

Cite this: *RSC Adv.*, 2018, **8**, 5652

## Black Si-doped $\text{TiO}_2$ nanotube photoanode for high-efficiency photoelectrochemical water splitting

Zhenbiao Dong,<sup>a</sup> Dongyan Ding,  \*<sup>a</sup> Ting Li<sup>a</sup> and Congqin Ning<sup>b</sup>

Black Si-doped  $\text{TiO}_2$  ( $\text{Ti-Si-O}$ ) nanotubes were fabricated through Zn metal reduction of the  $\text{Ti-Si-O}$  nanotubes on  $\text{Ti-Si}$  alloy in an argon atmosphere. The nanotubes were used as a photoanode for photoelectrochemical (PEC) water splitting. Both Si element and  $\text{Ti}^{3+}$ /oxygen vacancies were introduced into the black  $\text{Ti-Si-O}$  nanotubes, which improved optical absorption and facilitated the separation of the photogenerated electron–hole pairs. The photoconversion efficiency could reach 1.22%, which was 7.18 times the efficiency of undoped  $\text{TiO}_2$ . It demonstrated that a Si element and  $\text{Ti}^{3+}$ /oxygen vacancy co-doping strategy could offer an effective method for fabricating a high-performance  $\text{TiO}_2$ -based nanostructure photoanode for improving PEC water splitting.

Received 2nd January 2018  
Accepted 23rd January 2018DOI: 10.1039/c8ra00021b  
[rsc.li/rsc-advances](http://rsc.li/rsc-advances)

## 1 Introduction

PEC solar-driven water splitting hydrogen production based on semiconductor photocatalysts has been regarded as a promising method to decrease fossil fuel consumption and solve environmental problems.<sup>1</sup> Among many photocatalysts, titanium dioxide ( $\text{TiO}_2$ ) has attracted wide interest since the first report of PEC water splitting in 1972 because of its low cost, favorable conduction band edge and chemical inertness.<sup>2–5</sup> However, it suffers from two major drawbacks. One is the short diffusion paths of the charge carriers, and another is its large band gap. This leads to rapid recombination of photoinduced electrons–holes and little absorption in the visible light region.<sup>6,7</sup> Therefore, great efforts have been made to improve the photoelectrical properties through ion doping,<sup>8,9</sup> narrow band gap semiconductor coupling,<sup>10,11</sup> noble metal deposition,<sup>12,13</sup> quantum dot sensitization, *etc.*<sup>14,15</sup>

In recent years, black  $\text{TiO}_2$  through self-doped  $\text{Ti}^{3+}$  and/or oxygen vacancies has emerged as an effective approach to improve the photocatalytic performance.<sup>16</sup> The defect states in bulk  $\text{TiO}_2$  can change the transmission path of photogenerated carriers and promote the separation-transfer process of electrons and holes. The semiconductor conductivity and charge transport were effectively promoted due to the introduction of  $\text{Ti}^{3+}$ .<sup>17</sup> Therefore, the photocatalytic efficiency of black  $\text{TiO}_2$  can be improved. In addition, Si element doping has been also

proved to be a feasible method to improve the photoelectrical performance of  $\text{TiO}_2$ .<sup>18–20</sup> To date, several synthesis methods such as hydrogen thermal treatment,<sup>17</sup> hydrogen plasma,<sup>21</sup> chemical reduction,<sup>22,23</sup> chemical oxidation,<sup>24</sup> electrochemical reduction<sup>25</sup> and anodization-annealing<sup>26</sup> have been proposed to prepare  $\text{Ti}^{3+}$  self-doped  $\text{TiO}_2$ -based photocatalysts. Meanwhile, according to published literatures, Si-doped  $\text{TiO}_2$  photocatalysts could be synthesized by chemical vapor deposition,<sup>20</sup> hydro-thermal process,<sup>27</sup> or anodizing Ti plates in  $\text{HF}/\text{Na}_2\text{SiF}_6$  solution.<sup>28</sup> However, most of these methods are generally impractical for large-scale commercial production due to high cost or complicated conditions. Up to now, there was no literature report about the preparation of Si element and  $\text{Ti}^{3+}$  co-doping  $\text{TiO}_2$  nanotubes on  $\text{Ti-Si}$  alloy for PEC water splitting.

Inspired by thermal reduction reaction and ions-doping modification methods,<sup>8,9,22</sup> we fabricated black  $\text{Ti-Si-O}$  nanotubes through Zn-reduction of the as-annealed  $\text{Ti-Si-O}$  nanotubes on  $\text{Ti-Si}$  alloy. The microstructures, optical property and the PEC water splitting properties were investigated. Electronic structures were also calculated, and possible mechanism of the black  $\text{Ti-Si-O}$  for PEC water splitting was proposed. As expected, the black  $\text{Ti-Si-O}$  nanotubes photoanode co-doped with Si element and  $\text{Ti}^{3+}$  exhibited remarkable PEC water splitting properties in comparison with undoped  $\text{TiO}_2$ .

## 2 Experimental

### 2.1 Preparation of the black $\text{Ti-Si-O}$ nanotubes

A schematic illustration of fabrication procedures for black  $\text{Ti-Si-O}$  nanotubes is shown in Fig. 1. Before the fabrication of black  $\text{Ti-Si-O}$  nanotubes,  $\text{Ti-Si-O}$  nanotube arrays were obtained through anodizing  $\text{Ti-5 wt\% Si}$  alloy ( $\text{Ti-5Si}$ ) at 40 V for 20 min in ethylene glycol-based electrolytes (mixed by 0.5 wt%

<sup>a</sup>Institute of Electronic Materials and Technology, School of Materials Science and Engineering, Shanghai Jiao Tong University, Shanghai 200240, China. E-mail: [dyding@sjtu.edu.cn](mailto:dyding@sjtu.edu.cn)

<sup>b</sup>State Key Laboratory of High Performance Ceramics and Superfine Microstructure, Shanghai Institute of Ceramics, Chinese Academy of Sciences, Shanghai 200050, China



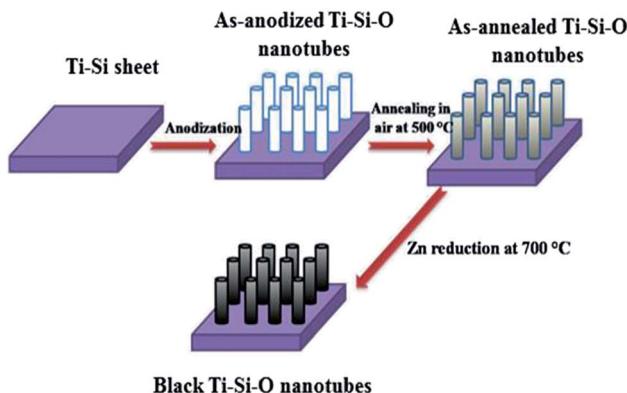


Fig. 1 Schematic diagram fabrication processes of the black Ti-Si-O nanotubes.

