3+}$ , the decrease in emission intensity with increasing temperature is less pronounced than for  $YAG:Ce^{3+}$  and has been also attributed to a change of the 4f–5d absorption strength, partly related to absorption by impurity phases ( $\beta-Ca_2SiO_4$ ,  $CaSiO_3$ ,  $CeO_2$ ), whereas the luminescence decay time is nearly constant ( $\approx 71$  ns) up to  $T = 700$  K.<sup>24</sup> The higher intrinsic thermal stability with respect to luminescence of  $CSSO:Ce^{3+}$  as compared to  $YAG:Ce^{3+}$  not only makes it a promising phosphor for application in pc-WLEDs, but it is also highly interesting from a more fundamental point of view. For the structurally similar phosphor  $SYG:Ce^{3+}$ , there is a lack of information as the thermal stability of luminescence has not been investigated, yet.

The aim of the present work is to compare and understand the emission and excitation characteristics, especially in regard to the intrinsic thermal stability, of  $SYG:Ce^{3+}$ ,  $CSSO:Ce^{3+}$ , and  $YAG:Ce^{3+}$  upon excitation at 420–470 nm, thus mimicking the experimental conditions for a pc-WLED using a blue LED. With a view toward practical applications, we also evaluate the performance of  $CSSO:Ce^{3+}$  phosphor in a pc-WLED prototype. Combined analyses of the emission and excitation spectra, the temperature dependence of the luminescence decay time of  $Ce^{3+}$  in  $SYG:Ce^{3+}$ ,  $CSSO:Ce^{3+}$ , and  $YAG:Ce^{3+}$ , and the vacuum referred binding energy (VRBE) diagram for  $Ce^{3+}$  in the respective host lattices, suggest that the separation between the host CB and the  $Ce^{3+}$  5d<sub>1</sub> level plays the most critical role in determining the thermal stability of the three garnet type phosphors.

## 2 Experimental

### 2.1 Sample preparation

$CSSO:Ce^{3+}$  was prepared through solid state sintering, by mixing stoichiometric amounts of  $CaCO_3$  (Aldrich,  $\geq 99\%$  purity),  $Sc_2O_3$  (Strem, 99.99% purity),  $SiO_2$  (Aldrich,  $\geq 99.995\%$  purity), and  $Ce(NO_3)_3 \cdot 6H_2O$  (Aldrich, 99.99% purity), followed by three repetitive heat treatments for 3 h at 1450 °C, with intermediate grindings and compacting of pellets under ambient conditions.  $SYG:Ce^{3+}$  was prepared by solid state sintering, by mixing stoichiometric amounts of AR grade  $SrCO_3$  (Aldrich,  $\geq 99\%$  purity),  $Y_2O_3$  (Aldrich, 99.99% purity),  $GeO_2$  (Aldrich, 99.998% purity), and  $Ce(NO_3)_3 \cdot 6H_2O$  (Aldrich, 99.99% purity), followed by three repetitive heat treatments at 1250 °C (24 h), with intermediate grindings and compacting of pellets.  $YAG:Ce^{3+}$  was prepared by co-precipitation. Aqueous solutions of  $Y(NO_3)_3 \cdot 6H_2O$  (Strem, 99.999% purity),  $Ce(NO_3)_3 \cdot 6H_2O$  (Aldrich, 99.99% purity), and  $NH_4Al(SO_4)_2 \cdot 12H_2O$  (Aldrich,  $> 99\%$  purity) were prepared from AR reagents. The mixed nitrate and sulfate solution was slowly added to a 3 M solution of ammonium hydroxide, keeping the pH  $> 10$ . The precipitate was washed with diluted ammonia and ethanol, filtered and dried overnight at 90 °C to obtain the precursor powder, which was subsequently annealed at 1100 °C for 2 h. The as-prepared samples were phase pure with no detectable traces of impurities as evaluated from powder X-ray diffraction patterns, collected on a Thermo ARL X'TRA powder diffractometer, operating in the Bragg–Brentano



geometry and equipped with a Cu-anode X-ray source ( $K_{\alpha}$ ,  $\lambda = 1.5418 \text{ \AA}$ ), using a Peltier Si (Li) cooled solid state detector. The patterns were collected with a scan rate of  $0.03^{\circ} \text{ s}^{-1}$  in the  $2\theta$  range  $5\text{--}90^{\circ}$ .

## 2.2 Photoluminescence spectroscopy

The PL excitation and emission spectra were measured using two different setups. Room temperature PL spectra were measured at the University of Verona, Italy, using a Fluorolog 3 spectrofluorometer (Horiba-Jobin Yvon), equipped with a Xe lamp, a double excitation monochromator, a single emission monochromator (mod. HR320) and a photomultiplier in photon counting mode for the detection of the emitted signal. Variable temperature emission spectra and luminescence decay curves were measured at Chalmers University of Technology, Sweden, over the temperature range  $T = 80\text{--}860 \text{ K}$ . The temperature was controlled using a Linkam THMS 600 heating stage. The excitation was achieved using a 454 nm laser (DeltaDiode DD-450L) connected to a DD-Ci picosecond diode controller from Horiba Scientific. Due to a very weak and relatively broad emission band (420–490 nm) from the laser, which overlaps partly with the shorter-wavelength part of the emission bands of the studied phosphors, the laser light was further monochromatized by the use of a monochromator (300–800 nm Manual Mini-Chrom, Edmund Optics). Upon the pulsed excitation with a repetition rate of 500 kHz, luminescence decay curves were then detected using a H10721-20 photosensor module (Hamamatsu) containing a photomultiplier and a high-voltage power supply circuit in the 230–920 nm wavelength range, which was coupled to a 500 nm longpass filter (A10033-62, Hamamatsu). The detected decay curves were recorded on an Agilent DSO-X 2022A oscilloscope. The emission spectra were measured as the samples were excited by the 454 nm pulsed laser with a repetition rate of 100 MHz, by using an Ocean Optics USB2000+ UV-Vis spectrometer coupled to an optical fiber with a 455 nm longpass colored glass filter (FGL455, Thorlabs) attached in front of the fiber inlet. Reference spectra were measured using the same setting as described above, but without the blue light irradiating on the samples, and subtracted from the spectra of the samples.

## 2.3 Phosphor-converted WLED prototype

In order to evaluate the actual performance of  $\text{CSSO}:\text{Ce}^{3+}$  phosphor in a pc-WLED device, we integrated this phosphor onto a commercial 420 nm emitting blue LED chip, using the novel pc-WLED design, as developed by Bin Im *et al.*<sup>25</sup> The phosphor was molded with silicon resin, put on a cap that was placed onto the LED, and left to solidify for a period over 24 h. A small air gap ( $\approx 5 \text{ mm}$ ; refractive index  $n = 1$ ) between the LED chip and phosphor cap ( $n \approx 1.7$ ) reduces the amount of light that is reflected back into the chip, due to the differences in refractive indices of the various layers. Further, the phosphor layer is insulated from the heat as generated by the LED chip, thus reducing a loss of emission intensity due to thermal quenching effects.

The spectral output of the pc-WLED prototype was monitored on a OL 770 multi-channel spectroradiometer from

Optronic Laboratories at Chalmers University of Technology, Göteborg. The input current was varied using a DC power supply from BK Precision.

# 3 Results

## 3.1 Photoluminescence spectra

The room temperature PL excitation and emission spectra of  $\text{CSSO}:\text{Ce}^{3+}$ ,  $\text{SYG}:\text{Ce}^{3+}$ , and  $\text{YAG}:\text{Ce}^{3+}$  are shown in Fig. 2. The excitation spectra (black color) were measured at a fixed emission wavelength of 504 nm ( $\text{CSSO}:\text{Ce}^{3+}$ ), 513 nm ( $\text{SYG}:\text{Ce}^{3+}$ ), and 525 nm ( $\text{YAG}:\text{Ce}^{3+}$ ), respectively. The spectra possess excitation maxima at approximately 440 nm ( $\text{CSSO}:\text{Ce}^{3+}$ ), 427 nm ( $\text{SYG}:\text{Ce}^{3+}$ ), and 470 nm ( $\text{YAG}:\text{Ce}^{3+}$ ), with a full width at half maximum (FWHM) of 73 nm, 53 nm, and 72 nm, respectively, for the three materials. The excitation bands correspond to  $\text{Ce}^{3+} 4f \rightarrow 5d_1$  transitions ( $\lambda_{5d_1}$ ) in the respective hosts. Higher-energy bands peaking at approximately 285 nm ( $\text{CSSO}:\text{Ce}^{3+}$ ), 297 nm ( $\text{SYG}:\text{Ce}^{3+}$ ), and 340 nm ( $\text{YAG}:\text{Ce}^{3+}$ ), which correspond to  $\text{Ce}^{3+} 4f \rightarrow 5d_2$  transitions ( $\lambda_{5d_2}$ ), can be also observed. The energy difference ( $\Delta E_{5d_{1-2}}$ ) between the band maxima of the excitation transitions to the  $5d_2$  and  $5d_1$  levels is  $12360 \text{ cm}^{-1}$ ,  $10251 \text{ cm}^{-1}$ , and  $8135 \text{ cm}^{-1}$ , for  $\text{CSSO}:\text{Ce}^{3+}$ ,  $\text{SYG}:\text{Ce}^{3+}$ , and  $\text{YAG}:\text{Ce}^{3+}$ , respectively, and can be used as a measure for the strength of “tetragonal” crystal-field splitting of the 8-fold coordinated  $\text{Ce}^{3+}$  ions.<sup>11,26</sup> The  $\Delta E_{5d_{1-2}}$  values are in good agreement with the literature: for  $\text{CSSO}:\text{Ce}^{3+}$ , the value ( $12360 \text{ cm}^{-1}$ ) is very similar to the one reported for  $\text{Ce}^{3+}$ -doped  $\text{Lu}_2\text{CaMg}_2\text{Si}_3\text{O}_{12}$  (a similar  $\text{Ce}^{3+}$ -doped silicate garnet) with

