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Epitaxial growth and interfacial property of monolayer MoS₂ on gallium nitride

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Two-dimensional (2D) transition-metal dichalcogenides (TMDCs) on semiconductor substrates are important for next-generation electronics and optoelectronics. In this study, we demonstrate the growth of monolayer MoS₂ on a lattice-matched gallium nitride (GaN) semiconductor substrate by chemical vapor deposition (CVD). The monolayer MoS₂ triangles exhibit optical properties similar to that of typical single-crystal MoS₂ sheets, as verified by the Raman, photoluminescence, and morphological characterizations. The Raman and PL features and their intensity mappings suggest that the as-grown MoS₂ on GaN substrate can achieve high quality and uniformity, demonstrating that GaN substrate is favorable for 2D MoS₂ growth. Moreover, the interfacial property and stacking structure were investigated by first-principles density functional theory (DFT) calculations to confirm the interlayer interactions of monolayer MoS₂ on GaN. Accordingly, the ability to grow high quality monolayer MoS₂ on semiconductor GaN substrate would open a new route toward the synthesis of hetero and composite structures for promising electronic and optoelectronic device applications.

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1. Introduction

The most investigated member of the two-dimensional (2D) transition-metal dichalcogenide (TMD) materials, monolayer molybdenum disulfide (MoS₂), has attracted widespread attention for a variety of next-generation electrical and optoelectronic devices applications due to its atomically sharp interfaces, ultrathin dimensions, flexibility and large optical effects.^{1,2} The proof-of-concept ultrafast field-effect transistors (FETs), optical devices, and flexible electronics^{3–5} based on 2D MoS₂ have been investigated. In general, at the beginning, the 2D MoS₂ applied on these devices was obtained by exfoliation of geological MoS₂ crystals and mechanically transferred to the desired substrates. However, small specimen sizes, uncertainty in placement location, and lack of composition control are the main drawbacks for the applications of geological MoS₂.⁶ In order to overcome these limitations, many methods have been used to grow the 2D materials with engineered composition. 2D MoS₂ has been synthesized by chemical vapor deposition (CVD),⁷ metal-organic chemical vapor deposition (MOCVD),⁸ and pulsed laser deposition (PLD).⁹ In most of these reports, monolayer MoS₂ has been primarily grown on SiO₂ substrate, which resulted in the random orientation of MoS₂ domains and grain boundaries

because of the amorphous nature of the substrate and its relatively high surface roughness.¹⁰

Moreover, the growth and fabrication of thin-film devices on semiconductor wafers constitute the bulk of the semiconductor industry.¹¹ Thus, the growth of high-quality 2D materials on semiconductor substrates would have technological advantages and broad applications. Considerable progress has been made on the growth of 2D MoS₂,¹² MoSe₂,¹³ WS₂,¹⁴ WSe₂,¹⁵ and black phosphorene¹⁶ on insulating substrates of SiO₂ and other compounds. Dumcenco *et al.* reported the successful growth of single-domain and large monolayer MoS₂ films on c-sapphire substrates.¹⁷ Growth of monolayer MoS₂ on an inert and nearly lattice-matching mica substrate by using a low-pressure chemical vapor deposition method was successfully demonstrated.¹⁸ Although these growth techniques are well established, the resulting 2D materials still need to be transferred to some unique semiconducting substrate surfaces to be integrated into the microelectronic or optoelectronic technologies for special functions. In fact, transfer processes¹⁹ would consume additional time and cause severe interfacial defects. Hence, the direct growth of 2D materials on semiconducting wafers or films can overcome these problems and result in ideal quality films. Many attempts have been made to grow 2D materials directly on semiconducting substrates, including Si,²⁰ GaAs,²¹ SiC,²² and others,²³ and preliminary hetero-junction devices have been fabricated. However, due to the unfavorable lattice mismatches, the domain orientation and size uniformity are not yet satisfactory. In contrast, the 2D material growth on more suitable semiconductor substrates may potentially offer

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better crystal quality. For example, the growth of MoS₂ on *c*-plane GaN substrate offers better crystal quality because the hexagonally symmetric GaN surface matches well with MoS₂. The *c*-plane lattice constants of GaN and MoS₂ are $a = 3.18 \text{ \AA}$ and 3.16 \AA , respectively, corresponding to a small mismatch of 2%. Moreover, the commonly used Ga-polar GaN substrate is uniformly terminated with Ga²⁺ cations, which facilitate the adsorption of S anions and the growth of smooth S–Mo–S layers. Therefore, it is particularly interesting to grow 2D MoS₂ on GaN substrate due to their close lattice match and favorable electronic band structure alignment for electronic and optoelectronic device requirements.