$\text{NH}_4\text{F}$  and 3 vol%  $\text{H}_2\text{O}$  in ethylene glycol). The as-anodized Ti-Si-O sample was heat-treated at 500 °C in an electrical furnace for 2 h to induce crystalline phase. The as-annealed Ti-Si-O sample and Zn particles were put in two ceramic boats and sent into tube furnace. Finally, black Ti-Si-O nanotubes were fabricated through Zn reduction of the as-annealed Ti-Si-O samples at 700 °C in argon atmosphere for 4 h. For reference, undoped  $\text{TiO}_2$ , Si-doped  $\text{TiO}_2$  and black undoped  $\text{TiO}_2$  nanotubes were also fabricated.

## 2.2 Material characterization

Morphology of the black Ti-Si-O nanotubes was observed with a scanning electron microscope (SEM, FEISIRION 200). Typical nanotubes morphology and corresponding element composition were characterized with a transmission electron microscope (TEM, JEOL JEM-2100) and the attached energy-dispersive X-ray (EDX) spectrometer. Phase composition and crystal structure was analyzed by Raman spectra with a 532 nm laser excitation wavelength on a thermal dispersive spectrometer. X-ray photoemission spectra (XPS) were obtained on a Kratos Axis Ultra DLD system to investigate the surface chemical states. The optical absorption property was evaluated by UV-visible absorption spectra measured on a Lambda 750S UV-vis spectrophotometer.

## 2.3 PEC measurement

The PEC water splitting properties were measured based on an electrochemical working station (CHI Instruments, model CHI660C). The conventional three-electrode system consists of black Ti-Si-O nanotubes working electrode, Pt foil counter electrode and Ag/AgCl reference electrode was put into 1.0 M KOH electrolyte solution during PEC measurement. A xenon-lamp with calibrated intensity 100 mW  $\text{cm}^{-2}$  provides the simulated light source. The photocurrent response was recorded by linear sweep voltammetry curve and amperometric  $I-t$  curve. Electrochemical impedance spectroscopy was obtained under dark condition at 0 V vs. Ag/AgCl. Mott-Schottky plots measurement was also conducted at 1000 Hz in dark.

## 2.4 Electronic structure calculation

Electronic structure calculation was carried out using the Cambridge Serial Total Energy Package (CASTEP) module based on the density functional theory (DFT) first-principles.<sup>18</sup> Si atom and oxygen vacancy were built based on  $2 \times 2 \times 1$  anatase  $\text{TiO}_2$  supercell as the computational model of black Ti-Si-O. During the simulation calculations, the Perdew-Burke-Ernzerhof (PBE) exchange-correlation function under generalized gradient approximation (GGA) was adopted.<sup>19</sup> The cut-off energy set as 340 eV, with Brillouin zone  $k$ -point mesh of  $3 \times 3 \times 3$ . The convergence criterion of electronic self-consistent energy and atom residual force was  $5 \times 10^{-5}$  eV per atom and  $10^{-3}$  eV  $\text{\AA}^{-1}$ , respectively. Electronic structure calculation was performed under these conditions after finishing geometry optimization.

## 3 Results and discussion

Fig. 2a shows top view SEM image of the Ti-Si-O nanotubes. The average nanotubes diameter was around 55 nm. From cross-sectional view SEM image (inset of Fig. 2a), it presented regular and well-aligned nanotube arrays with average nanotubes length of about 2.0  $\mu\text{m}$ . The SEM image of black Ti-Si-O nanotubes is shown in Fig. 2b. It still kept nanotubular morphology after Zn reduction at 700 °C, demonstrating that the Ti-Si-O nanotubes had good thermal stability. Typical TEM image of the black Ti-Si-O nanotubes can be found in Fig. 2c. Corresponding EDX mapping images and composition are shown in Fig. 2d-g. It revealed that Ti, Si and O elements were distributed in the nanotubes, suggesting that Si element existed in the black Ti-Si-O nanotubes.

Raman spectra were obtained to analyze the phase structure of the black Ti-Si-O nanotubes (Fig. 3). The undoped, Si-doped and black undoped  $\text{TiO}_2$  nanotubes fabricated at 700 °C were also characterized in order to identify the difference of their phase structures. For undoped  $\text{TiO}_2$ , typical Raman shift ( $\text{cm}^{-1}$ ) detected at around 142 ( $\text{B}_{1g}$ ), 234 ( $\text{E}_g$ ), 444 ( $\text{E}_g$ ) and 607 ( $\text{A}_{1g}$ ) corresponded to rutile peaks,<sup>29</sup> while the Raman modes detected at around 395 ( $\text{B}_{1g}$ ) and 513 ( $\text{A}_{1g}$ ) corresponded to anatase  $\text{TiO}_2$ .<sup>29</sup> Thus, we confirmed that rutile phase mainly existed in the undoped  $\text{TiO}_2$ . It shifted to a slight higher Raman shift and no impurity phase was found for the Si-doped system, which was caused by Si-doping. Also, the intensity of rutile Raman peaks (at around 234 ( $\text{E}_g$ ), 444 ( $\text{E}_g$ ) and 607 ( $\text{A}_{1g}$ )  $\text{cm}^{-1}$ ) decreased and it mainly consisted of anatase phase from the remnant typical anatase Raman peaks.<sup>29</sup> This indicated that Si-doping could suppress the phase transformation from anatase to rutile, the result was in agreement with previous report.<sup>30</sup> After Zn reduction, no obvious phase composition altered for the black undoped  $\text{TiO}_2$  and black Ti-Si-O. However, it is worth noting that the strongest  $\text{E}_g$  mode at 142  $\text{cm}^{-1}$  of undoped  $\text{TiO}_2$  and 144  $\text{cm}^{-1}$  of Ti-Si-O showed a blue shift after Zn reduction. In addition, the peaks intensity decreased with broadened full-width half-maximum (FWHM) (inset of Fig. 3). This could be analyzed by the induced oxygen vacancies according to previous studies, as it could cause crystal domain size and non-stoichiometry effects.<sup>22,31</sup>