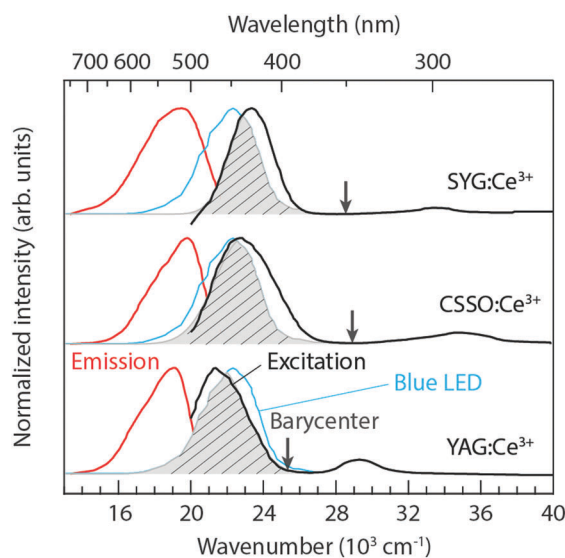


Fig. 2 PL excitation (black color) and emission (red color) spectra for  $\text{SYG}:\text{Ce}^{3+}$ ,  $\text{CSSO}:\text{Ce}^{3+}$ , and  $\text{YAG}:\text{Ce}^{3+}$ , as measured at room temperature. The arrows indicate the  ${}^2E_g$  barycenter of the  $\text{Ce}^{3+} 5d$  levels. The emission spectrum (blue color) corresponding to an InGaN blue LED is included for comparison. The PL ( $4f \rightarrow 5d_1$  and  $5d_1 \rightarrow 4f$ ) spectra of the three materials are normalized to the emission maximum of the blue LED. The spectral overlaps between the materials (excitation) and the blue LED (emission) are marked by the shaded grey areas.



**Table 1** Luminescence data for SYG:Ce<sup>3+</sup>, CSSO:Ce<sup>3+</sup>, and YAG:Ce<sup>3+</sup>, as described in the text

Sample	$\lambda_{5d_2}$ (nm)	$\lambda_{5d_1}$ (nm)	$\lambda_{4f}$ (nm)	$\Delta E_{5d_{1-2}}$ (cm <sup>-1</sup> )	$\Delta S$ (cm <sup>-1</sup> )	(x, y)	CCT (K)	$T_{80\%}^a$ (K)	$E_{5d_1-CB}^b/E_a^c$ (eV)
SYG:Ce <sup>3+</sup>	297	427	513	10 251	3926	(0.314, 0.509)	5911	280	0.45/0.33
CSSO:Ce <sup>3+</sup>	285	440	504	12 360	2886	(0.273, 0.542)	6749	> 860	1.36/0.5–1.0 <sup>d</sup>
YAG:Ce <sup>3+</sup>	340	470	525	8135	2229	(0.355, 0.576)	5130	550	1.17/0.45

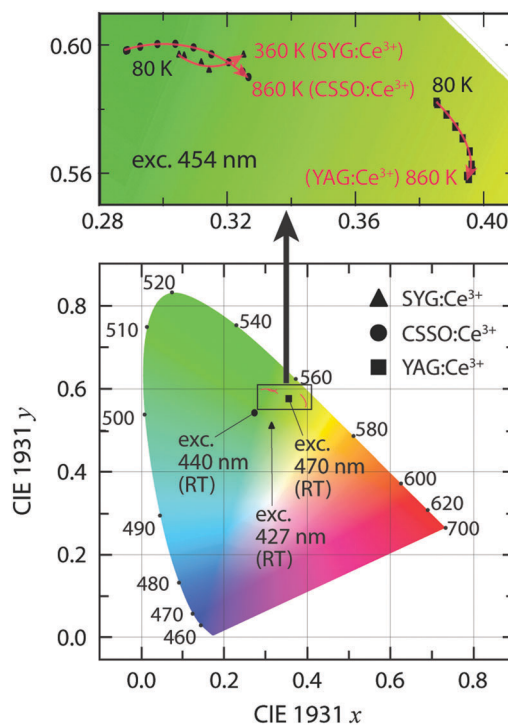
<sup>a</sup>  $T_{80\%}$  is the temperature at which the luminescence decay time has decreased to 80% with respect to the value at low temperature (here 80 K), *i.e.* with respect to that of the “intrinsic” decay time also known as lifetime. <sup>b</sup> Estimated from the VRBE diagrams (Fig. 6). <sup>c</sup>  $E_a$  is obtained from the fits of the temperature dependent luminescence decay time [Fig. 5(d)] to a single-barrier quenching model. <sup>d</sup> Estimated from  $E_{5d_1-CB}$  of CSSO:Ce<sup>3+</sup>, using the relation between  $E_{5d_1-CB}$  and  $E_a$  of SYG:Ce<sup>3+</sup> and YAG:Ce<sup>3+</sup>, respectively.

$\Delta E_{5d_{1-2}} = 11\,000\text{--}12\,000\text{ cm}^{-1}$ .<sup>27</sup> In the case of YAG:Ce<sup>3+</sup>, it is also comparable with previously reported ones, *e.g.* 8135 cm<sup>-1</sup> (this work) *vs.* 7600 cm<sup>-1</sup> by Setlur *et al.*<sup>27</sup> The spectroscopic data of  $\lambda_{5d_1}$ ,  $\lambda_{5d_2}$ , and  $\Delta E_{5d_{1-2}}$  are summarized in Table 1.

The Ce<sup>3+</sup> 4f → 5d<sub>1</sub> excitation band of YAG:Ce<sup>3+</sup> is well located within the InGaN blue emission. The spectral overlap between the excitation band of YAG:Ce<sup>3+</sup> and the emission band of an InGaN blue LED is about 85% under the circumstance that the maxima of the 4f → 5d<sub>1</sub> excitation bands are normalized to the emission maximum of the blue LED, see the shaded grey area in Fig. 2. This suggests a high QE of the conversion of blue light into green-yellow light, which is an important criterion for most practical applications.<sup>28</sup> In comparison, the spectral overlaps are around 83% and 58%, for CSSO:Ce<sup>3+</sup> and SYG:Ce<sup>3+</sup>, respectively, suggesting a lower QE for these materials.

The emission spectra (red color) were measured by fixing the excitation wavelength at 440 nm (CSSO:Ce<sup>3+</sup>), 427 nm (SYG:Ce<sup>3+</sup>), and 470 nm (YAG:Ce<sup>3+</sup>). The spectral maxima of the 5d<sub>1</sub> → 4f emission taking place after a relaxation (Stokes shift,  $\Delta S$ ) from higher to lower vibrational states of the Ce<sup>3+</sup> 5d<sub>1</sub> level, as denoted by  $\lambda_{4f}$  in Table 1, are shifted with respect to each other, with the spectrum peaking at 504 nm for CSSO:Ce<sup>3+</sup> ( $\Delta S = 2886\text{ cm}^{-1}$ ), at around 513 nm for SYG:Ce<sup>3+</sup> ( $\Delta S = 3926\text{ cm}^{-1}$ ), and at around 525 nm for YAG:Ce<sup>3+</sup> ( $\Delta S = 2229\text{ cm}^{-1}$ ). A trend of shifting the emission color from green for CSSO:Ce<sup>3+</sup> and SYG:Ce<sup>3+</sup> to greenish-yellow light for YAG:Ce<sup>3+</sup> is evident from the CIE 1931 diagram (Fig. 3). The corresponding color coordinates (x and y) are shown in Table 1. Included in Table 1 are also the correlated color temperatures (CCTs) at  $T = 300\text{ K}$ , as determined from the respective color coordinates.

