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Introduction

The ever-growing researches on ceramic-based electromagnetic interference (EMI) shielding materials have been generally motivated by their specific applications as light-weight and highly-efficient structural EMI shielding components in modern spacecraft which are commonly exposed to thermally harsh environments.^{1–4} Ceramics exhibit the unique combination of excellent mechanical strength and thermal stability at high temperatures which usually cannot be achieved in traditional EMI shielding materials such as metallic or electrically conductive polymer composites. Besides, ceramics usually exhibit the merits of light weight, good corrosion resistance and high dielectric loss at high temperatures.^{4–7}

One of the biggest obstacles for common ceramics as effective structural EMI shielding materials is the relatively low shielding effectiveness (SE) which had been generally ascribed to their poor electrical conductivity.^{2,4,8} Similar to polymer composites,⁹⁻¹² the electrical conductivity and EMI SE of ceramics could be significantly enhanced by incorporating highly conductive nano-fillers such as nano-layered Ti_3SiC_2 ,^{7,13}

Boron carbide composites with highly aligned graphene nanoplatelets: light-weight and efficient electromagnetic interference shielding materials at high temperatures

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B₄C-based ceramic composites containing 0–2 vol% highly aligned graphene nanoplatelets (GNPs) are fabricated. The electromagnetic interference (EMI) shielding properties of the obtained composites are investigated at X-band (8.2–12.4 GHz) frequency range from room-temperature up to 800 °C. All composites exhibit outstanding EMI shielding properties with satisfactory frequency- and thermal-stability. The shielding effectiveness (SE) of GNP/B₄C composites increases monotonically with increasing GNP loading. Superior room-temperature SE close to 40 dB is achieved with only 2 vol% GNPs and high SE around 35 dB still persists at 800 °C. Considering their relatively low density, GNP/B₄C composites possess a high specific shielding effectiveness (SSE) of 16 dB cm³ g⁻¹ which is among the highest values in reported ceramic-based shielding composites. Especially, the GNP/B₄C composite with 2 vol% GNPs exhibits the highest SSE/t (SSE divided by thickness) values at temperatures above 200 °C for all reported shielding composites. The excellent shielding properties of GNP/B₄C composites arise mainly from the high electrical conductivity, high dielectric loss and the multiple reflections by the highly aligned and large-sized GNP layers.

carbon fiber (Cf),14,15 carbon nanotube (CNT)16,17 and graphene.¹⁻³ Specially, it is worth noting that composites containing two-dimensional graphene could achieve considerably high EMI SE at a much lower filler content compared to onedimensional or three-dimensional nano-fillers,18,19 and the high SE dominantly arose from the high absorption loss.¹⁹ Yuan et al. demonstrated that graphene nanoplatelets (GNPs) are more effective in enhancing the electrical conductivity and EMI SE of polymer composites by comparing the EMI shielding properties between single-wall carbon nanotube and graphene sheet/polyaniline composites.19 It can be speculated that the two-dimensional nature and high aspect ratio of GNPs played a significant role in enhancing the shielding of EMI and the total SE could be possibly further improved by optimizing the interfaces and microstructure of ceramic composites.^{11,20} Selfaligned reduced graphene oxide (rGO) in epoxy composites could give rises to a remarkable EMI SE of 38 dB.20 It was recently reported that the two-dimensional nano-layered MAX phase (a class of layered ternary transition-metal carbides and nitrides with a general formula of $M_{n+1}AX_n$, wherein M is an early transition metal, A is an A-group element, X is either C or N, and *n* varies from 1 to 3)²¹ Ti₃AlC₂ with higher aspect ratio grains and higher texture degree is more favourable for high EMI SE.²² Specially, a giant X-band (8.2-12.4 GHz) EMI SE around 92 dB was reported in well aligned MXene film which is

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another type of highly conductive two-dimensional materials exfoliated from MAX phases.²³

In addition to the total SE, the specific SE (SSE) defined as the ratio between the total SE and the density of composite is considered as another important criterion to evaluate the EMI shielding materials especially for aerospace applications.^{23,24} In an effort to reduce the weights of the shielding materials, foamstructured composites were produced with the aid of expanding agents.24-26 Although ceramic-based composites could achieve high EMI SE by incorporating highly conductive nanofillers, their SSE was generally lower especially than those of conductive polymer-based composites owning to their relatively higher density. Consequently, it is of vital importance to choose suitable ceramic matrix with much lower density. Meanwhile, the chosen ceramic matrix should exhibit satisfactory mechanical properties and thermal stability, especially at high temperatures.