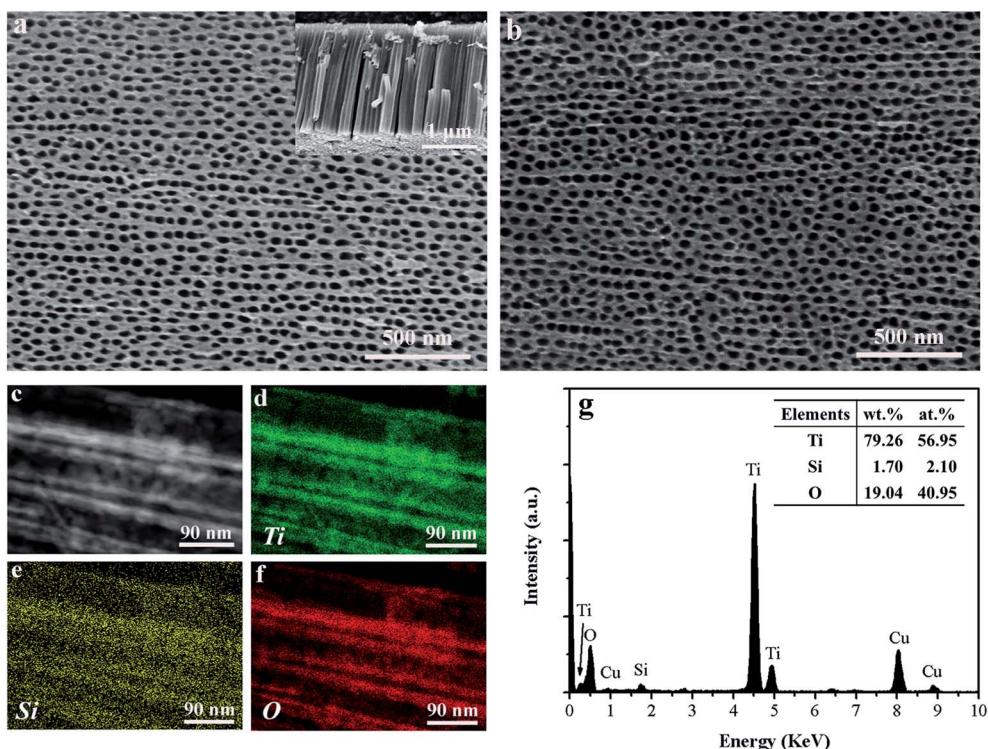


Fig. 2 (a) Surface morphology of the Ti–Si–O nanotubes, the inset shows cross-sectional view of Ti–Si–O nanotubes. (b) Top view SEM image of the black Ti–Si–O nanotubes. (c) Typical TEM image of the black Ti–Si–O nanotubes and corresponding EDX mapping images of (d) Ti, (e) Si and (f) O. (g) EDS pattern and chemical composition of the black Ti–Si–O nanotubes.

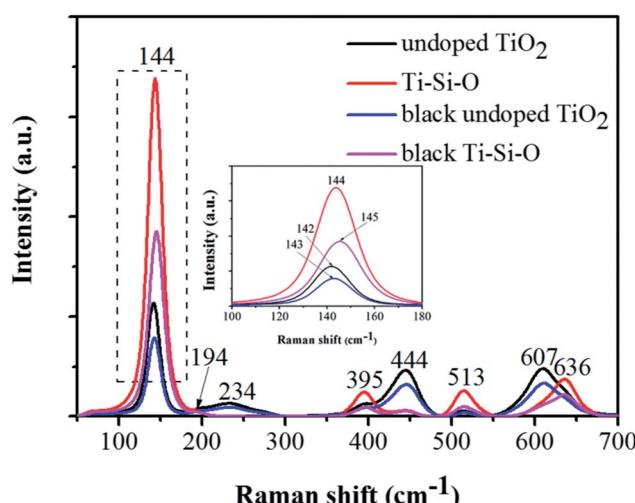


Fig. 3 Raman spectra of the undoped  $\text{TiO}_2$ , Ti–Si–O, black undoped  $\text{TiO}_2$  and black Ti–Si–O nanotubes. The inset of this figure shows enlargement peaks in the dashed line region.

Composition and chemical states of surface elements of the black Ti–Si–O nanotubes were investigated with XPS. Fig. 4a shows full XPS spectrum of the black Ti–Si–O nanotubes. Si, Ti and O elements were detected, suggesting that they existed in the black Ti–Si–O nanotubes. In addition, no evident Zn 2p peaks were found from the full spectrum, suggesting that Zn

metal may only help to reduce Ti–Si–O nanotubes and little Zn residues left on the surface of black Ti–Si–O nanotubes. This result was in accordance with our TEM EDS analysis, in which no obvious Zn signals were detected. High resolution spectra of the Si 2p, Ti 3d and O 1s of the sample surface were recorded (Fig. 4b–d). The binding energy located at around 101.9 eV for Si 2p (Fig. 4b). It was lower than the Si 2p binding energy of  $\text{SiO}_2$  (103.4 eV), indicating effective Si-doping due to the electronegativity difference between Ti (1.54) and Si (1.90).<sup>20,32,33</sup> The Ti 2p high resolution spectrum of black Ti–Si–O nanotubes was shown in Fig. 4c. Two peaks at around 458.6 eV and 464.3 eV were assigned to  $\text{Ti}^{4+}$   $2p_{3/2}$  and  $\text{Ti}^{4+}$   $2p_{1/2}$  species, respectively. While another two peaks at 457.5 eV and 463.2 eV were attributed to  $\text{Ti}^{3+}$  states.<sup>34–36</sup> The O 1s spectrum was shown in Fig. 4d. It could be fitted into Ti–O (529.6 eV), Si–O (530.4 eV) and oxygen vacancy (531.6 eV), respectively.<sup>36,37</sup> The result further confirmed that oxygen vacancies formed together with  $\text{Ti}^{3+}$  states, *i.e.*, both of them coexisted in the black Ti–Si–O nanotubes.

UV-vis absorption spectra of different  $\text{TiO}_2$  samples were measured to study the optical absorption property. As shown in Fig. 5, the absorption edge of undoped  $\text{TiO}_2$  showed a red shift after Si-doping. After Zn reduction, the absorption intensity of black undoped  $\text{TiO}_2$  intensified and the absorption edge appeared red shift compared to undoped  $\text{TiO}_2$ . Meanwhile, black Ti–Si–O nanotubes presented improved absorption intensity and the absorption edge showed a further red-shift compared to those of undoped  $\text{TiO}_2$ , Ti–Si–O and black

