

Cite this: *RSC Adv.*, 2018, **8**, 25881

## New insights into the surface plasmon resonance (SPR) driven photocatalytic H<sub>2</sub> production of Au–TiO<sub>2</sub>†

Jinlin Nie,<sup>ab</sup> Jenny Schneider,<sup>b</sup> Fabian Sieland,<sup>ID b</sup> Long Zhou,<sup>a</sup> Shuwei Xia<sup>ID \*a</sup> and Detlef W. Bahnemann<sup>ID \*bc</sup>

The Surface Plasmon Resonance (SPR) driven photocatalytic H<sub>2</sub> production upon visible light illumination ( $\geq 500$  nm) was investigated on gold-loaded TiO<sub>2</sub> (Au–TiO<sub>2</sub>). It has been clearly shown that the Au-SPR can directly lead to photocatalytic H<sub>2</sub> evolution under illumination ( $\geq 500$  nm). However, there are still some open issues about the underlying mechanism for the SPR-driven photocatalytic H<sub>2</sub> production, especially the explanation of the resonance energy transfer (RET) theory and the direct electron transfer (DET) theory. In this contribution, by means of the EPR and laser flash photolysis spectroscopy, we clearly showed the signals for different species formed by trapped electrons and holes in TiO<sub>2</sub> upon visible light illumination ( $\geq 500$  nm). However, the energy of the Au-SPR is insufficient to overcome the bandgap of TiO<sub>2</sub>. The signals of the trapped electrons and holes originate from two distinct processes, rather than the simple electron–hole pair excitation. Results obtained by Laser Flash Photolysis spectroscopy evidenced that, due to the Au-SPR effect, Au NPs can inject electrons to the conduction band of TiO<sub>2</sub> and the Au-SPR can also initiate e<sup>–</sup>/h<sup>+</sup> pair generation (interfacial charge transfer process) upon visible light illumination ( $\geq 500$  nm). Moreover, the Density Functional Theory (DFT) calculation provided direct evidence that, due to the Au-SPR, new impurity energy levels occurred, thus further theoretically elaborating the proposed mechanisms.

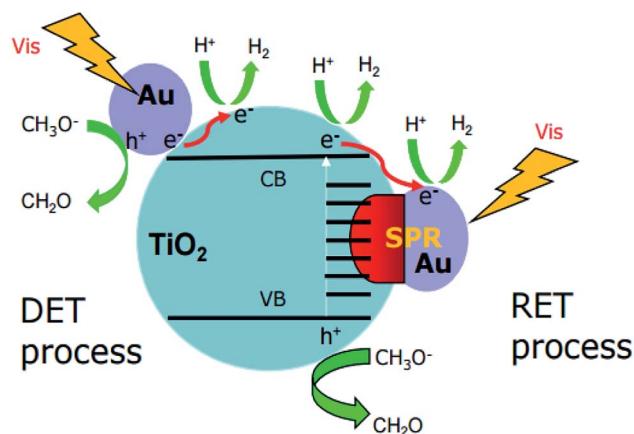
Received 25th June 2018  
Accepted 9th July 2018

DOI: 10.1039/c8ra05450a  
[rsc.li/rsc-advances](http://rsc.li/rsc-advances)

## 1. Introduction

Photocatalytic H<sub>2</sub> evolution from water utilizing sun light is gaining importance for providing clean and sustainable energy.<sup>1–3</sup> However, for practical applications, photocatalysts are desired to work efficiently under visible light which accounts for nearly half of the solar light.<sup>4</sup> Recently, the rapid advances in plasmonic photocatalysis have accelerated the progress in the visible-light harvesting ability of TiO<sub>2</sub>-based photocatalysts. For instance, it has been reported that the surface-loading of Au nanoparticles can greatly increase the photocatalytic H<sub>2</sub> production ability of Au–TiO<sub>2</sub> upon visible light illumination.<sup>4–9</sup>

However, the underlying mechanism remains unclear in the respective literatures.<sup>4a,10–14</sup> Generally, two main mechanisms have been proposed to explain the visible light activity of Au–TiO<sub>2</sub> photocatalysts (summarized in Fig. 1). The first mechanism known as resonance energy transfer (RET) is



<sup>a</sup>Key Laboratory of Marine Chemistry Theory and Technology, Ministry of Education, College of Chemistry and Chemical Engineering, Ocean University of China, Songling Road 238, 266100 Qingdao, China. E-mail: shuweixia@ouc.edu.cn

<sup>b</sup>Institut für Technische Chemie, Leibniz Universität Hannover, Callinstr. 3, D-30167 Hannover, Germany. E-mail: bahnemann@iftc.uni-hannover.de

<sup>c</sup>Laboratory "Photoactive Nanocomposite Materials", Saint-Petersburg State University, Ulyanovskaya Str. 1, Peterhof, 198504 Saint-Petersburg, Russia

† Electronic supplementary information (ESI) available: Spectra of the employed filters and EPR spectra of bare TiO<sub>2</sub> obtained under different conditions as well as the assignments of the EPR signals. XRD pattern of Au–TiO<sub>2</sub>. Side and top view of Au–TiO<sub>2</sub> model. Band structure of bulk anatase, (101) surface and partial density of states of TiO<sub>2</sub> (101) surface. See DOI: 10.1039/c8ra05450a

Fig. 1 Two main proposed mechanisms for H<sub>2</sub> production using Au–TiO<sub>2</sub> in water/methanol mixtures under visible light illumination: direct electron transfer (DET) process (left); resonance energy transfer (RET) process (right).