In this study, monolayer MoS₂ has been grown on a lattice-matched GaN semiconductor substrate by CVD method. Accordingly, triangle single-layer MoS₂ domains were commonly observed on the GaN substrate, and the excitation laser power dependent PL and Raman spectra were characterized to investigate the optical property of monolayer MoS₂ on GaN. Moreover, uniform color contrast of the PL and Raman intensity mappings suggest that the as-grown triangular monolayer MoS₂ on GaN substrate are of high quality and uniformity. Finally, first-principles density functional theory (DFT) calculations were employed to investigate the interfacial property and interactions of monolayer MoS₂ on semiconductor GaN.

2. Methods

Growth procedure

Monolayer MoS₂ was grown by CVD on the (0001) plane of GaN substrate. The substrate was cleaned by an ultrasonic cleaner with acetone, absolute ethyl alcohol, and deionized water. Following this, they were placed face-down on a crucible containing ~0.1 g of MoO₃ (99.9%) and loaded into a CVD heating furnace (downstream) with a 60 mm outer diameter quartz tube. CVD growth was performed at atmospheric pressure using ultra-high-purity N₂ as the carrier gas. A second crucible containing 1 g of sulfur powder (99.9% purity) was located in another heating zone, which was upstream from the growth substrate. Pure N₂ gas was first introduced into the CVD furnace with a flow rate of 100 sccm for 20 min to purge the tube and then maintained at a flow rate of 50 sccm during the entire growth process. During the heating process, the downstream heating zone was initially heated to 800 °C in 120 min, while the upstream heating zone was first maintained at room temperature for 70 min and then heated to 200 °C in 50 min, ensuring that the sulfur powder was sufficiently evaporated when the MoO₃ powder was heated to 800 °C. Consequently, the growth temperature was maintained at 800 °C for 15 min, and then naturally cooled down to room temperature in 2 h.

Characterization

The surface morphologies and thicknesses of the as-grown MoS₂ on GaN substrate were characterized by optical microscope, field emission scanning electron microscope (FE-SEM: JEOL JSM-7000F) and atomic force microscope (AFM: Bruker

Dimension ICON). The Raman and photoluminescence (PL) characterizations were performed using a Renishaw LabRAM Inviacmicro-Raman system using a 532 nm excitation laser at room temperature, with laser spot size set at about 1 μm using a 50× objective.

Density functional theory modelling

First-principles calculations based on density functional theory (DFT) were employed to study the interfacial and electronic structure properties of monolayer MoS₂ on GaN (0001) plane. ATOMISTIX TOOLKIT (ATK) code was used in implementing the calculations.²⁴ It is well known that the band gaps calculated at generalized gradient approximation (GGA) level are usually underestimated compared with the experimental values for some materials. To ensure the accuracy of the simulations, we optimized the structures of the MoS₂–GaN system by using Perdew–Burke–Ernzerhof (PBE) of GGA²⁵ for the exchange–correlation functional together with the SG15 pseudopotentials. The structure was optimized using the limited-memory Broyden–Fletcher–Goldfarb–Shanno (LBFGS) scheme till the forces on individual atoms were mitigated to less than 0.05 eV \AA^{-1} and the stress error tolerance was less than 0.1 GPa. To take the van der Waals (vdW) interaction into account, the Grimme's DFT-D2 dispersion-correction approach was adopted for all exchange correlation energies due to its higher accuracy and less empiricism.²⁶ A well-conserved $9 \times 9 \times 1$ *k*-point mesh was used for the structure optimization and electronic structure calculation with a density mesh cutoff energy of 75 Hartree. The buffer space of a vacuum of 20 Å was set to eliminate the coupling between two adjacent units. In the calculation, six atomic layers of GaN were chosen for the interface simulations since the obtained results do not change appreciably beyond this thickness, where the method has been widely used previously.

