



 Cite this: *RSC Adv.*, 2018, 8, 22998

## Graphynes: indispensable nanoporous architectures in carbon flatland

 Anto James, Chris John, Cheriya cheruvakkara Owais, Stephen Nagaraju Myakala, Sarap Chandra Shekar,<sup>†</sup> Jyoti Roy Choudhuri<sup>‡</sup> and Rotti Srinivasamurthy Swathi \*

Theoretical design and experimental realization of novel nanoporous architectures in carbon membranes has been a success story in recent times. Research on graphynes, an interesting class of materials in carbon flatland, has contributed immensely to this success story. Graphyne frameworks possessing sp and sp<sup>2</sup> hybridized carbon atoms offer a variety of uniformly distributed nanoporous architectures for applications ranging from water desalination, gas separation, and energy storage to catalysis. Theory has played a pivotal role in research on graphynes, starting from the prediction of various structural forms to the emergence of their remarkable applications. Herein, we attempt to provide an up-to-date account of research on graphynes, highlighting contributions from numerous theoretical investigations that have led to the current status of graphynes as indispensable materials in carbon flatland. Despite unsolved challenges in large-scale synthesis, the future appears bright for graphynes in present theoretical and experimental research scenarios.

 Received 30th April 2018  
Accepted 7th June 2018

DOI: 10.1039/c8ra03715a

rsc.li/rsc-advances

### 1 Introduction to the family of graphynes

“Carbon seems the most favourable case, because it has the richest chemistry” is a remark made by Stephen Hawking about the chemical basis of alien life in the universe.<sup>1</sup> Carbon is the second most abundant element in the human body, due to its catenation property, although it is the fifteenth most abundant

element in the Earth’s crust. Carbon has many allotropic forms: the natural ones are diamond and graphite, and the important man-made allotropes are fullerenes, carbon nanotubes (CNTs), graphene and carbon nanofoams. Interestingly, most of these materials essentially consist of carbon atoms in the sp<sup>2</sup> state of hybridization. They have expanded the scope of available zero-, one-, and two-dimensional materials (0D, 1D, and 2D, respectively) with their wide range of mechanical, electronic, optical, chemical and electrochemical properties leading to many advanced applications in photovoltaic cells, organic light emitting devices, field-effect transistors and chemical sensors.<sup>2–4</sup> Among these, one-atom-thick graphene is known to be one of the most important allotropes exhibiting a variety of intriguing electronic, thermal,<sup>5</sup> optical,<sup>6,7</sup> and mechanical properties.<sup>8</sup> It has potential applications in electronic

*School of Chemistry, Indian Institute of Science Education and Research Thiruvananthapuram (IISER-TVM), Vithura, Kerala, India-695551. E-mail: swathi@iisertvm.ac.in*

<sup>†</sup> Present address: Institute for Computational Physics, University of Stuttgart, Allmandring 3, 70569 Stuttgart, Germany.

<sup>‡</sup> Present address: Presidency University, Bangalore, India-560064.



*Anto James is currently pursuing his doctoral studies under the supervision of Dr R. S. Swathi at the School of Chemistry, Indian Institute of Science Education and Research Thiruvananthapuram (IISER-TVM), India. His research is aimed at designing nano-oscillators by investigating the permeation of fullerenes through the pores of higher order graphynes.*



*Chris John is an Integrated PhD student pursuing her doctoral studies under the supervision of Dr R. S. Swathi at IISER-TVM. Her research focuses on developing optimal parameters for model potentials governing the interactions of atoms and molecules with graphene, carbon nanotubes, as well as boron nitride sheets and tubes.*



and optoelectronic devices,<sup>9–11</sup> chemical sensors,<sup>12</sup> nanocomposites<sup>13,14</sup> and energy storage.<sup>15–20</sup> First scientific study of graphene was carried out by Wallace in 1947,<sup>21</sup> although it was first observed by Boehm *et al.*<sup>22</sup> in 1962 and isolated by Geim and Novoselov<sup>23</sup> in 2004. Wallace theoretically predicted the band structure of graphene considering it as the monolayer model structure for graphite. Other forms of carbon networks such as graphene oxide, graphane, graphone and fluorographane can be derived by chemical functionalization of graphene. Hydrogenation and fluorination of graphene yield graphane and fluorographane, respectively. Partial hydrogenation (on alternate carbons) of graphene gives graphone. The introduction of epoxy, carboxyl and hydroxyl linkages on graphene leads to the formation of graphene oxide. Of these range of systems, one can notice that graphene has  $sp^2$  hybridized

carbons, graphane and fluorographane have  $sp^3$  hybridized carbons and graphone and graphene oxide have both  $sp^2$  and  $sp^3$  hybridized carbons in their networks.

In the last two decades, new members of the carbon family referred to as graphynes, one-atom-thick two-dimensional networks similar to graphene have gained much importance.<sup>24,25</sup> Graphynes (GYs) were first theoretically proposed by Baughman and co-workers in 1987. Indeed, theory preceded and guided the synthesis and fabrication of most of the new forms of carbon-based materials known today. For instance, graphane and graphone were also first predicted theoretically and subsequently realized in experiments.<sup>26–29</sup> Theoretically, GYs are constructed by replacing some carbon-carbon  $sp^2$  bonds in graphene with acetylenic (*i.e.*, consisting of single- and triple-bonded) linkages.<sup>30</sup> Various types of GYs are proposed based on the variation of positions at which the acetylenic linkers are added. Depending on the structures, they are classified into various members known as  $\alpha$ -GY,  $\beta$ -GY,  $\gamma$ -GY,  $\delta$ -GY, rhombic-GY, 6,6,12-GY *etc.*,<sup>31</sup> as shown in Fig. 1. GYs possess intrinsically porous structures and the porosity can be tuned by changing the number of acetylenic linkers in the architectures. According to the number of acetylenic linker groups ( $N$ ), they are named as GY- $N$  with  $N$  equal to zero yielding the graphene.



*Cheriya Chervakkara Owais is an Integrated PhD student pursuing his doctoral studies under the supervision of Dr R. S. Swathi at IISER-TVM. His research focuses on employing continuum approaches for modeling carbon-based nanostructures such as graphene, carbon nanotubes and graphynes.*



*Stephen Nagaraju Myakala is a BS-MS student at IISER-TVM working under the supervision of Dr R. S. Swathi. He is currently a recipient of the INSPIRE fellowship.*



*S. Chandra Shekar completed his PhD under the supervision of Dr. R. S. Swathi at IISER-TVM. He is currently pursuing his post-doctoral studies at the Institute for Computational Physics, University of Stuttgart, Germany. His research is mainly focused on the two-dimensional carbon networks such as graphene, graphyne and graphdiyne and their utility in the context of lithium ion batteries, DNA sensing and photoswitching.*



*Jyoti Roy Choudhuri completed his PhD under the supervision of Prof. Amalendu Chandra at IIT, Kanpur. Subsequently, he had a postdoctoral stint at the Virginia Commonwealth University under the supervision of Prof. Alenka Luzar. He followed this with a short stint in the research group of Dr R. S. Swathi at IISER-TVM. Jyoti is currently an Assistant Professor at the Presidency University, Bangalore, India.*



*R S Swathi obtained PhD in theoretical chemistry from Indian Institute of Science, Bangalore. Subsequently, she joined as an Assistant Professor in School of Chemistry, IISER-TVM. Her research group employs analytical and computational approaches for modeling chemical phenomena involving carbon-based and metal-based nanostructures. She is a recipient of the Young Scientist Awards from Indian National Science Academy, National Academy of Sciences, Indian Academy of Sciences, Kerala State Council for Science, Technology and Environment and Distinguished Lectureship Award from the Chemical Society of Japan. She is currently on the Editorial Advisory Board of ACS Applied Materials & Interfaces.*



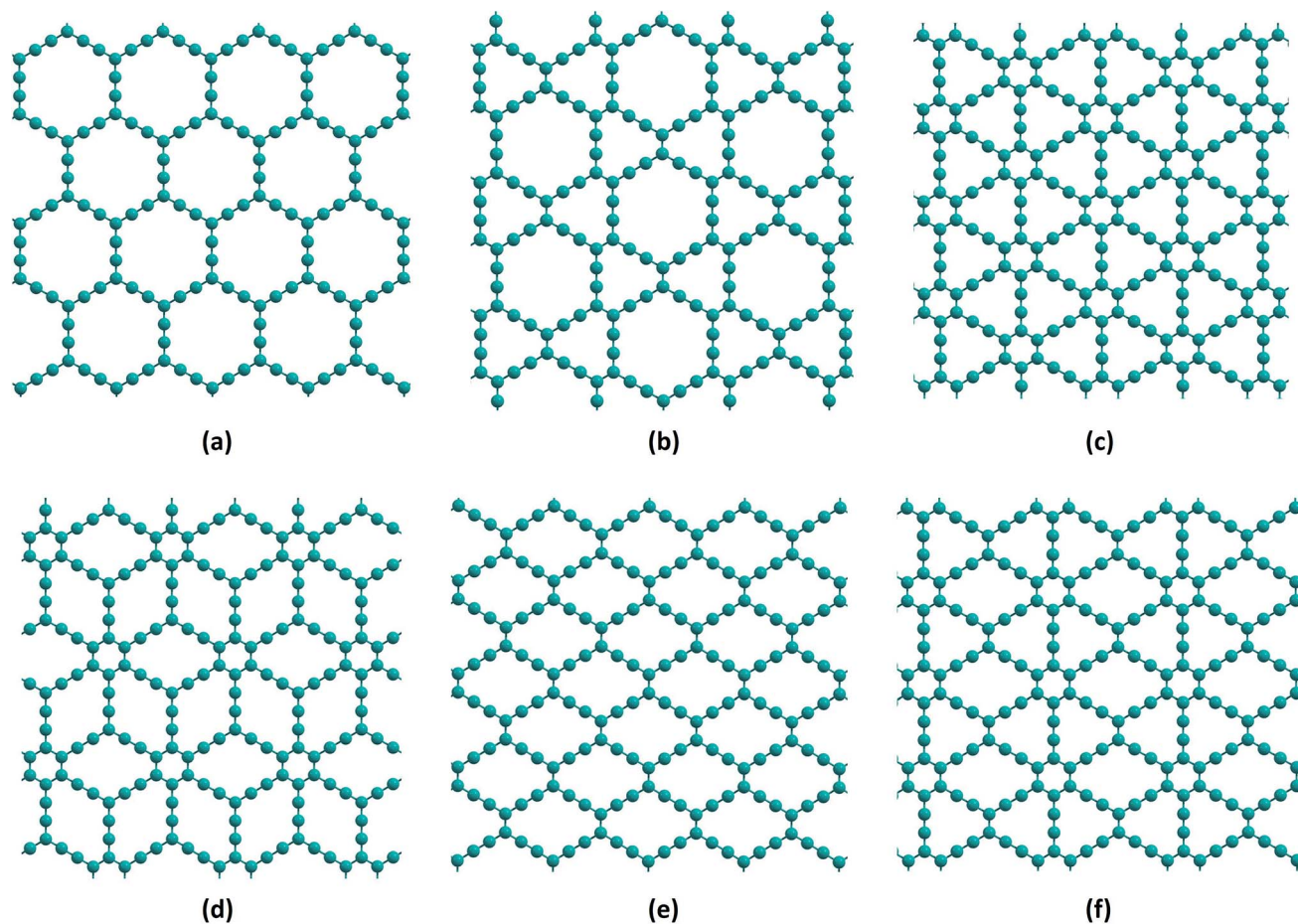


Fig. 1 Representative structural models for (a)  $\alpha$ -GY (b)  $\beta$ -GY (c)  $\gamma$ -GY (d)  $\delta$ -GY (e) rhombic-GY and (f) 6,6,12-GY.

GY-2 of  $\gamma$ -GY (known as graphdiyne, GDY) and rhombic-GY (known as carbon ene-yne, CEY) are the only structures reported to be successfully synthesised in large quantities.<sup>32,33</sup> Thin films of GDY have been synthesised on top of copper<sup>34</sup> and silver substrates.<sup>35</sup> A recent report shows that thin films of rhombic-GY-2 can also be synthesised on copper substrates.<sup>36</sup> However, small annulenic units of various GYs are prepared from the benzannulated and perethynylated dehydroannulene-derived substructures. A recent review on GDY gives a detailed picture of various significant applications as well as importance of GDY.<sup>37</sup> Unique structural and electronic properties of various GYs make them even more imperative. Alike graphene, GYs also show some unique chemical, electronic, mechanical and structural properties<sup>38–41</sup> that have important applications ranging from energy storage to nanoelectronics.<sup>42</sup> Nanoporous membranes of graphene and its analogues have been employed for achieving selective permeation of gas molecules, ions, hydrocarbons, biomolecules *etc.* However, attaining uniform porosity in graphene is a difficult task and hence carbon membranes like GYs which possess intrinsic pores are of great interest.<sup>43</sup> Theoretical studies have predicted the chemical inertness of free-standing graphyne layers and their stability under ambient temperature.<sup>44–46</sup>  $\gamma$ -GYs exhibit non-zero band gap in contrast to the zero band gap of graphene. The distortion

in the Dirac cone of 6,6,12-GY arising due to the presence of sp and sp<sup>2</sup> carbons results in the directionality of electric conductance.<sup>47</sup> Graphene is one of the stiffest materials known to exist after carbyne and diamond. Addition of acetylenic linkers in graphene results in reduction of stiffness. From  $\gamma$ -GY-1 to  $\gamma$ -GY-5, increase in sp carbon atom content results in 10 to 50% reduction in Young's modulus compared to graphene.<sup>39</sup> The decrease in carbon atom density in GYs compared to graphene also results in lower values of elastic modulus,<sup>48</sup> as shown in Table 1. The velocity of sound waves in  $\gamma$ -GY is predicted to vary with external stress, indicating the application of GYs as surface acoustic sensors.<sup>49</sup> Researchers suggest that these materials can surpass graphene in the near future.<sup>47</sup>

Similar to graphene which can be rolled into CNTs, GYs also form nanotubes. Graphyne nanotubes (GNTs) were proposed in 2003 by Coluci and co-workers.<sup>50</sup> The key difference between a CNT and a GNT is that the latter has porous walls, thus allowing material transport through the sidewalls. Coluci *et al.* classified the GNTs into three families,  $\alpha$ -GNT,  $\beta$ -GNT and  $\gamma$ -GNT, each derived from their corresponding parent graphyne sheets, respectively. Furthermore, the metallic or semi-conducting nature of GNTs is governed by the same rule as in case of CNTs. ( $n,0$ ) and ( $n,n$ )  $\alpha$ -GNTs are referred to as zigzag and armchair, respectively, whereas, ( $2m,m$ ) and ( $n,0$ )  $\beta$ -GNTs



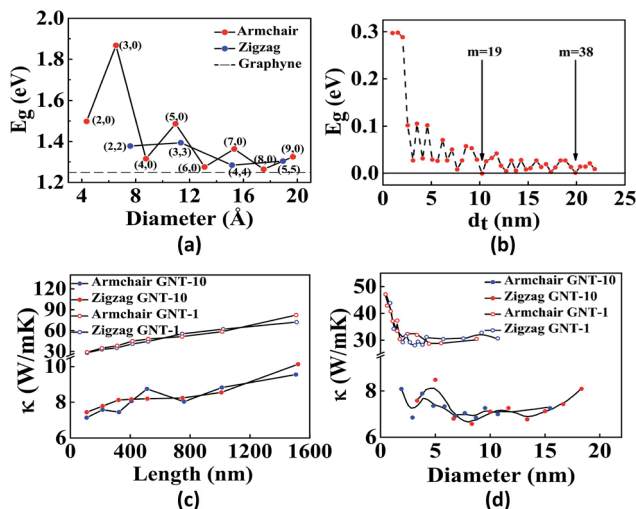
**Table 1** Structural and mechanical properties of various graphynes and graphene. Reproduced from ref. 48 with permission from the AIP Publishing

Model system	Percentage of acetylenic linkages	Atom density (atoms per nm <sup>2</sup> )	Young's Modulus (GPa)
$\alpha$ -GY	100	18.92	120
$\beta$ -GY	66.67	23.13	261
$\gamma$ -GY	33.33	29.61	505
6,6,12-GY	41.67	28.02	445
Graphene	0	39.95	995

and  $\gamma$ -GNTs are referred to as zigzag and armchair, respectively.<sup>50</sup> Comparing the band structures of the three, it was found that the band gap has an oscillatory behaviour for  $\alpha$ -GNT and  $\beta$ -GNT as a function of the GNT diameter, while  $\gamma$ -GNT has diameter-independent band gap. Contrary to this, Wang *et al.* have found that  $\gamma$ -GNTs also show oscillatory behaviour<sup>51</sup> (Fig. 2a and b). The key issues at present are the stability of the different forms of GNTs and the feasibility of the different synthetic methods. Although several reports predict that  $\gamma$ -GNTs are likely the first ones to be synthesised because of their lowest energy (and hence the most stable amongst the three forms), there have been successful attempts at the synthesis of other types of GNTs. Li and co-workers successfully demonstrated synthesis of high-quality GDY nanotubes with a wall thickness of 15 nm and nanotube length of about 40 microns obtained after annealing of the GNTs at 650 °C.<sup>52</sup> Alaei *et al.* studied the interaction of transition metals with  $\gamma$ -GNTs. Complexes of  $\gamma$ -GNTs with Co and Fe yielded net non-zero magnetic moments, opening up several applications in electronics and spintronics.<sup>53</sup> In addition to being nonmagnetic,

pristine  $\gamma$ -GNTs possess zero electric dipole moment as well. Deb *et al.* have shown that functionalization of H<sub>2</sub>O molecule on  $\gamma$ -GNTs can cause an increase in dipole moment.<sup>54</sup> Thus,  $\gamma$ -GNTs act as electron acceptors and therefore can be used as n-type semiconductors. Hu *et al.* found that  $\gamma$ -GNTs have remarkably low thermal conductivity (10 W mK<sup>-1</sup>), about two orders of magnitude lower when compared to any of the pristine or chemically functionalized CNTs.<sup>55</sup> Direct proportionality of thermal conductivity was observed to the length of the tube in the range of 0.1 to 1.5  $\mu$ m, beyond which it showed negligible change. Likewise, there was steep increase of thermal conductivity for tubes with a diameter less than 2 nm but was independent for all diameters greater than 2 nm (Fig. 2c and d). In addition to their array of distinct properties, de Sousa *et al.* have found that  $\gamma$ -GNTs also exhibit superplasticity.<sup>56</sup> Studies of the newly synthesised GNTs and nanowires with their field emission properties, high conductivity and mobility open numerous opportunities in the fabrication of graphyne nanodevices.<sup>52,57</sup> Graphyne nanoscrolls are structures obtained by rolling nanosheets in a papyrus-like topology. The nanoscroll diameter can be easily tuned when compared to the GNT diameters. Solis *et al.* have predicted that their stability depends on the critical value of the ratio between the length and the height of the GY sheets used.<sup>58</sup>

In this review, we describe some of the potential chemical applications of GYs on the basis of the latest theoretical and experimental results. We have attempted to launch an up-to-date account of most topics relevant in this field. This review is organised as follows. In Section 2, we have discussed the different approaches that have been used for the synthesis of GYs. This is followed by Sections 3 to 8, where various potential applications of GYs ranging from water desalination, gas separation, gas sensing, energy storage, catalysis to optoelectronics are discussed in detail. Finally, the review is concluded by giving a brief exposition to the future of GYs in Section 9.