based on the idea that the SPR enhances the local electromagnetic field which in turn facilitates the generation of  $e^-/h^+$  pairs near the semiconductor surface.<sup>3,11,12</sup> D. B. Ingram<sup>11</sup> and Z. W. Seh<sup>12</sup> proposed this mechanism according to their results based on finite-difference time-domain (FDTD) simulations and discrete-dipole approximation (DDA) simulations, respectively, yet without exhibiting any experimental evidence. In addition, K. Awazu<sup>15</sup> also obtained results supporting this mechanism when adding an insulating layer between the noble metal nanoparticles and the semiconductor. In this study, the plasmonic metal nanoparticles were homogeneously covered with a  $\text{SiO}_2$  shell thus preventing the direct electron transfer process. However, due to the energy transfer the enhanced photocatalytic activity was still achieved. The second mechanism explains the visible activity of  $\text{Au}-\text{TiO}_2$  by the SPR-excited electron transfer from the Au nanoparticles to the conduction band of  $\text{TiO}_2$ , known as direct electron transfer (DET).<sup>4a,10,13,16</sup> C. Gomes Silva<sup>10</sup> and coworkers proposed this mechanism based on the observation that  $\text{Au}-\text{TiO}_2$  photocatalysts exhibit photocatalytic activity for  $\text{H}_2$  evolution upon 532 nm laser illumination. Recently, J. B. Priebe *et al.*<sup>4a</sup> provided experimental evidence for the DET mechanism based on results obtained employing electron paramagnetic resonance (EPR) spectroscopy. In the EPR measurements, the authors detected the signals for electrons trapped at oxygen vacancies and for  $\text{Ti}^{3+}$  ions, which was considered as evidence for the electron transfer process from Au to  $\text{TiO}_2$ . Herein, it was emphasized that the absorption of  $\text{TiO}_2$  is situated absolutely below 420 nm, therefore, the signals occurred due to the action of the Au nanoparticles rather than of  $\text{TiO}_2$ . Moreover, it was excluded that the signals originated from electrons located in the Au nanoparticles themselves, as much higher *g* values and line widths are the requisites for these electrons to yield such signals.<sup>4</sup> In conclusion, J. B. Priebe and co-worker argued that the obtained EPR-signals provide evidence for trapped electrons in  $\text{TiO}_2$  thus indicating a hot electron injection from the Au nanoparticles to  $\text{TiO}_2$ . However, it should be mentioned here that the results obtained by these EPR measurements can also be explained by means of the RET mechanism, since according to this mechanism the SPR-enhanced electromagnetic field can initiate electron-hole pair excitation close to the  $\text{TiO}_2$  surface and thus the formation of the trapped electrons at this surface. Besides, due to pre-existing defects in  $\text{TiO}_2$ ,<sup>11b,16</sup> such as oxygen vacancies,  $\text{TiO}_2$  exhibits also a small absorption shoulder above 400 nm indicative for surface defects. Actually, by means of EPR and laser flash photolysis spectroscopy, the obtained results clearly showed that bare  $\text{TiO}_2$  can be activated by the visible light illumination ( $\geq 420$  nm) or using the 420 nm laser beam, which has long been neglected in the previous studies.<sup>4a,11,12,17</sup>

Therefore, the origin of the photoinduced electrons and the underlying mechanism upon visible light illumination remains unclear. Herein, we directly focus on the Au plasmon band, that is employing a 500 nm cutoff filter ( $\geq 500$  nm) or using the 532 nm laser beam to totally get rid of the confusion caused by the direct excitation of  $\text{TiO}_2$ .

## 2. Experimental section

### Photocatalyst preparation

$\text{Au}-\text{TiO}_2$  photocatalysts consisting of 1 wt% gold on  $\text{TiO}_2$  (Evonik Aeroxide P25) were prepared by the typical sol immobilization (labeled as SIM) method,<sup>4</sup> which was conducted by adding an aqueous solution of poly vinyl alcohol (PVA) (1.2 mL, 1 wt% sol, Aldrich, >99%) to 1 mL  $\text{HAuCl}_4 \cdot 3\text{H}_2\text{O}$  (0.05 M). A freshly prepared  $\text{NaBH}_4$  solution (2.5 mL, 0.1 M, Aldrich, >96%) was added dropwise to the above solution, with the color turning dark. After 30 min, the  $\text{TiO}_2$  support (1.0 g) was added and the suspension was further stirred for 12 h at 25 °C. Then, the suspension was centrifuged, washed three times with distilled water and dried for 12 h at 70 °C. The as-prepared samples were used for catalytic and spectroscopic experiments without any other treatment.

### Material characterization

The UV-Vis absorption spectra were recorded employing a UV-visible spectrophotometer (Varian Cary 100 Bio). TEM images were characterized by Transmission Electron Microscopy (TEM) (FEI Tecnai G2 F20 TMP) using a 200 kV field emission gun (FEG). XRD patterns were monitored using a Bruker D8 Advance X-ray Diffraction system.

### Photocatalytic $\text{H}_2$ evolution test

The photocatalytic hydrogen production experiments were conducted under argon atmosphere in a double-wall quartz glass reactor, with a cooling system (Julabo) maintaining the temperature at 20 °C. The 1 g L<sup>-1</sup> photocatalyst suspensions were prepared by suspending Au (1 wt%)- $\text{TiO}_2$  in methanol/water mixtures (volume ratio 1 : 5). The suspension was sonicated for 30 min before being transferred into the reactor and flushed with argon for 15 min to remove gases dissolved in the suspension and to make sure that Ar remained in the headspace of the reactor. The argon gas rate was kept constant at 10 cm<sup>3</sup> min<sup>-1</sup> during the photocatalytic measurement and the inlet gas flow rate for QMS was 1 cm<sup>3</sup> min<sup>-1</sup>. Before illumination, the system was stabilized in the dark for 1 h to get a stable baseline. Afterwards, the suspension was illuminated employing an Osram XBO 1000 W Xe-Lamp (Müller) equipped with a 420 nm/500 nm cutoff filter.