Boron carbide (B_4C) is one of the most light-weight ceramics with a density of only 2.52 g cm^{-3} which is comparable to polymer matrix.27 In addition, the excellent hardness, mechanical strength and thermal stability render B4C ceramics promising candidates of high-temperature structural EMI shielding ceramic matrix.27 The semiconducting characteristic of B₄C could also largely benefit the EMI shielding in comparison to other insulating ceramics.28 Accordingly, it can be expected that B₄C will be one of the most suitable ceramic matrices for high-temperature EMI shielding. Our previous study has demonstrated that the room-temperature EMI SE of B₄C could be significantly enhanced by adding GNPs.²⁹ Nevertheless, the EMI shielding properties of B4C-based composites at elevated temperatures still needs to be evaluated and the underlying mechanisms for the high shielding properties needs further clarification.

In this study, B_4C was chosen as the ceramic matrix with 0– 2 vol% GNPs being uniformly dispersed. The electrical, dielectric and EMI shielding properties of the obtained GNP/ B_4C composites were systematically investigated in the temperature range from 25 to 800 °C at the frequency range of 8.2–12.4 GHz (X-band). The origins of the high EMI shielding properties were explored and related mechanisms were presented.

Experimental

Composites preparation

GNPs were purchased from Deyang Carbonene Technology Co.,Ltd. The GNPs are stacks of graphene sheets about 6–8 nm in thickness and 5–10 μ m in level dimensions. GNPs were first dispersed in isopropanol and sonicated for 1 h. B₄C powders (Grade HS, H. C. Starck GmbH, Germany) mixed with different amount of GNPs dispersion (0–2 vol%) and 10 vol% Ti₃AlC₂ powders (purity > 98%, Forsman, China) as sintering aid were homogenized by attrition milling in isopropanol media for 10 h. Afterwards, the slurry was dried and then sieved. The mixed powders were placed in a 50 mm diameter graphite die and hotpressed in vacuum under 1900 °C for 30 min. A uniaxial pressure of 30 MPa was applied during sintering. The bulk density of the sintered samples was measured by the Archimedes' method.

The microstructural characterizations of the samples were determined by scanning electron microscopy (SEM, JEOL JSM 6300) and transmission electron microscopy (TEM, JEM-2100). The crystal structure of the samples was studied by X-ray diffraction (XRD). The crystal structure refinement was carried out using the FullProf program based on a triphase system with structural models of TiB₂, B_4C and Al_2O_3 . The diffraction lines were modelled by a pseudo-Voigt function convoluted with axial divergence asymmetry function and the background was modelled by a linear interpolation between a set of fixed points. The following parameters were also refined for each phase: the scale factor, the zero point of detector, the unit-cell parameters, the atomic site coordinates, and the overall displacement parameter.

Evaluation of EMI properties

The room-temperature electrical conductivity was measured with four-probe (dc power supply Agilent E364×A) on samples with a dimension of $13.0 \times 2.5 \times 2.5$ mm³ using a tubular support of alumina as holder. For EMI SE characterization, specimens with dimensions of 22.86 \times 10.16 \times 1.50 mm³ were cut and polished. The magnitudes of complex scattering parameters (S-parameters) that correspond to reflection $(S_{11} \text{ or } S_{22})$ and transmission $(S_{21} \text{ or }$ S_{12} in the X-band frequency were determined through waveguide method with a vector network analyzer (Agilent N5230A). For accuracy of measurement, the device is carefully calibrated with Through-Reflect-Line (TRL) approach. The hightemperature measurement of SE was performed in a waveguide heated by an inner heater at a rate of 10°C min⁻¹. The temperature range is 25-800 °C with a temperature span of 100 °C, and each temperature spot was stabilized for 10 min in order to ensure the accuracy of measurements.