3. Results and discussion

Fig. 1(a) shows the optical microscopy image of MoS₂ grown on GaN substrate. As the GaN substrate is grown by MOCVD, the dislocation density on the substrate surface is as low as 10^{-7} cm^{-3} ,²⁷ which can be neglected between the interface of GaN and as-grown MoS₂ domain. It can be clearly seen that some of the as-grown MoS₂ domains are triangular with uniform color contrast, where the lateral sizes have a certain limit due to the high energy caused by the lattice mismatch between the MoS₂ and the substrates. To confirm the thickness of the obtained MoS₂ films, a typical AFM image shown in Fig. 1(b) presents the edge of a triangular MoS₂ domain on the GaN substrate. The height profile taken from the solid line in Fig. 1(b) reveals a film thickness of 0.8 nm, which is consistent with the range of a single-layered MoS₂ film on bare substrates (other than GaN substrate, 0.6–0.9 nm) grown by CVD,^{28,29} demonstrating that the monolayer MoS₂ has been successfully obtained by the growth procedure on GaN substrate. Fig. 1(d) presents the SEM images of the as-grown monolayer MoS₂ on GaN, which indicate that most of the MoS₂ domains are monolayer single crystal, with large coverage ratio on the substrate. Additionally, it can be



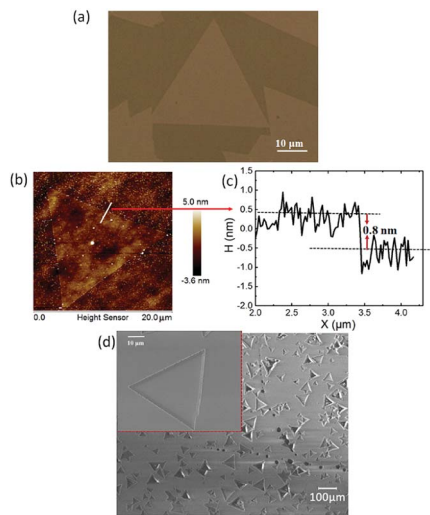


Fig. 1 (a) Optical microscopy image, (b) AFM image, (c) the corresponding height profile, and (d) SEM images of the as-grown MoS₂ on GaN substrate.

observed from the inset of the enlarged SEM image in Fig. 1(d) that the lateral size of the largest single triangular MoS₂ domain reaches as high as 50 μm on the GaN substrate.

Fig. 2(a) shows the Raman spectra of triangular MoS₂ domain at room temperature. Typical peak positions of in-plane E_{2g}¹ (~383 cm⁻¹) and out-of-plane A_{1g} (~402 cm⁻¹) Raman modes of the vibrations in MoS₂ can be clearly observed. The wave number difference between the E_{2g}¹ and A_{1g} peaks estimated from the Raman spectra is less than 19 cm⁻¹, identifying that the thickness of the as-grown MoS₂ nanosheet is single layer.³⁰ The behaviors as a function of excitation laser power (with the maximum excitation laser power denoted as P₀) imply that the E_{2g}¹ vibration changes slightly, while the A_{1g} vibration softens (red shifts) from 385.7 cm⁻¹ to 403.25 cm⁻¹ with the increase in the excitation power. The powers of the excitation laser lines are maintained well below 1 mW to avoid heating effects. The red-shift suggests that van der Waals force plays a minor role, while stacking-induced structure changes and long-range coulombic interlayer interactions may dominate the change in atomic vibrations within MoS₂ on GaN substrate.³⁰ Fig. 2(b) depicts the room-temperature PL spectra of monolayer

