



Cite this: *RSC Adv.*, 2018, 8, 4372

# Switching properties of epitaxial $\text{La}_{0.5}\text{Sr}_{0.5}\text{CoO}_3/\text{Na}_{0.5}\text{Bi}_{0.5}\text{TiO}_3/\text{La}_{0.5}\text{Sr}_{0.5}\text{CoO}_3$ ferroelectric capacitor

J. M. Song,<sup>ab</sup> L. H. Luo,<sup>c</sup> X. H. Dai,<sup>a</sup> A. Y. Song,<sup>a</sup> Y. Zhou,<sup>a</sup> Z. N. Li,<sup>a</sup> J. T. Liang<sup>a</sup> and B. T. Liu<sup>\*a</sup>

$\text{La}_{0.5}\text{Sr}_{0.5}\text{CoO}_3/\text{Na}_{0.5}\text{Bi}_{0.5}\text{TiO}_3/\text{La}_{0.5}\text{Sr}_{0.5}\text{CoO}_3$  (LSCO/NBT/LSCO) ferroelectric capacitors have been successfully fabricated on (001)  $\text{SrTiO}_3$  substrate, in which the LSCO film is prepared by magnetron sputtering and the NBT film by pulsed laser deposition. Both X-ray diffraction and transmission electron microscopy techniques confirm that the (001) oriented LSCO/NBT/LSCO heterostructure is epitaxially grown on  $\text{SrTiO}_3$  substrate. The remnant polarization, coercive field and relative dielectric constant of the LSCO/NBT/LSCO capacitor, measured at  $250 \text{ kV cm}^{-1}$ , are  $15.6 \mu\text{C cm}^{-2}$ ,  $47 \text{ kV cm}^{-1}$  and 559, respectively. Moreover, the capacitor possesses very good fatigue-resistance, and less pulse width dependence as well as piezoelectric properties ( $d_{33} = 145 \text{ pm V}^{-1}$ ). It is found that the leakage current density of the LSCO/NBT/LSCO capacitor meets well with ohmic conduction behavior at applied fields lower than  $55 \text{ kV cm}^{-1}$  and bulk-limited space charge-limited conduction at the fields higher than  $55 \text{ kV cm}^{-1}$ .

Received 20th November 2017

Accepted 17th January 2018

DOI: 10.1039/c7ra12575e

rsc.li/rsc-advances

## 1. Introduction

Ferroelectric thin film materials have been extensively investigated due to their good ferroelectric, dielectric, piezoelectric and optical properties for various kinds of applications, *e.g.* phase shifters, microelectromechanical systems, sensors, and ferroelectric random access memory (FeRAM).<sup>1–3</sup> Lead based  $\text{Pb}(\text{Zr}_{1-x}\text{Ti}_x)\text{O}_3$  (PZT) films are intensively studied worldwide including by our group due to their large remnant polarization and small coercive voltage.<sup>4–6</sup> However, PZT films are not environmentally friendly materials, thus many efforts have been made to develop lead-free new type ferroelectric materials.<sup>7–8</sup> Sodium bismuth titanate,  $\text{Na}_{0.5}\text{Bi}_{0.5}\text{TiO}_3$  (NBT), with A-site composite perovskite structure is considered as an excellent candidate to replace lead based ferroelectric materials.<sup>9,10</sup> Recently, we have successfully prepared a single-phase high quality NBT bulk material by solid state reaction method, and investigated the optical properties and ferroelectric properties.<sup>11,12</sup> The  $\text{Na}_{0.5}\text{Bi}_{0.5}\text{TiO}_3$  ceramics have good ferroelectric properties at room temperature and relatively high Curie temperature ( $320 \text{ }^\circ\text{C}$ ), which makes ferroelectric NBT film based device very promising.<sup>13</sup> Polycrystalline NBT films with good ferroelectric remnant polarization ( $\sim 11.9 \mu\text{C cm}^{-2}$ ) have been

grown on Pt/Ti/SiO<sub>2</sub>/Si substrates using various methods such as radio-frequency magnetron sputtering and metal organic decomposition process.<sup>14,15</sup> Compared to the polycrystalline films, it is believed that the highly oriented ferroelectric films (especially the epitaxial films) can possess higher polarization and less leakage current due to the lack of grain boundaries in the epitaxial films. Highly (111) oriented NBT film was prepared on Pt/Ti/SiO<sub>2</sub>/Si substrate by a sol-gel process, it is found that the remnant polarization and coercive field are  $20.6 \mu\text{C cm}^{-2}$  and  $112 \text{ kV cm}^{-1}$ , respectively.<sup>16</sup> (001) and (011) oriented epitaxial NBT films have been successfully fabricated on platinum coated MgO and  $\text{SrTiO}_3$  substrates by pulsed laser deposition, and very good dielectric and ferroelectric properties have been reported.<sup>17,18</sup> Compared to Pt electrode, the cheaper LSCO electrode can provide oxide/oxide (LSCO/NBT) interface, different from the metal/oxide (Pt/NBT) interface, which may further impact the NBT physical properties. From the viewpoint of our best knowledge, no results have been reported on the fabrication and characterization of epitaxial NBT capacitor sandwiched by oxide electrodes. In this paper, we report on the fabrication of the epitaxial LSCO/NBT/LSCO/ $\text{SrTiO}_3$  (STO) heterostructure, and the structural and switching properties of the LSCO/NBT/LSCO ferroelectric capacitors have been systematically investigated.