The total SE (SE $_{\rm T})$ can be applied and characterized as follows: 30

$$SE_{T} = 10 \log(P_{i}/P_{t}) = -10 \log|S_{12}|^{2}$$
(1)

in which P_i and P_t represent the incident power and transmitted power respectively, SE due to reflection (SE_R) and absorption loss (SE_A) were calculated from the *S*-parameters as follows,³⁰

$$SE_{R} = -10 \log(1 - |S_{11}|^{2})$$
(2)

$$SE_{A} = -10 \log(|S_{21}|^{2}/(1 - |S_{11}|^{2}))$$
(3)

Results and discussion

Fig. 1 shows the XRD profiles of the obtained GNP/B₄C composites with 0 and 2 vol% GNPs. Besides the main diffraction peaks corresponding to B_4C , TiB_2 and Al_2O_3 phases were also detected in all composites. The presence of TiB_2 and Al_2O_3 phases can be ascribed to the reactions between B_4C and Ti_3AlC_2 at high temperatures, as suggested in our previous study.³¹ The diffraction peak at ~26° reveals the presence of carbon, which is related to the residual carbon from reactions



Fig. 1 Typical XRD patterns and XRD refinements of GNP/B_4C composites with (a) 0 vol% and (b) 2 vol% GNPs.

between B4C and Ti3AlC2.32 The XRD refinements were performed in order to obtained the amount of different phases. The amorphous carbon the residual carbon from reactions between B4C and Ti3AlC2 was not considered during the refinement. According to the refinement results, the molar ratios of TiB₂, B₄C and Al₂O₃ phases are around 13%, 83% and 4% respectively. Table 1 summarizes the phase composition, density and electrical conductivity of various GNP/B4C composites. All composites exhibit similar density which is slightly higher than that of monolithic B₄C ceramics due to the presence of TiB₂. The existence of highly conductive TiB₂ phase also enhances the electrical conductivity of the composites. The composites without GNPs exhibit a moderate electrical conductivity of 27 S m⁻¹, which is higher than that of monolithic B₄C ceramics.²⁸ With the increase of GNPs loading, the electrical conductivity increases monotonically. A jump of electrical conductivity occurs at 1 vol% GNPs loading and a maximum value of 1850 S m⁻¹ was obtained in composites with 2 vol% GNPs. The enhancement of electrical conductivity is generally ascribed to the formation of numerous conducting paths by GNPs layers and TiB₂ particles.³²

Fig. 2 displays the fracture surfaces of the GNP/B₄C composites with different GNP loadings. All the B₄C composites show dense microstructures, demonstrating the effective promotion of densification of B₄C by adding Ti₃AlC₂. The B₄C exhibits a uniform size distribution with a fine grain size around 0.5 μ m. The presence of TiB₂ particles can be observed in all composites, as indicated in the red circles in Fig. 2(a). The

Table 1 Density and electrical conductivity of GNP/B ₄ C composites							
GNPs content (vol%)	Density (g cm ⁻³)	Electrical conductivity (S m^{-1})					
0.0	2.63	27 ± 5					
0.5	2.60	35 ± 4					
1.0	2.58	980 ± 12					
2.0	2.57	1850 ± 17					



Fig. 2 Fracture surfaces of GNP/B_4C composites with different GNP loading. (a) 0 vol%; (b) 0.5 vol%; (c-1) and (c-2) 1 vol%; (d-1) and (d-2) 2 vol%.