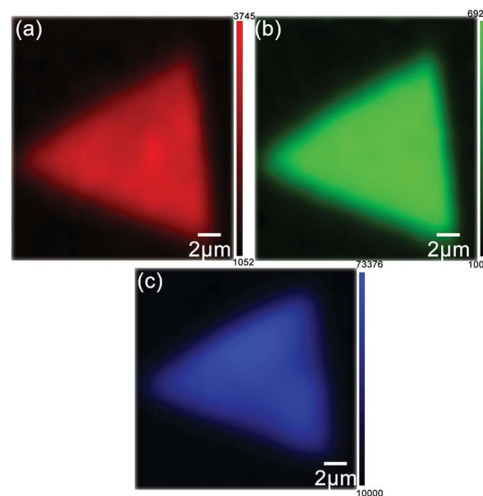


Fig. 3 Raman peak intensity mapping at 383 cm⁻¹ (a) and 402 cm⁻¹ (b), and PL peak intensity mapping at 674 nm (c) of monolayer MoS₂ on GaN substrate.

MoS₂ grown on GaN as a function of excitation laser power. It is notable that the emission peak excited by the lowest laser power is located at ~674 nm, further demonstrating the single-layer feature of the as-grown MoS₂ on GaN.¹⁸ Additionally, only strikingly strong PL of A peak exciton, with no shoulder peak of B peak exciton is observed from Fig. 2(b). Such a single PL peak feature has been reported in exfoliated monolayer MoS₂ suspended on holed Si substrate.³¹ Hence, the observation of a single (A excitonic emission) peak is a strong evidence that the monolayer MoS₂ grown herein on GaN substrate is undoubtedly of high quality.¹⁸ Moreover, it is observed that the PL peaks of MoS₂ on GaN substrate exhibit slight shift as the excitation laser intensity increases, indicating that possible heating effect caused by increasing laser power can be excluded. It can also be derived that the full width at half-maximum (FWHM) of the PL spectra for MoS₂ on GaN exhibits no apparent broadening, demonstrating that GaN substrate is favorable for 2D MoS₂ growth due to its lattice-match feature and the appropriate thermal conductivity of the GaN substrate.

In order to further investigate the uniformity of monolayer triangular MoS₂ domains on the GaN substrate, Fig. 3(a) and (b) present the Raman intensity mappings at 383 cm⁻¹ and

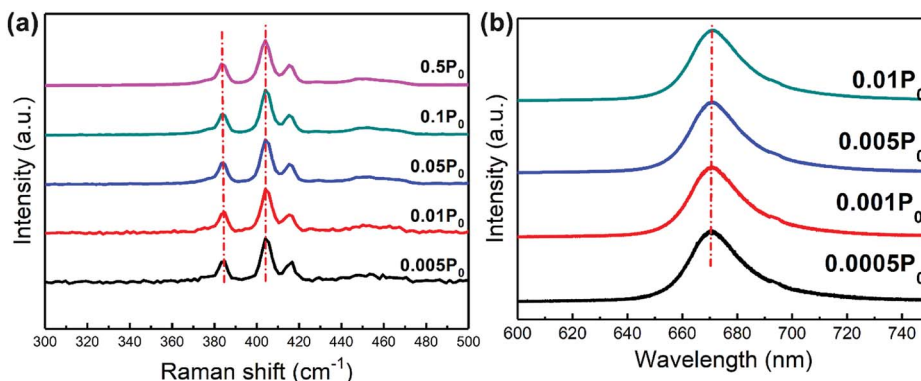


Fig. 2 Normalized Raman spectra (a) and PL spectra (b) of monolayer MoS₂ on GaN as a function of excitation laser power.



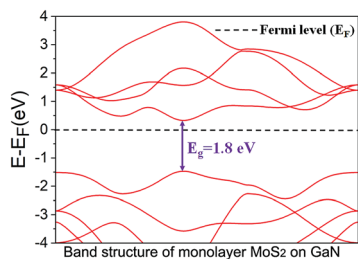


Fig. 4 The band structure of bare monolayer MoS₂ calculated in a primitive unit cell.