**Fig. 2** Representative properties of GNTs: effect of tube diameter on the band gap of (a)  $\gamma$ -GNTs and (b) zigzag  $\beta$ -GNTs. Dependence of the thermal conductivity on (c) tube length and (d) tube diameter of various GNTs. Part a is reproduced from ref. 51 with permission from the American Chemical Society. Part (b) is reproduced from ref. 59 and parts (c) and (d) are reproduced from ref. 55 with permissions from the American Physical Society.

## 2 Approaches for the synthesis of graphynes

Ever since they were theoretically predicted in 1987, GYs in various forms and structures have been modelled. Extensive studies of their unique properties have been carried out, indicating a plethora of applications. However, the central challenge of devising efficient industrially realizable synthetic methods still remains. Recently, Li *et al.* synthesised large area of GDY films *via* the homocoupling of hexaethynylbenzene on a Cu foil in the presence of pyridine.<sup>61</sup> As predicted by Baughman and co-workers, the obtained films were highly crystalline in nature, as confirmed by their XRD patterns. Qian *et al.* synthesised GDY films on ZnO nanorod arrays by using a self-catalysed vapor–liquid–solid growth process. GDY films were fabricated by vapour deposition on ZnO nanorod arrays from the vapour source of GDY powder. These GDY films exhibit very low turn on field of 4.2 V  $\mu\text{m}^{-1}$ .<sup>62</sup> Despite the large area and high crystallinity obtained in the early methods of GDY synthesis, these methods involve use of hazardous solvents and strong



reaction conditions. However, Nishihara and co-workers overcame this limitation by synthesising few layers of GDY of higher quality using a bottom-up approach (Fig. 3a).<sup>60</sup> They carried out catalytic polymerization of hexaethynylbenzene in a mixture of dichloromethane and toluene with Cu as a catalyst. GDY sheets thus synthesised are characterized by transmission electron microscope and atomic force microscope (Fig. 3b and c).

Another example of bottom-up approach for the synthesis of  $\gamma$ -GYs is oligotrimerization and cyclotrimerization of graphyne subunits. In recent advances, the hexakis-benzenes are proven to be the best graphyne subunits. While they seem to be the easy route, they have certain drawbacks such as additional cyclo-oligomeric and/or polymeric products often resulting in a low yield of the desired materials.<sup>63</sup> Intramolecular cyclization is yet another powerful method for the construction of more complex structures of the GY networks. However, Haley showed that it produces a meagre yield. To overcome this problem, considering the high efficiency of alkyne metathesis, Haley combined it with pre-organized propynyl groups obtaining a higher yield of 20%.<sup>63</sup> Wu *et al.* also developed a similar alkyne metathesis, synthesising graphyne-like porous networks whose thickness can be tuned by adjusting the reaction time.<sup>64</sup> To minimise other side products, various methods such as intramolecular alkylation with suitable

molecules like bis(dibenzylideneacetone)palladium(0) ( $\text{Pd}(\text{dba})_2$ ) under high dilution conditions were also discovered.

Ding's group found that small carbyne chains on nano-active materials have a curved polynic structure on less active metal surfaces (Cu), whereas, they possess a linear configuration on active metals like Ni.<sup>65</sup> Hence, a synthesis method of GY sheets by self-assembly of the carbyne chains at low temperature *via* carbon cluster sputtering was suggested.<sup>65</sup> It has been demonstrated that certain functionalization leads to high selectivity in alkyne homocoupling forming polymer strands which exhibit mutual lateral attraction yielding multiple stranded assemblies that are stable at room temperature. Klappenberger *et al.* constructed high quality, flexible CN-functionalized 3–2 GDY nanowires (with three phenyl and two alkyne (terminal) moieties) *via* on-surface covalent coupling of terminal alkyne building blocks.<sup>66</sup> From these results, the authors also concluded that secondary functional groups drastically increased the coupling efficiency. Jia and co-workers synthesised thin films of rhombic-GY-2 on copper substrates using a solvent-phase reaction.<sup>36</sup> Thin films were synthesised using tetraethynylethene molecules in pyridine solvent under  $\text{N}_2$  protected atmosphere. Apart from these successful syntheses, many dehydroannulene and radialene molecules are reported to be synthesised in literature, which can be considered as the molecular sub-units of GYs. A few representative molecular model systems of GYs are shown in Fig. 3d–f.

### 3 Graphynes for water desalination

Graphynes are intrinsically nanoporous one-atom-thick membranes with good mechanical and chemical stability. Unlike graphene, GYs can be used for molecular filtration without any post-synthetic modifications. Well-defined porous structures of GYs prevent defect formation over the membranes. Water desalination, one of the key aspects of molecular filtration is becoming increasingly important with the depletion of natural freshwater resources. Recent molecular simulation studies and experiments have demonstrated that GYs are much better membranes for water desalination compared to the commercially available forward osmosis (FO) and reverse osmosis (RO) membranes. Liu and co-workers have used the  $\gamma$ -GY-2 membranes synthesised over Cu foils to remove lead ions from water.<sup>67</sup> Acetylenic linkages in  $\gamma$ -GY-2 were shown to interact strongly with metal ions and their adsorption helps in removal of heavy metal ions.<sup>67</sup> However, most of the studies related to water desalination are carried out theoretically and the results are yet to be realized experimentally.

In 2014, Bartolomei *et al.* studied the interactions between  $\gamma$ -GYs and water molecules using second-order Moller-Plesset perturbation theory calculations (MP2C) with extended basis set.<sup>68</sup> Calculations were carried out on the smallest precursors of GYs, namely annulenes. Lengths of the sides of the triangular pores of  $\gamma$ -GY-1,  $\gamma$ -GY-2 and  $\gamma$ -GY-3, were estimated to be 1.3, 3.9 and 6.4 Å respectively, while the van der Waals diameter of water is in the range of 3.15–3.28 Å. There was a finite barrier for the penetration of water through  $\gamma$ -GY-1 and  $\gamma$ -GY-2 while  $\gamma$ -GY-3 allowed the barrierless flow of water. However, hydrogen

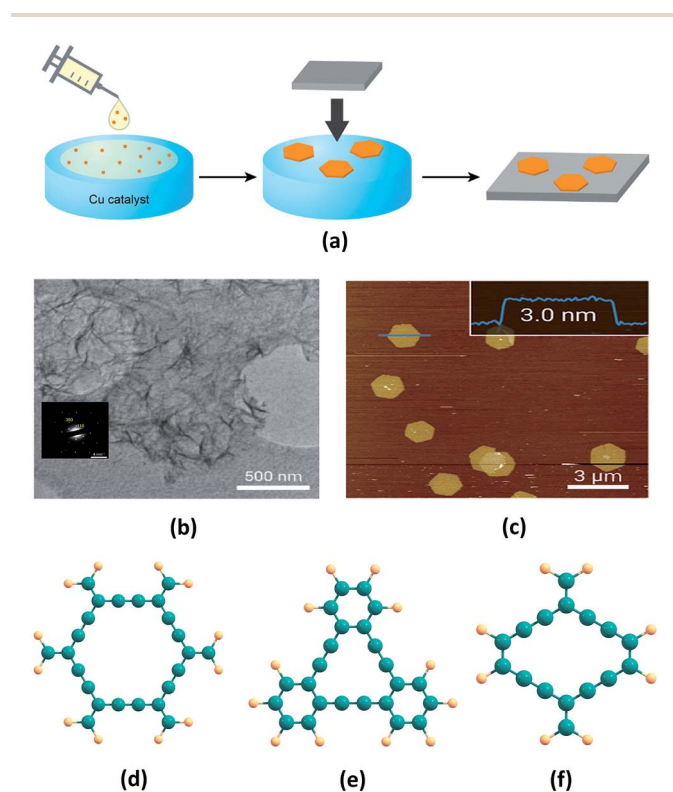


Fig. 3 (a) Illustration of the synthetic procedure and transfer of GDY thin films on copper substrates as adopted by Nishihara and co-workers. (b) Transmission electron microscopy image and selected-area electron diffraction pattern (inset) at the liquid–liquid interface, and (c) atomic force microscopy topographical image at the gas–liquid interface of GDY nanosheets synthesised by a bottom-up approach. Molecular sub-units of (d)  $\alpha$ -GY, (e)  $\gamma$ -GY and (f) rhombic-GY. Parts (a–c) are adapted from ref. 60 with permissions from the American Chemical Society.



bonding interaction from a second water molecule on the other side of the graphyne pore can reduce the barrier for water penetration, as illustrated in Fig. 4a. A comparative study of MP2C calculations with Lennard-Jones (LJ) and improved Lennard-Jones (ILJ) pair potentials was carried out by Bartolomei and co-workers.<sup>68</sup> ILJ force field (FF) optimization curves serve as a good approximation for water-graphyne interactions (Fig. 4b). The ILJ pair potential is given by

$$V(x) = \varepsilon \left[ \frac{6}{n(x) - 6} \left( \frac{1}{x} \right)^{n(x)} - \frac{n(x)}{n(x) - 6} \left( \frac{1}{x} \right) \right], \quad (1)$$

where  $n(x) = \beta + 4.0x^2$  and  $x = R/R_m$ .  $\varepsilon$ ,  $R_m$  and  $\beta$  are parameters depending on the nature of the interacting particles.

Employing molecular dynamics (MD) simulations, Zhang *et al.* demonstrated that  $\gamma$ -GY membranes could be used as FO membranes to purify water.<sup>70</sup> Water fluxes through a series of GY- $N$  membranes ( $N = 3, 4, 5$  and  $6$ ) were studied and they were all found to be of the same order of magnitude. In the simulation, no linear correlation between structural properties like pore size and water flux during FO was observed. Water flux through the membrane was influenced by hydrogen bonding, electrostatic and van der Waals interactions between water molecules, solutes and membranes. An increase in the number of average hydrogen bonds was observed with the increase in pore size of the membrane, which resulted in the slowing of water flow through pores. Similar to RO, the salt rejection efficiency of GY as FO membrane is dependent on the size

exclusion effect.  $\gamma$ -GY-3 gives a complete salt rejection while  $\gamma$ -GY- $N$  membranes with  $N > 3$  showed decreased salt rejection. Water flux through GY membranes is found to be around  $10 \text{ l cm}^{-1} \text{ h}^{-1}$ , which is three to four orders of magnitude higher than the commercial RO and FO membranes. As GYs are one-atom-thick membranes, concentration polarization phenomenon, which results in a reduction of water flux through FO membranes is minimal. Water flux through each pore and pore density on the sheet are two important factors that determine water permeability across GY membranes.<sup>69</sup> Water flow per unit pore for GY- $N$ s increases stepwise from  $N = 3$  to  $N = 6$ . As the pore size of the GY sheet increases, the pore density starts to decrease. As a result, effective water flow per unit area is the highest for  $\gamma$ -GY-4. Qin *et al.* have observed that the water molecules near the interface of GY show peculiar periodicity and higher mass density compared to the bulk.<sup>71</sup> The higher water viscosity can be attributed to the formation of layers of in-plane hydrogen-bonded networks near the membrane. Passage of water molecules through the pores requires the breaking of this ordered structural arrangement. The water molecule at the centre of the triangular pore of  $\gamma$ -GY forms a hydrogen bond with another water molecule across the membrane helping in permeation. As soon as the water permeates through the membrane, water molecule rotates in such a way that it forms two hydrogen bonds with the water layer and a third hydrogen bond with the next incoming water molecule. Thus, at a time, a single water molecule passes through the membrane resulting in quantized water transport through GYs.<sup>69</sup> Hydrogen bond formation between water molecules across the membrane helps in the formation of short single-file water structure, promoting fast water transmission as well as an increase in the concentration of water molecules close to the membrane surface.<sup>72</sup> Due to the rotation, there exists no consistency in the orientation of the water dipole while permeating through GY membranes. However, other membrane channels like CNTs retain the dipole orientations during permeation.<sup>72</sup>

Kou and co-workers carried out MD simulations to study permeation of water through graphyne membranes.<sup>72</sup> Water molecules were able to permeate through  $\gamma$ -GY-3 while the smaller pores of  $\gamma$ -GY-2 do not allow water molecules to pass through. Later, Kou *et al.* demonstrated that the  $\gamma$ -GY-3 membrane shows complete salt rejection for hydrostatic pressure varying from 0 to 350 MPa, while other graphyne analogues with higher pore size show decrease in salt rejection with an increase in pressure.<sup>73</sup> It was also observed that an increase in salt concentration decreases water permeability through GYs. With the increase in salt concentration, more water molecules become part of hydration structure and thus prevent random motion of molecules and permeability. Similar to other nanoporous membranes, the passage of ions through GY membranes is dependent on the hydrodynamic radii of the ions. Using MD simulations, Xue *et al.* have shown that the  $\alpha$ -GY-1 and  $\beta$ -GY-1 membranes give 100% salt rejection.<sup>74</sup>  $\gamma$ -GY-4 pores allow the permeation of  $\text{K}^+$  and  $\text{Na}^+$  ions. Water flux across the membranes for different GYs,  $\alpha$ -GY-1,  $\beta$ -GY-1,  $\gamma$ -GY-3 and  $\gamma$ -GY-4 was estimated to be 5.5, 8.11, 11.7 and  $14.31 \text{ l cm}^{-2}$  per day per MPa, respectively. Xue and co-workers mapped the

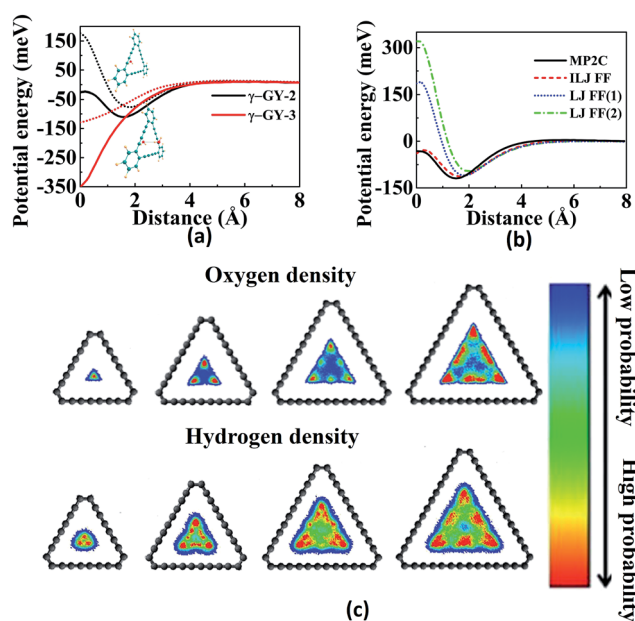


Fig. 4 Water permeation through GYs: (a) effect of hydrogen bonding on the energy barrier for water penetration through  $\gamma$ -GYs. (b) Comparison of potential energy scans of water penetration through nanopores of  $\gamma$ -GY-2 calculated using MP2C as well as ILJ and LJ FFs. (c) Oxygen atom and hydrogen atom densities inside the pores of various  $\gamma$ -GYs. Parts (a) and (b) are reproduced from ref. 68 with permissions from the American Chemical Society. Part (c) is reproduced from ref. 69 with permission from the Nature Publishing Group.



distribution characteristics of oxygen density while water molecules passed through the GY membranes (Fig. 4c). The density distributions for  $\alpha$ -GY and  $\beta$ -GY are observed to be round in shape, with the highest probability of water density at the centre of the pore. For  $\gamma$ -GY-3 and  $\gamma$ -GY-4, the density distribution map is triangular.  $\gamma$ -GY-3 has one channel for water passage while  $\gamma$ -GY-4 has three apparent channels for water to pass through.<sup>74</sup> Energy profiles for the passage of water molecules and ions are different across  $\alpha$ -GY-1. Minor valleys present near the graphyne pore help in the formation of the hydrogen-bonded single-file configuration of water molecules on both sides of the membrane, which in turn helps in higher water flux. GYs are found to be more efficient than conventional RO membranes in both water flux as well as salt rejection.