GNPs could be easily distinguished as the thin layers indicated by yellow arrows in Fig. 2(c-1) and (d-1). It can be seen that the GNPs were distributed uniformly in all GNP/B₄C composites. The GNPs are more uniformly distributed compared to our previous study due to the utilization of isopropanol as dispersion media.32 Correspondingly, the electrical conductivity becomes slightly higher as demonstrated in Table 1. The level dimension of GNPs layers varies from several micrometres to tens of micrometres, which is consistent with the raw GNPs layers. GNP layers with a level dimension over 20 µm were occasionally seen as indicated in Fig. 2(d-2). The level dimension of GNPs in this study is apparently larger than most studies in literature.1-3 Accordingly, the GNP layers show certain degree of parallel alignment under high applied pressure during sintering, as demonstrated in Fig. 2(c-1) and (d-1). The large level dimensions and texture-like alignment of GNPs are highly desirable and were supposed to induce high EMI shielding efficiency for electromagnetic waves that emanate through the thickness direction.^{20,33} Fig. 3(a) and (b) display the TEM images of the GNP/B₄C composites at lower and higher resolutions respectively. GNP was found to locate well along the grain boundaries. A simple estimate taking into account the thickness of the single graphene sheets ($\sim 2 \text{ nm}$) gives a proportion of \sim 30 layers of graphene in the observed GNP sheet. The relatively higher observed thickness is a consequence of two GNPs stacking closely together, as shown in Fig. 3(b).

Paper



Fig. 3 TEM images of B_4C/GNP composites with 2 vol% GNPs. (a) Lower magnification; (b) higher magnification.

Fig. 4(a) compares the total room-temperature EMI SE (SE_T) of GNP/B₄C composites with different GNPs content measured at the X-band frequency range. The SE_T of B₄C composite without GNPs shows a relatively high value (\sim 32 dB) with a tiny variation at the measured frequency range. The high SE_T could be mainly attributed to the semiconducting nature of B₄C matrix and the presence of highly conductive TiB₂ particles. The addition of 0.5 vol% GNPs gives rise to a slight enhancement of the SE_T to \sim 33 dB. With further increase of GNPs, SE_T increased remarkably and a maximum SE_T close to 40 dB was obtained in GNP/B₄C composite with only 2 vol% GNPs which is almost the lowest amount of nano-fillers required to achieve such a high EMI SE_T in reported ceramic- or even polymer-based composites.1-3,7,20 Correspondingly, 99.99% of the incident EMI energy could be effectively shielded at room-temperature for GNP/B4C composite with 2 vol% GNPs. The overall EMI SE_T could be



Fig. 4 (a)–(c) represent the frequency dependencies of SE_T, SE_R and SE_A respectively of GNP/B₄C composites at room temperature; (d) frequency dependency of SE_A/SE_R; (e) temperature dependency of SE_T for GNP/B₄C composites with different GNP loading.

further divided into three parts which correspond to reflection (SE_R) , absorption (SE_A) and the multiple reflections (SE_M) respectively.³⁴ The multiple reflection effect, nevertheless, is included in the absorption because the re-reflected waves could get absorbed or dissipated within the composite material.²³ Accordingly, the SE_T of GNP/B₄C composites could be expressed as follows for simplicity,

$$SE_{T} = SE_{R} + SE_{A}$$
(4)

The frequency dependences of the SE_R and SE_A of GNP/B₄C composites were exhibited in Fig. 4(b) and (c) respectively. The SE_R exhibits a significant enhancement when the GNPs loading is higher than 1 vol%, which is in consistency with the dependence of electrical conductivity on GNPs loading. The composite with 2 vol% GNPs shows high SE_R in the range of 11–15 dB. The high SE_R can be ascribed to the giant impedance mismatch between GNP/B₄C composites and the free air. In theory, for a plane wave radiation, the far field reflection loss (SE_R) is given by⁷

$$SE_R = 39.5 + 10 \log \sigma / (2\pi f \mu)$$
 (5)