The relation of the emission wavelengths between the three materials exhibits an opposite trend as one could expect based on the strength of their tetragonal crystal-field splitting (see  $\Delta E_{5d_{1-2}}$  in Table 1). A larger  $\Delta E_{5d_{1-2}}$  correlates with a more blue-shifted spectrum (smaller  $\lambda_{4f}$ ), which we interpret as the <sup>2</sup>E<sub>g</sub> barycenter of the lower 5d levels is ordered in energy as follows: CSSO:Ce<sup>3+</sup> > SYG:Ce<sup>3+</sup> > YAG:Ce<sup>3+</sup>, as indicated by the arrows in Fig. 2. The difference in the <sup>2</sup>E<sub>g</sub> barycenter may be related to the size-mismatch between Ce<sup>3+</sup> (1.143 Å) and the substituting host atom, Ca<sup>2+</sup> (1.12 Å), Y<sup>3+</sup> (1.019 Å), and Sr<sup>2+</sup> (1.27 Å), in 8-fold coordination.<sup>15</sup> The size-mismatch ratios are approximately 2%, 10%, and 12% for CSSO:Ce<sup>3+</sup>, SYG:Ce<sup>3+</sup>, and YAG:Ce<sup>3+</sup>, respectively. Therefore, a larger size-mismatch appears to lower the energy level of the <sup>2</sup>E<sub>g</sub> barycenter, even



**Fig. 3** CIE 1931 color coordinate diagram for SYG:Ce<sup>3+</sup>, CSSO:Ce<sup>3+</sup> and YAG:Ce<sup>3+</sup>, upon excitation at 427, 440, and 470 nm, respectively, at room temperature (RT). The variable temperature CIE coordinates of the three materials upon excitation at 454 nm are presented in detail in the upper image.

though the local 8-fold environment surrounding the Ce<sup>3+</sup> ions is less tetragonally distorted.

### 3.2 Temperature dependence of the emission spectra

To investigate the thermal stability of luminescence of the three garnet phosphors, both the emission intensity and the decay time of the Ce<sup>3+</sup> emission were evaluated. It should be noted that because of the complexity of the various factors that can contribute to the temperature dependence of the luminescence intensity (*e.g.*, temperature dependence of the absorption strength and temperature dependence of energy migration and reabsorption), an absolute, quantitative analysis of the (intrinsic) temperature quenching is difficult from the emission intensity data alone.<sup>22</sup> Hence, for an evaluation of the intrinsic quenching process, luminescence decay time data are expected to provide a better, more accurate description.



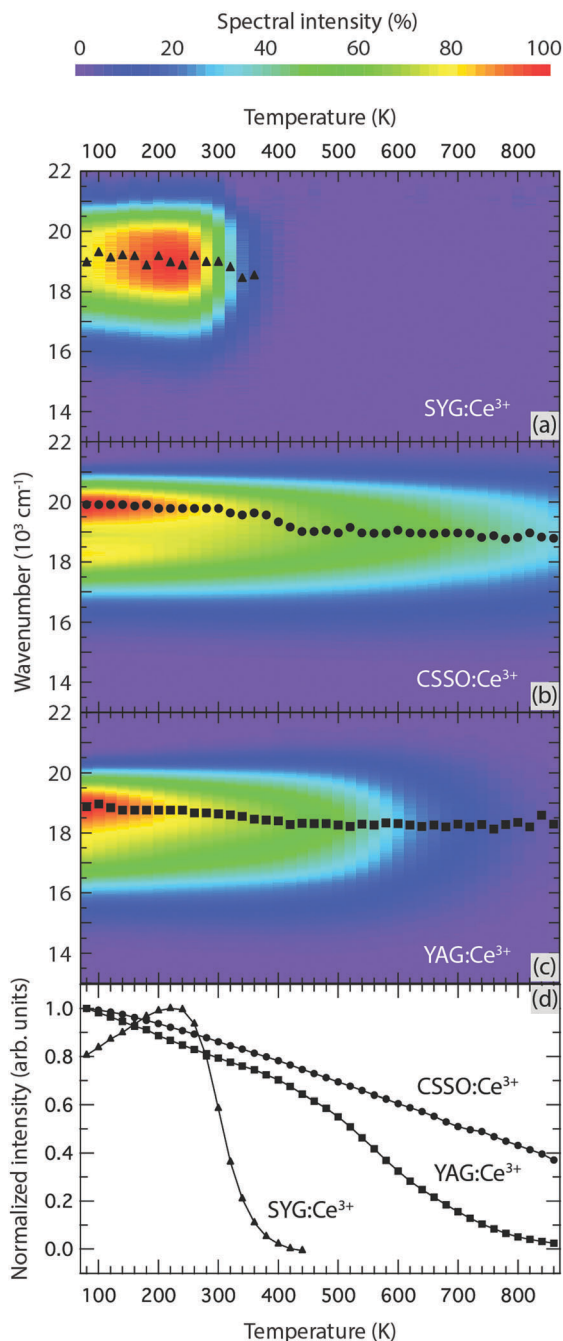


Fig. 4 Temperature dependence of the PL emission spectra and emission maxima [triangles, circles, and squares in (a–c)] of (a) SYG:Ce<sup>3+</sup>, (b) CSSO:Ce<sup>3+</sup>, and (c) YAG:Ce<sup>3+</sup>, upon excitation at 454 nm. (d) Temperature dependence of the normalized emission intensity for SYG:Ce<sup>3+</sup>, CSSO:Ce<sup>3+</sup>, and YAG:Ce<sup>3+</sup>.

Fig. 4(a–c) shows the emission spectra for excitation at 454 nm for all three samples over the temperature range of  $T = 80$ – $860$  K. Fig. 4(d) compares the temperature dependence of the emission intensity, as integrated over the energy region of  $13\,000$ – $22\,000$   $\text{cm}^{-1}$  and normalized to the maximum of the integrated intensity at  $T = 80$ – $860$  K. For CSSO:Ce<sup>3+</sup>, the integrated emission intensity exhibits a quite monotonic decrease upon increase of the temperature from  $T = 80$  K up

to  $T = 860$  K. For SYG:Ce<sup>3+</sup>, the intensity increases from  $T = 80$  K to approximately  $T = 240$  K, which is followed by an abrupt decrease at  $T > 240$  K until the luminescence intensity is virtually completely quenched at  $T \approx 380$ – $400$  K. For YAG:Ce<sup>3+</sup>, the intensity shows an almost linear decrease between  $T = 80$  K and  $T = 500$  K, followed by a more pronounced decrease upon further temperature increase. Similar observations have been done in previous studies.<sup>22</sup> The slow decrease of the luminescence intensity at  $T < 500$  K is caused mainly by a reduction of the absorption strength of the  $4f \rightarrow 5d_1$  transition. This is because the second level of the  $^2F_{5/2}$  level is more thermally populated and the transition from this level to the lowest  $5d$  level is symmetry forbidden.<sup>22,29</sup>

In contrast to the broadband emission spectra at RT (Fig. 2), the emission spectra at  $T = 80$  K show a distinguished double peak signature [especially for CSSO:Ce<sup>3+</sup>, cf. Fig. 4(b)] corresponding to radiative  $5d_1 \rightarrow ^2F_{5/2}$  (higher energy) and  $5d_1 \rightarrow ^2F_{7/2}$  (lower energy) transitions, respectively. To determine the spectra of these two radiative transitions, two Gaussian functions were used to fit the doublet band in each spectrum, which reveals that the energy separation of the maxima of the  $^2F_{5/2}$  and  $^2F_{7/2}$  bands is approximately  $1710$   $\text{cm}^{-1}$ ,  $1410$   $\text{cm}^{-1}$ , and  $1560$   $\text{cm}^{-1}$  at  $T = 80$  K for CSSO:Ce<sup>3+</sup>, SYG:Ce<sup>3+</sup>, and YAG:Ce<sup>3+</sup>, respectively. The energy difference between the  $^2F_{5/2}$  and  $^2F_{7/2}$  levels is comparable to the expected value of  $2000 \pm 250$   $\text{cm}^{-1}$  for Ce<sup>3+</sup>-doped garnet phosphors.<sup>11</sup> The different values between CSSO:Ce<sup>3+</sup>, SYG:Ce<sup>3+</sup>, and YAG:Ce<sup>3+</sup> is due to a difference in the spin–orbit coupling strength as Ce<sup>3+</sup> ions interact with the neighboring atoms in the respective hosts.<sup>29</sup> Furthermore, the emission maximum exhibits a shift toward longer wavelengths (red-shift) upon increasing temperature for all phosphors in the measured temperature range, see Fig. 4(a–c). Note that the emission maximum of SYG:Ce<sup>3+</sup> is only shown up to  $T = 360$  K since the spectra at  $T > 360$  K are too weak to determine the emission maximum. The red-shift trend of the emission maximum upon temperature increase is in accordance with the temperature behavior of the color coordinates (Fig. 3). For YAG:Ce<sup>3+</sup>, the red-shift of the color coordinates persists up to around  $T = 600$  K, which is in accordance with the results reported in ref. 19 and 30, whereas it turns to a blue-shift at  $T > 600$  K. The blue-shift for YAG:Ce<sup>3+</sup> at  $T > 600$  K was not observed in previous studies due to limited highest temperatures ( $T \leq 573$  K).<sup>19,30</sup> It is worth noting that the blue-shift of the color coordinates of YAG:Ce<sup>3+</sup> from  $T = 600$  K to  $T = 860$  K (Fig. 3) is not evident for the emission maximum in the same temperature range [Fig. 4(c)]. The discrepancy between the color coordinates and the emission maximum originates from the way in which they have been obtained. For the calculation of the color coordinates  $x$  and  $y$ , the whole emission spectrum, which is determined by all spectral parameters such as the position, width and intensity of the  $^2F_{5/2}$  and  $^2F_{7/2}$  bands, is used. However, the emission maximum considers only the intensity of the superposition of the two bands. The latter implies that the emission maximum only contains partial information of the spectrum. This information depends on the extent of the spectral overlap between the  $^2F_{5/2}$  and  $^2F_{7/2}$  bands



as explained by the following example: the emission maximum for  $\text{CSSO}:\text{Ce}^{3+}$  at  $T = 80 \text{ K}$  [Fig. 4(b)] is mainly determined by the  $^2\text{F}_{5/2}$  band. Unlike the color coordinates, the emission maximum only reflects the emission color of part of the spectrum, particularly when the separation between the  $^2\text{F}_{5/2}$  and  $^2\text{F}_{7/2}$  bands is large. Therefore, to investigate the shift of the emission color as a function of temperature, the variation of the color coordinates, as shown in the upper image in Fig. 3, reflects the effect in a more accurate manner. It is also noted that the RT color coordinates (Fig. 3) are different between the ones derived from the emission spectra at RT in Fig. 2 and 4(a–c). This is likely an effect of that the measurements were performed with different excitation wavelengths.