Graphyne membranes possess acetylenic linkages which can get hydrogenated or hydroxylated in the presence of protons or radicals in water. The evaluation of water desalination performance of functionalized GYs is also important. Raju *et al.* studied the desalination performance of bare as well as hydrogenated GYs.<sup>75</sup> Hydrogenation of GYs results in 20–40% reduction in pore area. Water flux through the GY membranes decreases in the following order:  $\gamma$ -GY-4 > hydrogenated  $\gamma$ -GY-4 (H  $\gamma$ -GY-4) >  $\gamma$ -GY-3 >  $\alpha$ -GY-1 >  $\gamma$ -GY-2 > hydrogenated  $\gamma$ -GY-3 (H  $\gamma$ -GY-3) > hydrogenated  $\alpha$ -GY-1 (H  $\alpha$ -GY-1) > hydrogenated  $\gamma$ -GY-2 (H  $\gamma$ -GY-2). Hydrogenated  $\gamma$ -GY-2 does not allow any water passage up to a pressure of 100 MPa. Except for  $\gamma$ -GY-4, all membranes show high salt rejection efficiency (greater than 75%) for pressures up to 2 GPa.  $\gamma$ -GY-2,  $\gamma$ -GY-3,  $\alpha$ -GY-1, H  $\alpha$ -GY-1 and H  $\gamma$ -GY-3 show a significant water flux as well as salt rejection efficiency that is required for an efficient RO membrane. Hydrogenation of GYs does not change the shape of probability density map of oxygen passing through membranes, but the area of oxygen density distribution is significantly reduced. Similar to graphene nanopores, pore functionalization can introduce selectivity for the passage of the ions. H  $\gamma$ -GY-4 membrane allows selective permeation of negatively charged ions, while hydroxylation of the  $\gamma$ -GY-4 pore results in the passage of positively charged ions and rejection of negatively charged ions.

## 4 Graphynes for gas separation

The separation of gases is considered to play a vital role in various processes, from industrial-scale to small-scale applications, such as H<sub>2</sub> production from syngas, separation of atmospheric gases for medical and industrial utilization, and isotope separation for usage in nuclear power.<sup>76–79</sup> In recent times, people are searching for clean and renewable energy due to the depletion of fossil fuels and increased environmental problems. Owing to natural abundance and environmental friendliness, hydrogen has been identified for this purpose. Therefore, it is very essential to separate hydrogen from other undesirable species in an efficient and cost-effective way. Among various processes of gas separation, membrane technology proves to be one of the well accepted approaches due to its several advantages including facile operation, low energy consumption, and easy maintenance.<sup>80</sup> Among diverse membranes, graphene-

based membranes are of profound interest and have a great potential for gas separation due to their one-atom-thickness, and consequent high efficiency.<sup>81,82</sup> Even small atoms like He cannot permeate through graphene sheet.<sup>83,84</sup> It is therefore imperative to create pores using experimental techniques<sup>85,86</sup> to achieve gas permeability. However, the creation of pores in graphene often leads to defects in the lattice and is not cost effective. Thus, GYs are proposed as alternative materials for gas separation as they possess large natural pores formed by the replacement of some carbon-carbon bonds in graphene with acetylenic linkages.

Jiao *et al.* suggested GDY as the membrane for achieving the separation of H<sub>2</sub> from CH<sub>4</sub> and CO using density functional theory (DFT) and transition state theory (TST).<sup>87</sup> The uniformly distributed pores, with sizes in between the van der Waals diameter of H<sub>2</sub> and CH<sub>4</sub>/CO make GDY a good option for hydrogen sieving from syngas. The selectivity for the separation of binary mixtures of gases (H<sub>2</sub>/O<sub>2</sub>, H<sub>2</sub>/N<sub>2</sub>, H<sub>2</sub>/CO<sub>2</sub>, H<sub>2</sub>/CH<sub>4</sub>) was studied by Zhao and co-workers using MD simulations.<sup>88</sup> Permeance of H<sub>2</sub> was found to increase as the pressure increases from 47 MPa to 1.5 GPa and decrease with increasing pressure after 1.5 GPa. Cranford *et al.* investigated temperature and force dependence of the hydrogen diffusion and its separation from syngas through GDY pores using MD simulations.<sup>89</sup> The energy barrier for H<sub>2</sub> molecule is determined to be 0.11 eV. High temperature is needed for the permeation of CO and CH<sub>4</sub>, whereas H<sub>2</sub> can pass through at room temperature. A driving force can enhance the selectivity as well as permeation of H<sub>2</sub> up to a limit of 40 pN. Increasing the force above this limit results in permeation of CO and further increase results in CH<sub>4</sub> permeation. The introduction of nitrogen atoms in GDY by replacing three sp<sup>2</sup> hybridized carbon atoms enhances the hydrogen selectivity over syngas.<sup>90</sup> The energy barrier for the permeance of hydrogen decreases whereas the barrier height increases for other gases. These methods for hydrogen purification involve complexity due to chemical modification of the nanopores. Tan *et al.* investigated whether the selectivity for hydrogen purification can be enhanced by the introduction of positive charge in GDY.<sup>91</sup> GDY displays a good performance for the separation of H<sub>2</sub> from mixtures of large molecules like CH<sub>4</sub>. However, the pore structure appears to be too large to accomplish ideal H<sub>2</sub> separation from small molecules such as CO and N<sub>2</sub>. Zhang *et al.* explored the use of rhombic-GY-1, whose pore size is in between that of  $\gamma$ -GY-1 and GDY for the advancement of selectivity of hydrogen purification from CO<sub>2</sub>/N<sub>2</sub>/CH<sub>4</sub>.<sup>92</sup> The separation is not based on the kinetic diameter but according to the physical and chemical interactions between the pores and the penetrating molecules.<sup>43</sup>

The purification of H<sub>2</sub> is also investigated using  $\alpha$ -GY and  $\gamma$ -GY. Sang *et al.* modified the triangular pores in  $\gamma$ -GY-1 by rupturing one-third of acetylenic bonds and replacing the unsatisfied valences using nitrogen and hydrogen atoms to obtain  $\gamma$ -GYN and  $\gamma$ -GYH, respectively.<sup>93</sup> The increased pore size which is in between that of  $\gamma$ -GY-1 and  $\gamma$ -GY-2 results in a decrease of the energy barrier for the selective separation of H<sub>2</sub> from other gases using both  $\gamma$ -GYN and  $\gamma$ -GYH. It is found that there is no electron density overlap between H<sub>2</sub> and  $\gamma$ -GYX



(X = N, H). MD and DFT results show that both modified membranes are acceptable candidates for the sieving of H<sub>2</sub> even though  $\gamma$ -GYN is the best among them. As temperature increases, the selectivity was found to decrease as shown in Fig. 5a. Alaghemandi *et al.* used hydrogenated  $\alpha$ -GY with and without defects as membranes for molecular sieving of pairs of H<sub>2</sub>/N<sub>2</sub> and H<sub>2</sub>/CH<sub>4</sub>.<sup>94</sup> Since CH<sub>4</sub> molecule is impermeable, the selectivity of H<sub>2</sub> is infinity and is independent of temperature, whereas in case of H<sub>2</sub>/N<sub>2</sub>, the selectivity of H<sub>2</sub> decreases with increase in temperature due to high kinetic energy and permeation probability of N<sub>2</sub>. In case of hydrogenated  $\alpha$ -GY with a defect (which is obtained by the removal of a carbon atom that is surrounded by three carbon atoms), H<sub>2</sub> selectivity increases with increase in temperature due to rearrangements of sp<sup>3</sup> free arms at the defect position.

Oxygen purification from toxic gases is essential since it is of great importance in medical field. Meng *et al.* studied the separation of O<sub>2</sub> gas from a mixture of toxic gases using  $\gamma$ -GY-2.<sup>95</sup> The oxidation of acetylenic bond does not occur since the process is associated with a large energy barrier. The electron density overlap of atoms of the triangular pore rims of  $\gamma$ -GY-2 and O<sub>2</sub> is very less and hence the barrier is easily surmountable. Separation of CO<sub>2</sub> and N<sub>2</sub> is a challenge as it decreases the energy content of natural gas as well as corrodes the pipelines. Zhao *et al.* investigated the separation of CO<sub>2</sub> and N<sub>2</sub> from CH<sub>4</sub> using modified GDY obtained by removing one-third of acetylenic linkages and replacing with H, F and O atoms forming GDY-H, GDY-F and GDY-O, respectively.<sup>96</sup> GDY-F and GDY-O are the best candidates for the separation of CO<sub>2</sub>/CH<sub>4</sub> and N<sub>2</sub>/CH<sub>4</sub>

due to the strong electronic interactions between CO<sub>2</sub> and the membranes, which decrease the barrier height for transmission. GDY-O can separate CO<sub>2</sub>/N<sub>2</sub> since the permeance of CO<sub>2</sub> is higher by two orders of magnitude. The permeances of CO<sub>2</sub> and N<sub>2</sub> through GDY-F and GDY-O are higher than in case of nanoporous graphene, making them a good choice for the separation process. Increasing the temperature increases the kinetic energy leading to the permeation of both gas molecules, thereby causing a decrease in selectivity.

Lei *et al.* explored the adsorption and diffusion of H<sub>2</sub>S and CH<sub>4</sub> through multilayer  $\gamma$ -GY-*N* (*N* = 1–3) (MGN-*N*) at various temperatures and pressures.<sup>97</sup> The adsorption of H<sub>2</sub>S is much higher than that of CH<sub>4</sub> since its polarity results in strong interactions between the adsorbent and the adsorbate species. Under low pressure, MGN-1 exhibits the best performance for the adsorption of both the species, whereas at high pressure, adsorption of H<sub>2</sub>S is higher in case of MGN-3 due to the formation of molecular clusters. At low pressure, H<sub>2</sub>S selectivity increases but at higher values of pressure, the selectivity decreases. Increasing temperature decreases the selectivity of H<sub>2</sub>S due to the increased permeance of CH<sub>4</sub>.

<sup>3</sup>He is known to be a rare and highly precious gas, obtained as a byproduct during the radioactive decay of tritium. It has high demand in low-temperature research institutes and cryogenic industries.<sup>98</sup> It is heavily used in dilution refrigerators. <sup>3</sup>He also has applications in large neutron scattering facilities.<sup>99</sup> Recently, investigation of the quantum effects during the transmission of <sup>3</sup>He and <sup>4</sup>He through a single GDY sheet has demonstrated GDY to be a potential He isotope separator.<sup>100</sup> A comparison between the zero-point energy and tunneling effects depicts that both effects are highly relevant at low temperatures (20–30 K). The zero-point energy of <sup>3</sup>He is higher than that of <sup>4</sup>He, allowing the faster transmission of <sup>4</sup>He. At low temperature, tunneling effect is found to be slightly higher than the contribution from zero-point energy which favours the selective permeation of the lighter isotope. Bartolomei *et al.* employed quantum dynamics simulations for probing the efficiency of GDY and 2D-polyphenylene towards achieving He separation from natural gas.<sup>101</sup> Since the pore size of GDY is larger than the van der Waals diameter of He, it has better efficiency for the separation of He. Improved Lennard-Jones potential is used to describe the interactions between He atoms and the membrane. As the temperature increases, the selectivity for He/Ne, Ne/CH<sub>4</sub> and He/CH<sub>4</sub> decreases.

## 5 Graphynes for molecular adsorption and gas sensing

Gas sensors have a wide range of applications such as monitoring air pollution, medical diagnosis, leakage detection *etc.*<sup>102–104</sup> The efficiency of gas sensors depends on selectivity as well as sensitivity of the electronic properties of the substrates. Two-dimensional nanoporous materials such as graphene and GYs can be used for gas detection (Fig. 5). Due to uniformly distributed pores and high sensitivity, GYs have the potential to be used in gas sensing technology. The adsorption of gas

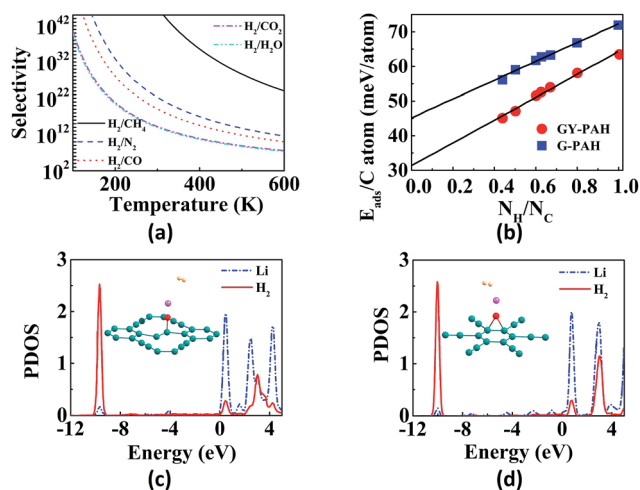


Fig. 5 Illustrations of separation, adsorption and sensing using GYs: (a) H<sub>2</sub> selectivity over H<sub>2</sub>O, CO<sub>2</sub>, N<sub>2</sub>, CO and CH<sub>4</sub> as a function of temperature for  $\gamma$ -GYN. (b) Adsorption energy per carbon atom as a function of H : C ratio of PAH molecules for binding on graphene (G-PAH) and  $\gamma$ -GY-1 (GY-PAH). Partial density of states (PDOS) of Li 2s and H<sub>2</sub> 1s orbitals suggesting the strong interaction between H<sub>2</sub> and Li-decorated oxidised GY at (c) T1 site and (d) B1 site. Part (a) is reproduced from ref. 93 with permission from the Elsevier. Part (b) is reproduced from ref. 110 with permission from the Wiley Periodicals, Inc. Parts (c) and (d) are reproduced from ref. 114 with permission from the AIP Publishing LLC.





molecules leads to charge transfer, which causes a change in conductance of the system. The sensitivity as well as the adsorption propensity can further be enhanced by metal decoration, defects and doping.

Pristine GYs are of great interest in the detection of various gases due to their unique properties. Formaldehyde is a colourless gas which causes environmental pollution and is even considered to be carcinogenic. Majidi *et al.* investigated whether GYs can be used for the detection of HCHO using DFT.<sup>105</sup> The small binding energy and large equilibrium distance imply the physisorption of HCHO on GYs ( $\alpha$ ,  $\beta$  and  $\gamma$ -GYs). To increase the efficiency of detection of HCHO, Chen *et al.* used Sc and Ti decorated GDY.<sup>106</sup> GDY, which is a metal, transforms into an n-type semiconductor upon the approach of HCHO. The binding energy for adsorption is higher in case of pristine GDY compared to pristine  $\gamma$ -GY-1 due to increased number of  $\pi$  electrons. It is found that decoration with Sc and Ti enhances the binding energy, with Sc resulting in higher binding strength. The change in the electronic property of  $\gamma$ -GY-1 upon adsorption of H<sub>2</sub>O<sub>2</sub> was studied by Majidi and co-workers.<sup>107</sup> The preferential orientation of H<sub>2</sub>O<sub>2</sub> is the configuration in which the O–O bond is perpendicular to the sheet and it gets adsorbed at the triangular pore. Increasing the number of H<sub>2</sub>O<sub>2</sub> molecules decreases the band gap and increases the conductivity and hence  $\gamma$ -GY-1 transforms to an n-type semiconductor. Deb *et al.* studied the adsorption of boron halides (BF<sub>3</sub>, BCl<sub>3</sub>, BI<sub>3</sub>) on  $\gamma$ -GY and monitored the change in electronic properties on adsorption.<sup>108</sup> It was found that the order of binding strength is BF<sub>3</sub> < BCl<sub>3</sub> < BI<sub>3</sub> for the adsorption process. The GYs transform to n-type semiconductors on the adsorption of BCl<sub>3</sub> and BI<sub>3</sub>, whereas in case of BF<sub>3</sub> adsorption, they transform as p-type semiconductors. Maximum electron transfer occurs in case of BI<sub>3</sub>, indicating that GY senses BI<sub>3</sub> much better than the other two halides. The dipole moment of  $\gamma$ -GY is high in the case of adsorption of BI<sub>3</sub> due to the strong electron transfer. Majidi *et al.* studied the electronic properties and the binding energies for the adsorption of tetracyanoethylene molecules on GYs.<sup>109</sup> The small binding energy and the large adsorption distance imply physisorption of molecules on the sheet.  $\alpha$  and  $\beta$ -GYs which are semi-metals transform to semiconductors due to the adsorption of tetracyanoethylene that accepts electrons from the sheets.

PAHs (polycyclic aromatic hydrocarbons) are formed during the incomplete ignition of organic molecules and many of them are considered to be toxic. Arriagada *et al.* investigated the adsorption of PAHs on graphene and GYs.<sup>110</sup> The binding energies for the adsorption of PAHs on GYs are low due to the porous structures and decreased  $\pi$  system of GYs compared to graphene, as shown in Fig. 5b. Smaller PAHs prefer stacked orientation due to increased  $\pi$ – $\pi$  interactions, whereas for larger PAHs, the preferred orientation is bridge or rotated. The electron transfer from PAHs to GYs decreases the band gap and hence increases the conductivity. Ozmaian *et al.* studied the diffusion and controlling motion of C<sub>60</sub> on  $\gamma$ -GY-*N* (*N* = 1–5) sheets.<sup>111</sup>  $\gamma$ -GY-1 has high binding energy due to high carbon atom density. The diffusion of C<sub>60</sub> on the sheets can be controlled by using various GYs. In case of  $\gamma$ -GY-1, mobile

assemblies are formed whereas stable assemblies are formed in case of  $\gamma$ -GY-2 and  $\gamma$ -GY-5.

Zhang *et al.* investigated the adsorption of gas molecules on  $\gamma$ -GY-1.<sup>112</sup> It is found that CO, NH<sub>3</sub> and SO<sub>2</sub> bind weakly, whereas H<sub>2</sub>S, NO<sub>2</sub> and NO bind strongly through charge transfer and change  $\gamma$ -GY-1 from semiconductor to metal. The strong binding of NO and NO<sub>2</sub> indicates that  $\gamma$ -GY-1 can be used for their detection. The competitive adsorption of CO<sub>2</sub> and H<sub>2</sub> on  $\gamma$ -GY-1 was studied by Kwon and co-workers using DFT.<sup>113</sup> The gas molecules can be adsorbed on three sites (i) triangular (ii) bridge and (iii) hexagonal. Considering the adsorption properties, both CO<sub>2</sub> and H<sub>2</sub> preferably adsorb on the triangular pores (hollow sites), however due to the electrostatic field caused by atomic charges on the sheet, the preference of binding is for CO<sub>2</sub>.