where σ is the electrical conductivity, *f* is the frequency and μ is the permeability. The SE_R tends to decrease with increasing frequency according to the above equation which is well consistent with the results shown in Fig. 4(b). If we substitute the electrical conductivity (shown in Table 1), frequency (8.2-12.4 GHz) and permeability (assumed to be 1 for all composites) into eqn (5), the SE_R was calculated to be around 20 dB for composite with 2 vol% GNPs. The relatively lower experimentally obtained values could be ascribed to the secondary reflections of reflected radiation due to the presence of GNP layers. Compared to SE_R, all GNP/B₄C composites exhibit considerably higher absorption loss SEA. The composites with 0 and 0.5 vol% GNPs exhibit similar SE_A between 22 and 24 dB. The SE_A shows a step increase at 1 vol% GNPs loading and no further significant improvement was observed with further increasing GNP loading to 2 vol%. High SEA values between 25 and 28 dB were obtained in composite with 2 vol% GNPs. The SEA of all composites tend to increase with increasing frequency, which is attributed to the increased effective thickness of the composites with decreasing wavelength. The ratios between SE_A and SE_R were calculated and shown in Fig. 4(d). The SE_A/SE_R ratios exhibit high values between 1.5 and 3.5 indicating that the SE_A made a dominant contribution to the SE_T for all composites. The high-temperature EMI shielding of GNP/B4C composites was evaluated up to 800 °C and the temperature dependences of SE_T for different GNP/B₄C composites at 10 GHz were shown in Fig. 4(e). The EMI SE_T of all GNP/B₄C composites exhibit satisfactory thermal stability and they experienced only a slight decrease with increasing temperature. This is consistent with the results of graphene oxides (GO) composites for which the SE_T decreases with increasing temperature when the GO mass is below 8 vol%.3 The slight decrease of SE_T with increasing temperature can be

Table 2 Electromagnetic shielding performance of representative ceramic-based composites

Matrix	Filler and loading	Thickness (mm)	Frequency range	Temperature range	EMI SE (dB)		
					RT	HT	Ref.
SiO ₂	r-GO 20 wt%	1.5	X band	25–200 °C	~ 36	~ 37	3
SiO_2	MWCNT 10 wt%	2.5	X band	100–500 °C	${\sim}21$	${\sim}23$	17
SiO_2	C _f 20 wt%	2.5	X band	25–600 °C	${\sim}12$	${\sim}12$	15
SiO ₂	OMC^a 10 vol%	5.0	X band	25 °C	${\sim}40$		37
SiC	C _f 40 vol%	3.0	X band	25 °C	~ 31	_	14
SiC _f /SiC	Ti ₃ SiC ₂ 14.4 wt%	3.0	X band	25–600 °C	${\sim}20$	${\sim}45$	7
SiC _f /SiC	PyC^b 3.3 vol%	2.0	X band	25 °C	${\sim}26$		35
Al_2O_3	GNP 2 vol%	1.5	X band	25–400 °C	~ 23	~ 37	2
Al_2O_3	Ti ₃ SiC ₂ 25 vol%	1.0	Ku-band	25 °C	~ 32		13
BaTiO ₃	GNP 4 wt%	1.5	X band	25 °C	${\sim}42$		1
Ti ₃ AlC ₂	_	1.5	X band	25–800 °C	~ 34	~ 31	7
3Y-TZP	MWCNT 9 wt%	1.0	Ku band	25 °C	25-30		8
Si ₃ N ₄	PyC 11.7 vol%	2.8	X band	25 °C	$\sim \! 43$		36
B ₄ C	GNP 2 vol%	1.5	X band	25-800 °C	$\sim \! 40$	~ 36	

ascribed to the higher defect mobility and therefore the decreased dielectric loss at high temperatures. Although with a slight decrease, the GNP/B₄C composites still show high EMI SE_T even at temperature as high as 800 °C. Especially, the GNP/B₄C composites with 2 vol% GNPs exhibit high SE_T larger than 35 dB at 800 °C, demonstrating that GNP/B₄C composites are efficient high-temperature EMI shielding materials.