On a fundamental level, the color shifting effect may be associated with a change in the tetragonal crystal field acting on the  $\text{Ce}^{3+}$  ions upon heating, *e.g.*, through the increase of vibrationally induced tetragonal distortions of the local  $\text{CeO}_8$  moieties (red-shift), in competition with local thermal lattice expansion that, conversely, reduces the tetragonal crystal field (blue-shift).<sup>31</sup> As the lattice size increases with elevating temperature, the local symmetry of  $\text{CeO}_8$  becomes more cubic-like due to the elongation of the Ce–O bonds.<sup>9</sup> The red-shift may, however, be also explained by thermally activated energy migration and energy transfer to  $\text{Ce}^{3+}$  ions emitting at longer wavelength as the excitation and emission bands are thermally broadened.

### 3.3 Luminescence decay time measurements

Fig. 5(a–c) shows the luminescence decay curves for  $\text{SYG}:\text{Ce}^{3+}$ ,  $\text{CSSO}:\text{Ce}^{3+}$ , and  $\text{YAG}:\text{Ce}^{3+}$  for excitation at 454 nm, at a selection of temperatures. Data for all temperatures are given in Fig. S1 (ESI<sup>†</sup>). All decay curves can be adequately approximated by single exponential fits. Fig. 5(d) compares the luminescence decay time,  $\tau(T)$ , over the measured temperature ranges, as extracted from the single exponential fits to the  $T$  dependent data in (a–c). In the low-temperature regime ( $T = 80\text{--}200 \text{ K}$ ), the luminescence decay time remains nearly constant, *i.e.* 52, 65, and 94 ns for  $\text{SYG}:\text{Ce}^{3+}$ ,  $\text{CSSO}:\text{Ce}^{3+}$ , and  $\text{YAG}:\text{Ce}^{3+}$ , respectively. The “intrinsic” decay time (lifetime) of the three  $\text{Ce}^{3+}$ -doped garnet phosphors is within the range of the lifetime of  $\text{Ce}^{3+}$  luminescence ( $<100 \text{ ns}$ ) in many host materials.<sup>32</sup> However, the lifetime of  $\text{YAG}:\text{Ce}^{3+}$  (94 ns) is higher than the values (60–67 ns) as reported in the literature by Bachmann *et al.*<sup>22</sup> for  $\text{YAG}:\text{Ce}^{3+}$  for the  $\text{Ce}^{3+}$  dopant concentrations of 0.033–3.333%. Fig. 5(d) also compares the temperature dependent decay time of  $\text{YAG}:\text{Ce}^{3+}$  between this work and the one reported in ref. 22. The longer lifetime of  $\text{YAG}:\text{Ce}^{3+}$  observed here may be attributed to enhanced re-absorption processes due to a larger overlap between the excitation and emission spectra (Fig. 2), compared to the spectra reported by Bachmann *et al.*<sup>22</sup> The difference in spectral overlap may stem from different synthesis details such as starting reagents and heat treatments, which may, for example, affect the distribution (homogeneity) of  $\text{Ce}^{3+}$  ions in the YAG host lattice, and lead to differences in the distance between  $\text{Ce}^{3+}$  ions. Such differences in the Ce–Ce distance should affect the energy transfer among  $\text{Ce}^{3+}$  ions, which shifts the position of the

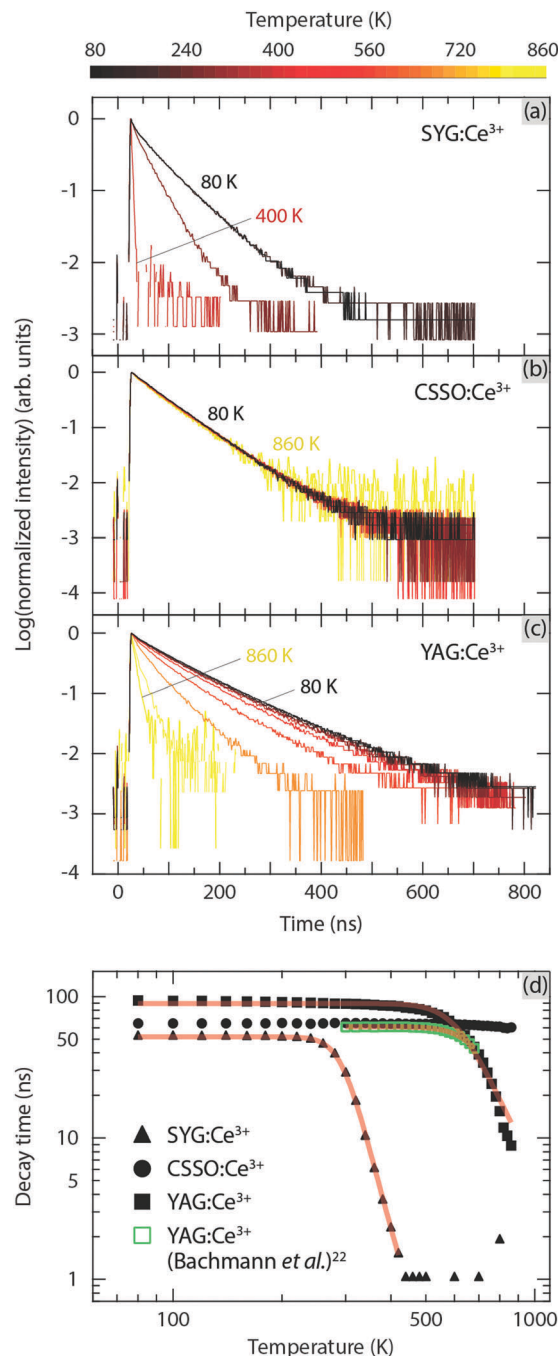


Fig. 5 (a) Temperature dependence of the luminescence decay curves for (a)  $\text{SYG}:\text{Ce}^{3+}$ , (b)  $\text{CSSO}:\text{Ce}^{3+}$ , and (c)  $\text{YAG}:\text{Ce}^{3+}$ . (d) Luminescence decay time determined from the single-exponential fits in (a–c) as well as the decay time of  $\text{YAG}:\text{Ce}^{3+}$  from Bachmann *et al.*,<sup>22</sup> as a function of temperature. Red solid curves are fits using a single-barrier quenching model.<sup>33,34</sup>

emission spectrum<sup>22</sup> and hence alters the spectral overlap, *i.e.* the probability of re-absorption processes.

Upon further temperature increase ( $T > 200 \text{ K}$ ), the three materials show different behaviors. The onset temperature of thermal quenching for  $\text{SYG}:\text{Ce}^{3+}$ ,  $\text{CSSO}:\text{Ce}^{3+}$ , and  $\text{YAG}:\text{Ce}^{3+}$  are at approximately  $T = 220 \text{ K}$ ,  $600 \text{ K}$ , and  $400 \text{ K}$ , respectively.



In agreement with the recent study on CSSO:Ce<sup>3+</sup>,<sup>24</sup> its luminescence decay time is approximately constant (71 ns) up to  $T = 600\text{--}700$  K, which confirms an excellent internal thermal stability up to at least  $T = 600$  K, although we observe a slightly different lifetime (65 ns) as compared to previous work.<sup>24</sup> At  $T > 600$  K, the decay time of CSSO:Ce<sup>3+</sup> exhibits a relatively slow decrease from 64 ns at  $T = 600$  K to 60 ns at  $T = 860$  K, which further confirms that CSSO:Ce<sup>3+</sup> has strong resistance against thermal quenching at high temperatures (up to at least  $T = 860$  K). In contrast, the luminescence decay time of SYG:Ce<sup>3+</sup> decreases abruptly from about 51 ns at  $T = 220$  K to 29 ns at  $T = 300$  K, and to 2 ns at  $T = 400$  K at which the luminescence is nearly completely quenched. For YAG:Ce<sup>3+</sup>, the luminescence decay time decreases abruptly from  $\tau = 86$  ns at  $T = 400$  K to  $\tau = 63$  ns at  $T = 600$  K, and to  $\tau = 9$  ns at  $T = 860$  K. The temperature at which the luminescence decay time has dropped to 80% of its “intrinsic” value,  $T_{80\%}$ , is approximately  $T = 280$  K,  $T > 860$  K, and  $T = 550$  K, for SYG:Ce<sup>3+</sup>, CSSO:Ce<sup>3+</sup>, and YAG:Ce<sup>3+</sup>, respectively. The resistance to thermal quenching of luminescence for  $5d_1 \rightarrow 4f$  emission for the three materials is hence ranked as follows, CSSO:Ce<sup>3+</sup> > YAG:Ce<sup>3+</sup> > SYG:Ce<sup>3+</sup>. It is worth noting that the  $T_{80\%}$  of YAG:Ce<sup>3+</sup> with 1% doping concentration ( $T = 550$  K) is almost 100 K lower than that ( $T = 645$  K) of the same material synthesized with different starting reagents and heat treatments.<sup>22</sup> This is in accordance with the longer lifetime obtained here (94 ns) compared to the one ( $\approx 61$  ns) reported in ref. 22.