The effect of adsorption of O<sub>2</sub> on the electronic properties of GYs is theoretically studied by Kang and co-workers<sup>115</sup> and it was observed that the binding energy of O<sub>2</sub> with various GYs is in the order:  $\alpha$ -GY >  $\beta$ -GY >  $\gamma$ -GY. Adsorption of oxygen atoms over  $\alpha$  and  $\beta$ -GYs results in the formation of oxides of GYs opening up the band gap in these materials. Omidvar *et al.* investigated the enhancement of CO sensing ability of  $\gamma$ -GY upon metal decoration.<sup>116</sup> Pristine GY is not suitable for the sensing of CO due to the weak interaction between the sheet and CO, resulting in only slight variation in the electronic properties. However, there is considerable variation in the band gap of the system upon the adsorption of CO in case of metal-decorated GY. Beheshtian *et al.* explored the change in electronic properties upon the adsorption of HCN on pristine  $\gamma$ -GY-1 and Si-doped  $\gamma$ -GY-1.<sup>117</sup> Upon HCN adsorption on pristine GY, the decrease in band gap is less, implying weak adsorption thereby making  $\gamma$ -GY-1 unsuitable for HCN detection. The adsorption of HCN on Si-doped GY decreases the band gap due to the donation of electrons from the nitrogen atom to the LUMO of Si and hence increases the conductivity. The electronic properties of HCN/Si-doped  $\gamma$ -GY show that doped GYs can be used for HCN detection. Phosgene is a pale yellow poisonous gas used as a chemical weapon during World War II. Felegari *et al.* explored the effect of adsorption of phosgene on boron, nitrogen and Si-doped  $\gamma$ -GYs.<sup>118</sup> Phosgene molecule binds to electron deficient Si sites in such a way that the electron-rich oxygen atoms face the GY sheets. The charge transfer from phosgene to doped GY sheet increases its conductivity.

Peyghan *et al.* investigated the effect of doping  $\gamma$ -GYs with Ni and Si as the dopants on gas sensing.<sup>119</sup> The resulting GYs were used to study the variation in electronic properties upon the adsorption of NH<sub>3</sub>. In the case of pristine  $\gamma$ -GYs, the adsorption was very weak, which indicates small variation in electronic properties. Doped GYs (designated as MGY-*N* where M = Ni, Si; *N* = 1, 2) are obtained by replacing carbon atoms from hexagonal and triangular pores of  $\gamma$ -GYs with Ni/Si, respectively. The density of states (DOS) plot indicates that doping enhances the adsorption strength of NH<sub>3</sub> with maximum enhancement reported for Si-doped GYs. Lu *et al.* studied the effect of doping on the gas sensing ability of pristine  $\gamma$ -GY-1 using DFT calculations.<sup>120</sup> Except O<sub>2</sub>, all other gas molecules (CO, CH<sub>4</sub>, CO<sub>2</sub>, NH<sub>3</sub> and NO) undergo physisorption over  $\gamma$ -GY-1, causing negligible



change in band gap. Thus,  $\gamma$ -GY-1 cannot be used as a gas sensor for other gases. The Mn atom can get adsorbed at the triangular pore of  $\gamma$ -GY-1. The decrease in the band gap, due to the introduction of Mn results in enhancement of the gas sensing ability. The selectivity as well as sensitivity are enhanced by the doping of Mn over  $\gamma$ -GY, since the gas molecules are adsorbed *via* chemisorption.

## 6 Graphynes for energy storage

### 6.1 Hydrogen storage

Hydrogen is the cleanest fuel and is considered to be the future replacement for natural fuels. Currently, storing of this potential energy source is a challenge due to its low critical temperature.<sup>121</sup> Many research initiatives are directed towards exploring an efficient way for the storage and transportation of hydrogen. The techniques for the storage include liquefaction, compression, chemical hydrides and physisorption.<sup>121,122</sup> The liquefaction as well as compression do not guarantee safety since high pressure is required, whereas use of appropriate adsorption materials assure storage at low pressures.<sup>123</sup> GYs, which are 2D nanoporous materials with uniform pores and high pore densities can be used for hydrogen storage. Zhang *et al.* investigated the hydrogen storage capacity of bilayer  $\gamma$ -GY-2 using DFT calculations.<sup>129</sup> Bilayer  $\gamma$ -GY-2 has overlapping triangular pores that can trap H<sub>2</sub> molecules. The H<sub>2</sub> molecules, upon permeation through the first triangular pore get adsorbed in between the two layers in bilayers, resulting in hydrogen storage. The small barrier height for desorption can be easily surmounted, ensuring the unloading of trapped hydrogen molecules. The limited hydrogen storage capacity of pristine GYs can be enhanced through metal decoration. Metal-decorated GYs are widely reported, but the challenge in fabrication is the higher cohesive energies of metals compared to the binding energies of metals with GYs.

Li-decorated GYs have been used as storage media for hydrogen in numerous studies. A comparative study of the hydrogen storage capacities of  $\gamma$ -GY-1 and  $\gamma$ -GY-2 was performed using the first-principles calculations by Srinivasu and co-workers.<sup>124</sup> The first Li atom adsorbed preferably occupies triangular pores by interacting with the three acetylenic bonds and the binding strength is higher for  $\gamma$ -GY-1 compared to  $\gamma$ -GY-2. As the number of hydrogen molecules adsorbed per Li increases, interaction energy decreases as shown in Fig. 6a. The maximum number of H<sub>2</sub> molecules adsorbed per cationic Li site is three. Even though the adsorption energy is higher in case of  $\gamma$ -GY-2 than  $\gamma$ -GY-1, the adsorption capacity is in reverse order since the number of carbon atoms per unit cell is higher for  $\gamma$ -GY-2. First-principles study of the hydrogen storage in Li-decorated  $\gamma$ -GY-1 was also performed by Guo and co-workers.<sup>130</sup> The  $\gamma$ -GY-1 system is found to have high binding energy with Li due to the low ionization potential of Li which donates its s electron to the GY sheet. The hydrogen storage capacity is found to be 18.6 wt%, and each Li adsorbs four H<sub>2</sub> molecules. The interaction between H<sub>2</sub> and sheet is due to the electric field induced by the ionic Li. The adsorption of the fourth H<sub>2</sub> molecule is due to the charge transfer through the

polarized H<sub>2</sub> molecules that are already adsorbed. In a study carried out by Zhang *et al.*, it is found that double-sided Li-decorated GY adsorbs seven molecules per Li atom with a capacity of 15.15 wt%, wherein the adsorption is attributed to the polarization interaction mechanism.<sup>131</sup> Lu and co-workers used Li-decorated 6,6,12-GY, which consists of rhombus-like acetylenic (SA) rings apart from triangular acetylenic (TA) rings for the hydrogen storage.<sup>132</sup> Li occupies the centre of SA rings since number of sp hybridized carbons is more in SA rings. The adsorption of hydrogen is stronger at the TA site even though the electron donation of Li to 6,6,12-GY is facile at the SA site. It is found that storage capacity increases to 19.3 wt% upon Li decoration, where the enhanced adsorption is due to the electric field produced by the polarized H<sub>2</sub> molecules. Kumar *et al.* used Li-decorated GY networks for hydrogen storage.<sup>133</sup> Each GY linker binds two Li atoms forming Li-GY complexes. It was found that each Li atom adsorbs three H<sub>2</sub> molecules weakly through the polarisation of Li forming metal-GY framework-Li<sub>8</sub>-24H<sub>2</sub> with a storage capacity of 6.4 wt%. The adsorption energy decreases with increasing number of H<sub>2</sub>. Since adsorption of H<sub>2</sub> is weak, increasing the temperature leads to desorption. Yan *et al.* used Li-decorated oxidised GY (obtained by oxidation of the acetylenic bond and binding Li with the oxygen atom) as a hydrogen sensor.<sup>114</sup> The oxygen atom forms (i) epoxy bond with two C atoms of hexagonal pore (B1 site) or triangular pore (B2 site), (ii) bond with one C atom of triangular pore (T1 site) or hexagonal pore (T2 site). It is found that Li atom binds strongly to T1 site followed by B1 site. Partial density of states (PDOS) indicates the overlap of the 2s orbital of Li with the  $\sigma$  bond of H<sub>2</sub> in both sites as shown in Fig. 5c and d. Each Li atom adsorbs four or five H<sub>2</sub> molecules with a storage capacity of 12.03 wt%. Lu *et al.* investigated the storage capacity of metal-decorated boron-doped  $\gamma$ -GY.<sup>134</sup> The storage capacity is low for pristine GYs due to weak interaction between H<sub>2</sub> and sheet, which can be enhanced by the adsorption of Li on GY sheets. Since the cohesive energy of Li is higher than the interaction energy, clustering occurs resulting in low adsorption. This can be solved by boron doping of the  $\gamma$ -GY sheet which transforms it to a p-type semiconductor. The binding of H<sub>2</sub> with the hexagonal pore is favoured since at the triangular pore, strong charge transfer from Li to the GY sheet occurs, resulting in low adsorption energy.

Apart from Li, other alkali metals and alkaline earth metals are also used for hydrogen storage. Liu *et al.* used Na-decorated single-sided and double-sided  $\gamma$ -GY and BN-yne (GY doped with BN) for the storage of hydrogen by DFT calculations.<sup>138</sup> Decorating GY with Na leads to the transfer of s electrons from Na to the sheets. In case of Na-decorated single-sided and double-sided GY, each Na atom adsorbs three hydrogen molecules and each of them has a capacity of 3.49 and 5.98 wt%, respectively. The BN-yne decorated with Na on both sides also adsorbs three hydrogens per Na atom and each of them has a capacity of 5.48 wt%. Hwang *et al.* studied the storage of H<sub>2</sub> molecules in Ca-decorated  $\alpha$ ,  $\beta$ ,  $\gamma$ -GYs.<sup>139</sup> It was found that Ca occupies hexagonal pores in  $\alpha$  and triangular pores in  $\beta$  and  $\gamma$ -GYs. As the number of H<sub>2</sub> increases, adsorption energy decreases in all these cases. They found that transition metals undergo



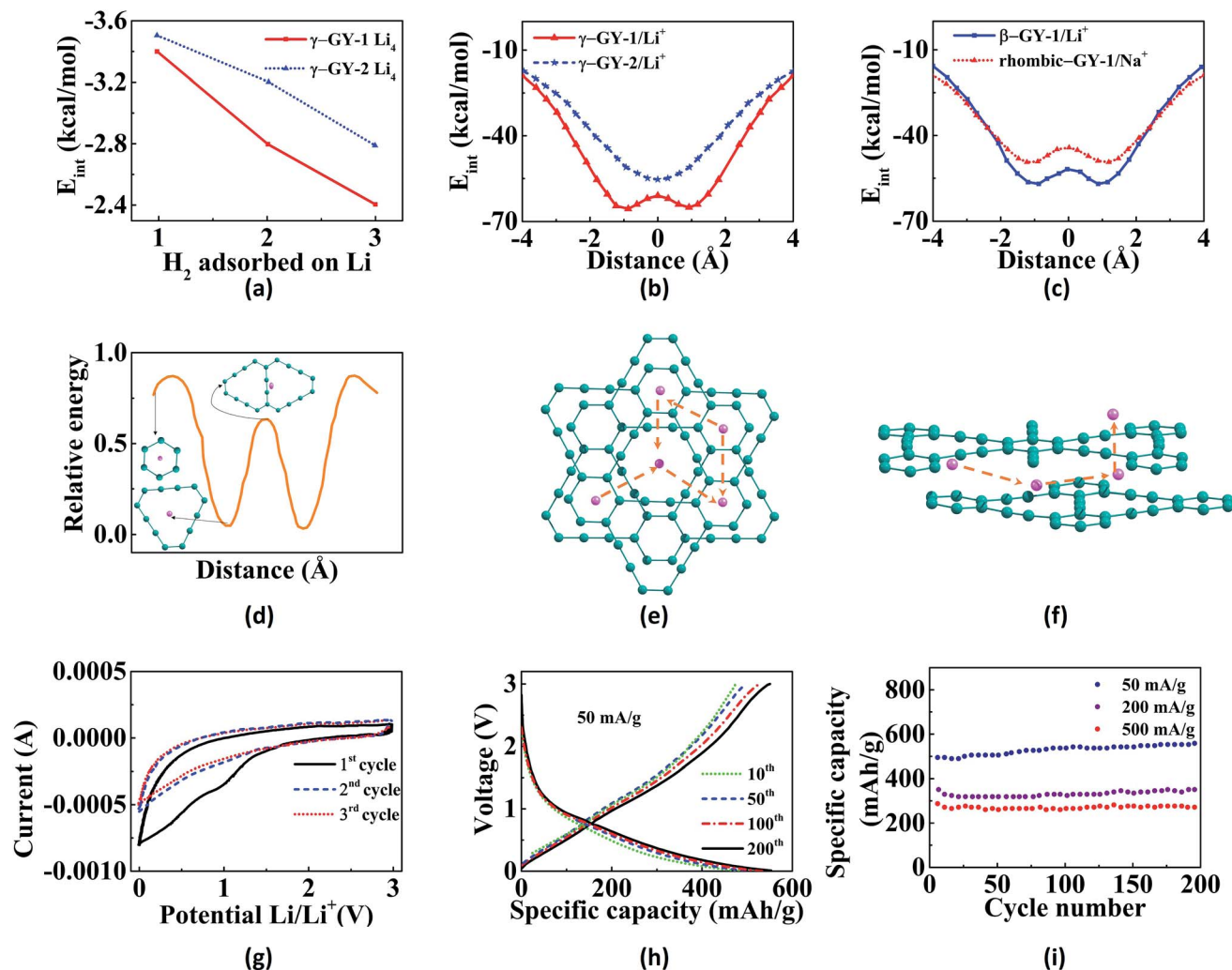


Fig. 6 Applications of GYs in energy storage: (a) variation in interaction energy with increase in number of  $\text{H}_2$  adsorbed per Li. Potential energy scans for the passage of (b)  $\text{Li}^+$  through  $\gamma$ -GY-1 and  $\gamma$ -GY-2 and (c)  $\text{Li}^+$  and  $\text{Na}^+$  through  $\beta$ -GY-1 (triangular pore) and rhombic-GY-1, respectively. (d) Diffusion pathway of Li atom as it moves through different adsorption sites on  $\gamma$ -GY-1. Schematics of the (e) in-plane and (f) out-of-plane diffusion of Li through various sites of bulk  $\gamma$ -GY-1. (g) Cyclic voltammetric curves at initial three cycles with scan rate as  $1 \text{ mV s}^{-1}$ . (h) galvanostatic charge/discharge profiles with current density as  $50 \text{ mA g}^{-1}$  and (i) cycle performance of lithium-ion batteries with GDY as the anode. Part (a) is reproduced from ref. 124, parts (b) and (c) are reproduced from ref. 125 and parts (d–f) are reproduced from ref. 126 with permissions from the American Chemical Society. Part (g) is reproduced from ref. 127 with permission from the Elsevier. Parts (h) and (i) are reproduced from ref. 128 with permission from the Royal Society of Chemistry.

clustering on graphyne sheet, implying that they would be poor choice for metal decoration. Li *et al.* also found that triangular pores are stable adsorption sites for Ca.<sup>140</sup> The hybridization of the  $\sigma$  states of  $\text{H}_2$  and 3d orbitals of Ca indicates weak adsorption. In case of Ca-decorated single-sided and double sided  $\gamma$ -GY-1, six  $\text{H}_2$  molecules can be trapped per Ca atom. The storage capacity of Ca-decorated double-sided graphyne was found to be 9.6 wt%. Guo *et al.* investigated the hydrogen storage capacity of different metal-decorated GYs.<sup>141</sup> The metal adsorption is found to be preferable at triangular pores of  $\gamma$ -GYs. The metal donates  $s$  electrons to GY sheet and GY back-donates electrons to 3d orbitals of metal. Sc and Ti bind strongly resulting in higher adsorption strength when compared to other metal-decorated GYs. The hydrogen storage capacity is found to be higher for Li and Ca-decorated GYs.

GNTs are also used to study the trapping of  $\text{H}_2$ . First-principles calculations were employed to study the storage of  $\text{H}_2$  in Ca-decorated GNTs.<sup>142</sup> Since the binding energy of Ca with GNT is higher than the cohesive energy of Ca and the Ca–Ca distance is large, clustering of Ca on GNTs is ruled out. The hydrogen molecules bind to Ca due to the polarization and hybridization between the  $\sigma$  state of  $\text{H}_2$  and 3d orbitals of Ca. The hydrogen storage capacity was found to be 7.47 wt%. Due to curvature effect, the metal atoms bind strongly with GNTs preventing clustering.<sup>143</sup>

The application of external field prevents metal aggregation and enhances the hydrogen storage capacity. Employing DFT calculations, Liu *et al.* explored the enhancement of hydrogen storage of Mg-decorated  $\gamma$ -GY in presence of applied electric field.<sup>144</sup> The binding energy exceeds cohesive energy on applying



**Table 2** The binding energies, adsorption heights and barrier heights (for the out-of-plane diffusion) for various metal ions on graphyne sheets. Reproduced from ref. 125 with permission from the American Chemical Society

Graphynes	Metal ions	Binding energy (kcal mol <sup>-1</sup> )	Adsorption height (Å)	Barrier height (kcal mol <sup>-1</sup> )
γ-GY-1	Li <sup>+</sup>	64.22	1.10	3.54
	Na <sup>+</sup>	46.36	1.70	35.60
γ-GY-2	Li <sup>+</sup>	58.01	0.0	0.0
	Na <sup>+</sup>	50.14	0.0	0.0
α-GY-1	Li <sup>+</sup>	49.42	0.0	0.0
	Na <sup>+</sup>	40.10	0.0	0.0
β-GY-1 (hexagonal pore)	Li <sup>+</sup>	53.62	0.0	0.0
	Na <sup>+</sup>	45.34	0.0	0.0
Rhombic-GY-1	Li <sup>+</sup>	62.57	0.0	0.0
	Na <sup>+</sup>	45.71	1.01	1.96

an external electric field, which prevents the clustering of Mg. The polarization and hybridization between H<sub>2</sub> and Mg lead to strong adsorption. Zhang *et al.* studied the hydrogen storage in Ti-decorated GY in presence of external electric field of varying strengths by first-principles calculations.<sup>145</sup> The Ti atom donates electron to the γ-GY sheet, which in turn back-donates electrons to the 3d orbital of Ti. The adsorption of H<sub>2</sub> is due to electrostatic interactions between the sheet and H<sub>2</sub> molecules, which is the result of polarization effect caused by the applied electric field.