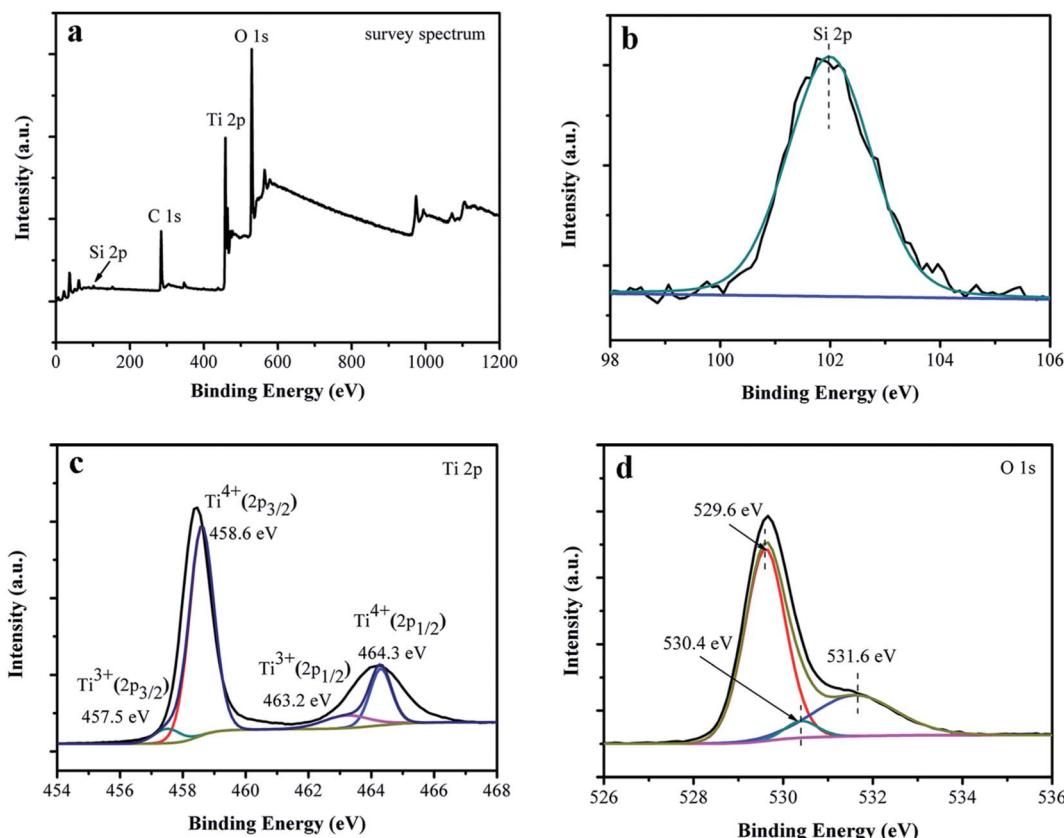


Fig. 4 XPS results of the black Ti-Si-O nanotubes. (a) Full spectrum. High resolution spectra of (b) Si 2p, (c) Ti 3d and (d) O 1s.

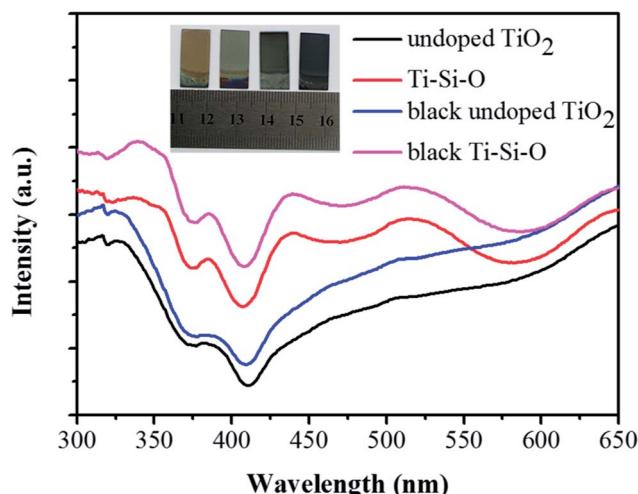


Fig. 5 UV-vis absorption spectra of the undoped  $\text{TiO}_2$ ,  $\text{Ti-Si-O}$ , black undoped  $\text{TiO}_2$  and black  $\text{Ti-Si-O}$  nanotubes. The inset shows digital photographs of different  $\text{TiO}_2$  samples.

undoped  $\text{TiO}_2$ . According to the Kubelka-Munk equation,<sup>38</sup> the extended red-shift absorption edge corresponded to a narrowed band gap for the black  $\text{Ti-Si-O}$ . Surface colors of different  $\text{TiO}_2$  nanotubes photoanodes were also observed and the digital photographs were shown in the inset of Fig. 5. The undoped  $\text{TiO}_2$  and  $\text{Ti-Si-O}$  shows pale yellow and pale blue, respectively.

While the black undoped  $\text{TiO}_2$  and black  $\text{Ti-Si-O}$  shows black color. This could also be an indirect reflection of  $\text{Ti}^{3+}$ /oxygen vacancies self-doping in the black  $\text{Ti-Si-O}$  nanotubes, as it could form color center and be favorable for visible light absorption.<sup>39,40</sup>

Transient linear sweep voltammetry ( $I-V$ ) curve and potentiostatic transient  $I-t$  response was recorded to evaluate PEC properties of the black  $\text{Ti-Si-O}$  nanotubes photoanode. Fig. 6a and b shows transient  $I-V$  curves and  $I-t$  responses of the different  $\text{TiO}_2$  nanotubes photoanodes. At the same applied bias (0 V relative to the reference electrode), the photocurrent density was  $0.35 \text{ mA cm}^{-2}$  for undoped  $\text{TiO}_2$ . It was  $0.76 \text{ mA cm}^{-2}$  for  $\text{Ti-Si-O}$ , indicating that Si-doping facilitated the separation of photogenerated carriers.<sup>27,28</sup> After Zn reduction, the photocurrent density of black undoped  $\text{TiO}_2$  increased to  $0.92 \text{ mA cm}^{-2}$  and the black  $\text{Ti-Si-O}$  nanotubes photoanode exhibited much better photocurrent response with a remarkable photocurrent density of  $1.78 \text{ mA cm}^{-2}$ . This suggested that  $\text{Ti}^{3+}$  exhibited very efficient charge separation capability of photogenerated electrons and holes. Apparently, the PEC property of the undoped  $\text{TiO}_2$  was greatly improved due to Si-doping and Zn reduction process. In other words, the synergistic effects of Si-doping and self-doped  $\text{Ti}^{3+}$  in the black  $\text{Ti-Si-O}$  nanotubes played a significant role in promoting the PEC water splitting performance.

Photoelectrode stability is an important indicator for PEC water splitting.<sup>13</sup> It can be reflected by the change of