Table 2 compares the EMI SE of various advanced ceramic composites reported in literatures.1-3,8,13-17,22,35-37 It can be observed that considerably high EMI SE over 30 dB can be achieved in many ceramic composites by incorporating highly conductive nanofillers such as Cf, CNT, GNP and Ti₃AlC₂ particles. More importantly, their high EMI SE could be well maintained or even enhanced at high temperatures up to 600 °C. However, the high EMI SE of most ceramic composites was achieved either by high filler loading or large thickness. Moreover, the preparation of some ceramics such as SiC/SiC composites with high SE is costly and time-consuming compared to the preparation of GNP/B4C composite in this study. It seems that only Al₂O₃ composites with 2 vol% GNPs exhibit similar EMI shielding properties with B4C composites in this study. However, the density of Al₂O₃ is much higher than B₄C. Regarding the aerospace applications of EMI shielding materials, the specific EMI shielding effectiveness SSE (SE divided by density of composite) is more appropriate to judge the shielding performances of different shielding composites.23,24 Nevertheless, SSE alone is not a sufficient parameter for understanding overall effectiveness, as a higher SSE can simply be achieved at a larger thickness, which directly increases the weight of the final product. Therefore, a more realistic parameter is to divide SSE by the material thickness (SSE/t). Fig. 5 shows the temperature dependent SSE/t of various EMI shielding composites reported in literature and this study. It can be observed that metals exhibit both lower SSE/t values and lower application temperature; carbon/polymer composites show extremely high SSE/t values while lower application temperature

(generally lower than 300 °C); Traditional ceramic-based composites exhibit the highest application temperature (\geq 400 °C) while relatively lower SSE/*t* values which are comparable to metals such as copper. GNP/B₄C composite with 2 vol% GNPs in this study exhibits simultaneously high SSE/*t* values and application temperature. The high SSE/*t* around 100 dB cm² g⁻¹ were well maintained up to 800 °C. It can be concluded that GNP/B₄C composite with 2 vol% GNPs in this study exhibit the highest EMI shielding efficiency both at temperatures higher than 400 °C in all EMI shielding composites investigated so far, indicating that the application of B₄C/GNP composites as light-weight structural EMI shielding materials at high temperatures is highly promising.

Generally, the high shielding properties of GNP/B₄C composites can be attributed to the high reflection loss and absorption loss. The high reflection loss comes from the highly conductive networks formed by GNPs and TiB₂. The absorption of the radiation could arise from different energy dissipation processes as ohmic loss, dielectric loss and magnetic loss.³⁵ In GNP/B₄C composites magnetic loss can be generally neglected.



Fig. 5 Comparison of SSE/t for typical EMI shielding materials reported in literature and this study.

Paper

The ohmic loss is mainly due to the high electrical conductivity of GNP/B₄C composites. Compared to ohmic loss, the dielectric loss is expected to make a dominant contribution to the high attenuation of microwave in GNP/B₄C composites. Fig. 6(a) illustrates the real part of the complex permittivity of GNP/B₄C composites with different GNP loadings. Composite without GNPs exhibit relatively high permittivity between 200 and 400 in the X-band frequency range and it decreases gradually with increasing frequency, which is higher than the reported values in literatures. The high permittivity of B₄C composites mainly arises from the hopping electrons of B₄C and free electrons in TiB₂ phase. Besides, the defects produced during hightemperature sintering and reactions could also contribute to the high permittivity. With the increase of GNPs content, the real permittivity keeps decreasing to \sim 50 at 1 vol% GNPs. With further increasing GNPs content, no more decrease of permittivity was observed. The significantly decreased real permittivity can be ascribed to the weaken ability to storage electrical charges by the formation of highly conductive GNPs networks in B₄C composites with high GNPs loading. Being opposite to the real part of permittivity, the dielectric loss tan δ defined as the ratio between the imaginary and the real part of the complex permittivity of GNP/B₄C composites increases monotonically with increasing GNP content, as exhibited in Fig. 6(b). When the GNP content is below 1 vol%, the dielectric loss kept around 1. GNP/B₄C composite with 2 vol% GNPs exhibit a highly lossy characteristic with extremely high dielectric loss between 1 and 9. The high dielectric loss of GNP/B₄C composites can be ascribed to the high imaginary permittivity which was generally described as below,7