The overall thermal quenching behavior can be adequately fitted to a single-barrier quenching model [Fig. 5(d)], *i.e.*  $\tau(T) = 1/[\Gamma_\nu + \Gamma_0 \exp(-E_a/k_B T)]$ , where  $\Gamma_\nu$  is the radiative rate,  $\Gamma_0$  is the attempt rate for non-radiative processes,  $E_a$  is the activation energy for the thermal quenching processes, and  $k_B$  is the

Boltzmann constant.<sup>33,34</sup> From a free fit, we obtain  $E_a$  of  $0.33 \pm 0.005$  eV (SYG:Ce<sup>3+</sup>), and  $0.45 \pm 0.002$  eV (YAG:Ce<sup>3+</sup>), see Table 1. In comparison, the  $E_a$  of thermal quenching of YAG:1%Ce<sup>3+</sup> estimated from the data of Bachmann *et al.*<sup>22</sup> is  $0.53 \pm 0.007$  eV [Fig. 5(d)]. This is in good agreement with the one obtained here (0.45 eV), which suggests that the thermal quenching of luminescence of the two YAG:Ce<sup>3+</sup> samples is driven by similar quenching process(es). As for CSSO:Ce<sup>3+</sup>, the decay time is almost constant up to  $T = 860$  K (*i.e.*, the highest applied  $T$ ), which makes a free fit to the data unreliable. In order to estimate  $E_a$  for CSSO:Ce<sup>3+</sup>, we employed an empirical relation based on the results of SYG:Ce<sup>3+</sup> and YAG:Ce<sup>3+</sup>, which shows that  $E_a$  is lowered by 26–62% with respect to  $E_{5d_1\text{-CB}}$  (Table 1), whose details are described in the next section. With this method, we estimate  $E_a$  to 0.5–1.0 eV for CSSO:Ce<sup>3+</sup>. In order to interpret in more detail these experimental findings, especially in regard to the positions of the ground and excited configurations of the Ce<sup>3+</sup> dopants with respect to the valence and conduction bands of the respective hosts, we have constructed the VRBE diagrams of Ce<sup>3+</sup> within SYG, CSSO, and YAG, as described in the following.

### 3.4 Vacuum referred binding energy diagrams

The way in which we have built the VRBE diagrams for the different hosts is based on published methods.<sup>11,35–43</sup> The model, called the chemical shift model, provides electronic structure with absolute binding energies relative to the energies of the electron at rest, in vacuum.

Using the spectroscopic data of the Ce<sup>3+</sup> luminescence from the present work (Fig. 2 and Table 1) in comparison with previous studies, and the data for the Eu<sup>2+</sup> and Eu<sup>3+</sup> dopants in the studied hosts,<sup>11,44–46</sup> the VRBE diagrams were constructed

Table 2 Compilation of VRBE parameters and spectroscopic data for the phosphors, SYG, CSSO, and YAG, as doped with Eu<sup>2+</sup>, Eu<sup>3+</sup>, or Ce<sup>3+</sup>

VRBE parameters	SYG	CSSO	YAG
• Input data, obtained from literature and Table 1			
$E_V$ (eV), top of valence band	−8.43 <sup>44</sup>	−8.90 <sup>45</sup>	−9.38 <sup>11</sup>
$E^{ex}$ (eV), exciton creation energy	5.90 <sup>44a</sup>	6.81 <sup>45a</sup>	7.10 <sup>11</sup>
$E_{VC}$ (eV), band gap energy	6.37 <sup>44</sup>	7.35 <sup>45</sup>	7.67 <sup>11a</sup>
$E^{CT}$ (eV), charge transfer energy of $O^{2-} \rightarrow Eu^{3+}$	4.43 <sup>44</sup>	4.86 <sup>45</sup>	5.42 <sup>11</sup>
$U$ (eV), 4f-electron Coulomb repulsion energy of $Eu^{3+}$	6.66 <sup>44b</sup>	6.93 <sup>45</sup>	6.80 <sup>11</sup>
$\lambda_{5d_1}$ (nm), band maximum of $4f \rightarrow 5d_1$ excitation (this work/reference)	427/435 <sup>44</sup>	440/450 <sup>46</sup>	470/457 <sup>11</sup>
$\lambda_{5d_2}$ (nm), band maximum of $4f \rightarrow 5d_2$ excitation (this work/reference)	297/300 <sup>44</sup>	285/—	340/340 <sup>11</sup>
• Output data, calculated using the input data above			
$E_C$ (eV), bottom of conduction band <sup>c</sup>	−2.06	−1.55	−1.71
$E_{Eu^{2+}4f}$ (eV), 4f ground level of $Eu^{2+}$ <sup>d</sup>	−4.00	−4.04	−3.96
$E_{Eu^{3+}4f}$ (eV), 4f ground level of $Eu^{3+}$ <sup>e</sup>	−10.66	−10.97	−10.76
$E_{Ce^{3+}4f}$ (eV), 4f ground level of $Ce^{3+}$ <sup>f</sup>	−5.42	−5.73	−5.52
$E_{Ce^{3+}5d_1}$ (eV), $5d_1$ excited level of $Ce^{3+}$ (this work/reference) <sup>g</sup>	−2.51/−2.57	−2.91/−2.97	−2.88/−2.81
$E_{Ce^{3+}5d_2}$ (eV), $5d_2$ excited level of $Ce^{3+}$ (this work/reference) <sup>g</sup>	−1.24/−1.29	−1.38/—	−1.87/−1.87
$E_{5d_1\text{-CB}}$ (eV), difference between $E_{Ce^{3+}5d_1}$ and $E_C$ (this work/reference)	0.45/0.51	1.36/1.42	1.17/1.10
$E_{5d_2\text{-CB}}$ (eV), difference between $E_{Ce^{3+}5d_2}$ and $E_C$ (this work/reference)	−0.82/−0.77	−0.17/—	0.16/0.16
$D$ (eV), red-shift of the lowest 5d level of $Ce^{3+}$ (this work/reference) <sup>h</sup>	3.22/3.27	3.30/3.36	3.48/3.41

<sup>a</sup> Derived from the relation  $E_{VC} = (1 + 8\%)E^{ex}$ .<sup>11b</sup> Derived from the known  $E_{Ce^{3+}4f} = -5.42$  eV,<sup>44</sup> and using the relation  $E_{Ce^{3+}4f} = E_{Eu^{2+}4f} - U + \Delta E_{4f}^{3+}$ , where  $\Delta E_{4f}^{3+} (=5.24$  eV) is the difference between the 4f binding energies of  $Eu^{3+}$  and  $Ce^{3+}$ .<sup>11c</sup> Derived from the relation  $E_C = E_V + E_{VC}$ ,<sup>11</sup> see arrow 1 in Fig. 6. <sup>d</sup> Derived from the relation  $E_{Eu^{2+}4f} = E_V + E^{CT}$ ,<sup>11</sup> see arrow 2 in Fig. 6. <sup>e</sup> Derived from the relation  $E_{Eu^{3+}4f} = E_{Eu^{2+}4f} - U$ ,<sup>11</sup> see arrow 3 in Fig. 6. <sup>f</sup> Derived from the relation  $E_{Ce^{3+}4f} = E_{Eu^{3+}4f} + \Delta E_{4f}^{3+}$ , where  $\Delta E_{4f}^{3+}$  is defined in b,<sup>11</sup> see arrow 4 in Fig. 6. <sup>g</sup> Derived from the relation  $E_{Ce^{3+}5d_i} = E_{Ce^{3+}4f} + 1240/\lambda_{5d_i}$ , where  $i = 1$  and  $2$ ,<sup>11</sup> see arrow 5 in Fig. 6. <sup>h</sup> Derived from the relation  $D = E_{Ce^{3+}free} - 1240/\lambda_{5d_1}$ , where  $E_{Ce^{3+}free} = 6.12$  eV.<sup>11</sup>