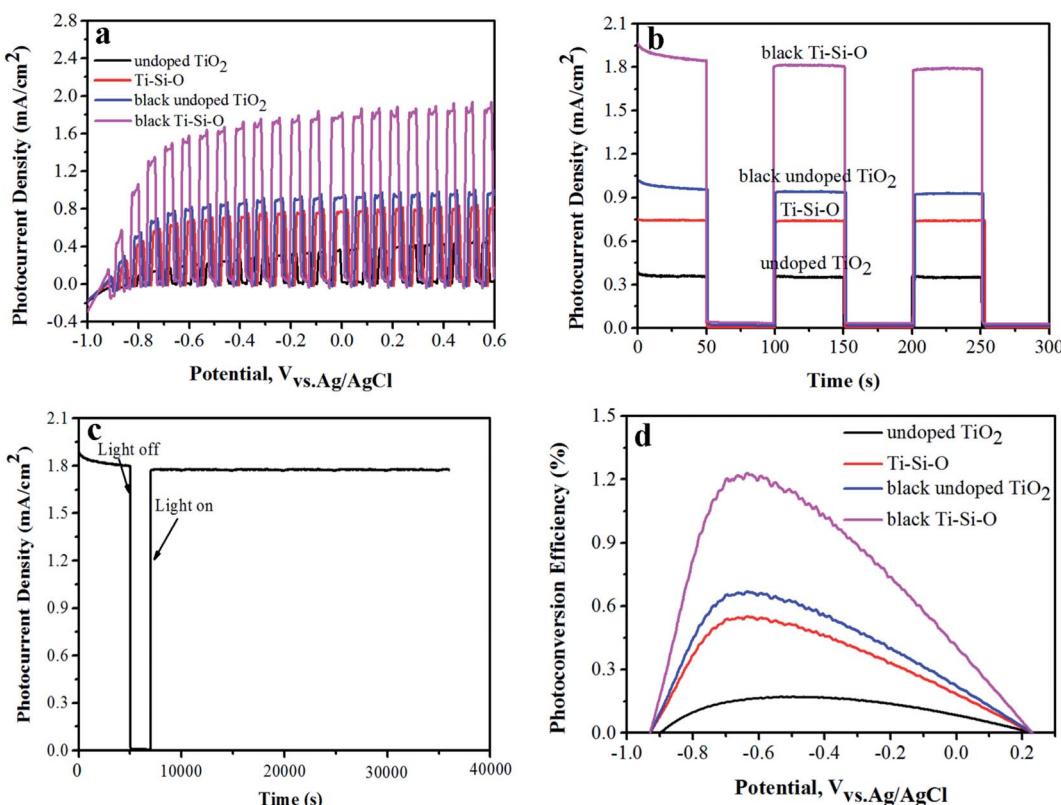


Fig. 6 (a) Transient  $I$ - $V$  curves and (b) potentiostatic transient  $I$ - $t$  response of the different TiO<sub>2</sub> nanotubes photoanodes. (c) PEC stability of the black Ti-Si-O nanotubes photoanode. (d) Photoconversion efficiency of the different TiO<sub>2</sub> nanotubes photoanodes.

photocurrent density *versus* irradiation time. The photocurrent density of black Ti-Si-O nanotubes photoanode under nearly 10 h continuous illumination was measured. It showed no obvious decay during long-time irradiation (Fig. 6c), suggesting that the black Ti-Si-O nanotubes photoanode exhibited good PEC stability and thus it was feasible used for high-efficiency PEC water splitting.

The solar-to-hydrogen photoconversion efficiency ( $\eta$ ) meets the eqn (1):<sup>41</sup>

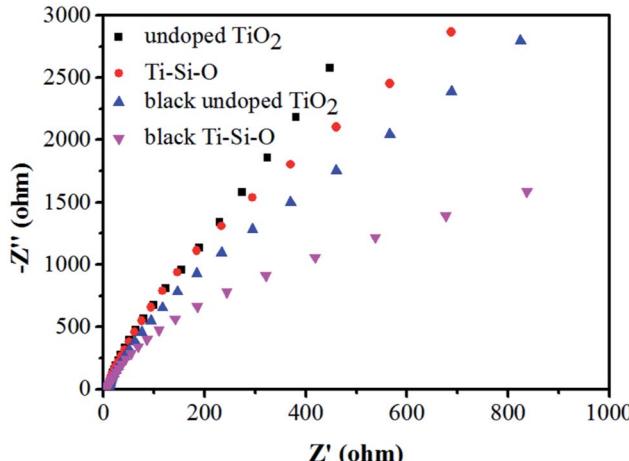


Fig. 7 Nyquist plots of the different TiO<sub>2</sub> nanotubes photoanodes.

$$\eta = I(1.23 - V_{\text{RHE}})/J_{\text{light}} \quad (1)$$

where  $J_{\text{light}}$  is the illumination intensity (100 mW cm $^{-2}$ ), and  $I$  is the photocurrent density. The reversible hydrogen potential  $V_{\text{RHE}}$  could be obtained according to the equation of  $V_{\text{RHE}} = V_{\text{Ag/AgCl}} + 0.059\text{pH} + 0.199$ ,<sup>29</sup> where  $V_{\text{Ag/AgCl}}$  represents the applied bias, and pH value is 13.6 for the KOH electrolyte.

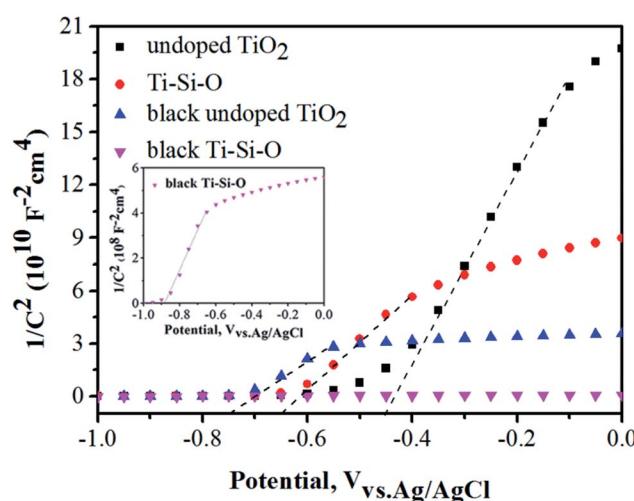


Fig. 8 MS plots of the different TiO<sub>2</sub> samples. The inset of this figure shows Mott-Schottky plots of black Ti-Si-O nanotubes photoanode.

As shown in Fig. 6d, the black Ti–Si–O nanotubes photoanode showed a peak value of 1.22%, which was about 7.18 times that of undoped  $\text{TiO}_2$  (0.17%). The black Ti–Si–O nanotubes photoanodes exhibited a higher photoconversion efficiency, which was mainly caused by  $\text{Ti}^{3+}$  species. The  $\text{Ti}^{3+}$  could capture electrons effectively and prevent the photogenerated electron–hole from rapid recombination,<sup>42,43</sup> which made the photon utilization rate increase. Therefore, the PEC property of the black Ti–Si–O photoanode was improved.

EIS spectra of the different photoanodes were also recorded to evaluate the charge transfer characters between the semiconductor electrode and the interface. Fig. 7 shows Nyquist plots of the different photoanodes. Obviously, the black Ti–Si–O exhibited a smaller impedance arc diameter compared with undoped  $\text{TiO}_2$ , Ti–Si–O and black undoped  $\text{TiO}_2$ . The smaller arc diameter indicated that it was favorable for charge transfer across the interface of semiconductor electrode and electrolyte.<sup>23</sup>