## 6.2 Lithium-ion batteries

Graphite rods and graphene sheets are extensively explored materials as anodes for lithium-ion batteries. Lithium ions can get intercalated in between the sheets in graphite. Graphyne sheets containing both sp<sup>2</sup> and sp hybridized carbons can potentially adsorb more metal atoms. The porous nature of GYs provides additional tools to finetune their properties thereby making them better anodic materials in batteries. γ-GY has two types of adsorption sites for metal atoms. Site with the smaller pore (site A) is made up of six-membered phenyl ring and the site with the larger pore (site B) is made up of twelve-membered ring composed of both sp<sup>2</sup> and sp hybridized carbons. Compared to site A, lithium is adsorbed strongly over site B. Binding of the first two lithium atoms on a unit cell of γ-GY-1 occurs in-plane to the carbon surface and they occupy the two

B sites.<sup>124</sup> Further addition of metal atoms results in occupation of positions above and below the C<sub>6</sub> rings. Binding energies for first lithium atom are estimated to be 64.5 kcal mol<sup>-1</sup> and 61.9 kcal mol<sup>-1</sup> for γ-GY-1 and γ-GY-2, respectively, for binding at site B.<sup>124</sup> γ-GY-1 nanopores can only adsorb one atom in site B, while γ-GY-2 pores can adsorb up to three lithium atoms.<sup>146</sup>

In graphite and multilayer graphenes, diffusion of Li is confined to the interlayer spacing.<sup>147</sup> However, both in-plane and out-of-plane diffusion of ions are allowed in intrinsically nanoporous graphyne membranes (Fig. 6).<sup>126,146</sup> γ-GY-1 and γ-GY-2 provide high in-plane mobility and the barriers for hopping between adsorption sites are less than 20 kcal mol<sup>-1</sup> as shown in Fig. 6d and e. The barrier for the out-of-plane diffusion of Li ion through the C<sub>6</sub> ring is estimated to be 190 kcal mol<sup>-1</sup>. Thus, out-of-plane diffusion through C<sub>6</sub> ring is not possible under normal conditions. The larger triangular pore in γ-GY-1 provides a barrier of approximately 4 kcal mol<sup>-1</sup>, allowing fast diffusion (Fig. 6). Higher mobility of metal ions helps in achieving higher current density as well as complete reversibility during multiple charging–discharging cycles. Chandra Shekar *et al.* studied the rattling motion or out-of-plane mobility of different alkali and alkaline-earth metal ions through various GY sheets.<sup>125,148</sup> The binding energies, adsorption heights and barrier heights for passage through the pores of various GYs for lithium and sodium are provided in Table 2. Potential energy curves for the out-of-plane rattling motion of

**Table 3** Theoretically predicted maximum capacities and lithiation potentials for different anodic materials in lithium-ion batteries

System	Composition	Specific capacity (mA h g <sup>-1</sup> )	Volumetric capacity (mA h cm <sup>-3</sup> )	Lithiation potential (V)
Graphite	C <sub>6</sub> Li	372 <sup>a,b</sup>	818 <sup>b</sup>	~0.5 <sup>c</sup>
γ-GY-1	C <sub>3</sub> Li	623.4 <sup>a</sup> (1117 <sup>c</sup> )	1589 <sup>b</sup>	2.8–1.7 <sup>a</sup>
γ-GY-2	C <sub>2,25</sub> Li	788.4 <sup>a</sup>	—	2.7–1.9 <sup>a</sup>
α-GY-1	C <sub>6</sub> Li <sub>3</sub>	1117 <sup>c</sup>	1364 <sup>c</sup>	—
α-GY-2	C <sub>6</sub> Li <sub>7,31</sub>	2719 <sup>c</sup>	2032 <sup>c</sup>	0.85–0.35 <sup>c</sup>
Rhombic-GY-2	C <sub>5</sub> Li <sub>6</sub>	2680 <sup>d</sup>	—	0.75–0.23 <sup>d</sup>
Boron-doped γ-GY-1 (3B)	—	1144 <sup>e</sup>	—	4.0–2.97 <sup>e</sup>
Boron-doped γ-GY-1 (1B)	—	1125 <sup>e</sup>	—	3.5–2.66 <sup>e</sup>

<sup>a</sup> Values are taken from ref. 124. <sup>b</sup> Values are taken from ref. 135. <sup>c</sup> Values are taken from ref. 136. <sup>d</sup> Values are taken from ref. 137. <sup>e</sup> Values are taken from ref. 134.



metal ions through various GYs are shown in Fig. 6b and c. Lower diffusion barriers for metal ions through GYs help in fast charging of battery. However, in the case of discharging, the process becomes difficult energetically and requires application of a voltage across anode and cathode.<sup>126</sup> Adsorption of lithium atoms on GY is followed by charge transfer from the metal to the sheet. These electrons occupy the conduction band and result in electrostatic interaction between metal adatoms and GY sheet. Zhang *et al.* attribute the site preference to the charge transfer from Li to the larger triangular pore since the conduction band lies along this  $C_{12}$  ring.<sup>126</sup>

Adsorption sites of lithium on  $\alpha$ -GY are in-plane and slightly off-centre towards the edge of hexagon.<sup>135</sup> Lithium adsorption over  $\alpha$ -GY results in in-plane distortion of carbon atoms in the sheet. It is predicted that lithium can be dispersed on  $\alpha$ -GY-2 to give a maximum composition of up to  $C_6Li_{1,18}$  and corresponding specific capacity is  $4259 \text{ mA h g}^{-1}$ . One of the important factors affecting the applicability of GYs as anode materials is the structural stability of sheets during intercalation and de-intercalation. Calculated cohesive energies for  $\alpha$ -GY and  $\gamma$ -GY are  $189.1$  and  $196 \text{ kcal mol}^{-1}$ , respectively, while the cohesive energy for graphite is  $209.9 \text{ kcal mol}^{-1}$ .<sup>136</sup> Detailed study of structural stability and phase transformation (from GY to graphite) is required to address this issue.<sup>136</sup> Theoretical predictions for boron-doped  $\gamma$ -GY were made by Lu *et al.* for the storage of lithium and dihydrogen.<sup>134</sup> Calculations were carried out using Grimme DFT-D2 dispersion correction method. Up to three boron atoms were doped at the phenyl ring of  $\gamma$ -GY (1BG, 2BG, 3BG respectively) and the adsorption energy of lithium adatom increased with the number of borons. The lithiation potentials ( $\text{Li}/\text{Li}^+$ ) were  $3.50$ ,  $2.66$ ,  $4.00$  and  $2.97 \text{ V}$  in  $1\text{Li}@1\text{BG}$  (one lithium adatom adsorbed on 1BG),  $6\text{Li}@1\text{BG}$ ,  $1\text{Li}@3\text{BG}$  and  $6\text{Li}@3\text{BG}$ , respectively, while lithiation potentials in graphite and  $4\text{Li}-\text{GY}$  are  $0.5$  and  $1.7 \text{ V}$ , respectively. Specific capacity of lithium battery is predicted to be  $1125 \text{ V}$  for  $6\text{Li}@1\text{BG}$  and increases nominally on further doping. The theoretically predicted maximum energy storage capacities and lithiation potentials for different types of graphynes as anodic material have been summarised in Table 3.

Xu *et al.* have proposed the usage of GYs as anode materials in sodium-ion battery.<sup>149</sup> Storage capacity of  $\gamma$ -GY-1 and  $\gamma$ -GY-2 was predicted to be  $\text{NaC}_4$  and  $\text{NaC}_3$ , respectively. One sodium can bind with the hexagonal  $C_{12}$  ring in  $\gamma$ -GY-1 sheets, while three atoms can bind in the  $C_{18}$  ring of  $\gamma$ -GY-2. However, the binding of three atoms results in blockage of out-of-plane diffusion through the pore. Sodium atom can also bind to the  $\gamma$ -GY sheets by positioning above the  $C_6$  ring. However, combined experimental and theoretical study carried out by Zhang *et al.* has shown that sodium ions bind only to the  $C_{18}$  ring in well-stacked  $\gamma$ -GY-2 sheets except at defective sites in stacking.<sup>150</sup> Observed capacity of GDY-based sodium-ion battery is  $261 \text{ mA h g}^{-1}$  at  $50 \text{ mA g}^{-1}$  current density.

Experimental studies on the lithium-ion batteries have proved the superior quality of  $\gamma$ -GY-2 as an anode material.<sup>151,152</sup> Wang *et al.* have synthesised both lithium-ion batteries and capacitors with excellent cyclic stability and good capacitance or energy storage performance.<sup>127</sup> They observed reversible

capacities of  $908 \text{ mA h g}^{-1}$  and  $526 \text{ mA h g}^{-1}$  at  $0.05 \text{ A g}^{-1}$  and  $1 \text{ A g}^{-1}$  current densities, respectively. Specific capacitance observed was  $208 \text{ F A g}^{-1}$  at  $1 \text{ A g}^{-1}$  current density. Cyclic voltammetric curves (Fig. 6g) do not show any distinguishable reduction and oxidation peaks, except in the first cycle. Low initial coulombic efficiency and irreversibility in the first cathodic scan as shown in cyclic voltammogram trace is due to formation of a stable solid electrolyte interface film at the electrode.<sup>127,128,152</sup> Due to the unique porous nature of GYs, the charging and discharging cycles are almost reversible after few initial cycles.<sup>128</sup> As depicted in Fig. 6h, the charge and discharge capacities in Li-ion batteries with  $\gamma$ -GY-2 anodes remain almost constant after  $10^{\text{th}}$  cycle. Similarly, a high specific capacity is retained over many cycles reversibly (Fig. 6i).  $\gamma$ -GY-2 based battery achieved an energy density of around  $100 \text{ W h kg}^{-1}$  at both low and high power densities.<sup>152</sup> The battery also showed 94.7% retention of energy density after 1000 cycles. Hydrogen-substituted  $\gamma$ -GY-2 is fabricated using coupling of triethynylbenzene. The flexible electrode based on hydrogen-substituted GDY is reported to have a specific capacity of  $650 \text{ mA h g}^{-1}$  after 100 cycles. Nitrogen-doped  $\gamma$ -GY-2 (N-GDY) based lithium-ion batteries were synthesised by Yang *et al.* and Zhang and co-workers.<sup>153,154</sup> Nitrogen doping was carried out by Zhang *et al.* using the prolonged exposure of ammonia at  $873 \text{ K}$ . While, Yang *et al.* fabricated N-GDY using derivatized pyrimidine and pyridine monomers. Anodes based on N-GDY showed a better electrochemical performance compared to the pristine ones. Yang and co-workers observed the specific capacities of batteries to be  $1168$  and  $1165 \text{ mA h g}^{-1}$ , respectively for anodes with pyridine and pyrimidine-based GDYs at a current density of  $100 \text{ mA g}^{-1}$ .

## 7 Graphynes for catalysis

We now review some of the applications of GYs in the area of catalysis. First-principles calculations have been performed on various pristine and doped GYs to design optimal catalysts for reactions such as oxygen reduction reaction, oxygen evolution reaction, CO oxidation *etc.* Experiments along these lines are only beginning to emerge. The rate of an oxygen reduction reaction (ORR) in fuel cells is usually improved by using Pt-based catalysts. However, there are several limitations for Pt catalysts and hence other alternatives are widely being researched.<sup>155,156</sup> Out of these, carbon-based catalysts are of greater interest because of their ease of availability. Doped graphene and CNTs have shown good ORR electrocatalytic activity.<sup>157-160</sup> The improved catalytic activity is due to the positive charges on the doped systems. Since pristine graphene has no such positive charges, it does not show any catalytic activity. In contrast to graphene, GYs have uneven charge distribution, making them an excellent choice for the catalysis of ORR.<sup>161</sup> GYs are expected to show better catalytic activity than doped graphene and CNTs because they have a large number of positively charged sites for catalysis. Using DFT calculations, Wu *et al.* reported the electrocatalytic activity of  $\alpha$ -GY-1 for ORRs in acidic fuel cells.<sup>161</sup> In an acidic fuel cell, there are two possible pathways for the first electron transfer in the ORR process. One is the



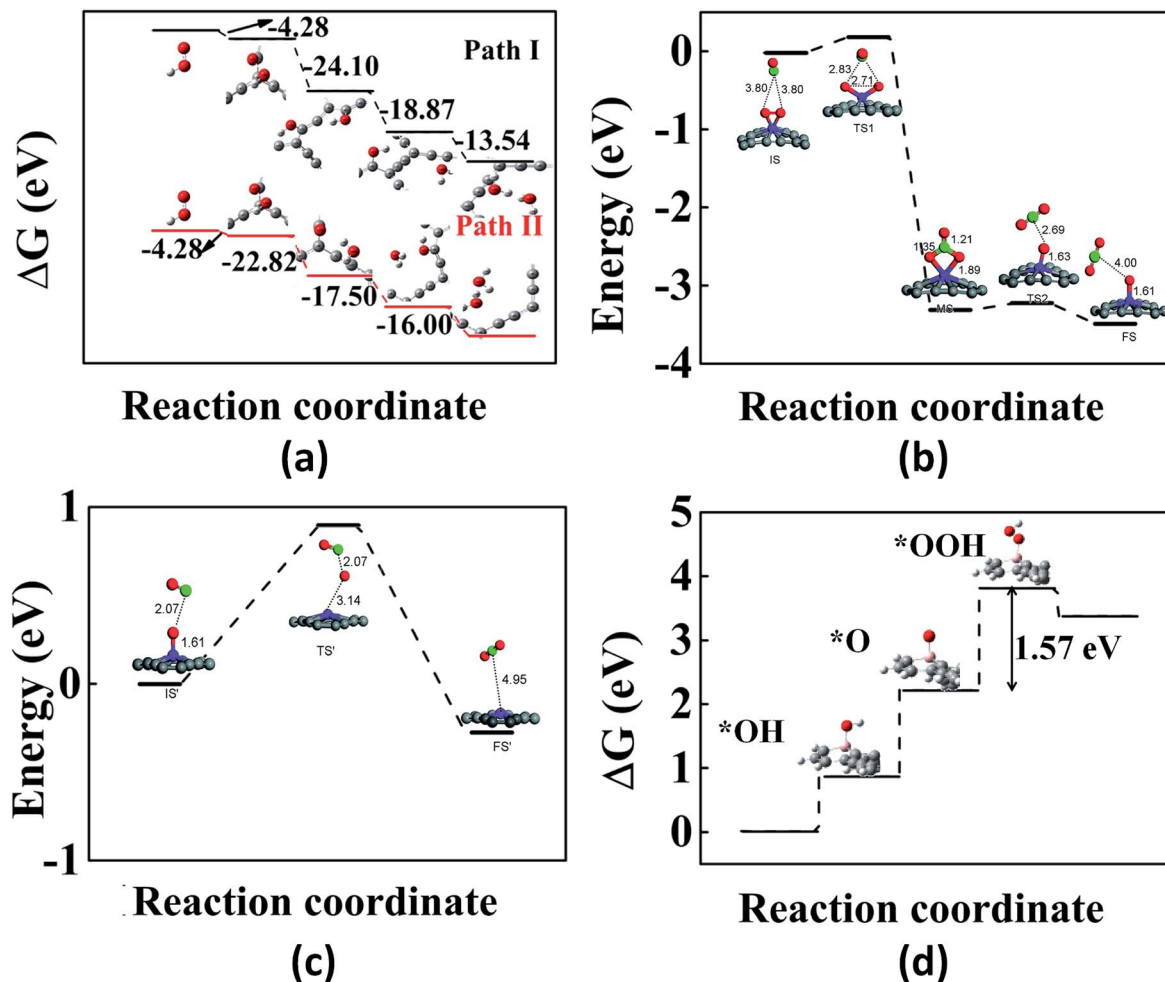


Fig. 7 Illustrations of GY-catalyzed reactions: (a) energy differences between reaction intermediates (*via* path I and path II) of the  $\alpha$ -GY-catalyzed ORR. (b and c) The first and second steps of CO oxidation catalyzed by  $\gamma$ -GY-supported Fe atom. (d) Free energy plot for the OER on B-doped  $\gamma$ -GY. Part (a) is reproduced from ref. 161 with permission from the American Chemical Society. Parts (b) and (c) are reproduced from ref. 162 with permissions from the Royal Society of Chemistry. Part (d) is reproduced from ref. 163 with permission from the Elsevier.

adsorption of  $\text{OOH}^+$  on  $\alpha$ -GY-1 and the other is the direct adsorption of  $\text{O}_2$  on  $\alpha$ -GY-1, which then reacts with an  $\text{H}^+$  to form  $\text{OOH}^+$ . Adsorption of  $\text{OOH}^+$  is a more favourable process than that of  $\text{O}_2$ . The next  $\text{H}^+$  may move to the O atom bound to the C atom (path I) or the oxygen bound to the H atom (path II). Fig. 7a shows the adsorption energy differences between the reactants and the products for each step and as we can see from the figure, path I is more favourable compared to path II. Thus,  $\alpha$ -GY-1 proves to be an excellent catalyst for the ORR process that happens *via* a 4-electron pathway, which includes an O–O bond breakage and formation of a  $\text{H}_2\text{O}$  molecule.