### Electron paramagnetic resonance (EPR) spectroscopy

EPR measurements were carried out with a MiniScope MS400 spectrometer (Magnettech GmbH, Germany). A 500 W Mercury-Lamp (Müller) equipped with an optical cutoff filter (GC420/GC500) enables the illumination in the UV-Vis and visible light region. For low temperature measurements (90 K), the samples were placed in X-Band standard EPR tubes (Wilmad, 2 mm, O.D.) and cooled with a commercial EPR cold finger quartz Dewar by applying liquid nitrogen. The *g* values were calculated with the formula  $h\nu = g\beta B_0$  ( $B_0$  – external magnetic field,  $\beta$  – Bohr magneton,  $g_e$  – Landé *g*-factor) as



described in the literature.<sup>4a</sup> The frequency was at around 9.42 GHz, and the  $B_0$  was varied between 305–365 mT.

### Laser flash photolysis spectroscopy

During the laser flash photolysis measurements the samples were excited with an Nd-YAG laser (Brilliant B, Quantel) and the spectra were monitored by a laser flash photolysis spectrometer (LKS 80, Applied Photophysics). The excitation wavelength was tuned to 420 nm by employing an Optical Parametric Oscillator (OPO) (MagicPRISM, OPOTEK Inc.). Together with the OPO the 3rd harmonic (355 nm) of the Brilliant B laser was used to generate the visible light pulses with 6 mm diameter and 6 ns pulse length. The 532 nm laser was supplied by the 2nd harmonic (532 nm) of the Brilliant B laser. All experiments were performed in diffuse reflectance mode, where the analyzing light (supplied by a 150 W xenon arc lamp) was focused onto the sample and the reflected light was guided through the monochromator into the detector (Hamamatsu R928 photomultiplier). The light level detected by the oscilloscope was kept at 100 mV for all measurements to compensate the wavelength dependant sensitivity. Furthermore, all examinations were conducted under  $N_2$  atmosphere and 12 shots were averaged for every transient. The energy density of the 420 nm and 532 nm laser beam was around 5 mJ  $cm^{-2}$  per shot, respectively. The transient absorption change in reflectance  $\Delta J$  was calculated from the absorbance values according to the literature,<sup>18</sup> which also provides more detailed information about the set-up.

### Theoretical calculation details

All spin-polarized calculations of optimization and energy bands were performed by Cambridge Sequential Total Energy Package (CASTEP). Projected augmented wave (PAW) pseudo-potential was used to describe core electrons and plane wave basis set with 380 eV to represent valance electrons. Exchange-correlation energy was approximated by Perdew–Burke–Ernzerhof (PBE) functional of the generalized gradient approximation (GGA). Since the majority of P25 is anatase, we chose it for calculation. The model for anatase  $TiO_2$  dominant (101) surface was represented with a 12 monolayer unit cell with a 15 Å vacuum slab. Au– $TiO_2$  model was constructed with a two gold atom cluster, to get rid of multi isomers of Au clusters,<sup>19–21</sup> and  $TiO_2$  (101) surface, with bottom six atomic layers fixed and upper six layers relax (ESI Fig. S4†). A  $4 \times 4 \times 2$   $k$ -points girds was used to optimize anatase cell,  $2 \times 2 \times 1$   $k$ -points girds was used for  $TiO_2$  (101) surface and Au– $TiO_2$  optimization and energy bands calculation both. The optimized lattice values of anatase were 3.78 and 9.49 Å, close to experimental values (3.78 and 9.5 Å). To overcome the self-interaction error of GGA method, a Hubbard correction (DFT+U) of 8.3 eV was imposed on Ti d-states; a good qualitative description of the electronic was obtained (Fig. S5†) with 3.23 eV band gap for bulk and 3.43 eV for (101) surface. No correction was applied to gold.

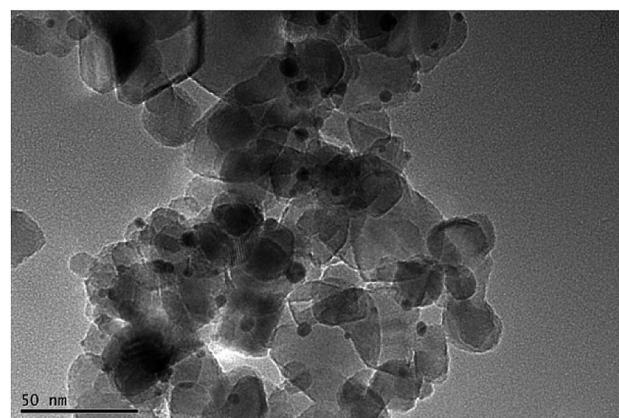


Fig. 2 TEM image of Au (1 wt%)- $TiO_2$ .

## 3 Results

### Material characterization

Fig. 2 shows the transmission electron microscopy (TEM) image of Au– $TiO_2$ . As can be seen, the lighter particles are the  $TiO_2$  matrix and the darker particles are the Au nanoparticles. The TEM images of Au– $TiO_2$  reveal that the Au nanoparticles are well deposited on the surface of  $TiO_2$  matrix exhibiting particle sizes between 5 nm and 10 nm. Besides, the contents of Au– $TiO_2$  photocatalyst were analyzed by means of X-Ray Diffraction (XRD). Accordingly, in addition to the peaks of anatase  $TiO_2$  (JCPDS # 01-070-7347) and rutile  $TiO_2$  (JCPDS # 03-065-5714), the XRD pattern exhibited peaks for Au (JCPDS # 03-065-2870), as shown in Fig. S1.† Due to the low concentration, the peaks of Au are weak. However, no peaks for other impurities (e.g. carbon or PVA) were detected within the detection limit.