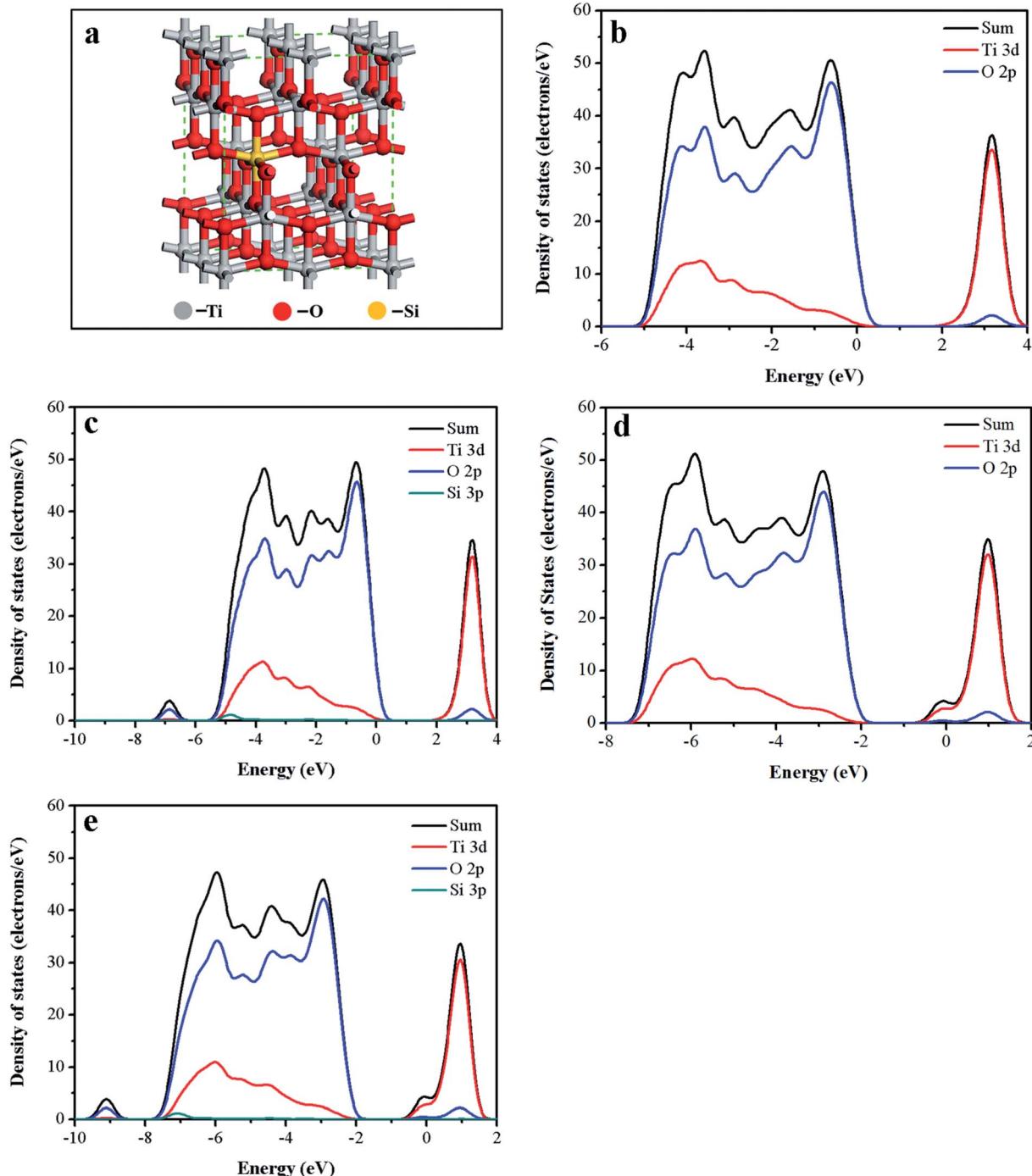


Fig. 9 (a) The computational model of black Ti–Si–O. DOS of (b) undoped  $\text{TiO}_2$ , (c) Ti–Si–O, (d) black undoped  $\text{TiO}_2$  and (e) black Ti–Si–O.

Mott-Schottky (MS) plots of the different  $\text{TiO}_2$  photoanodes were measured and shown in Fig. 8. The flat band potential ( $V_{FB}$ ) and conducting type can be acquired according to eqn (2):<sup>29,34</sup>

$$1/C^2 = (2/e_0\epsilon\epsilon_0N_d)[(V - V_{FB}) - kT/e_0] \quad (2)$$

where  $e_0$  is the fundamental electron charge,  $\epsilon_0$  and  $\epsilon$  is the permittivity of vacuum and dielectric constant of  $\text{TiO}_2$ , respectively.  $V$  is the applied potential at the electrode. As can be seen from Fig. 8, all the  $\text{TiO}_2$  samples present n-type semiconductor property according to positive slope in the MS plots.<sup>42</sup> The acquired  $V_{FB}$  of undoped  $\text{TiO}_2$ ,  $\text{Ti-Si-O}$ , black undoped  $\text{TiO}_2$  and black  $\text{Ti-Si-O}$  nanotubes photoanodes are  $-0.45$  V,  $-0.64$  V,  $-0.75$  V and  $-0.88$  V, respectively. More negative  $V_{FB}$  implied more favorable charge separation and transport efficiency when used as photoelectrode.<sup>42,44</sup>

Meanwhile, the carrier density ( $N_d$ ) can be estimated by the eqn (3):<sup>29,45</sup>

$$N_d = (2/e_0\epsilon\epsilon_0)[d(1/C^2)/dV]^{-1} \quad (3)$$

A smaller slope of the linear part in MS plots corresponded to large donor density according to previous reports.<sup>17,46</sup> Apparently, the black  $\text{Ti-Si-O}$  nanotubes photoanode presented a smaller slope than that of undoped  $\text{TiO}_2$ ,  $\text{Ti-Si-O}$ , black undoped  $\text{TiO}_2$  samples. This indicated that the as-annealed  $\text{Ti-Si-O}$  nanotubes reduced by Zn induced an increase of donor density. The donor density of black  $\text{Ti-Si-O}$  was calculated to be  $1.66 \times 10^{21} \text{ cm}^{-3}$ , which was higher than that of  $2.28 \times 10^{19} \text{ cm}^{-3}$  for black undoped  $\text{TiO}_2$ ,  $1.23 \times 10^{19} \text{ cm}^{-3}$  for  $\text{Ti-Si-O}$  and  $6.15 \times 10^{18} \text{ cm}^{-3}$  for undoped  $\text{TiO}_2$ . As a result, it could improve the conductivity and charge transport of black  $\text{Ti-Si-O}$  nanotubes photoanode. In addition, the increased donor density could raise the Fermi energy toward to conduction band, which could cause the band bending at semiconductor photoelectrode/electrolyte interface and facilitate the separation-transfer process of charge carriers.<sup>22,23</sup>

Electronic structure of the black  $\text{Ti-Si-O}$  was calculated to better understand the roles of Si-doping and  $\text{Ti}^{3+}$  self-doping. Electronic structures of the undoped  $\text{TiO}_2$ ,  $\text{Ti-Si-O}$  and black undoped  $\text{TiO}_2$  samples were also calculated as comparisons. Optimized computational model of the black  $\text{Ti-Si-O}$  was shown in Fig. 9a. Fig. 9b shows density of states (DOS) of the undoped  $\text{TiO}_2$ . Ti 3d and O 2p states constitute conduction band (CB) and valence band (VB), respectively. The calculated band gap (2.19 eV) of the undoped  $\text{TiO}_2$  was smaller than the experimental value due to self-shortcoming of the exchange-change function.<sup>47,48</sup> No foreign levels appeared in the band gap of  $\text{Ti-Si-O}$  (Fig. 9c). But the VB was broadened due to Si doping. This facilitated separation and transportation process of the photogenerated electron-hole.<sup>18</sup> Compared to undoped  $\text{TiO}_2$  and  $\text{Ti-Si-O}$ , the CB bottom and VB maximum of the black undoped  $\text{TiO}_2$  and black  $\text{Ti-Si-O}$  (Fig. 9d and e) shifted to a more negative potential. This was probably caused by  $\text{Ti}^{3+}$  self-doping, as the hybridization of Ti 3d states with O 2p and Si 3p states might be responsible for the shift of VB and CB.<sup>40,48</sup> The