402 cm⁻¹, respectively, from which uniform color distribution and strong intensity can be observed in all the positions of the single triangular domain. Furthermore, the PL intensity mappings of MoS₂ on GaN at a wavelength of 674 nm (shown in Fig. 3(c)) also exhibits very uniform color contrast across the entire surface, with a small discrepancy at the very edge of the triangle domain, which would have been caused by the partial absorption of the laser spot energy. Such a uniformity of the PL and Raman intensity mappings over the entire surface further suggests that the as-grown triangular monolayer MoS₂ on GaN substrate is of high quality with a low presence of defects. The low presence of defects is beneficial because local defects may easily absorb water vapor and oxygen molecules in the air and thus affect the PL and Raman intensities locally, which might lead to inhomogeneity of PL and Raman intensities, as proved by previous studies.³²

First-principles calculations based on density functional theory (DFT) were performed to investigate the interfacial property and stacking structure and confirm the interlayer interactions of monolayer MoS₂ on GaN. Fig. 4 shows the electronic band structure of the bare monolayer MoS₂ calculated in a primitive unit cell. It can be seen that the maximum of the valence band and the minimum of the conduction band appear at the same point, indicating a direct band gap $E_g = 1.8$ eV, which is consistent with the PL emission peak value in our experimental results, as mentioned above and in other reported results,^{28,33} thus verifying the accuracy of the optimized atomic configurations.

Fig. 5(a) shows the top view and side view of the optimized structure of monolayer MoS₂ and GaN substrate, and the electronic band structure of the interfacial system is shown in Fig. 5(b). It can be clearly seen that the growth of MoS₂ on GaN does not significantly perturb the band structure and only

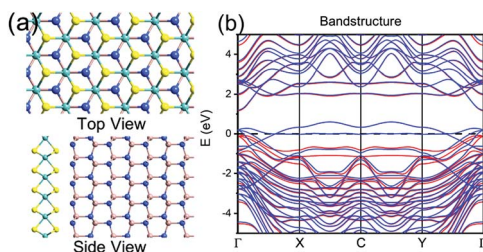


Fig. 5 (a) Top view and side view of the atomic structure of MoS₂ on GaN. (b) The electronic band structure of the interfacial MoS₂-GaN system.

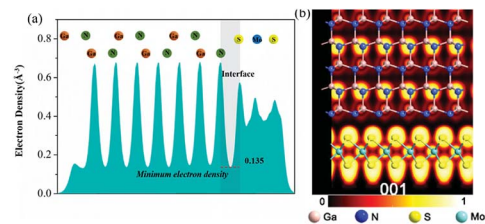


Fig. 6 (a) The average electron density in planes normal to the monolayer MoS₂-GaN interface. (b) The electron localization function (ELF) of the monolayer MoS₂-GaN geometrical structure.

brings about a very small band gap change. New electronic bands below -0.5 eV are introduced after the monolayer MoS₂ adsorption, while most bands of the interfacial MoS₂-GaN system retain the nature of original GaN bands, and the band gap does not change considerably. These results demonstrate that the growth of monolayer MoS₂ on the GaN surface would not make significant changes on the interfacial property of GaN, and the calculated binding energy of monolayer MoS₂ and (0001) GaN is as low as 0.0993 eV.

Furthermore, the average electron density plot along the monolayer MoS₂-GaN system is depicted in Fig. 6(a). The value of the electron density provides a visual and quantitative evaluation of the interfacial interactions. Higher electron density at the interface indicates higher interfacial interactions. The minimum value of electron density at the interface is 0.135 \AA^{-3} for the MoS₂-GaN system, indicating a low van der Waals force at the interface. Thus, the MoS₂ overlayers are not expected to show much strain because of their weak vdW interactions with the substrate. Furthermore, the electron localization function of the monolayer MoS₂-GaN system is displayed in Fig. 6(b) in order to estimate the bonding nature between MoS₂ and GaN substrate. It can be noted that the electron localization function indicates that electrons localize inside the GaN and MoS₂ layers instead of the interlayer space, suggesting the weak van der Waals force between the two layers. From this finding, we can conclude that the weak van der Waals force plays a minor role on the atomic vibration at the interface between monolayer MoS₂ and GaN, which is in agreement with the abovementioned experimental findings.