In this study, the optical properties of Au nanoparticles, bare  $TiO_2$  and Au– $TiO_2$  were characterized by UV-Vis spectroscopy. As Fig. 3 shows Au-related samples exhibit plasmon band maxima located above 500 nm, namely Au nanoparticles with  $\lambda_{max}$  at

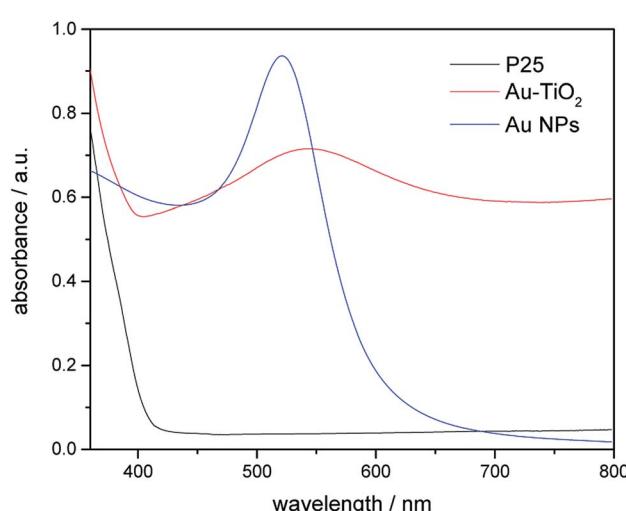


Fig. 3 Diffuse-reflectance UV-Vis spectra of bare  $TiO_2$  (black), Au NPs (blue) and Au– $TiO_2$  (red).



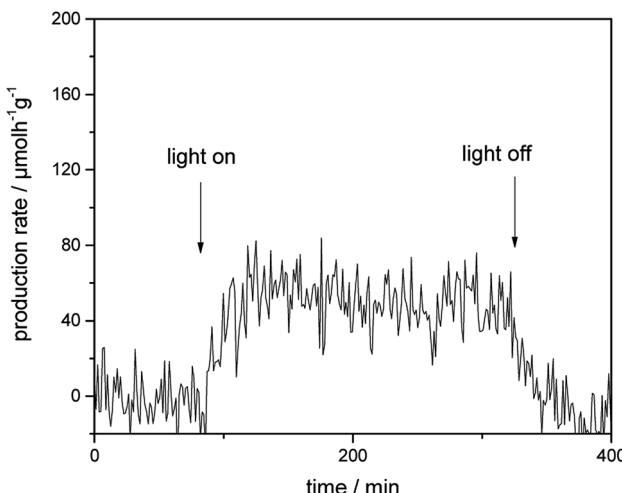


Fig. 4 Photocatalytic  $\text{H}_2$  evolution rate of  $\text{Au}-\text{TiO}_2$  obtained from  $\text{CH}_3\text{OH}/\text{H}_2\text{O}$  mixtures upon visible light illumination in the presence of a 500 nm cutoff filter.

520 nm,  $\text{Au}-\text{TiO}_2$  with  $\lambda_{\text{max}}$  at 543 nm, respectively. The absorption edge of bare  $\text{TiO}_2$  is located around 400 nm, which is in good agreement with previous studies.<sup>4b</sup> Spectra of bare  $\text{TiO}_2$  and the employed filters are presented in the ESI Fig. S2.†

### Photocatalytic $\text{H}_2$ evolution test

Fig. 4 presents the photocatalytic  $\text{H}_2$  evolution rates from  $\text{CH}_3\text{OH}/\text{H}_2\text{O}$  mixtures obtained employing Au-loaded  $\text{TiO}_2$  under visible-light illumination in the presence of a 500 nm cutoff filter. However, no formation of  $\text{H}_2$  was detected for bare  $\text{TiO}_2$  photocatalyst within the detection limit under the same reaction condition. In comparison to bare  $\text{TiO}_2$ ,  $\text{Au}-\text{TiO}_2$  exhibited  $\text{H}_2$  production ability upon visible light illumination in the presence of a 500 nm cutoff filter, indicating that Au-SPR can directly lead to photocatalytic  $\text{H}_2$  production.

### EPR analysis

The photoinduced EPR signals of trapped electrons and holes for bare  $\text{TiO}_2$  and  $\text{Au}-\text{TiO}_2$  were recorded in Ar at 90 K by (applying liquid nitrogen). The EPR parameters of the detected signals and their assignments were presented (ESI Fig. S3 and Table S1†), which are in good agreement with the previous literatures.<sup>4,22–24</sup>

In the presence of a 500 nm cutoff filter, when comparing the EPR signals of bare  $\text{Au}-\text{TiO}_2$  obtained in the dark with those observed upon visible light illumination ( $\geq 500$  nm), small but clear signal of trapped electrons (anatase  $\text{Ti}^{3+}$  at  $g = 1.990$ ) was observed (see Fig. 5), indicating the photo-induced electrons were trapped in  $\text{TiO}_2$ .

### Laser flash photolysis analysis

Fig. 6 shows the transient absorption signals measured at 400 nm and 680 nm for bare  $\text{TiO}_2$  and  $\text{Au}-\text{TiO}_2$  in the absence of any electron acceptor or donor. From the literature, it is known that  $\text{Ti}^{3+}$  species absorb in the wavelength range around

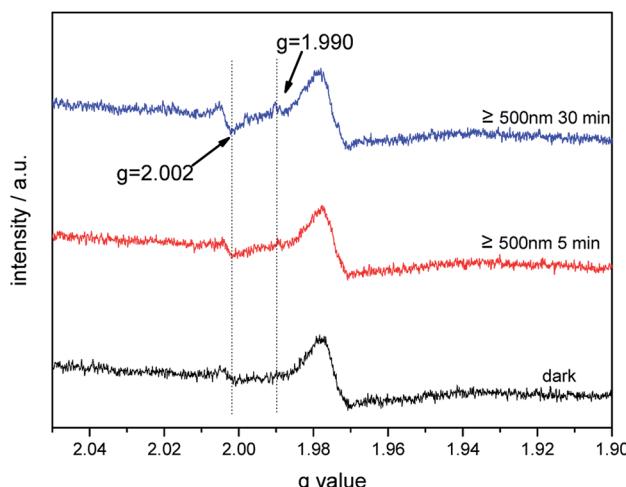


Fig. 5 EPR spectra of  $\text{Au}-\text{TiO}_2$  obtained in the dark and under visible light illumination ( $\geq 500$  nm) at 90 K.