Incorporation of transition metals (TM) on GYs changes the electronic structure and magnetic properties in such a way that the ORR catalytic activity is improved. Srinivasu *et al.* studied Fe, Co and Ni decorated GYs as catalysts for ORR.<sup>164</sup> DFT studies showed that the most favourable position for metal binding is above the triangular pores of  $\gamma$ -GY-1.  $\gamma$ -GY-TM systems have higher binding energies than the corresponding graphene-TM systems due to the presence of sp hybridized carbon atoms. The most preferred site for  $\text{O}_2$  adsorption is just above the TM.  $\text{O}_2$

adsorbs *via* the side-on mode in  $\gamma$ -GY-Fe case and *via* the end-on mode for Ni and Co systems. The order of reactivity of the metals is given by  $\text{Fe} > \text{Co} > \text{Ni}$ , indicating that  $\gamma$ -GY-Fe is the best ORR catalyst among these systems. In acid medium,  $\gamma$ -GY-TM systems proceed through a 4-electron process and the  $\gamma$ -GY-Co system shows a higher over-potential than  $\gamma$ -GY-Fe, making the  $\gamma$ -GY-Fe system a better catalyst. In alkaline medium,  $\gamma$ -GY-Co proceeds through a 2-electron mechanism, whereas the  $\gamma$ -GY-Fe system proceeds *via* a 4-electron mechanism. Similar to acid medium,  $\gamma$ -GY-Fe is a better catalyst than  $\gamma$ -GY-Co in alkaline medium too. From the computed free energy profiles, it was concluded that  $\gamma$ -GY-Fe shows enhanced catalytic efficiency in the alkaline medium compared to the acid medium.

Metal-free ORR catalysts with good stability and tolerance to CO and methanol are of growing interest and the heteroatom-doped GYs (excluding transition metal atoms as the heteroatoms) seem to exhibit these desired properties due to the electronegativity differences between C atoms and the heteroatoms. Kong *et al.* investigated the effect of B doping on  $\gamma$ -GY-1 for its use as a catalyst.<sup>165</sup> B doping changes the electronic



structure of  $\gamma$ -GY-1 and introduces positive charge density on the sheet which enhances the catalytic ability. Using first-principles study, Chen *et al.* studied the mechanism of ORR catalysis by B, N single-doped  $\alpha$ -GY-1,  $\gamma$ -GY-1 and co-doped  $\alpha$ -GY-1.<sup>166</sup> B and N single-doped  $\alpha$ -GY-1 showed low catalytic efficiency due to high over-potential in the reduction steps, whereas the single-doped  $\gamma$ -GY-1 showed much better catalytic efficiency. Thus, the configuration of the doped GY also plays an important role in deciding the catalytic ability. Initially, the authors studied three co-doped systems –  $\alpha$ -B<sub>1</sub>N<sub>2</sub>G,  $\alpha$ -B<sub>1</sub>N<sub>3</sub>G,  $\alpha$ -B<sub>1</sub>N<sub>4</sub>G and except for  $\alpha$ -B<sub>1</sub>N<sub>2</sub>G, all steps for all the systems had positive reversible potential. These systems did not show much improvement in catalytic activity and therefore  $\alpha$ -B<sub>1</sub>(N<sub>3</sub>)<sub>4</sub>G with increased N doping was studied. It showed a higher efficiency than  $\alpha$ -B<sub>1</sub>N<sub>4</sub>G due to improved onset potential. N-doped  $\beta$ -GY-1 could also efficiently catalyze ORR *via* a 4-electron process.<sup>167</sup> The  $\beta$ -GY-1 with an sp<sup>2</sup> C replaced by N showed higher efficiency than the GY with an sp<sup>2</sup> C replaced by N because of a higher positive charge created in the former case. Das and co-workers compared the influence of B doping on  $\alpha$ -GY,  $\beta$ -GY,  $\gamma$ -GY,  $\delta$ -GY, 6,6,12-GY, and rhombic-GY using first-principles calculations.<sup>168</sup> All systems show a 4-electron pathway for ORR, but only the Rsp<sup>2</sup>-rhombic system (Rsp<sup>2</sup> represents replacing sp<sup>2</sup> C at the corner of the rectangular or square ring of rhombic-GY with the dopant) showed a monotonically exothermic ORR pathway. All other systems have an endothermic step in between. Thus, it is the B-doped rhombic-GY that is of greater interest for future ORR applications. N-doped GDY shows a catalytic efficiency comparable to the Pt/C system with improved stability and tolerance to methanol.<sup>169</sup> The imine N enhanced the positive effect on the sheet to a larger extent than the pyridinic N atom. Recently, Lv *et al.* treated the N-doped GDY to high temperature to get porous N-doped GDY with superior catalytic activity.<sup>170</sup> N-GDY was further calcined in an ammonia atmosphere to increase the nitrogen content. The N,N'-GDY thus prepared showed an even higher catalytic activity. The effect of co-doping GDY with multi-elements for improving the ORR rate was investigated by Zhang and co-workers.<sup>171</sup> The GDY was co-doped with N and S, N and B, N and F successfully and out of these, N and F co-doped GDY showed the best catalytic performance with good stability and improved CO and methanol tolerance.

Another interesting and emerging class of carbon-based catalysts for ORR is GY nanotubes. Recently, Chen studied the effect of the doped  $\gamma$ -GNT on ORR catalytic efficiency.<sup>172</sup> The GNTs were doped with N, B, P and S atoms and their catalytic activities were studied using DFT. S-GNT is a poor ORR catalyst and the efficiencies of the other three catalysts follow the order: N-GNT > B-GNT > P-GNT. The O<sub>2</sub> molecule is adsorbed in the end-on mode for N and B-GNTs, whereas ORR is initiated by the bridge adsorption of O<sub>2</sub> in P-GNT.

GDYs also act as metal-free catalysts for CO oxidation.<sup>173</sup> It is again the positively charged C atoms that facilitate the CO oxidation *via* the Eley-Rideal (ER) mechanism. The adsorption energy of O<sub>2</sub> is higher than that of CO and even in a mixture of CO and O<sub>2</sub>, the GDY system will be majorly covered by O<sub>2</sub> because of its higher binding energy. Initially, the O<sub>2</sub> adsorbs in a parallel fashion to the GDY (side-on mode). Next, there is an

O–O bond cleavage with a carbonate-like intermediate formation which is finally followed by a CO<sub>2</sub> formation. CO oxidation is also catalyzed by GY-supported single Fe atom.<sup>162</sup> The non-uniform charge distribution favours the adsorption of single Fe atom on  $\gamma$ -GY-1 and results in a better distribution of the metal catalyst unlike in the case of pristine graphene, where the metal atoms move to the edges. The metal atom shows preferential adsorption over the hollow site above the centre of the acetylenic ring. As in the previous case, O<sub>2</sub> adsorption is preferred over CO and the CO oxidation proceeds *via* the ER mechanism as shown in Fig. 7b and c.

In a similar manner, GDY acts as an excellent support for Co nanoparticles to catalyze oxygen evolution reaction (OER).<sup>174</sup> A support material like GY helps in preventing the decrease in catalytic efficiency caused by the aggregation of the nanoparticles. Li *et al.* successfully synthesised a 3D Cu@GDY/Co electrode that shows good OER catalytic activity. The bare Cu foam or the Cu foam with GDY does not show any catalytic activity on its own. It is the combination of Cu foam, GDY and Co nanoparticle that shows high catalytic activity. The GDY layer supported on the Cu foam improves the conductivity and prevents the aggregation of the Co catalyst. Using DFT calculations, it was shown that 2D B-doped  $\gamma$ -GY-1 is a promising metal-free catalyst for OER (Fig. 7d).<sup>163</sup> The effect of N and B doping at different sites of  $\gamma$ -GY-1 on the catalytic activity was investigated and it was found that B doping of  $\gamma$ -GY-1 resulted in better catalytic performance than N doping and within B doping, replacing sp<sup>2</sup> carbons by B atoms gives lowest over-potential and hence results in better OER catalysts.

GYs were recently investigated as potential candidates to support Pt for the oxidation of ethanol.<sup>175</sup> The binding energies of Pt on  $\gamma$ -GY,  $\beta$ -GY, GDY and graphene were compared and they follow the order: graphene < GDY <  $\gamma$ -GY <  $\beta$ -GY. Binding on GYs leads to higher binding energies compared to pristine graphene due to more reactive sp carbon atoms present in these systems. Adsorption of a Pt cluster (Pt<sub>4</sub>) was also studied and in this case the binding energy followed a different order: graphene <  $\beta$ -GY <  $\gamma$ -GY < GDY. However, Pt-supported GYs did not show any remarkable enhancement of ethanol oxidation.

GYs also show interesting catalytic properties for the dehydrogenation of light metal complex hydrides.<sup>176</sup> The decrease in the energy required for the removal of H atoms from three systems, namely LiAlH<sub>4</sub>, LiBH<sub>4</sub> and NaAlH<sub>4</sub> was studied using DFT calculations. All the three systems showed positive binding energies with GDY and the greatest stabilization on hydrogen removal was achieved by LiBH<sub>x</sub> (x = 2, 3 and 4). The catalytic enhancement observed on using GDY can be ascribed to three major reasons: (i) the sideways orientation of the alkali metal and hydride results in interaction between B/Al and C atoms, (ii) the smaller distance between the metal and GDY helps in the removal of H atom from the complex and (iii) the interaction between alkali metal and C increases the electron density on the C atom on removal of the H atom, which causes a decrease in the energy required to remove H from hydride.

The use of GDY in the field of photocatalysis is another emerging aspect that has many applications in the field of environmental remediation. Wang and co-workers studied the



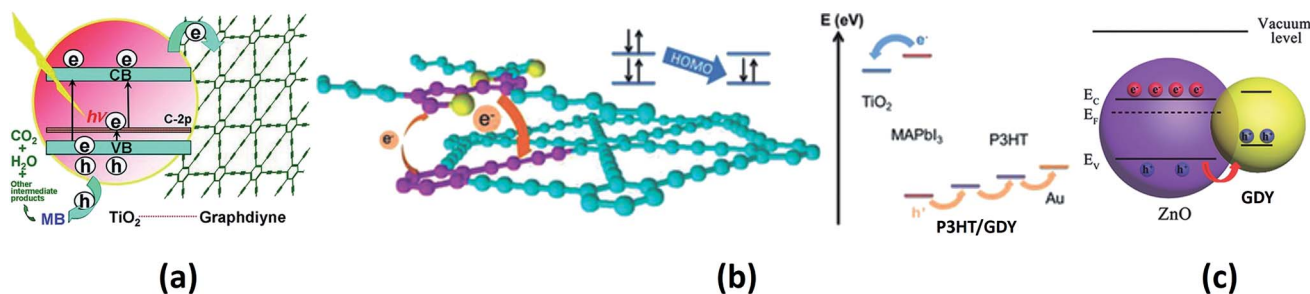


Fig. 8 Schematics of various photo-processes assisted by GYs: (a) photocatalytic degradation of methylene blue over P25–GDY. (b) Illustration of energy levels in P3HT/GDY-based perovskite solar cell. (c) Illustration of hole transfer from ZnO NP to GDY NP that enhances UV photo-detection. Parts (a), (b) and (c) are reproduced from ref. 177, 182 and 186, respectively, with permissions from the John Wiley and Sons.

activity of titania–GDY nanocomposites (P25–GDY) of different compositions as photocatalysts for the degradation of methylene blue.<sup>177</sup> P25–GDY showed improved activity than bare P25 due to the conjugated nature of GDY. The excited electron moves from titania to GDY by a percolation mechanism, thereby suppressing charge recombination. A schematic representation of this process is shown in Fig. 8a. Further, experimental studies were carried out to understand the photocatalytic properties of GDY-modified TiO<sub>2</sub>.<sup>178</sup> Yang *et al.* compared the catalytic ability of TiO<sub>2</sub>–GDY and TiO<sub>2</sub>–graphene composites on the photodegradation of methylene blue and it was seen that the TiO<sub>2</sub>–GDY composite showed higher catalytic performance. The enhanced charge separation, higher oxidation ability and the presence of larger impurity levels in TiO<sub>2</sub>–GDY composite are the major reasons attributed to its better performance. Another GDY hybrid, Ag/AgBr/GO/GDY has been successfully implemented as a photocatalyst for the degradation of methyl orange.<sup>179</sup> The GO/GDY part of the hybrid suppresses the charge recombination and the hybrid shows 100% methyl orange degradation. A novel GDY–ZnO hybrid was investigated as a catalyst for the photodegradation of azo dyes.<sup>180</sup> GDY–ZnO nanohybrid showed improved photocatalytic activity towards the degradation of methylene blue and rhodamine B. It also showed a superior catalytic efficiency towards the photodegradation of phenol compared to bare ZnO nanoparticles. Thus, we see that, a majority of the suggested applications of various forms of GYs in catalysis have emerged out of first-principles calculations and are yet to be realized in experiments. The authors believe that recent experiments demonstrating the role of GDY in photocatalysis would provide an impetus to researchers towards probing the other forms of GYs for related applications.

## 8 Graphynes for optoelectronics

Recent research has indicated that GDY can be employed to achieve improved performance of optoelectronic devices such as solar cells and photodetectors. In this section, we describe some of the applications of GDYs in optoelectronics. The effect of GDY doping on the efficiency of polymer solar cells was investigated by Du and co-workers.<sup>181</sup> Different ratios of GDY was doped into poly(3-hexathiophene) (P3HT):[6,6]-phenyl-C<sub>61</sub>-

butyric acid methyl ester (PCBM) active layer and the short-circuit current,  $J_{sc}$  and power conversion efficiency (PCE) were calculated. The doped cells showed superior PCE and higher  $J_{sc}$  due to better electron transport *via* the GDY percolation paths. Perovskite-based solar cells may or may not contain a hole transporting material (HTM) layer. However, the HTM layer is usually necessary to improve the efficiency as it suppresses charge recombination and helps in hole transportation/separation. Inspired by the improved efficiency of GDY-doped P3HT in organic solar cells, Xiao *et al.* employed GDY doping on P3HT hole transporting material of perovskite solar cell and observed a higher PCE.<sup>182</sup> The strong  $\pi$ – $\pi$  stacking between GDY and P3HT enhances the hole transporting capacity (Fig. 8b). The GDY aggregates also enhance the light absorption of the solar cell. The doping of GDY in the electron transport layer (ETL) of perovskite has also been investigated. A work by Kuang *et al.* showed a remarkable PCE enhancement when the ETL component of PCBM was doped with GDY.<sup>183</sup> GDY improved electron mobility, conductivity, charge extraction as well as the ETL coverage on the cell. The enhancement of  $J_{sc}$  is attributed to the high conductivity of GDY. Very recently, improved electron transport was achieved in CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> perovskite solar cell by doping both PCBM and ZnO film of the cell with GDY.<sup>184</sup> The dual-doped solar cell showed a remarkable enhancement of PCE and stability. The dual doping with GDY decreased the surface defects of PCBM and ZnO, thereby suppressing the charge recombination and current leakage. GDY has also been used to enhance the catalytic activity of dye-sensitized solar cells (DSSC).<sup>185</sup> GDY nanosheets with Pt nanoparticles were shown to improve electron transfer and catalytic activity when used as the counter electrode of a DSSC. The GDY-based substrates showed better performance than bare Pt and Pt nanoparticle-reduced graphene oxide counter electrodes.

A photoelectrochemical cell (PEC) generates hydrogen by conversion of the solar energy into a photoelectrode. The higher surface area and light harvesting abilities make quantum dots (QDs) an excellent material for photoelectrodes. They are mainly used as photoanodes because electron transfer through the QDs can take place easily.<sup>187,188</sup> However, the transport of holes through QDs is difficult and hence QD-based photocathodes are limited. Therefore, researchers are trying to develop materials that can facilitate the hole transportation through





photocathodes. Li *et al.* investigated the use of GDY as the HTM of the QD-sensitized photocathode for hydrogen production in a PEC water splitting cell and it was seen that the high hole mobility through GDY and the strong  $\pi$ - $\pi$  interaction between GDY and mercaptopyrindine surface-functionalized CdSe QDs enhance the performance of the PEC.<sup>189</sup> Another study by Gao *et al.* utilizes the high carrier mobility through GDY to enhance the PEC behaviour of BiVO<sub>4</sub> by synthesising a GDY/BiVO<sub>4</sub> composite photoanode for the PEC.<sup>190</sup> The faster extraction of the carriers from BiVO<sub>4</sub> by GDY helps in reducing unwanted carrier recombinations and prevents the oxidation of the BiVO<sub>4</sub> layer by the holes, thereby increasing the stability of the electrode.

GDY is also being used to improve the performance of photodetectors (PDs). ZnO is one of the most common materials used to fabricate ultraviolet (UV) PDs. However, there are still efforts to improve the performance of UV PDs. Recently, GD:ZnO nanocomposites have been synthesised and they exhibit enhanced photoresponse.<sup>186</sup> The GDY layer, as explained before, reduces charge recombination by trapping the holes that are generated by UV irradiation. Fig. 8c is an illustration of the hole transfer from ZnO NP to GDY NP. The PN junction depletion region width quickly decreases (increases) upon light illumination on (off) and this results in improved photoresponse. Semiconducting single-walled carbon nanotubes (s-SWNTs) are generally used for IR detection, but the poor separation of the excitons by s-SWNTs and exciton quenching by small quantities of metallic nanotubes present in the s-SWNTs raise the demand for discovering new materials that can combine with s-SWNTs to overcome these problems. Zheng *et al.* found that the introduction of GDY layers can improve the photoelectric conversion of the s-SWNTs.<sup>191</sup> The GDY layer helps in faster dissociation of the excitons and increases carrier density, which in turn improves the photoresponse.