$$\varepsilon'' = \varepsilon''_{\text{Relax}} + \sigma/\omega\varepsilon_0 \tag{6}$$

where $\varepsilon_{\text{Relax}}''$ is the electron relaxation polarization, σ is the electrical conductivity, ε_0 is the dielectric constant in vacuum, and ω is the angular frequency. The first term in eqn (6) is related to the relaxation loss of microwave and arises mainly from interface polarization and defects produced during hightemperature sintering. The last term corresponds to the conductance loss resulted from the formation of conductive nanofiller networks acting as dissipating mobile charge carriers in the composites. The high conductivity gives rise to high imaginary permittivity and consequently high electromagnetic radiation dissipation. Besides the high dielectric loss, the high absorption loss of GNP/B4C composites partially arises from the presence of aligned two-dimensional nano-layered



Fig. 6 (a) The real permittivity and (b) the dielectric loss of various GNP/B_4C composites.

architecture.^{20,33} The parallel aligned GNPs layers could induce multiple reflections inside the composites. The multiple reflections prolong the EM wave propagation path and enhance the absorption accordingly. The multiple reflections are speculated to be responsible for the different EMI shielding efficiency in composites containing GNP and CNT. With increasing temperature, the dielectric loss decreased due to the lower internal friction during the rotation of the dipoles, thereby leads to a slight decrease of SE_T.

Conclusions

Light-weight B_4C composites with high EMI shielding properties were successfully obtained by incorporating a small amount of GNPs. The composites containing only 2 vol% GNPs exhibit a high room-temperature shielding effectiveness around 40 dB and a satisfactory shielding effectiveness over 35 dB at 800 °C in the whole X-band frequency range. A high SSE/*t* around 100 dB cm² g⁻¹ which was almost the highest value reported in ceramic-based shielding composites was achieved in GNP/B₄C composites containing 2 vol% GNPs. The outstanding shielding performance of GNP/B₄C composites was ascribed to: (1) high reflection loss and ohmic loss due to high electrical conductivity; (2) high dielectric loss resulting from the interfacial polarization and large amounts of defects; (3) multiple reflections by parallel-aligned GNP layers.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- 1 Y. Qing, Q. Wen, F. Luo, W. Zhou and D. Zhu, *J. Mater. Chem. C*, 2016, 4, 371–375.
- 2 Y. Qing, Q. Wen, F. Luo and W. Zhou, *J. Mater. Chem. C*, 2016, 4, 4853–4862.
- 3 B. Wen, M. Cao, M. Lu, W. Cao, H. Shi, J. Liu, X. Wang, H. Jin, X. Fang, W. Wang and J. Yuan, *Adv. Mater.*, 2014, 26, 3484– 3489.
- 4 X. Yin, L. Kong, L. Zhang, L. Cheng, N. Travitzky and P. Greil, *Int. Mater. Rev.*, 2014, **59**, 326–355.
- 5 B. Zhang, J. Li, J. Sun, S. Zhang, H. Zhai and Z. Du, *J. Eur. Ceram. Soc.*, 2002, **22**, 93–99.
- 6 X. Li, L. Zhang, X. Yin, L. Feng and Q. Li, *Scr. Mater.*, 2010, **63**, 657–660.