680 nm,<sup>25a</sup> while the signal of the trapped holes can be monitored at 400–550 nm.<sup>25b</sup>

As shown in Fig. 6a, the transient signal of the trapped electron in  $\text{TiO}_2$  was detected indicating that bare  $\text{TiO}_2$  can be directly excited by 420 nm laser beam. The decay of the transient signals can be related to the internal electron/hole recombination.<sup>18,25</sup> Upon 532 nm laser excitation, which is the plasmon band of the Au NPs, in addition to the transient absorption signal of the trapped electrons at 680 nm, the signal of the trapped holes in  $\text{TiO}_2$  (see Fig. 6b) was monitored at 400 nm. As expected, no transient signals were observed for bare  $\text{TiO}_2$  in the blank experiment (see Fig. 6b inset). The trapped holes observed at 400 nm exhibit nearly the same transient signal intensity and kinetic like the trapped electrons detected at 680 nm.

### DFT calculation analysis

To further analyze the enhanced photocatalytic activities of  $\text{Au}-\text{TiO}_2$ , the electronic structure of  $\text{Au}-\text{TiO}_2$  system were calculated, showed in Fig. 7. As shown in Fig. 7a, energy levels of the system are apparently split due to the reduction of symmetry and the existence of surface dangling bonds. Moreover, new energy levels are appeared, two near the top of valence band and bottom of conduction band each, and four near Fermi level. Meanwhile, Fig. 7 presented clear evidence that new isolated impurity energy levels are mainly attributed to d orbitals of Au clusters, and the top of valence band by the hybrid orbitals of Au and  $\text{TiO}_2$  surface.

## 4. Discussion

Obviously,  $\text{Au}-\text{TiO}_2$  can produce  $\text{H}_2$  upon visible light illumination (employing a 500 nm cutoff filter). However, bare  $\text{TiO}_2$  absorbs no light in the wavelength range above 500 nm and no  $\text{H}_2$  can be formed ( $\geq 500$  nm) as evidenced by blank experiments. Besides, the XRD pattern showed no other impurities



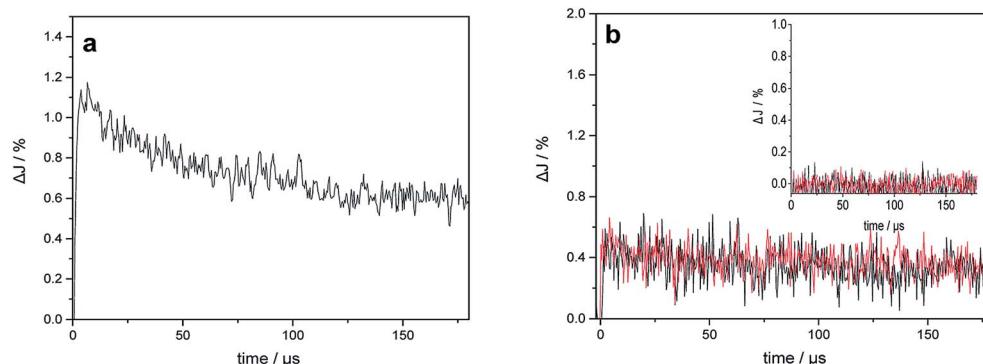


Fig. 6 (a) Transient absorption signals of bare TiO<sub>2</sub> observed at 680 nm upon 420 nm laser excitation; (b) transient absorption signals of bare TiO<sub>2</sub> (inset) and Au-TiO<sub>2</sub> observed at 680 nm (black) and 400 nm (red) upon 532 nm laser excitation in a N<sub>2</sub> atmosphere.

(e.g. carbon or PVA) were detected within the detection limit. Hence, it can be excluded that the visible-activity of Au-TiO<sub>2</sub> above 500 nm is caused by a carbon radical stem from polyvinyl alcohol remaining in the catalyst from the synthesis. The observation of any H<sub>2</sub> and signals of trapped electrons and holes can only be related to the Au-SPR effect. As shown in Fig. S3,† the appearance of new EPR signals can be observed upon illumination ( $\geq 500$  nm) which can be attributed to trapped electrons and trapped holes according to the literature.<sup>4,22–24</sup> The new formed EPR signals indicate the excitation of Au-TiO<sub>2</sub> photocatalyst upon visible light illumination above 500 nm. However, as new formed EPR signals are indistinctive (especially the signals at  $g = 2.002$ – $2.005$  (ref. 4 and 22–24)) laser excitation was employed to totally get rid of the confusion caused by the proximity of  $g$  values for signals of trapped electrons and trapped holes.