4. Conclusion

In conclusion, we have achieved epitaxial chemical vapor deposition growth of monolayer MoS₂ on a GaN semiconductor substrate. High crystalline quality and uniformity of the as-grown monolayer MoS₂ is demonstrated by their Raman and PL features and their intensity mappings, suggesting that GaN substrate is favorable for 2D MoS₂ growth. Moreover, laser power dependent Raman and PL characterizations were performed to investigate the optical and atomic vibration properties of MoS₂ on GaN. The relatively weak van der Waals interaction controls the crystalline growth and plays a minor role on the atomic vibration, which was explored and confirmed by interfacial structure and interaction investigations between monolayer MoS₂ and GaN substrate based on the first-



principles density functional theory calculations. As a result, the ability to grow high quality monolayer MoS₂ on semiconductor GaN substrate would open a route toward the synthesis of hetero and composite structures for electronic and optoelectronic device applications.

Conflicts of interest

The authors declare no competing financial interest.

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Notes and references

- 1 A. K. Geim and I. V. Grigorieva, *Nature*, 2013, **499**, 419–425.
- 2 H. Jeong, S. Bang, H. M. Oh, H. J. Jeong, S.-J. An, G. H. Han, H. Kim, K. K. Kim, J. C. Park, Y. H. Lee, G. Lerondel and M. S. Jeong, *ACS Nano*, 2015, **9**, 10032–10038.
- 3 D. Krasnozhan, D. Lembke, C. Nyffeler, Y. Leblebici and A. Kis, *Nano Lett.*, 2014, **14**, 5905–5911.
- 4 K. F. Mak, K. He, J. Shan and T. F. Heinz, *Nat. Nanotechnol.*, 2012, **7**, 494–498.
- 5 G.-H. Lee, Y.-J. Yu, X. Cui, N. Petrone, C.-H. Lee, M. S. Choi, D.-Y. Lee, C. Lee, W. J. Yoo, K. Watanabe, T. Taniguchi, C. Nuckolls, P. Kim and J. Hone, *ACS Nano*, 2013, **7**, 7931–7936.
- 6 E. Varrla, C. Backes, K. R. Paton, A. Harvey, Z. Gholamvand, J. McCauley and J. N. Coleman, *Chem. Mater.*, 2015, **27**, 1129–1139.
- 7 H. Liu, M. W. Si, S. Najmaei, A. T. Neal, Y. C. Du, P. M. Ajayan, J. Lou and P. D. D. Ye, *Nano Lett.*, 2013, **13**, 2640–2646.
- 8 K. Kang, S. Xie, L. Huang, Y. Han, P. Y. Huang, K. F. Mak, C.-J. Kim, D. Muller and J. Park, *Nature*, 2015, **520**, 656–660.
- 9 C. R. Serrao, A. M. Diamond, S.-L. Hsu, L. You, S. Gadgil, J. Clarkson, C. Carraro, R. Maboudian, C. Hu and S. Salahuddin, *Appl. Phys. Lett.*, 2015, **106**, 052101.
- 10 O. V. Yazyev and Y. P. Chen, *Nat. Nanotechnol.*, 2014, **9**, 755–767.
- 11 L. J. Van Ruyven, *J. Lumin.*, 1984, **29**, 123–161.
- 12 Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano, *Nat. Nanotechnol.*, 2012, **7**, 699–712.
- 13 H. Heo, J. H. Sung, G. Jin, J. H. Ahn, K. Kim, M. J. Lee, S. Cha, H. Choi and M. Jo, *Adv. Mater.*, 2015, **27**, 3803–3810.