using input data for  $E_V$ ,  $E^{ex}$ ,  $E_{VC}$ ,  $E^{CT}$ ,  $U$ ,  $\lambda_{5d_1}$ , and  $\lambda_{5d_2}$ , see Table 2. Fig. 6 compares the VRBE diagrams for the three hosts, as doped with  $\text{Eu}^{2+}$ ,  $\text{Eu}^{3+}$ , or  $\text{Ce}^{3+}$ . As seen in Table 2, the 4f-electron Coulomb repulsion energies  $U$  of  $\text{Eu}^{3+}$  in SYG (6.66 eV), CSSO (6.93 eV), and YAG (6.80 eV) are almost the same and are located in the 6.4–7.2 eV range, as typically observed for oxide materials.<sup>43</sup> The major difference between the three phosphors lays in the other parameters, *i.e.*  $E_V$ ,  $E^{ex}$ ,  $E_{VC}$ ,  $E^{CT}$ ,  $\lambda_{5d_1}$  (or  $D$ ), and  $\lambda_{5d_2}$  of the respective hosts, which affect the absolute location of the 5d levels with respect to the bottom of the CB ( $E_C$ ), see Fig. 6 and Table 2. The obtained energies ( $E_{5d_1-CB}$ ) of the  $\text{Ce}^{3+}$  5d<sub>1</sub> level is 0.45 eV (SYG: $\text{Ce}^{3+}$ ), 1.36 eV (CSSO: $\text{Ce}^{3+}$ ), and 1.17 eV (YAG: $\text{Ce}^{3+}$ ) below the CB, which are in good agreement with the results derived from the  $\lambda_{5d_1}$  values, as reported previously [dashed lines in Fig. 6 and  $E_{5d_1-CB}$  (reference) in Table 2]. The larger  $E_{5d_1-CB}$  value means that the electron at the  $\text{Ce}^{3+}$  5d<sub>1</sub> level requires more thermal energy (higher temperature) to be promoted into the CB and the delocalized electron in the CB may be followed by charge trapping (luminescence is “killed”) at defect levels located within the band gap. This type of quenching process has been confirmed by measuring the photoconductivity of  $\text{Y}_3\text{Al}_2\text{Ga}_3\text{O}_{12}:\text{Ce}^{3+}$  and  $\text{Y}_3\text{Ga}_5\text{O}_{12}:\text{Ce}^{3+}$  garnet phosphors when exciting the  $\text{Ce}^{3+}$  4f electron to the 5d<sub>1</sub> level as a function of temperature.<sup>47</sup> In this context, a phosphor with a larger  $E_{5d_1-CB}$  is expected to exhibit higher stability against thermal quenching of luminescence, which is in good agreement with our results, *cf.* CSSO: $\text{Ce}^{3+}$  ( $E_{5d_1-CB} = 1.36$  eV,  $T_{80\%} > 860$  K), YAG: $\text{Ce}^{3+}$  ( $E_{5d_1-CB} = 1.17$  eV,  $T_{80\%} = 550$  K), and SYG: $\text{Ce}^{3+}$  ( $E_{5d_1-CB} = 0.45$  eV,  $T_{80\%} = 280$  K), see Table 1. Similarly, for YAG: $\text{Ce}^{3+}$ , the separation between the  $\text{Ce}^{3+}$  5d<sub>2</sub> level and the bottom of the CB, as denoted by  $E_{5d_2-CB}$ , is much smaller than  $E_{5d_1-CB}$ . This suggests that, at a given temperature, more luminescence of  $\text{Ce}^{3+}$  ions with the excitation at  $\lambda_{5d_2}$  is killed due to charge trapping at defects, as

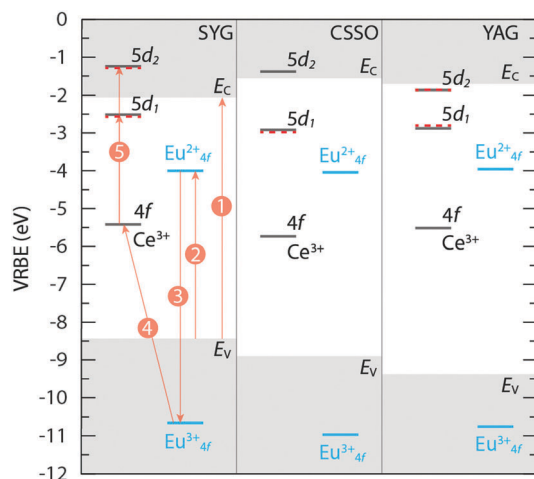


Fig. 6 VRBE diagrams for the three hosts, left: SYG, middle: CSSO, and right: YAG, as doped with  $\text{Eu}^{2+}$ ,  $\text{Eu}^{3+}$ , or  $\text{Ce}^{3+}$ . The arrows 1, 2, 3, 4, and 5 represent the shifts of energy by amounts of  $E_{VC}$ ,  $E^{CT}$ ,  $U$ ,  $\Delta E_{4f}^{3+}$ , and energy of light with wavelength  $\lambda_{5d_1}$  or  $\lambda_{5d_2}$ , respectively. Horizontal dashed lines (red color) represent the 5d<sub>1</sub> and 5d<sub>2</sub> levels of  $\text{Ce}^{3+}$  derived from the reported values of  $\lambda_{5d_1}$  and  $\lambda_{5d_2}$ . Details of all the energy levels and VRBE parameters are listed in Table 2.

compared to the excitation at  $\lambda_{5d_1}$ . This is evident from a comparison of the intensity of the excitation bands centered at  $\lambda_{5d_1}$  (stronger) and  $\lambda_{5d_2}$  (weaker), see Fig. 2. The  $\text{Ce}^{3+}$  5d<sub>2</sub> levels of CSSO: $\text{Ce}^{3+}$  and SYG: $\text{Ce}^{3+}$  are located within the CB, which implies that the luminescence is completely quenched upon excitation at  $\lambda_{5d_2}$ . However, weak  $\lambda_{5d_2}$  bands for CSSO: $\text{Ce}^{3+}$  and SYG: $\text{Ce}^{3+}$  can be still observed (Fig. 2). This may have its origin in a non-negligible amount of electrons in the CB returning to the 5d<sub>1</sub> level, followed by 5d<sub>1</sub> → 4f emission.

## 4 Discussion

The three green-to-yellow-emitting phosphors SYG: $\text{Ce}^{3+}$ , CSSO: $\text{Ce}^{3+}$ , and YAG: $\text{Ce}^{3+}$  have been shown to exhibit quite remarkably different luminescence properties, thus highlighting the important role of garnet host in regard to their optical performance. The nearly constant luminescence decay time ( $\tau \approx 60$  ns) up to very high temperature,  $T = 860$  K, for CSSO: $\text{Ce}^{3+}$  is a unique feature as shown by the present host and is even better than the  $\text{Na}_{3-2x}\text{Sc}_2(\text{PO}_4)_2:x\text{Eu}^{2+}$  blue phosphor, which was recently claimed to be a zero-thermal-quenching phosphor up to  $T = 473$  K.<sup>48</sup> Moreover, the 4f–5d<sub>1</sub> excitation and emission spectral characteristics of CSSO: $\text{Ce}^{3+}$  make it a very good candidate for pc-WLEDs based on a blue LED, in particular, under a high-power operation which generates significant heat in pc-WLEDs, including the phosphor materials.

The difference in intrinsic thermal stability of luminescence in SYG: $\text{Ce}^{3+}$ , CSSO: $\text{Ce}^{3+}$ , and YAG: $\text{Ce}^{3+}$  may have its origin in different quenching processes in the three materials. On a fundamental level, the three primary processes that are thought to play a role in the thermal quenching of luminescence in  $\text{Ce}^{3+}$ -doped materials exhibiting 5d–4f luminescence are (1) thermal ionization of the  $\text{Ce}^{3+}$  5d electrons into the CB of the host crystal, followed by charge trapping at defects, (2) thermally activated nonradiative energy migration among  $\text{Ce}^{3+}$  ions to killer centers (generally known as concentration quenching), and (3) thermally activated crossover from the 5d excited state to the 4f ground state *via* electron–phonon coupling mechanisms.<sup>6,22,31,49</sup> The fact that the  $\text{Ce}^{3+}$  concentration is identical (1%) in the three materials, but their stability against the thermal quenching of luminescence is remarkably different, suggests that concentration quenching is not responsible for the observed difference in the thermal quenching between the three phosphors. On the contrary, any of the other two processes are likely to play a significant role in determining the overall quenching behavior. The trend of increasing the thermal stability of luminescence (*e.g.* higher  $T_{80\%}$ ) with increasing  $E_{5d_1-CB}$  values, see Table 1, would suggest that the quenching behavior is primarily determined by thermal ionization. However, the relatively large difference between  $E_{5d_1-CB}$  and  $E_a$  shows that thermal ionization alone cannot determine the overall quenching behavior. Rather, these results point towards the contribution of both thermal ionization of the  $\text{Ce}^{3+}$  5d electrons into the CB of the host crystal, followed by charge trapping, and thermally activated crossover from the 5d





excited state to the 4f ground state *via* electron–phonon coupling mechanisms. In this context, we have previously shown that the excitation of the high-frequency phonon modes in the range of 600–900  $\text{cm}^{-1}$  for YAG:Ce<sup>3+</sup> promotes crossover non-radiative relaxation.<sup>31</sup> For CSSO, this set of phonon modes are found at somewhat higher vibrational frequencies (800–900  $\text{cm}^{-1}$ ),<sup>50</sup> meaning that a higher temperature is needed for them to be thermally excited. Under the assumption that thermally activated 5d<sub>1</sub> → 4f crossover contributes to the thermal quenching of luminescence in YAG:Ce<sup>3+</sup> and CSSO:Ce<sup>3+</sup>, this would be in line with the observed difference in phonon frequencies. In comparison, the same set of phonon modes for SYG is located in the same frequency range as for YAG, *i.e.* 600–900  $\text{cm}^{-1}$  (see Fig. S2, ESI<sup>†</sup>).<sup>51</sup> This implies that the lower quenching temperature of SYG:Ce<sup>3+</sup>, see  $T_{80\%}$  in Table 1, is primarily attributed to its relatively smaller  $E_{5d_1\text{-CB}}$  (0.45 eV) as compared to that of YAG:Ce<sup>3+</sup> (1.17 eV). In other words, this suggests that thermal quenching of luminescence in SYG:Ce<sup>3+</sup> is driven primarily by thermal ionization rather than by 5d<sub>1</sub> → 4f crossover relaxation.