Upon 532 nm laser ( $\sim 2.34$  eV) excitation, the transient signals of the trapped electrons and holes in TiO<sub>2</sub> were also observed, which further confirmed the presence of photo-generated electrons and holes. It should be noted here that, the transient signals of the Au-SPR induced electrons can be detected in ps scale<sup>14</sup> and most charges recombine within 1 ns.<sup>26,27</sup> Therefore, the significant decay process of the transient signals in Fig. 6b was not observed within the detection time scale. In the absence of any electron acceptor or donor, transient signals of the trapped electrons and holes in TiO<sub>2</sub> can be observed on longer time scale ( $\mu$ s), as presented. The long-lived charges can be correlated with the photocatalytic activity of the photocatalysts.<sup>28,29</sup>

It should be emphasized here that the energy of the incoming photons ( $\lambda \geq 500$  nm,  $E_{hv} \leq 2.48$  eV) is insufficient to excite the electrons from the valence band to the conduction band of TiO<sub>2</sub> and leave the holes in the valence band of TiO<sub>2</sub>. However, due to the Au-SPR, many electrons in Au have higher energy than the conduction band of TiO<sub>2</sub>, facilitating the electron injection from Au to the conduction band of TiO<sub>2</sub>.<sup>13,30</sup> This provided a reasonable explanation for the observation of the Ti<sup>3+</sup> in both EPR spectroscopy and laser flash photolysis spectroscopy. As reported,<sup>13,30</sup> the direct electron injection process is thermodynamically favorable, since the photon energy ( $h\nu$ ) only needs to overcome the gap between the Fermi level ( $E_f$ ) of Au

and the top of TiO<sub>2</sub> conduction band ( $E_{CB}$ ), that is:  $h\nu \geq E_{CB0} - E_f$ . Moreover, after the contact of Au and TiO<sub>2</sub>, the Fermi level of the Au nanoparticles will be equalized to the work function and the conduction band of TiO<sub>2</sub> will be bent down.<sup>30</sup> Therefore, the trapped electrons leading to the formation of Ti<sup>3+</sup> are from the “hot electron injection” of the Au nanoparticles, process II (a) in Fig. 8.

On the other hand, since no holes can be transferred from the Au nanoparticles to the valence band of TiO<sub>2</sub>, the EPR signals of the trapped holes in TiO<sub>2</sub> (signal D) can only originate from the Au-SPR initiated electron–hole pair generation in TiO<sub>2</sub>. However, it should be noted here again that the energy of the visible light ( $\geq 500$  nm, *i.e.*,  $\leq 2.5$  eV) is not sufficient to overcome the band gap energy of TiO<sub>2</sub>. In this case, it is assumed that the Au-SPR initiates the Interfacial Charge Transfer (IFCT), as proposed by Creutz.<sup>31</sup> Accordingly, electrons in the valence band of TiO<sub>2</sub> directly transfer to the surface Au clusters, that is not *via* the excited state,<sup>32</sup> while the holes are trapped in the valence band of TiO<sub>2</sub>, process II (b) in Fig. 8. Besides, after electrons injection from Au to the conduction band of TiO<sub>2</sub>, monovalent of Au might be generated. Therefore, the DET process also provides the possibility of interfacial charge transfer from valence band of TiO<sub>2</sub> to monovalent of Au.

Moreover, the obtained theoretical calculation results provide further evidences supporting the above proposed mechanisms. Firstly, a good qualitative description of the electronic was obtained (Fig. S5†) with 3.23 eV band gap for bulk and 3.43 eV for (101) surface. As presented above (Fig. 7), with the deposition of Au clusters on the surface of TiO<sub>2</sub>, new isolated impurity energy levels appeared mainly due to d orbitals of Au clusters. Therefore, the possibility of electron transfer from Au to the conduction band of TiO<sub>2</sub> is demonstrated by DFT calculation. Besides, the existence of Au impurity energy levels in band gap also provide channels for electrons to be excited with lower energy ( $\lambda \geq 500$  nm,  $E_{hv} \leq 2.48$  eV). Moreover, the maxima absorption wavelength of Au-TiO<sub>2</sub> in Fig. 3 (543 nm, 2.288 eV) is in good accordance with the energy difference between an impurity level and the bottom of conduction band in DFT result (2.279 eV). Furthermore, due to the small energy gaps, the electron transfer from the valence band of TiO<sub>2</sub> to the surface loaded Au clusters is thermodynamically feasible, as

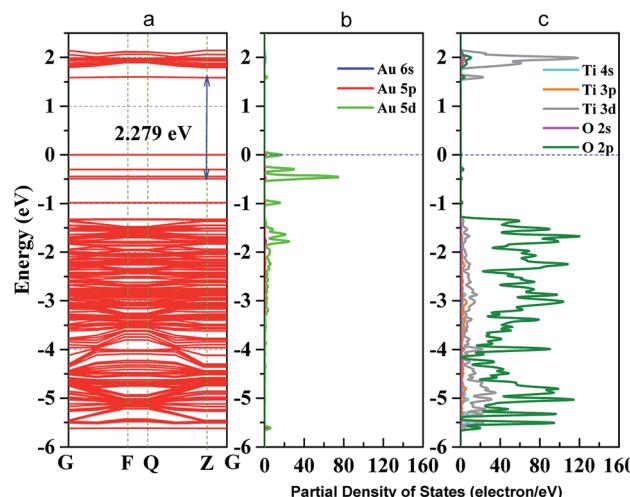


Fig. 7 (a) band structure of Au–TiO<sub>2</sub> slab; (b) partial density of states (PDOS) of Au cluster; (c) PDOS of TiO<sub>2</sub> (101) surface; blue dot line represents Fermi level.

proposed by IFCT process. Herein, the proposed mechanisms based on the experimental data are favored by the DFT calculation results. In case of the electron transfer processes can occur employing pure anatase as TiO<sub>2</sub> source, the same process should be easier in the case of using the mixture of anatase and rutile TiO<sub>2</sub>, for example, in Evonik Aerioxide P25 as TiO<sub>2</sub> source.