- 14 H. R. Gutiérrez, N. Perea-Lopez, A. L. Elias, A. Berkdemir, B. Wang, R. Lv, F. Lopez-Urias, V. H. Crespi, H. Terrones and M. Terrones, *Nano Lett.*, 2013, **13**, 3447–3454.
- 15 B. Liu, M. Fathi, L. Chen, A. Abbas, Y. Ma and C. Zhou, *ACS Nano*, 2015, **9**, 6119–6127.
- 16 L. K. Li, Y. Yu, G. J. Ye, Q. Ge, X. Ou, H. Wu, D. Feng, X. H. Chen and Y. Zhang, *Nat. Nanotechnol.*, 2014, **9**, 372–377.
- 17 D. Dumcenco, D. Ovchinnikov, K. Marinov, P. Lazic, M. Gilbertini, N. Marzari, O. L. Sanchez, Y. C. Kung, D. Krasnozhan, M. W. Chen, S. Bertolazzi, P. Gillet, A. F. I. Morral, A. Radenovic and A. Kis, *ACS Nano*, 2015, **9**, 4611–4620.
- 18 Q. Q. Ji, Y. F. Zhang, T. Gao, Y. Zhang, D. L. Ma, M. X. Liu, Y. B. Chen, X. F. Qiao, P. H. Tan, M. Kan, J. Feng, Q. Sun and Z. F. Liu, *Nano Lett.*, 2013, **13**, 3870–3877.
- 19 Z. Lin, Y. Zhao, C. Zhou, R. Zhong, X. Wang, Y. H. Tsang and Y. Chai, *Sci. Rep.*, 2015, **5**, 18596.
- 20 Y. Liu, L. Hao, W. Gao, Z. Wu, Y. Lin, G. Li, W. Guo, L. Yu, H. Zeng, J. Zhu and W. Zhang, *Sens. Actuators, B*, 2015, **211**, 537–543.
- 21 S. Lin, X. Li, P. Wang, Z. Xu, S. Zhang, H. Zhong, Z. Wu, W. Xu and H. Chen, *Sci. Rep.*, 2015, **5**, 15103.
- 22 E. W. Lee, L. Ma, D. N. Nath, C. H. Lee, A. Arehart, Y. Wu and S. Rajan, *Appl. Phys. Lett.*, 2014, **105**, 203504.
- 23 S. Vishwanath, X. Liu, S. Rouvimov, P. C. Mende, A. Azcatl, S. McDonnell, R. M. Wallace, R. M. Feenstra, J. K. Furdyna, D. Jena and H. G. Xing, *2D Mater.*, 2015, **2**, 024007.
- 24 J. Taylor, H. Guo and J. Wang, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2001, **63**, 121104.
- 25 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865–3868.
- 26 S. Grimme, J. Antony, S. Ehrlich and H. Krieg, *J. Chem. Phys.*, 2010, **132**, 154104.
- 27 G. F. Yang, P. Chen, Z. L. Wu, Z. G. Yu, H. Zhao, B. Liu, X. M. Hua, Z. L. Xie, X. Q. Xiu, P. Han, Y. Shi, R. Zhang and Y. D. Zheng, *J. Mater. Sci.: Mater. Electron.*, 2012, **23**, 1224–1228.
- 28 B. Radisavljevic, A. Radenovic, J. Brivio, V. Giacometti and A. Kis, *Nat. Nanotechnol.*, 2011, **6**, 147–150.
- 29 A. Splendiani, L. Sun, Y. B. Zhang, T. S. Li, J. Kim, C. Y. Chim, G. Galli and F. Wang, *Nano Lett.*, 2010, **10**, 1271–1275.
- 30 C. Lee, H. Yan, L. E. Brus, T. F. Heinz, J. Hone and S. Ryu, *ACS Nano*, 2010, **4**, 2695–2700.
- 31 K. F. Mak, C. Lee, J. Hone, J. Shan and T. F. Heinz, *Phys. Rev. Lett.*, 2010, **105**, 136805.
- 32 H. Nan, Z. L. Wang, W. H. Wang, Z. Liang, Y. Lu, Q. Chen, D. W. He, P. H. Tan, F. Miao, X. R. Wang, J. L. Wang and Z. H. Ni, *ACS Nano*, 2014, **8**, 5738–5745.
- 33 Y. H. Lee, L. L. Yu, H. Wang, W. J. Fang, X. Ling, Y. M. Shi, C. T. Lin, J. K. Huang, M. T. Chang, C. S. Chang, M. Dresselhaus, T. Palacios, L. J. Li and J. Kong, *Nano Lett.*, 2013, **13**, 1852–1857.