The excellent thermal stability of CSSO:Ce<sup>3+</sup> up to  $T = 860$  K, suggests that it is a promising phosphor for application in efficient pc-WLEDs, which require a high thermal stability of luminescence up to at least  $T = 450$  K (the temperature of the phosphors) in today's available LED packages.<sup>32</sup> Fig. 4(d) shows that the luminescence intensity of CSSO:Ce<sup>3+</sup> at  $T = 450$  K is about 26% lower with respect to that at  $T = 80$  K. Since the decay time of CSSO:Ce<sup>3+</sup> remains almost constant up to  $T = 860$  K, this reduction of the luminescence intensity is mainly attributed to the decrease of the absorption strength of the 4f → 5d<sub>1</sub> transition, as in the case for YAG:Ce<sup>3+</sup>.<sup>22</sup> In comparison to YAG:Ce<sup>3+</sup>, which shows a decrease of the luminescence intensity by  $\approx 37\%$  at  $T = 450$  K with respect to that at  $T = 80$  K [Fig. 4(d)], CSSO:Ce<sup>3+</sup> shows a better thermal stability in terms of luminescence intensity, in agreement with the literature.<sup>52</sup> The difference in thermal stability of luminescence intensity between CSSO:Ce<sup>3+</sup> and YAG:Ce<sup>3+</sup> tends to become larger at higher temperatures, *e.g.* differences of 11% at  $T = 450$  K, and 28% at  $T = 600$  K [Fig. 4(d)], which suggests that CSSO:Ce<sup>3+</sup> is a better candidate for high power pc-WLEDs.

Fig. 7(a) shows the PL emission spectra of the pc-WLED prototype. The prototype consists of the CSSO:Ce<sup>3+</sup> phosphor deposited on a commercial InGaN blue LED ( $\lambda \approx 420$  nm) under a voltage of 3.4, 3.6, 3.8, and 4.0 V, respectively. The integrated intensity of the emitted light of the pc-WLED prototype shows a linear increase with increasing applied voltage [Fig. 7(b)]. This suggests that the variation of the temperature of the blue LED chip upon applied voltage within the typical operating range for InGaN LEDs (3.4–4.0 V) has no significant impact on the strength of the 4f–5d transition (both blue absorption and green emission) of CSSO:Ce<sup>3+</sup> when coated on top of the LED chip. This agrees with the observed excellent thermal stability of the emission intensity of CSSO:Ce<sup>3+</sup> [Fig. 4(d)]. Under an applied voltage of 3.4 V, the prototype exhibits a luminous efficacy of optical radiation (LER) of 243  $\text{lm W}^{-1}$  ( $\approx 4.191 \times 10^{-3} \text{ lm}/1.725 \times 10^{-5} \text{ W}$ , defined as the ratio of the

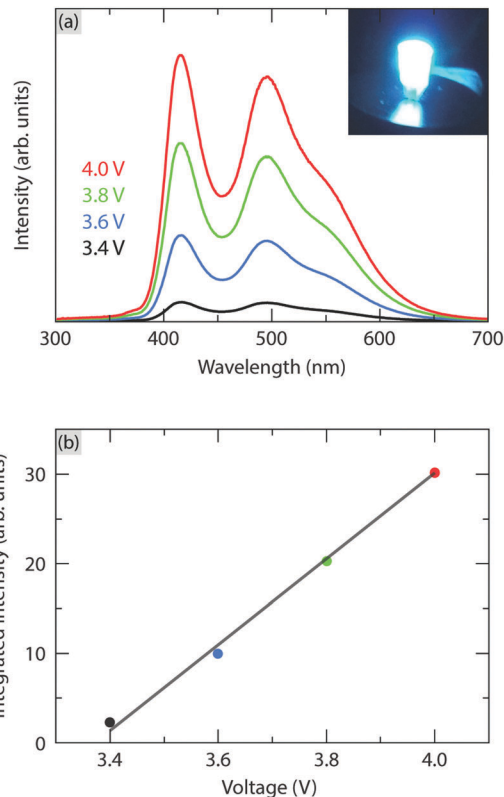


Fig. 7 (a) Emission spectra of the pc-WLED prototype coated with CSSO:Ce<sup>3+</sup> under operation with a voltage of 3.4 V, 3.6 V, 3.8 V, and 4.0 V, respectively. Inset: Bluish white emission from the pc-WLED operated with 3.4 V. (b) Integrated intensity of the emission spectrum of the pc-WLED prototype as a function of applied voltage.

total luminous flux to radiant power); the inset in Fig. 7(a) shows a photograph of the pc-WLED prototype under this bias. The prototype has a color rendering index (CRI) of 58, which is

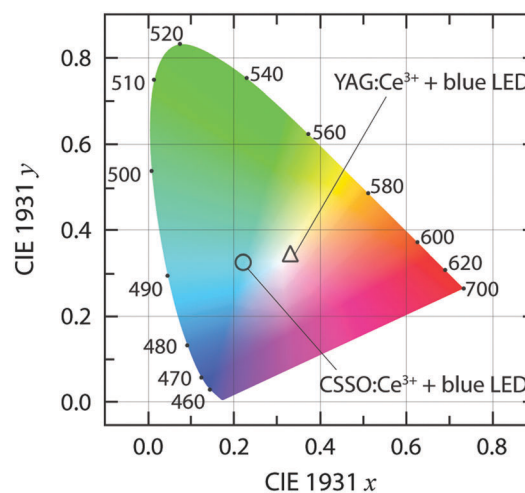


Fig. 8 CIE 1931 color coordinate diagram for the pc-WLED prototype based on CSSO:Ce<sup>3+</sup> and a blue LED [CIE 1931  $xy = (0.220, 0.324)$ , and CCT = 13 604 K] as derived from the spectrum of the prototype operated with 3.4 V in Fig. 7(a), which is compared to the one based on YAG:Ce<sup>3+</sup> and a blue LED [CIE 1931  $xy = (0.330, 0.339)$ , and CCT = 5600 K].



calculated using the Python colour science package called "Colour" with the measured spectral power distribution [Fig. 7(a)].<sup>53</sup> In comparison with other pc-WLEDs with similar CRIs (57–63), the LER obtained here is generally lower than the others (370–430 lm W<sup>-1</sup>),<sup>54</sup> which may be caused by a larger fraction of blue light leaking through the phosphor.<sup>28</sup> This relatively strong blue light is also evident from the CIE coordinates of the pc-WLED prototype that lay near the region of bluish white color or blue color (Fig. 8). Thus, the light emitted from the prototype is perceived as cold white light. The prototype exhibits a CCT = 13 604 K, which may be compared to that of the well known system comprising a blue LED and YAG:Ce<sup>3+</sup> (CCT = 5600 K),<sup>28</sup> see Fig. 8. The white balance and the CRI of the pc-WLED prototype may thus be improved by the addition of a red-emitting phosphor and/or red-emitting core/shell quantum dots to the present system, as implemented in other studies.<sup>28,52,55</sup>

## 5 Conclusions

To conclude, we have investigated the luminescence properties of the three garnet type phosphors CSSO:Ce<sup>3+</sup>, SYG:Ce<sup>3+</sup> and YAG:Ce<sup>3+</sup>. By monitoring the temperature dependence of the luminescence decay time, we establish an excellent, intrinsic, thermal stability of luminescence in CSSO:Ce<sup>3+</sup> with a nearly constant decay time ( $\approx 60$  ns) up to  $T = 860$  K. In comparison, SYG:Ce<sup>3+</sup> and YAG:Ce<sup>3+</sup> exhibit a significant reduction of the luminescence decay time, starting at around  $T = 280$  K and  $T = 550$  K, respectively, suggesting a lower internal thermal stability of luminescence in these materials. These experimental results are supported by the energy separation between the Ce<sup>3+</sup> 5d<sub>1</sub> level and the CB of the respective hosts, which are found at 1.36 eV (CSSO:Ce<sup>3+</sup>), 0.45 eV (SYG:Ce<sup>3+</sup>), and 1.17 eV (YAG:Ce<sup>3+</sup>), respectively, as predicted by their VRBE diagrams. The performance of CSSO:Ce<sup>3+</sup> was further evaluated by applying the phosphor on a blue InGaN LED. The system shows a luminous efficacy of optical radiation of 243 lm W<sup>-1</sup> and a linear response with increasing applied voltage.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

Financial support from the Swedish Research Council Formas (Grant No. 2013-1723), and the European Commission for funding through the Marie Curie Initial Training network LUMINET (Grant No. 316906) is gratefully acknowledged. We also thank Erica Viviani, University of Verona, for expert technical assistance, and Jumpei Ueda, Kyoto University, for fruitful discussions.