Besides, the EPR and the laser flash photolysis spectra (ESI Fig. S3† and 6a) show clear evidence that bare TiO<sub>2</sub> (P25) can be directly excited by visible light illumination ( $\lambda \geq 420$  nm) or 420 nm laser beam (process I in Fig. 8), which has long been ignored by some studies. Therefore, the investigation of SPR influence on the visible light photocatalytic activity of TiO<sub>2</sub> employing a 400/420 nm cutoff filter does not seem to be correct. Actually, upon visible light illumination ( $\lambda \geq 420$  nm), the direct excitation of TiO<sub>2</sub> and the SPR-induced excitation can both lead to the visible light photocatalytic activity of Au–TiO<sub>2</sub>, process I + II in Fig. 8.

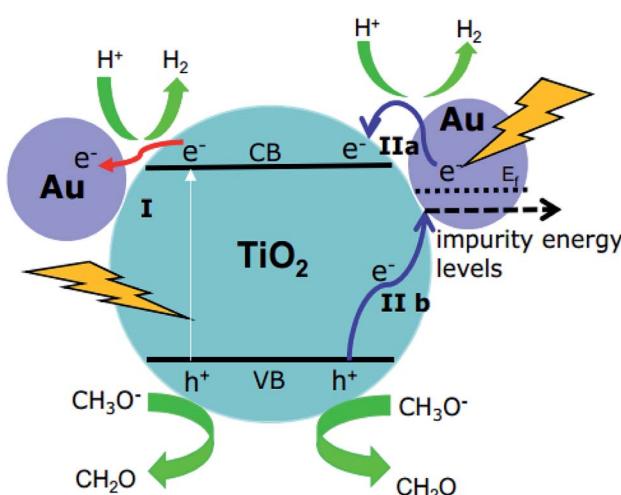


Fig. 8 Proposed mechanisms for H<sub>2</sub> production using Au–TiO<sub>2</sub> in water/methanol mixtures upon visible light illumination ( $\geq 420$  nm: process I + II;  $\geq 500$  nm: process II).

Therefore, upon illumination above 500 nm, due to the Au-SPR, electrons can be directly injected into the conduction band of TiO<sub>2</sub>, as reported previously.<sup>4a,10,13,14,26</sup> Simultaneously, the Au-SPR can also initiate electron hole pair generation in TiO<sub>2</sub>, while the electrons directly transfer from the valence band of TiO<sub>2</sub> to the surface Au-related species (IFCT process). The two pathways cooperatively contribute to the SPR-induced visible-driven photocatalytic H<sub>2</sub> production ability of Au–TiO<sub>2</sub> photocatalyst. This also provides a reasonable explanation for the photocatalytic activity when the plasmonic metal nanoparticles were homogeneously covered with an insulating SiO<sub>2</sub> shell.

## 5. Conclusions

In conclusion, Au-SPR driven photocatalytic H<sub>2</sub> production was observed upon visible light illumination ( $\geq 500$  nm). Direct experimental evidences were presented that Au nanoparticles can transfer the SPR-induced hot electrons into the conduction band of TiO<sub>2</sub> upon visible light illumination in the wavelength range above 500 nm. Simultaneously, the Au-SPR can also initiate electron–hole pair generation in TiO<sub>2</sub>. However, due to the photons energy/ $E_{\text{SPR}}$  is insufficient to overcome the band gap of TiO<sub>2</sub>, the initiated electrons directly transfer to the surface Au-related species through the IFCT process rather than *via* the excited state. The experimental results are in good agreement with the data obtained from DFT calculation. The DFT calculation analysis clearly shows how the d orbitals of Au clusters create impurity energy levels and decrease the band gap of Au–TiO<sub>2</sub>. We are optimistic that our contribution may contribute to the controversial discussion about the origin of the SPR-induced electron formation within Au–TiO<sub>2</sub> thus providing a new horizon for further investigations on exploring more effective visible light harvesting photocatalysts for solar energy conversion.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

Financial Support from the China Scholarship Council is gratefully acknowledged. The authors would like to thank the LNQE for TEM measurements. J. S. and D. B. kindly acknowledge financial support from the Federal Ministry of Education and Research (BMBF) for the project “DuaSol” (03SF0482C). F. S. and D. B. kindly acknowledge financial support from the Federal Ministry of Education and Research (BMBF) for the project “PureBau” (13N13350).

## References

- 1 J. A. Turner, *Science*, 1999, **285**, 687–689.
- 2 A. Kudo and Y. Miseki, *Chem. Soc. Rev.*, 2009, **38**, 253–278.
- 3 S. K. Cushing, J. Li, F. Meng, T. R. Senty, S. Suri, M. Zhi, M. Li, A. D. Bristow and N. Wu, *J. Am. Chem. Soc.*, 2012, **134**, 15033–15041.



4 (a) J. B. Priebe, M. Karnahl, H. Junge, M. Beller, D. Hollmann and A. Brückner, *Angew. Chem., Int. Ed.*, 2013, **52**, 11420–11424; (b) J. B. Priebe, J. Radnik, A. J. J. Lennox, M. Pohl, M. Karnahl, D. Hollmann, K. Grabow, U. Bentrup, H. Junge, M. Beller and A. Brückner, *ACS Catal.*, 2015, **5**, 2137–2148.

5 L. Liu, T. D. Dao, R. Kodiyath, Q. Kang, H. Abe, T. Nagao and J. Ye, *Adv. Funct. Mater.*, 2014, **24**, 7754–7762.

6 M. Murdoch, G. I. N. Waterhouse, M. A. Nadeem, J. B. Metson, M. A. Keane, R. F. Howe, J. Llorca and H. Idriss, *Nat. Chem.*, 2011, **3**(6), 489–492.

7 Y. C. Pu, G. M. Wang, K. D. Chang, Y. C. Ling, Y. K. Lin, B. C. Fitzmorris, C. M. Liu, X. H. Lu, Y. X. Tong, J. Z. Zhang, Y. J. Hsu and Y. Li, *Nano Lett.*, 2013, **13**, 3817–3823.