## 9 Conclusions and outlook

Overall, “graphynes” is the latest buzzword in carbon flatland. The family of graphynes is structurally diverse, offering researchers a variety of nanoporous architectures to play around with. The real challenge in this area of course lies in achieving successful large-scale synthesis of  $\alpha$ -,  $\beta$ -,  $\gamma$ -,  $\delta$ -, 6,6,12-GYs and rhombic-GYs. Bulk syntheses of  $\gamma$ -GY-2 and rhombic-GY-2 are only reported to date. Attempts are being made by several research groups to synthesise the other forms and the authors believe that they would soon be realised in experiments. Despite limited success in the synthetic protocols, theorists have taken strides towards understanding the fundamental properties and proposing newer applications of GYs. Some of the interesting properties of GYs include Dirac cone-like electronic structure, high carrier mobility, anisotropic electrical conductivity, stress or strain induced variation in band gap, mechanical robustness *etc.*, most of which are summarized in a recent review by Li and co-workers.<sup>24</sup> Structural diversity of the GY family has been showcased earlier by Ivanovskii.<sup>25</sup> The porous architectures of GNTs can result in potential applications such as drug delivery, molecular encapsulation *etc.* Their low thermal conductivity

and superplasticity in contrast to CNTs are attributed to the acetylenic linkages. In this review article, we focused on describing the theoretical as well as experimental aspects of the applications of GYs, emerging out of their interesting porous architectures. The applications considered range from water desalination, gas sensing and separation, molecular adsorption, energy storage, photovoltaics to catalysis. Theoretical studies have proposed that GYs could be excellent membranes for water desalination. The scope of GYs in forward osmosis is another exciting aspect that is being investigated in great detail. Gas separation using GYs is theoretically found to be a better option when compared to the conventional separation techniques. The variation in electrical conductivity of GYs on molecular adsorption contributes to their sensing ability. The uniform porous structures of GYs imply that they are a viable option for hydrogen storage. Doping of GY sheets and metal decoration have emerged as strategies for fine-tuning the properties like gas sensing and gas storage. Research has indicated improvement in storage capacity of Li-ion batteries on doping, fuelling interest in controlling the efficiencies of Na and Mg ion batteries as well. The unique electronic structure of GYs makes them promising materials for photovoltaics, eventually aiming at the fabrication of artificial photosynthetic devices. GYs, in pristine as well as doped form, are suggested as alternatives for metals in electrochemical actuators, molecular switches and drug delivery. Most of the efforts thus far have been in the direction of probing if GYs could emerge as viable alternatives to graphene for various purposes. It is interesting to note that, despite numerous theoretical and experimental studies on investigating the interactions between graphene and biopolymers such as DNA and proteins, a parallel effort involving GYs has hardly been undertaken.<sup>192,193</sup> We believe that studies probing the interfaces of GYs with biological systems will soon emerge. Therefore, it is clear that researchers have only scratched the surface of the numerous possibilities that the family of GYs has to offer and there are many more opportunities for theorists as well as experimentalists towards realizing the full potential of GYs.

## Conflicts of interest

The authors declare that there are no conflicts of interest.

## Acknowledgements

RSS acknowledges the Science and Engineering Research Board (SERB), Government of India for financial support through the SERB Women Excellence Award Scheme. AJ, CJ, CHO, and SC thank IISER-TVM for financial support. SM thanks the Department of Science and Technology (DST), Government of India for financial support through the INSPIRE fellowship.

## References

- 1 A. Coustenis and T. Encrenaz, *Life beyond Earth: the search for habitable worlds in the Universe*, Cambridge University Press, Cambridge, 2013, pp. 1–287.



- 2 H. B. Liu, J. L. Xu, Y. J. Li and Y. L. Li, *Acc. Chem. Res.*, 2010, **43**, 1496–1508.
- 3 H. Y. Zheng, Y. J. Li, H. B. Liu, X. D. Yin and Y. L. Li, *Chem. Soc. Rev.*, 2011, **40**, 4506–4524.
- 4 H. B. Liu, Y. L. Li, L. Jiang, H. Y. Luo, S. Q. Xiao, H. J. Fang, H. M. Li, D. B. Zhu, D. P. Yu, J. Xu and B. Xiang, *J. Am. Chem. Soc.*, 2002, **124**, 13370–13371.
- 5 A. A. Balandin, S. Ghosh, W. Bao, I. Calizo, D. Teweldebrhan, F. Miao and C. N. Lau, *Nano Lett.*, 2008, **8**, 902–907.
- 6 R. R. Nair, P. Blake, A. N. Grigorenko, K. S. Novoselov, T. J. Booth, T. Stauber, N. M. R. Peres and A. K. Geim, *Science*, 2008, **320**, 1308.
- 7 S. Bae, H. Kim, Y. Lee, X. Xu, J. Park, Y. Zheng, J. Balakrishnan, L. Tian, H. Ri Kim, Y. I. Song, Y. Kim, K. Kim, B. Ozyilmaz, J. Ahn, B. Hee Hong and S. Iijima, *Nat. Nanotechnol.*, 2010, **5**, 574–578.
- 8 C. Lee, X. Wei, J. W. Kysar and J. Hone, *Science*, 2008, **321**, 385–388.
- 9 K. A. Ritter and J. W. Lyding, *Nat. Mater.*, 2009, **8**, 235–242.
- 10 J. Wu, H. A. Becerril, Z. Bao, Z. Liu, Y. Chen and P. Peumans, *Appl. Phys. Lett.*, 2008, **92**, 263302.
- 11 X. Li, X. Wang, L. Zhang, S. Lee and H. Dai, *Science*, 2008, **319**, 1229–1232.
- 12 J. D. Fowler, M. J. Allen, V. C. Tung, Y. Yang, R. B. Kaner and B. H. Weiller, *ACS Nano*, 2009, **3**, 301–306.
- 13 C. Berger, Z. M. Song, T. B. Li, X. B. Li, A. Y. Ogbazghi, R. Feng, Z. Dai, A. N. Marchenkov, E. H. Conrad, P. N. First and W. A. de Heer, *J. Phys. Chem. B*, 2004, **108**, 19912–19916.
- 14 C. Berger, Z. Song, X. Li, X. Wu, N. Brown, C. Naud, D. Mayou, T. Li, J. Hass, E. H. Marchenkov, A. N. Conrad, P. N. First and W. A. de Heer, *Science*, 2006, **312**, 1191–1196.
- 15 S. Yang, X. Feng, S. Ivanovici and K. Müllen, *Angew. Chem., Int. Ed.*, 2010, **49**, 8408–8411.
- 16 H. Wang, L. F. Cui, Y. Yang, C. H. Sanchez, J. T. Robinson, Y. Liang, Y. Cui and H. Dai, *J. Am. Chem. Soc.*, 2010, **132**, 13978–13980.
- 17 S.-M. Paek, E. Yoo and I. Honma, *Nano Lett.*, 2008, **9**, 72–75.
- 18 L. Pan, X.-D. Zhu, K.-N. Sun, Y.-T. Liu, X.-M. Xie and X.-Y. Ye, *Nano Energy*, 2016, **30**, 347–354.
- 19 P. Long, Z. Xiao-Dong, X. Xu-Ming and L. Yi-Tao, *Adv. Funct. Mater.*, 2015, **25**, 3341–3350.
- 20 Y.-T. Liu, X.-D. Zhu, Z.-Q. Duan and X.-M. Xie, *Chem. Commun.*, 2013, **49**, 10305–10307.
- 21 P. R. Wallace, *Phys. Rev.*, 1947, **71**, 622–634.
- 22 H. P. Boehm, A. Clauss, G. O. Fischer and U. Hofmann, *Z. Anorg. Allg. Chem.*, 1962, **316**, 119–127.
- 23 K. S. Novoselov, A. K. Geim, S. V. Morozov, D. Jiang, Y. Zhang, S. V. Dubonos, I. V. Grigorieva and A. A. Firsov, *Science*, 2004, **306**, 666–669.
- 24 Y. J. Li, L. Xu, H. B. Liu and Y. L. Li, *Chem. Soc. Rev.*, 2014, **43**, 2572–2586.
- 25 A. L. Ivanovskii, *Prog. Solid State Chem.*, 2013, **41**, 1–19.
- 26 J. O. Sofo, A. S. Chaudhari and G. D. Barber, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2007, **75**, 153401.
- 27 D. C. Elias, R. R. Nair, T. M. G. Mohiuddin, S. V. Morozov, P. Blake, M. P. Halsall, A. C. Ferrari, D. W. Boukhvalov, M. I. Katsnelson, A. K. Geim and K. S. Novoselov, *Science*, 2009, **323**, 610–613.
- 28 J. Zhou, Q. Wang, Q. Sun, X. S. Chen, Y. Kawazoe and P. Jena, *Nano Lett.*, 2009, **9**, 3867–3870.
- 29 W. Zhao, J. Gebhardt, F. Späth, K. Gotterbarm, C. Gleichweit, H.-P. Steinrück, A. Görling and C. Papp, *Chem.–Eur. J.*, 2015, **21**, 3347–3358.
- 30 R. H. Baughman, H. Eckhardt and M. Kertesz, *J. Chem. Phys.*, 1987, **87**, 6687–6699.
- 31 B. G. Kim and H. J. Choi, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, **86**, 115435.
- 32 F. Diederich, *Nature*, 1994, **369**, 199–207.
- 33 U. H. F. Bunz, Y. Rubin and Y. Tobe, *Chem. Soc. Rev.*, 1999, **28**, 107–119.
- 34 G. Li, Y. Li, H. Liu, Y. Guo, Y. Li and D. Zhu, *Chem. Commun.*, 2010, **46**, 3256–3258.
- 35 Y.-Q. Zhang, N. Kepčija, M. Kleinschrodt, K. Diller, S. Fischer, A. C. Papageorgiou, F. Allegretti, J. Björk, S. Klyatskaya, F. Klappenberger, M. Ruben and J. V. Barth, *Nat. Commun.*, 2012, **3**, 1286.
- 36 Z. Jia, Z. Zuo, Y. Yi, H. Liu, D. Li, Y. Li and Y. Li, *Nano Energy*, 2017, **33**, 343–349.
- 37 Z. Jia, Y. Li, Z. Zuo, H. Liu, C. Huang and Y. Li, *Acc. Chem. Res.*, 2017, **50**, 2470–2478.
- 38 J. Zhong, J. Wang, J.-G. Zhou, B.-H. Mao, C.-H. Liu, H.-B. Liu, Y.-L. Li, T.-K. Sham, X.-H. Sun and S.-D. Wang, *J. Phys. Chem. C*, 2013, **117**, 5931–5936.
- 39 Y. Yang and X. Xu, *Comput. Mater. Sci.*, 2012, **61**, 83–88.
- 40 Y. Pei, *Physica B*, 2012, **407**, 4436–4439.
- 41 S. Lin and M. J. Buehler, *Nanoscale*, 2013, **5**, 11801–11807.
- 42 A. L. Ivanovskii, *Prog. Solid State Chem.*, 2013, **41**, 1–14.
- 43 C. Owais, A. James, C. John, R. Dhali and R. S. Swathi, *J. Phys. Chem. B*, 2018, **122**, 5127–5146.
- 44 S. W. Cranford and M. J. Buehler, *Carbon*, 2011, **49**, 4111–4121.
- 45 A. N. Enyashin and A. L. Ivanovskii, *Phys. Status Solidi B*, 2011, **248**, 1879–1883.
- 46 S. W. Cranford, D. B. Brommer and M. J. Buehler, *Nanoscale*, 2012, **4**, 7797–7809.
- 47 D. Malko, C. Neiss, F. Viñes and A. Görling, *Phys. Rev. Lett.*, 2012, **108**, 086804.
- 48 Y. Y. Zhang, Q. X. Pei and C. M. Wang, *Appl. Phys. Lett.*, 2012, **101**, 081909.
- 49 Q. Peng, W. Ji and S. De, *Phys. Chem. Chem. Phys.*, 2012, **14**, 13385–13391.
- 50 V. R. Coluci, S. F. Braga, S. B. Legoas, D. S. Galvão and R. H. Baughman, *Nanotechnology*, 2004, **15**, S142–S149.
- 51 X.-M. Wang and S.-S. Lu, *J. Phys. Chem. C*, 2013, **117**, 19740–19745.
- 52 G. Li, Y. Li, X. Qian, H. Liu, H. Lin, N. Chen and Y. Li, *J. Phys. Chem. C*, 2011, **115**, 2611–2615.
- 53 S. Alaei, S. Jalili and S. Erkoc, *Fullerenes, Nanotub. Carbon Nanostruct.*, 2014, **23**, 494–499.
- 54 J. Deb, B. Bhattacharya and U. Sarkar, *AIP Conf. Proc.*, 2016, **1731**, 050081.



- 55 M. Hu, Y. Jing and X. Zhang, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2015, **91**, 155408.
- 56 J. M. de Sousa, G. Brunetto, V. R. Coluci and D. S. Galvão, *Carbon*, 2016, **96**, 14–19.
- 57 X. Qian, Z. Ning, Y. Li, H. Liu, C. Ouyang, Q. Chen and Y. Li, *Dalton Trans.*, 2012, **41**, 730–733.
- 58 D. Solis, C. F. Woellner, D. D. Borges and D. S. Galvao, *MRS Adv.*, 2017, **2**, 129–134.
- 59 V. R. Coluci, S. F. Braga, S. B. Legoas, D. S. Galvão and R. H. Baughman, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2003, **68**, 035430.
- 60 R. Matsuoka, R. Sakamoto, K. Hoshiko, S. Sasaki, H. Masunaga, K. Nagashio and H. Nishihara, *J. Am. Chem. Soc.*, 2017, **139**, 3145–3152.
- 61 G. Li, Y. Li, H. Liu, Y. Guo, Y. Li and D. Zhu, *Chem. Commun.*, 2010, **46**, 3256–3258.
- 62 X. Qian, H. Liu, C. Huang, S. Chen, L. Zhang, Y. Li, J. Wang and Y. Li, *Sci. Rep.*, 2015, **5**, 7756.
- 63 M. M. Haley, *Pure Appl. Chem.*, 2008, **80**, 519–532.
- 64 B. Wu, M. Li, S. Xiao, Y. Qu, X. Qiu, T. Liu, F. Tian, H. Li and S. Xiao, *Nanoscale*, 2017, **9**, 11939–11943.
- 65 Q. Yuan and F. Ding, *Nanoscale*, 2014, **6**, 12727–12731.
- 66 F. Klappenberger, R. Hellwig, P. Du, T. Paintner, M. Uphoff, L. Zhang, T. Lin, B. A. Moghanaki, M. Paszkiewicz, I. Vobornik, J. Fujii, O. Fuhr, Y. Q. Zhang, F. Allegretti, M. Ruben and J. V. Barth, *Small*, 2018, **14**, 1704321.
- 67 R. Liu, J. Y. Zhou, X. Gao, J. Q. Li, Z. Q. Xie, Z. Z. Li, S. Q. Zhang, L. M. Tong, J. Zhang and Z. F. Liu, *Adv. Electron. Mater.*, 2017, **3**, 1700122.
- 68 M. Bartolomei, E. Carmona-Novillo, M. I. Hernandez, J. Campos-Martinez, F. Pirani, G. Giorgi and K. Yamashita, *J. Phys. Chem. Lett.*, 2014, **5**, 751–755.
- 69 C. Zhu, H. Li, X. C. Zeng, E. G. Wang and S. Meng, *Sci. Rep.*, 2013, **3**, 3163.
- 70 X. Zhang and J. G. Gai, *RSC Adv.*, 2015, **5**, 68109–68116.
- 71 Z. Qin and M. J. Buehler, *Nano Lett.*, 2015, **15**, 3939–3944.
- 72 J. L. Kou, X. Y. Zhou, Y. Y. Chen, H. J. Lu, F. M. Wu and J. T. Fan, *J. Chem. Phys.*, 2013, **139**, 064705.
- 73 J. L. Kou, X. Y. Zhou, H. J. Lu, F. M. Wu and J. T. Fan, *Nanoscale*, 2014, **6**, 1865–1870.
- 74 M. M. Xue, H. Qiu and W. L. Guo, *Nanotechnology*, 2013, **24**, 505720.
- 75 M. Raju, P. B. Govindaraju, A. C. T. van Duin and M. Ihme, *Nanoscale*, 2018, **10**, 3969–3980.
- 76 N. W. Ockwig and T. M. Nenoff, *Chem. Rev.*, 2007, **107**, 4078–4110.
- 77 G. Arora and S. I. Sandler, *Fluid Phase Equilib.*, 2007, **259**, 3–8.
- 78 M. Hankel, H. Zhang, T. X. Nguyen, S. K. Bhatia, S. K. Gray and S. C. Smith, *Phys. Chem. Chem. Phys.*, 2011, **13**, 7834–7844.
- 79 T. M. Nenoff, R. J. Spontak and C. M. Aberg, *MRS Bull.*, 2006, **31**, 735–741.
- 80 P. Bernardo, E. Drioli and G. Golemme, *Ind. Eng. Chem. Res.*, 2009, **48**, 4638–4663.
- 81 Y. Li, Z. Zhou, P. Shen and Z. Chen, *Chem. Commun.*, 2010, **46**, 3672–3674.
- 82 S. Blankenburg, M. Bieri, R. Fasel, K. Müllen, C. A. Pignedoli and D. Passerone, *Small*, 2010, **6**, 2266–2271.
- 83 J. S. Bunch, S. S. Verbridge, J. S. Alden, A. M. van der Zande, J. M. Parpia, H. G. Craighead and P. L. McEuen, *Nano Lett.*, 2010, **8**, 2458–2462.
- 84 O. Leenaerts, B. Partoens and F. M. Peeters, *Appl. Phys. Lett.*, 2010, **93**, 193107.
- 85 M. D. Fischbein and M. Drndć, *Appl. Phys. Lett.*, 2008, **93**, 113107.
- 86 G. Compagnini, F. Giannazzo, S. Sonde, V. Raineri and E. Rimini, *Carbon*, 2008, **47**, 3201–3207.
- 87 Y. Jiao, A. Du, M. Hankel, Z. Zhu, V. Rudolph and S. C. Smith, *Chem. Commun.*, 2011, **47**, 11843–11845.
- 88 W.-h. Zhao, L.-f. Yuan and J.-l. Yang, *Chin. J. Chem. Phys.*, 2012, **25**, 434–440.
- 89 S. W. Cranford and M. J. Buehler, *Nanoscale*, 2012, **4**, 4587–4593.
- 90 Y. Jiao, A. Du, S. C. Smith, Z. Zhu and S. Z. Qiao, *J. Mater. Chem. A*, 2015, **3**, 6767–6771.
- 91 X. Tan, L. Kou, H. A. Tahini and S. C. Smith, *Mol. Simul.*, 2015, **42**, 573–579.
- 92 H. Zhang, X. He, M. Zhao, M. Zhang, L. Zhao, X. Feng and Y. Luo, *J. Phys. Chem. C*, 2012, **116**, 16634–16638.
- 93 P. Sang, L. Zhao, J. Xu, Z. Shi, S. Guo, Y. Yu, H. Zhu, Z. Yan and W. Guo, *Int. J. Hydrogen Energy*, 2017, **42**, 5168–5176.
- 94 M. Alaghemandi, *Chem. Phys. Lett.*, 2015, **629**, 65–69.
- 95 Z. Meng, X. Zhang, Y. Zhang, H. Gao, Y. Wang, Q. Shi, D. Rao, Y. Liu, K. Deng and R. Lu, *ACS Appl. Mater. Interfaces*, 2016, **8**, 28166–28170.
- 96 L. Zhao, P. Sang, S. Guo, X. Liu, J. Li, H. Zhu and W. Guo, *Appl. Surf. Sci.*, 2017, **405**, 455–464.
- 97 G. Lei, C. Liu, Q. Li and X. Xu, *Fuel*, 2016, **182**, 210–219.
- 98 A. Cho, *Science*, 2009, **326**, 778–779.
- 99 W. P. Halperin, *Nat. Phys.*, 2014, **10**, 467–470.
- 100 M. I. Hernández, M. Bartolomei and J. Campos-Martínez, *J. Phys. Chem. A*, 2015, **119**, 10743–10749.
- 101 M. Bartolomei, E. Carmona-Novillo, M. I. Hernández, J. Campos-Martínez, F. Pirani and G. Giorgi, *J. Phys. Chem. C*, 2014, **118**, 29966–29972.
- 102 X. Bévenot, A. Trouillet, C. Veillas, H. Gagnaire and M. Clément, *Sens. Actuators, B*, 2000, **67**, 57–67.
- 103 S. Ryabtsev, A. Shaposhnick, A. Lukin and E. Domashevskaya, *Sens. Actuators, B*, 1999, **59**, 26–29.
- 104 W. Tsujita, A. Yoshino, H. Ishida and T. Moriizumi, *Sens. Actuators, B*, 2005, **110**, 304–311.
- 105 R. Majidi and A. R. Karami, *Phys. E*, 2014, **59**, 169–173.
- 106 X. Chen, P. Gao, L. Guo, Y. Wen, D. Fang, B. Gong, Y. Zhang and S. Zhang, *Phys. Lett. A*, 2017, **381**, 879–885.
- 107 R. Majidi and A. R. Karami, *Phys. E*, 2013, **54**, 177–180.
- 108 J. Deb, B. Bhattacharya, N. B. Singh and U. Sarkar, *Struct. Chem.*, 2016, **27**, 1221–1227.
- 109 R. Majidi and A. Karami, *Stud. Univ. Babeş-Bolyai, Chem.*, 2016, **61**, 177–184.
- 110 D. Cortés-Arriagada, *Int. J. Quantum Chem.*, 2017, **117**, 1–7.
- 111 M. Ozmaian, A. Fathizadeh, M. Jalalvand, M. R. Eftehadi and S. M. Allaei, *Sci. Rep.*, 2016, **6**, 21910.