## References

- 1 S. Pimputkar, J. S. Speck, S. P. DenBaars and S. Nakamura, *Nat. Photonics*, 2009, **3**, 180.
- 2 T. Suehiro, N. Hirotsuki and R.-J. Xie, *ACS Appl. Mater. Interfaces*, 2011, **3**, 811–816.
- 3 T.-C. Liu, B.-M. Cheng, S.-F. Hu and R.-S. Liu, *Chem. Mater.*, 2011, **23**, 3698–3705.
- 4 S. Nakamura, *MRS Bull.*, 2009, **34**, 101–107.
- 5 S. Ye, F. Xiao, Y. Pan, Y. Ma and Q. Zhang, *Mater. Sci. Eng., R*, 2010, **71**, 1–34.
- 6 Y.-C. Lin, M. Karlsson and M. Bettinelli, *Top. Curr. Chem.*, 2016, **374**(21), 1–47.
- 7 Z. Xia and A. Meijerink, *Chem. Soc. Rev.*, 2017, **46**, 275–299.
- 8 Z. Xia and Q. Liu, *Prog. Mater. Sci.*, 2016, **84**, 59–117.
- 9 J. L. Wu, G. Gundiah and A. Cheetham, *Chem. Phys. Lett.*, 2007, **441**, 250–254.
- 10 A. Kalaji, P. J. Saines, N. C. George and A. K. Cheetham, *Chem. Phys. Lett.*, 2013, **586**, 91–96.
- 11 P. Dorenbos, *J. Lumin.*, 2013, **134**, 310–318.
- 12 P. Dorenbos, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2002, **65**, 235110.
- 13 N. C. George, K. A. Denault and R. Seshadri, *Annu. Rev. Mater. Res.*, 2013, **43**, 481–501.
- 14 N. C. George, A. J. Pell, G. Dantelle, K. Page, A. Llobet, M. Balasubramanian, G. Pintacuda, B. F. Chmelka and R. Seshadri, *Chem. Mater.*, 2013, **25**, 3979–3995.
- 15 R. D. Shannon, *Acta Crystallogr., Sect. A: Cryst. Phys., Diffraction, Theor. Gen. Crystallogr.*, 1976, **32**, 751–767.
- 16 P. Ghigna, S. Pin, C. Ronda, A. Speghini, F. Piccinelli and M. Bettinelli, *Opt. Mater.*, 2011, **34**, 19–22.
- 17 F. Pan, M. Zhou, J. Zhang, X. Zhang, J. Wang, L. Huang, X. Kuang and M. Wu, *J. Mater. Chem. C*, 2016, **4**, 5671–5678.
- 18 S. J. Marin, M. O'Keeffe, V. G. Young and R. B. V. Dreele, *J. Solid State Chem.*, 1991, **91**, 173–175.
- 19 C. C. Lin and R.-S. Liu, *J. Phys. Chem. Lett.*, 2011, **2**, 1268–1277.
- 20 H. Daicho, T. Iwasaki, K. Enomoto, Y. Sasaki, Y. Maeno, Y. Shinomiya, S. Aoyagi, E. Nishibori, M. Sakata and H. Sawa, *et al.*, *Nat. Commun.*, 2012, **3**, 1132.
- 21 H. Zhu, C. C. Lin, W. Luo, S. Shu, Z. Liu, Y. Liu, J. Kong, E. Ma, Y. Cao and R.-S. Liu, *et al.*, *Nat. Commun.*, 2014, **5**, 4312.
- 22 V. Bachmann, C. Ronda and A. Meijerink, *Chem. Mater.*, 2009, **21**, 2077–2084.
- 23 P. F. Smet, A. B. Parmentier and D. Poelman, *J. Electrochem. Soc.*, 2011, **158**, R37–R54.
- 24 I. Berezovskaya, Z. Khapko, A. Voloshinovskii, N. Efyushina, S. Smola and V. Dotsenko, *J. Lumin.*, 2018, **195**, 24–30.
- 25 W. B. Im, N. George, J. Kurzman, S. Brinkley, A. Mikhailovsky, J. Hu, B. F. Chmelka, S. P. DenBaars and R. Seshadri, *Adv. Mater.*, 2011, **23**, 2300–2305.
- 26 A. Kaminska, A. Duzynska, M. Berkowski, S. Trushkin and A. Suchocki, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, **85**, 155111.
- 27 A. A. Setlur, W. J. Heward, Y. Gao, A. M. Srivastava, R. G. Chandran and M. V. Shankar, *Chem. Mater.*, 2006, **18**, 3314–3322.



- 28 M. R. Krames, O. B. Shchekin, R. Mueller-Mach, G. O. Mueller, L. Zhou, G. Harbers and M. G. Craford, *J. Disp. Technol.*, 2007, **3**, 160–175.
- 29 D. Robbins, *J. Electrochem. Soc.*, 1979, **126**, 1550–1555.
- 30 C.-C. Chiang, M.-S. Tsai and M.-H. Hon, *J. Electrochem. Soc.*, 2008, **155**, B517–B520.
- 31 Y.-C. Lin, P. Erhart, M. Bettinelli, N. C. George, S. F. Parker and M. Karlsson, *Chem. Mater.*, 2018, **30**, 1865–1877.
- 32 A. Setlur, *Electrochem. Soc. Interface*, 2009, **16**, 32.
- 33 S. Shionoya, W. M. Yen and H. Yamamoto, *Phosphor handbook*, CRC press, 2006.
- 34 B. Di Bartolo, *Advances in Nonradiative Processes in Solids*, Springer My Copy, UK, 2014.
- 35 P. Dorenbos, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, **85**, 165107.
- 36 P. Dorenbos, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2013, **87**, 035118.
- 37 P. Dorenbos, *J. Electrochem. Soc.*, 2005, **152**, H107–H110.
- 38 P. Dorenbos and E. G. Rogers, *ECS J. Solid State Sci. Technol.*, 2014, **3**, R150–R158.
- 39 E. Rogers and P. Dorenbos, *J. Lumin.*, 2014, **153**, 40–45.
- 40 E. Rogers and P. Dorenbos, *ECS J. Solid State Sci. Technol.*, 2014, **3**, R173–R184.
- 41 E. Rogers and P. Dorenbos, *J. Lumin.*, 2014, **155**, 135–140.
- 42 P. Dorenbos, *J. Lumin.*, 2013, **135**, 93–104.
- 43 P. Dorenbos, *ECS J. Solid State Sci. Technol.*, 2013, **2**, R3001–R3011.
- 44 H. Luo, L. Ning, Y. Dong, A. J. Bos and P. Dorenbos, *J. Phys. Chem. C*, 2016, **120**, 28743–28752.
- 45 L. Zhou, W. Zhou, F. Pan, R. Shi, L. Huang, H. Liang, P. A. Tanner, X. Du, Y. Huang and Y. Tao, *et al.*, *Chem. Mater.*, 2016, **28**, 2834–2843.
- 46 Y. Liu, X. Zhang, Z. Hao, Y. Luo, X. Wang and J. Zhang, *J. Mater. Chem.*, 2011, **21**, 16379–16384.
- 47 J. Ueda, S. Tanabe and T. Nakanishi, *J. Appl. Phys.*, 2011, **110**, 053102.
- 48 Y. H. Kim, P. Arunkumar, B. Y. Kim, S. Unithrattil, E. Kim, S.-H. Moon, J. Y. Hyun, K. H. Kim, D. Lee and J.-S. Lee, *et al.*, *Nat. Mater.*, 2017, **16**, 543.
- 49 J. Ueda, P. Dorenbos, A. J. Bos, A. Meijerink and S. Tanabe, *J. Phys. Chem. C*, 2015, **119**, 25003–25008.
- 50 F. Piccinelli, A. Speghini, G. Mariotto, L. Bovo and M. Bettinelli, *J. Rare Earths*, 2009, **27**, 555–559.
- 51 Y.-C. Lin, P. Erhart, M. Bettinelli and M. Karlsson, to be published.
- 52 Y. Shimomura, T. Honma, M. Shigeiwa, T. Akai, K. Okamoto and N. Kijima, *J. Electrochem. Soc.*, 2007, **154**, J35–J38.
- 53 Color rendering index was calculated using the Python colour science package called Colour developed by JetBrains PyCharm, <http://colour-science.org/>.
- 54 C.-C. Sun, Y.-Y. Chang, T.-H. Yang, T.-Y. Chung, C.-C. Chen, T.-X. Lee, D.-R. Li, C.-Y. Lu, Z.-Y. Ting and B. Glorieux, *et al.*, *J. Solid State Light.*, 2014, **1**, 19.
- 55 S. Abe, J. J. Joos, L. I. Martin, Z. Hens and P. F. Smet, *Light: Sci. Appl.*, 2017, **6**, e16271.