8 J. W. Hong, D. H. Wi, S.-U. Lee and S. W. Han, *J. Am. Chem. Soc.*, 2016, **138**, 15766–15773.

9 A. Stevanovic, M. Büttner, Z. Zhang and J. T. Yates, *J. Am. Chem. Soc.*, 2012, **134**(1), 324–332.

10 C. G. Silva, R. Juárez, T. Marino, R. Molinari and H. García, *J. Am. Chem. Soc.*, 2011, **133**, 595–602.

11 (a) D. B. Ingram and S. Linic, *J. Am. Chem. Soc.*, 2011, **133**, 5202–5205; (b) S. Linic, P. Christopher and D. B. Ingram, *Nat. Mater.*, 2011, **10**, 911–921.

12 Z. W. Seh, S. Liu, M. Low, S. Y. Zhang, Y. Liu, Z. Liu, A. Mlayah and M.-Y. Han, *Adv. Mater.*, 2012, **24**, 2310–2314.

13 K. Qian, B. C. Sweeny, A. C. Johnston-Peck, W. Niu, J. O. Graham, J. S. DuChene, J. Qiu, Y.-C. Wang, M. H. Engelhard, D. Su, E. A. Stach and W. D. Wei, *J. Am. Chem. Soc.*, 2014, **136**, 9842–9845.

14 Z. Bian, T. Tachikawa, P. Zhang, M. Fujitsuka and T. Majima, *J. Am. Chem. Soc.*, 2014, **136**, 458–465.

15 K. Awazu, M. Fujimaki, C. Rockstuhl, J. Tominaga, H. Murakami, Y. Ohki, N. Yoshida and T. Watanabe, *J. Am. Chem. Soc.*, 2008, **130**, 1676–1680.

16 Z. Liu, W. Hou, P. Pavaskar, M. Aykol and S. B. Cronin, *Nano Lett.*, 2011, **11**(3), 1111–1116.

17 S. Zhu, S. Liang, Q. Gu, L. Xie, J. Wang, Z. Ding and P. Liu, *Appl. Catal., B*, 2012, **119–120**, 146–155.

18 J. Schneider, K. Nikitin, M. Wark, D. W. Bahnemann and R. Marschall, *Phys. Chem. Chem. Phys.*, 2016, **18**, 10719–10726.

19 K. Z. Milowska and J. K. Stolarczyk, *Phys. Chem. Chem. Phys.*, 2016, **18**(18), 12716–12724.

20 L. A. Mancera and D. M. Benoit, *Phys. Chem. Chem. Phys.*, 2016, **18**(1), 529–549.

21 L. A. Mancera and D. M. Benoit, *Comput. Theor. Chem.*, 2015, **1067**, 24–32.

22 D. C. Hurum, A. G. Agrios, K. A. Gray, T. Rajh and M. C. Thurnauer, *J. Phys. Chem. B*, 2003, **107**, 4545–4549.

23 (a) T. Hirakawa, H. Kominami, B. Ohtani and Y. Nosaka, *J. Phys. Chem. B*, 2001, **105**, 6993–6999; (b) C. P. Kumar, N. O. Gopal, T. C. Wang, M.-S. Wong and S. C. Ke, *J. Phys. Chem. B*, 2006, **110**, 5223–5229.

24 V. Jovic, K. E. Smith, H. Idriss and G. I. N. Waterhouse, *ChemSusChem*, 2015, **8**, 2551–2559.

25 (a) D. W. Bahnemann, M. Hilgendorff and R. Memming, *J. Phys. Chem. B*, 1997, **101**, 4265–4275; (b) A. O. T. Patrocínio, J. Schneider, M. D. França, L. M. Santos, B. P. Caixeta, A. E. H. Machado and D. W. Bahnemann, *RSC Adv.*, 2015, **5**, 70536–70545.

26 A. Furube, L. Du, K. Hara, R. Katoh and M. Tachiya, *J. Am. Chem. Soc.*, 2007, **129**, 14852–14853.

27 L. Du, A. Furube, K. Yamamoto, K. Hara, R. Katoh and M. Tachiya, *J. Phys. Chem. C*, 2009, **113**(16), 6454–6462.

28 J. S. DuChene, B. C. Sweeny, A. C. Johnston-Peck, D. Su, E. A. Stach and W. D. Wei, *Angew. Chem., Int. Ed.*, 2014, **53**, 7887–7891.

29 (a) Z. Bian, T. Tachikawa, T. Majima and J. Phys. Chem. Lett., 2012, **3**, 1422–1427; (b) Z. Bian, T. Tachikawa, W. Kim, W. Choi and T. Majima, *J. Phys. Chem. C*, 2012, **116**, 25444–25453.

30 X. Zhang, Y. L. Chen, R. Liu and D. P. Tsai, *Plasmonic Photocatalysis*, *Rep. Prog. Phys.*, 2013, **76**, 046401–046441.

31 (a) C. Creutz, B. S. Brunschwig and N. Sutin, *J. Phys. Chem. B*, 2005, **109**, 10251–10260; (b) C. Creutz, B. S. Brunschwig and N. Sutin, *J. Phys. Chem. B*, 2006, **110**, 25181–25190.

32 (a) H. Irie, K. Kamiya, T. Shibanuma, S. Miura, D. A. Tryk, T. Yokoyama and K. Hashimoto, *J. Phys. Chem. C*, 2009, **113**, 10761–10766; (b) H. Yu, H. Irie, Y. Shimodaira, Y. Hosogi, Y. Kuroda, M. Miyauchi and K. Hashimoto, *J. Phys. Chem. C*, 2010, **114**, 16481–16487.