- 7 Y. Mu, W. Zhou, F. Wan, D. Ding, Y. Hu and F. Luo, *Composites, Part A*, 2015, 77, 195–203.
- 8 S. Shi and L. Ji, Nanotechnology, 2008, 19, 255707.
- 9 J.-M. Thomassin, C. Jérôme, T. Pardoen, C. Bailly, I. Huynen and C. Detrembleur, *Mater. Sci. Eng.*, *R*, 2013, 74, 211–232.
- 10 Y. Yang, M. C. Gupta, K. L. Dudley and R. W. Lawrence, *Nano Lett.*, 2005, **5**, 2131–2134.
- 11 D. Yan, H. Pang, B. Li, R. Vajtai, L. Xu, P. Ren, J. Wang and Z. Li, *Adv. Funct. Mater.*, 2015, **25**, 559–566.
- 12 M. H. Al-Saleh, W. H. Saadeh and U. Sundararaj, *Carbon*, 2013, **60**, 146–156.
- 13 S. Shi, L. Zhang and J. Li, J. Appl. Phys., 2008, 103, 124103.
- 14 L. Chen, X. Yin, X. Fan, M. Chen, X. Ma, L. Cheng and L. Zhang, *Carbon*, 2015, **95**, 10–19.
- 15 M. Cao, W. Song, Z. Hou, B. Wen and J. Yuan, *Carbon*, 2010, 48, 788–796.
- 16 Y. Qing, Y. Mu, Y. Zhou, F. Luo, D. Zhu and W. Zhou, *Eur. Ceram. Soc., J.*, 2014, **34**, 2229–2237.
- 17 B. Wen, M. Cao, Z. Hou, W. Song, L. Zhang, M. Lu, H. Jin, X. Fang, W. Wang and J. Yuan, *Carbon*, 2013, 65, 124–139.
- 18 S. H. Xie, Y. Y. Liu and J. Y. Li, *Appl. Phys. Lett.*, 2008, **92**, 243121.
- 19 B. Yuan, L. Yu, L. Sheng, K. An and X. Zhao, *J. Phys. D: Appl. Phys.*, 2012, **45**, 235108.
- 20 N. Yousefi, X. Sun, X. Lin, X. Shen, J. Jia, B. Zhang, B. Tang,
 M. Chan and J.-K. Kim, *Adv. Mater.*, 2014, 26, 5480–5487.
- 21 J. Wang and Y. Zhou, Annu. Rev. Mater. Res., 2009, **39**, 415-443.
- 22 Y. Tan, H. Luo, H. Zhang, X. Zhou and S. Peng, *Scr. Mater.*, 2017, **134**, 47–51.

- 23 F. Shahzad, M. Alhabeb, C. B. Hatter, B. Anasori, S. Man Hong, C. M. Koo and Y. Gogotsi, *Science*, 2016, **353**, 1137– 1140.
- 24 Z. Chen, C. Xu, C. Ma, W. Ren and H. Cheng, *Adv. Mater.*, 2013, **25**, 1296–1300.
- 25 B. Shen, Y. Li, D. Yi, W. Zhai, X. Wei and W. Zheng, *Carbon*, 2016, **102**, 154–160.
- 26 Y. Zhang, Y. Huang, T. Zhang, H. Chang, P. Xiao, H. Chen, Z. Huang and Y. Chen, *Adv. Mater.*, 2015, 27, 2049–2053.
- 27 F. Thévenot, J. Eur. Ceram. Soc., 1990, 6, 205-225.
- 28 T. L. Aselage, D. Emin and S. S. McCready, *Phys. Rev. B*, 2001, 64, 054302.
- 29 Y. Tan, H. Luo, H. Zhang, X. Zhou and S. Peng, *AIP Adv.*, 2016, **6**, 035208.
- 30 M. H. Al-Saleh and U. Sundararaj, *Carbon*, 2009, **47**, 1738–1746.
- 31 Y. Tan, H. Luo, H. Zhang, X. Zhou and S. Peng, *Ceram. Int.*, 2016, 42, 7347–7352.
- 32 Y. Tan, H. Luo, H. Zhang and S. Peng, *J. Eur. Ceram. Soc.*, 2016, **36**, 2679–2687.
- 33 B. Shen, W. Zhai and W. Zheng, *Adv. Funct. Mater.*, 2014, 24, 4542–4548.
- 34 D. D. L. Chung, Carbon, 2001, 39, 279-285.
- 35 D. Ding, Y. Shi, Z. Wu, W. Zhou, F. Luo and J. Chen, *Carbon*, 2013, **60**, 552–555.
- 36 X. Li, L. Zhang and X. Yin, *J. Eur. Ceram. Soc.*, 2013, **33**, 647–651.
- 37 J. Wang, C. Xiang, Q. Liu, Y. Pan and J. Guo, Adv. Funct. Mater., 2008, 18, 2995–3002.