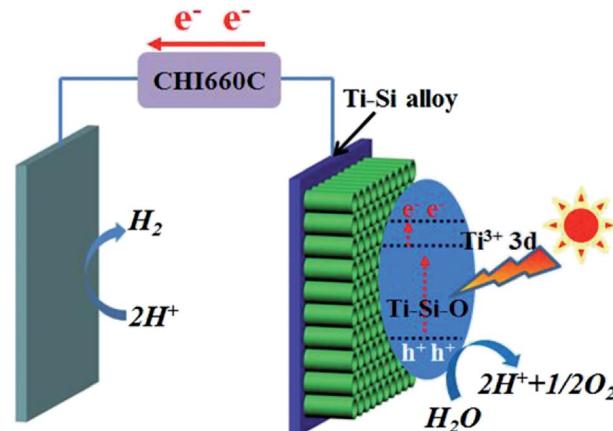


Fig. 10 Energy diagram and charge transfer in the black  $\text{Ti-Si-O}$  nanotubes photoanode for PEC water splitting.

negative shift of the energy band made the Fermi energy rise closely to conduction band. When the black  $\text{Ti-Si-O}$  nanotubes photoanode was in contact with electrolyte, the band bending degree would increase and thus enhance the separation process of the photoinduced electrons and holes. Therefore,  $\text{H}^+$  was prone to be reduced to  $\text{H}_2$  by the photogenerated electrons during PEC water splitting hydrogen evolution reaction process. In addition, for both the undoped  $\text{TiO}_2$  and Si-doped  $\text{TiO}_2$ , the black process could induce the formation of miniband rising up closely below the conduction band minimum. This could result in narrowed band gap according to previous reports,<sup>22,49</sup> which could enhance optical absorption shown by UV-vis absorption spectra.

Based on above experimental results and analyses, possible mechanism of the black  $\text{Ti-Si-O}$  nanotubes photoanode for PEC water splitting was proposed. Energy diagram and charge transfer in the black  $\text{Ti-Si-O}$  nanotubes photoanode was illustrated in Fig. 10.  $\text{Ti}^{3+}$  3d states were introduced into the  $\text{Ti-Si-O}$  nanotubes and thus produced donor states below the conduction band.<sup>39</sup> The  $\text{Ti}^{3+}$  could form conduction band tail, which could trap photogenerated electrons and decrease the recombination rate of the charge carriers.<sup>22,49</sup> Furthermore, the narrowed band gap could improve optical absorption.<sup>31</sup> Moreover, Si-doping favored the separation process of the photoinduced electrons and holes.<sup>19,50</sup> Once illuminated, the photogenerated holes would accumulate at the surface of the  $\text{Ti-Si-O}$  nanotubes for water oxidation ( $\text{H}_2\text{O} + 2\text{h}^+ \rightarrow 2\text{H}^+ + 1/2\text{O}_2$ ). Meanwhile, the photogenerated electrons quickly transferred to the cathode across the external circuit and participated in hydrogen evolution reaction ( $2\text{H}^+ + 2\text{e}^- \rightarrow \text{H}_2$ ). As a result, the PEC water splitting performance of black Si-doped  $\text{TiO}_2$  nanotubes photoanode was significantly improved.

## 4 Conclusions

In summary, black Si-doped  $\text{TiO}_2$  nanotubes photoanode was prepared by Zn metal reduction of the as-annealed  $\text{Ti-Si-O}$  nanotubes in argon atmosphere. Si element and  $\text{Ti}^{3+}$ /oxygen vacancies were introduced into the black  $\text{Ti-Si-O}$  nanotubes,



resulting in much better PEC properties than those of undoped  $\text{TiO}_2$  and Si-doped  $\text{TiO}_2$ . The improved PEC performance was mainly attributed to synergistic results of enhanced optical absorption, favorable charge separation and transfer property. A higher photocurrent density of  $1.78 \text{ mA cm}^{-2}$  at  $0 \text{ V}$  vs.  $\text{Ag}/\text{AgCl}$  was obtained. The maximum photoconversion efficiency was 1.22%, which was about 7.18 times the photoconversion efficiency of undoped  $\text{TiO}_2$ . We expected that these findings may contribute to develop high-performance  $\text{TiO}_2$ -based nanostructures photoanodes for improving PEC water splitting.

## Conflicts of interest

The authors declare no competing financial interest.

## Acknowledgements

This work was supported by National Natural Science Foundation of China (No. 51572170). We thank the contribution from SEM and XPS labs at Instrumental Analysis Center of SJTU.