- 112 P. Zhang, Q. Song, J. Zhuang and X.-J. Ning, *Chem. Phys. Lett.*, 2017, **689**, 185–189.
- 113 H. J. Kwon, Y. Kwon, T. Kim, Y. Jung, S. Lee, M. Cho and S. Kwon, *AIP Adv.*, 2017, **7**, 125013.
- 114 Z. Yan, L. Wang, J. Cheng, L. Huang, C. Zhu, C. Chen, L. Miao and J. Jiang, *J. Appl. Phys.*, 2014, **116**, 174304.
- 115 B. Kang, H. Liu and J. Y. Lee, *Phys. Chem. Chem. Phys.*, 2014, **16**, 974–980.
- 116 A. Omidvar and A. Mohajeri, *Mol. Phys.*, 2015, **113**, 3900–3908.
- 117 J. Beheshtian, A. A. Peyghan, Z. Bagheri and M. B. Tabar, *J. Struct. Chem.*, 2013, **25**, 1–7.
- 118 Z. Felegari and S. Hamedani, *Results Phys.*, 2017, **7**, 2626–2631.
- 119 A. A. Peyghan, S. F. Rastegar and N. L. Hadipour, *Phys. Lett. A*, 2014, **378**, 2184–2190.
- 120 Z. Lu, P. Lv, D. Ma, X. Yang, S. Li and Z. Yang, *J. Phys. D: Appl. Phys.*, 2018, **51**, 065109.
- 121 A. Zuttel, *Naturwissenschaften*, 2004, **91**, 157–172.
- 122 D. J. Durbin and C. Malardier-Jugroot, *Int. J. Hydrogen Energy*, 2013, **38**, 14595–14617.
- 123 A. Delahaye, A. Aoufi, A. Gicquel and I. Pentchev, *AIChE J.*, 2002, **48**, 2061–2073.
- 124 K. Srinivasu and S. K. Ghosh, *J. Phys. Chem. C*, 2012, **116**, 5951–5956.
- 125 S. C. Shekar and R. S. Swathi, *J. Phys. Chem. C*, 2015, **119**, 8912–8923.
- 126 H. Y. Zhang, M. W. Zhao, X. J. He, Z. H. Wang, X. J. Zhang and X. D. Liu, *J. Phys. Chem. C*, 2011, **115**, 8845–8850.
- 127 K. Wang, N. Wang, J. J. He, Z. Yang, X. Y. Shen and C. S. Huang, *Electrochim. Acta*, 2017, **253**, 506–516.
- 128 S. L. Zhang, H. B. Liu, C. S. Huang, G. L. Cui and Y. L. Li, *Chem. Commun.*, 2015, **51**, 1834–1837.
- 129 H. Zhang, X. Zhao, M. Zhang, Y. Luo, G. Li and M. Zhao, *J. Phys. D: Appl. Phys.*, 2013, **46**, 495307.
- 130 Y. Guo, K. Jiang, B. Xu, Y. Xia, J. Yin and Z. Liu, *J. Phys. Chem. C*, 2012, **116**, 13837–13841.
- 131 H. Zhang, M. Zhao, H. Bu, X. He, M. Zhang, L. Zhao and Y. Luo, *J. Appl. Phys.*, 2012, **112**, 084305.
- 132 J. Lu, Y. Guo, Y. Zhang and J. Cao, *Int. J. Hydrogen Energy*, 2014, **39**, 17112–17117.
- 133 S. Kumar and T. J. Dhillip Kumar, *ACS Appl. Mater. Interfaces*, 2017, **9**, 28659–28666.
- 134 R. Lu, D. Rao, Z. Meng, X. Zhang, G. Xu, Y. Liu, E. Kan, C. Xiao and K. Deng, *Phys. Chem. Chem. Phys.*, 2013, **15**, 16120–16126.
- 135 B. Jang, J. Koo, M. Park, H. Lee, J. Nam, Y. Kwon and H. Lee, *Appl. Phys. Lett.*, 2013, **103**, 263904.
- 136 H. J. Hwang, J. Koo, M. Park, N. Park, Y. Kwon and H. Lee, *J. Phys. Chem. C*, 2013, **117**, 6919–6923.
- 137 M. Makaremi, B. Mortazavi and C. V. Singh, *Appl. Mater. Today*, 2018, **10**, 115–121.
- 138 Y. Liu, W. Liu, R. Wang, L. Hao and W. Jiao, *Int. J. Hydrogen Energy*, 2014, **39**, 12757–12764.
- 139 H. J. Hwang, Y. Kwon and H. Lee, *J. Phys. Chem. C*, 2012, **116**, 20220–20224.
- 140 C. Li, J. Li, F. Wu, S.-S. Li, J.-B. Xia and L.-W. Wang, *J. Phys. Chem. C*, 2011, **115**, 23221–23225.
- 141 Y. Guo, X. Lan, J. Cao, B. Xu, Y. Xia, J. Yin and Z. Liu, *Int. J. Hydrogen Energy*, 2013, **38**, 3987–3993.
- 142 Y. S. Wang, P. Fei Yuan, M. Li, W. Fen Jiang, Q. Sun and Y. Jia, *J. Solid State Chem.*, 2013, **197**, 323–328.
- 143 J. Lu, Y. Guo, Y. Zhang, Y. Tang and J. Cao, *J. Solid State Chem.*, 2015, **231**, 53–57.
- 144 P.-P. Liu, H. Zhang, X.-L. Cheng and Y.-J. Tang, *Appl. Surf. Sci.*, 2016, **371**, 44–49.
- 145 L. Zhang, S. Zhang, P. Wang, C. Liu, S. Huang and H. Tian, *Comput. Theor. Chem.*, 2014, **1035**, 68–75.
- 146 C. H. Sun and D. J. Searles, *J. Phys. Chem. C*, 2012, **116**, 26222–26226.
- 147 K. Persson, Y. Hinuma, Y. S. Meng, A. Van der Ven and G. Ceder, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2010, **82**, 125416.
- 148 S. C. Shekar and R. S. Swathi, *J. Phys. Chem. A*, 2013, **117**, 8632–8641.
- 149 Z. M. Xu, X. Lv, J. Li, J. G. Chen and Q. S. Liu, *RSC Adv.*, 2016, **6**, 25594–25600.
- 150 S. L. Zhang, J. J. He, J. Zheng, C. S. Huang, Q. Lv, K. Wang, N. Wang and Z. G. Lan, *J. Mater. Chem. A*, 2017, **5**, 2045–2051.
- 151 C. S. Huang, S. L. Zhang, H. B. Liu, Y. J. Li, G. T. Cui and Y. L. Li, *Nano Energy*, 2015, **11**, 481–489.
- 152 H. P. Du, H. Yang, C. S. Huang, J. J. He, H. B. Liu and Y. L. Li, *Nano Energy*, 2016, **22**, 615–622.
- 153 Z. Yang, X. Shen, N. Wang, J. He, X. Li, X. Wang, Z. Hou, K. Wang, J. Gao, T. Jiu and C. Huang, *ACS Appl. Mater. Interfaces*, 2018, DOI: 10.1021/acsami.8b01823.
- 154 S. Zhang, H. Du, J. He, C. Huang, H. Liu, G. Cui and Y. Li, *ACS Appl. Mater. Interfaces*, 2016, **8**, 8467–8473.
- 155 V. R. Stamenkovic, B. Fowler, B. S. Mun, G. Wang, P. N. Ross, C. A. Lucas and N. M. Marković, *Science*, 2007, **315**, 493–497.
- 156 Z. Peng and H. Yang, *J. Am. Chem. Soc.*, 2009, **131**, 7542–7543.
- 157 K. Gong, F. Du, Z. Xia, M. Durstock and L. Dai, *Science*, 2009, **323**, 760–764.
- 158 L. Yang, S. Jiang, Y. Zhao, L. Zhu, S. Chen, X. Wang, Q. Wu, J. Ma, Y. Ma and Z. Hu, *Angew. Chem., Int. Ed.*, 2011, **50**, 7132–7135.
- 159 Z.-H. Sheng, L. Shao, J.-J. Chen, W.-J. Bao, F.-B. Wang and X.-H. Xia, *ACS Nano*, 2011, **5**, 4350–4358.
- 160 L. Qu, Y. Liu, J.-B. Baek and L. Dai, *ACS Nano*, 2010, **4**, 1321–1326.
- 161 P. Wu, P. Du, H. Zhang and C. Cai, *J. Phys. Chem. C*, 2012, **116**, 20472–20479.
- 162 P. Wu, P. Du, H. Zhang and C. Cai, *Phys. Chem. Chem. Phys.*, 2015, **17**, 1441–1449.
- 163 X. Kong, Y. Huang and Q. Liu, *Carbon*, 2017, **123**, 558–564.
- 164 K. Srinivasu and S. K. Ghosh, *J. Phys. Chem. C*, 2013, **117**, 26021–26028.
- 165 X.-k. Kong, Q.-w. Chen and Z. Sun, *RSC Adv.*, 2013, **3**, 4074–4080.



- 166 X. Chen, Q. Qiao, L. An and D. Xia, *J. Phys. Chem. C*, 2015, **119**, 11493–11498.
- 167 B. Kang, H. Shi, S. Wu, W. Zhao, H. Ai and J. Y. Lee, *Carbon*, 2017, **123**, 415–420.
- 168 B. K. Das, D. Sen and K. K. Chattopadhyay, *Phys. Chem. Chem. Phys.*, 2016, **18**, 2949–2958.
- 169 R. Liu, H. Liu, Y. Li, Y. Yi, X. Shang, S. Zhang, X. Yu, S. Zhang, H. Cao and G. Zhang, *Nanoscale*, 2014, **6**, 11336–11343.
- 170 Q. Lv, W. Si, Z. Yang, N. Wang, Z. Tu, Y. Yi, C. Huang, L. Jiang, M. Zhang, J. He and Y. Long, *ACS Appl. Mater. Interfaces*, 2017, **9**, 29744–29752.
- 171 S. Zhang, Y. Cai, H. He, Y. Zhang, R. Liu, H. Cao, M. Wang, J. Liu, G. Zhang, Y. Li, H. Liu and B. Li, *J. Mater. Chem. A*, 2016, **4**, 4738–4744.
- 172 X. Chen, *Phys. Chem. Chem. Phys.*, 2015, **17**, 29340–29343.
- 173 P. Wu, P. Du, H. Zhang and C. Cai, *Phys. Chem. Chem. Phys.*, 2014, **16**, 5640–5648.
- 174 J. Li, X. Gao, X. Jiang, X.-B. Li, Z. Liu, J. Zhang, C.-H. Tung and L.-Z. Wu, *ACS Catal.*, 2017, **7**, 5209–5213.
- 175 G. Tian, Z. Qi, W. Ma and Y. Wang, *ChemistrySelect*, 2017, **2**, 2311–2321.
- 176 H. Yu, A. Du, Y. Song and D. J. Searles, *J. Phys. Chem. C*, 2013, **117**, 21643–21650.
- 177 S. Wang, L. Yi, J. E. Halpert, X. Lai, Y. Liu, H. Cao, R. Yu, D. Wang and Y. Li, *Small*, 2012, **8**, 265–271.
- 178 N. Yang, Y. Liu, H. Wen, Z. Tang, H. Zhao, Y. Li and D. Wang, *ACS Nano*, 2013, **7**, 1504–1512.
- 179 X. Zhang, M. Zhu, P. Chen, Y. Li, H. Liu, Y. Li and M. Liu, *Phys. Chem. Chem. Phys.*, 2015, **17**, 1217–1225.
- 180 S. Thangavel, K. Krishnamoorthy, V. Krishnaswamy, N. Raju, S. J. Kim and G. Venugopal, *J. Phys. Chem. C*, 2015, **119**, 22057–22065.
- 181 H. Du, Z. Deng, Z. Lü, Y. Yin, L. Yu, H. Wu, Z. Chen, Y. Zou, Y. Wang, H. Liu and Y. Li, *Synth. Met.*, 2011, **161**, 2055–2057.
- 182 J. Xiao, J. Shi, H. Liu, Y. Xu, S. Lv, Y. Luo, D. Li, Q. Meng and Y. Li, *Adv. Energy Mater.*, 2015, **5**, 1401943.
- 183 C. Kuang, G. Tang, T. Jiu, H. Yang, H. Liu, B. Li, W. Luo, X. Li, W. Zhang, F. Lu, J. Fang and Y. Li, *Nano Lett.*, 2015, **15**, 2756–2762.
- 184 J. Li, T. Jiu, C. Duan, Y. Wang, H. Zhang, H. Jian, Y. Zhao, N. Wang, C. Huang and Y. Li, *Nano Energy*, 2018, **46**, 331–337.
- 185 H. Ren, H. Shao, L. Zhang, D. Guo, Q. Jin, R. Yu, L. Wang, Y. Li, Y. Wang, H. Zhao and D. Wang, *Adv. Energy Mater.*, 2015, **5**, 1500296.
- 186 Z. Jin, Q. Zhou, Y. Chen, P. Mao, H. Li, H. Liu, J. Wang and Y. Li, *Adv. Mater.*, 2016, **28**, 3697–3702.
- 187 K. Tvrđy, P. A. Frantsuzov and P. V. Kamat, *Proc. Natl. Acad. Sci. U. S. A.*, 2011, **108**, 29–34.
- 188 M. B. Wilker, K. E. Shinopoulos, K. A. Brown, D. W. Mulder, P. W. King and G. Dukovic, *J. Am. Chem. Soc.*, 2014, **136**, 4316–4324.
- 189 J. Li, X. Gao, B. Liu, Q. Feng, X. B. Li, M. Y. Huang, Z. Liu, J. Zhang, C. H. Tung and L. Z. Wu, *J. Am. Chem. Soc.*, 2016, **138**, 3954–3957.
- 190 X. Gao, J. Li, R. Du, J. Zhou, M. Y. Huang, R. Liu, J. Li, Z. Xie, L. Z. Wu, Z. Liu and J. Zhang, *Adv. Mater.*, 2017, **29**, 1605308.
- 191 Z. Zheng, H. Fang, D. Liu, Z. Tan, X. Gao, W. Hu, H. Peng, L. Tong, W. Hu and J. Zhang, *Adv. Sci.*, 2017, **4**, 1700472.
- 192 S. Chandra Shekar and R. S. Swathi, *J. Phys. Chem. C*, 2014, **118**, 4516–4528.
- 193 N. Parvin, Q. Jin, Y. Wei, R. Yu, B. Zheng, L. Huang, Y. Zhang, L. Wang, H. Zhang, M. Gao, H. Zhao, W. Hu, Y. Li and D. Wang, *Adv. Mater.*, 2017, **29**, 1606755.