## References

- 1 A. Wolcott, W. A. Smith, T. R. Kuykendall, Y. Zhao and J. Z. Zhang, *Small*, 2009, **5**, 104–111.
- 2 A. Fujishima and K. Honda, *Nature*, 1972, **238**, 37–38.
- 3 S. U. M. Khan, M. Al-Shahry and W. B. Ingler, *Science*, 2002, **297**, 2243–2245.
- 4 M. Ni, M. K. H. Leung, D. Y. C. Leung and K. Sumathy, *Renewable Sustainable Energy Rev.*, 2007, **11**, 401–425.
- 5 S. Martha, P. C. Sahoo and K. M. Parida, *RSC Adv.*, 2015, **5**, 61535–61553.
- 6 S. M. Gupta and M. Tripathi, *Chin. Sci. Bull.*, 2011, **56**, 1639–1657.
- 7 P. Roy, S. Berger and P. Schmuki, *Angew. Chem., Int. Ed.*, 2011, **50**, 2904–2939.
- 8 F. Dong, W. Zhao and Z. Wu, *Nanotechnology*, 2008, **19**, 365607–365616.
- 9 C. C. Chuang, C. K. Lin, T. T. Wang, V. Srinivasadesikan, P. Raghunath and M. C. Lin, *RSC Adv.*, 2015, **5**, 81371–81377.
- 10 S. B. Rawal, S. Bera, D. Lee, D. J. Jang and W. I. Lee, *Catal. Sci. Technol.*, 2013, **3**, 1822–1830.
- 11 M. Gholami, M. Qorbani, O. Moradlou, N. Naseri and A. Z. Moshfegh, *RSC Adv.*, 2014, **4**, 783–7844.
- 12 T. Sreethawong and S. Yoshikawa, *Chem. Eng. J.*, 2012, **197**, 272–282.
- 13 R. Liang, A. Hu, J. Persic and Y. N. Zhou, *Nano-Micro Lett.*, 2013, **5**, 202–212.
- 14 Y. X. Pan, T. Zhou, J. Han, J. Hong, Y. Wang, W. Zhang and R. Xu, *Catal. Sci. Technol.*, 2016, **6**, 2206–2213.
- 15 K. Shin, J. B. Yoo and J. H. Park, *J. Power Sources*, 2013, **225**, 263–268.
- 16 X. Chen, L. Liu and F. Huang, *Chem. Soc. Rev.*, 2015, **44**, 1861–1885.
- 17 G. Wang, H. Wang, Y. Ling, Y. Tang, X. Yang, R. C. Fitzmorris, C. Wang, J. Z. Zhang and Y. Li, *Nano Lett.*, 2011, **11**, 3026–3033.
- 18 W. Shi, Q. Chen, Y. Xu, D. Wu and C. Huo, *Appl. Surf. Sci.*, 2011, **257**, 3000–3006.
- 19 K. Yang, Y. Dai and B. Huang, *Chem. Phys. Lett.*, 2008, **456**, 71–75.
- 20 Y. Su, S. Chen, X. Quan, H. Zhao and Y. Zhang, *Appl. Surf. Sci.*, 2008, **255**, 2167–2172.
- 21 Z. Wang, C. Yang, T. Lin, H. Yin, P. Chen, D. Wan, F. Xu, F. Huang, J. Lin, X. Xie and M. Jiang, *Adv. Funct. Mater.*, 2013, **23**, 5444–5450.
- 22 Z. Wang, C. Yang, T. Lin, H. Yin, P. Chen, D. Wan, F. Xu, F. Huang, J. Lin, X. Xie and M. Jiang, *Energy Environ. Sci.*, 2013, **6**, 3007–3014.
- 23 Q. Kang, J. Cao, Y. Zhang, L. Liu, H. Xu and J. Ye, *J. Mater. Chem. A*, 2013, **1**, 5766–5774.
- 24 M. Xing, W. Fang, M. Nasir, Y. Ma, J. Zhang and M. Anpo, *J. Catal.*, 2013, **297**, 236–243.
- 25 C. Xu, Y. Song, L. Lu, C. Cheng, D. Liu, X. Fang, X. Chen, X. Zhu and D. Li, *Nanoscale Res. Lett.*, 2013, **8**, 391–397.
- 26 J. Dong, J. Han, Y. Liu, A. Nakajima, S. Matsushita, S. Wei and W. Gao, *ACS Appl. Mater. Interfaces*, 2014, **6**, 1385–1388.
- 27 R. Jin, Z. Wu, Y. Liu, B. Jiang and H. Wang, *J. Hazard. Mater.*, 2009, **161**, 42–48.
- 28 Y. Zhang, X. Li, D. Chen, N. Ma, X. Hua and H. Wang, *Scr. Mater.*, 2009, **60**, 543–546.
- 29 F. Cao, J. Xiong, F. Wu, Q. Liu, Z. Shi, Y. Yu, X. Wang and L. Li, *ACS Appl. Mater. Interfaces*, 2016, **8**, 12239–12245.
- 30 H. Zhang, X. Quan, S. Chen and H. Zhao, *Environ. Sci. Technol.*, 2006, **40**, 6104–6109.
- 31 L. Shen, Z. Xing, J. Zou, Z. Li, X. Wu, Y. Zhang, Q. Zhu, S. Yang and W. Zhou, *Sci. Rep.*, 2017, **7**, 41978–41988.
- 32 X. Zhang, B. Zhang, Z. Zuo, M. Wang and Y. Shen, *J. Mater. Chem. A*, 2015, **3**, 10020–10025.
- 33 J. Zhu, D. Zhang, Z. Bian, G. Li, Y. Huo, Y. Lu and H. Li, *Chem. Commun.*, 2009, 5394–5396.
- 34 H. Cui, W. Zhao, C. Yang, H. Yin, T. Lin, Y. Shan, Y. Xie, H. Gu and F. Huang, *J. Mater. Chem. A*, 2014, **2**, 8612–8616.
- 35 H. Khan and I. K. Swati, *Ind. Eng. Chem. Res.*, 2016, **55**, 6619–6633.
- 36 X. Zhang, H. Tian, X. Wang, G. Xue, Z. Tian, J. Zhang, S. Yuan, T. Yu and Z. Zou, *Mater. Lett.*, 2013, **100**, 51–53.
- 37 C. Xie, Q. Yang, Z. Xu, X. Liu and Y. Du, *J. Phys. Chem. B*, 2006, **110**, 8587–8592.
- 38 J. J. Sene, W. A. Zeltner and M. A. Anderson, *J. Phys. Chem. B*, 2003, **107**, 1597–1603.
- 39 W. Fang, M. Xing and J. Zhang, *Appl. Catal., B*, 2014, **160**, 240–246.
- 40 Z. Wang, B. Wen, Q. Hao, L. M. Liu, C. Zhou, X. Mao, X. Lang, W. J. Yin, D. Dai, A. Selloni and X. Yang, *J. Am. Chem. Soc.*, 2015, **137**, 9146–9152.
- 41 Z. Zhang and P. Wang, *Energy Environ. Sci.*, 2012, **5**, 6506–6512.
- 42 Q. Liu, D. Ding, C. Ning and X. Wang, *Int. J. Hydrogen Energy*, 2015, **40**, 2107–2114.
- 43 F. Teng, M. Li, C. Gao, G. Zhang, P. Zhang, Y. Wang, L. Chen and E. Xie, *Appl. Catal., B*, 2014, **148**, 339–343.
- 44 H. Cai, Q. Yang, Z. Hu, Z. Duan, Q. You, J. Sun, N. Xu and J. Wu, *Appl. Phys. Lett.*, 2014, **104**, 053114–053117.



- 45 L. Yu, J. He, C. Huang, M. Li, Y. Zhang, X. Zhou and H. Zhu, *RSC Adv.*, 2017, **7**, 54485–54490.
- 46 C. Mao, F. Zuo, Y. Hou, X. Bu and P. Feng, *Angew. Chem., Int. Ed.*, 2014, **53**, 10485–10489.
- 47 M. E. Kurtoglu, T. Longenbach, K. Sohlberg and Y. Gogotsi, *J. Phys. Chem. C*, 2011, **115**, 17392–17399.
- 48 T. Zhang, Z. Zhu, H. Chen, Y. Bai, S. Xiao, X. Zheng, Q. Xue and S. Yang, *Nanoscale*, 2015, **7**, 2933–2940.
- 49 X. Chen, L. Liu, P. Y. Yu and S. S. Mao, *Science*, 2011, **331**, 746–750.
- 50 W. Shi, Q. Chen, Y. Xu, D. Wu and C. F. Huo, *J. Solid State Chem.*, 2011, **184**, 1983–1988.