## 2. Experimental

Briefly, the Pt/LSCO/NBT/LSCO/STO heterostructure was prepared by a multi-step procedure. Step 1: 60 nm-thick LSCO

<sup>a</sup>Hebei Key Lab of Optic-Electronic Information and Materials, College of Physics Science & Technology, Hebei University, Baoding 071002, China. E-mail: btlui@hbu.cn

<sup>b</sup>College of Science, Agriculture University of Hebei, Baoding 071001, China

<sup>c</sup>College of Science, Ningbo University, Ningbo 315211, China



thin film was deposited on (001) STO single crystal substrate by magnetron sputtering at room temperature under the following conditions: Ar : O<sub>2</sub> = 3 : 1, power = 30 W. Post-annealing was conducted at 550 °C in a 1 atm oxygen-flowing tube furnace. Step 2: Na<sub>0.5</sub>Bi<sub>0.5</sub>TiO<sub>3</sub> target with 10% excessive bismuth and 10% sodium was used during deposition to compensate the loss of bismuth and sodium due to the strong volatility of bismuth and sodium. 250 nm-thick NBT film was further deposited on the LSCO/STO heterostructure by pulsed laser deposition at 550 °C and 7.5 Pa oxygen deposition pressure. The target and substrate distance was 4.5 cm; the laser energy density and repetition rate were 2 J cm<sup>-2</sup> and 5 Hz, respectively. Step 3: 60 nm-thick LSCO film and 70 nm thick Pt were sequentially deposited by magnetron sputtering on the surface of the NBT/LSCO/STO heterostructure through a shadow mask to get 7.85 × 10<sup>-5</sup> cm<sup>2</sup> circular pads as the top electrodes of the capacitors. The Pt film is used to improve the electrical contact between the measurement tip and LSCO. The Pt/LSCO/NBT/LSCO/STO heterostructure was further annealed at 550 °C in an oxygen-flowing tube furnace in order to make the as-grown top LSCO film crystallized. The crystallinity and phase of the sample was characterized by X-ray diffraction (XRD) and the cross-sectional interfaces of NBT/LSCO/STO heterostructure were investigated using transmission electron microscopy (TEM). The ferroelectric properties and leakage current behavior of the LSCO/NBT/LSCO capacitor were studied using a precision LC unit from Radiant Technologies. The dielectric and piezoelectric properties were measured using an Agilent 4294A and piezoresponse force microscope (PFM-Asylum Research), respectively.

### 3. Results and discussion

Fig. 1(a) presents the theta–2theta XRD spectrum of the NBT/LSCO/STO heterostructure. Besides the STO peaks, only (001), (002) diffraction peaks of NBT and LSCO are observed without any detectable secondary phase or diffraction peaks from other directions, indicating that both NBT and LSCO films are highly (001) oriented. In order to further prove the epitaxial property of NBT film on the LSCO/STO heterostructure, we performed the phi-scan of (110) plane of NBT film. As shown in the inset of Fig. 1, the four evenly spaced peaks with spacing of 90° reveal that the NBT film was epitaxially grown on NBT/LSCO/STO heterostructure. Fig. 1(b) presents the TEM image of the NBT/LSCO and LSCO/STO interfaces, no obvious interdiffusion can be found at the interfaces. Selected area diffraction (SAED) of NBT and LSCO films derived from fast Fourier transform (FFT) are respectively presented in the Fig. 1(b), indicating both NBT and LSCO are single crystalline films. Fig. 1(c) shows a typical high-resolution TEM image of the LSCO/STO heterostructure and single crystalline LSCO film is epitaxially grown on the surface of (001) STO substrate. The orientation relationship of these two layers is (001)LSCO//(001)STO and [100]LSCO//[100]STO. Fig. 1(d) shows a HRTEM image of the NBT/LSCO interface region viewed along the [100] zone axis of STO substrate. The interface is atomically sharp, indicating a very good epitaxial growth relationship exists between NBT and LSCO. The orientation relationship between the NBT film and the

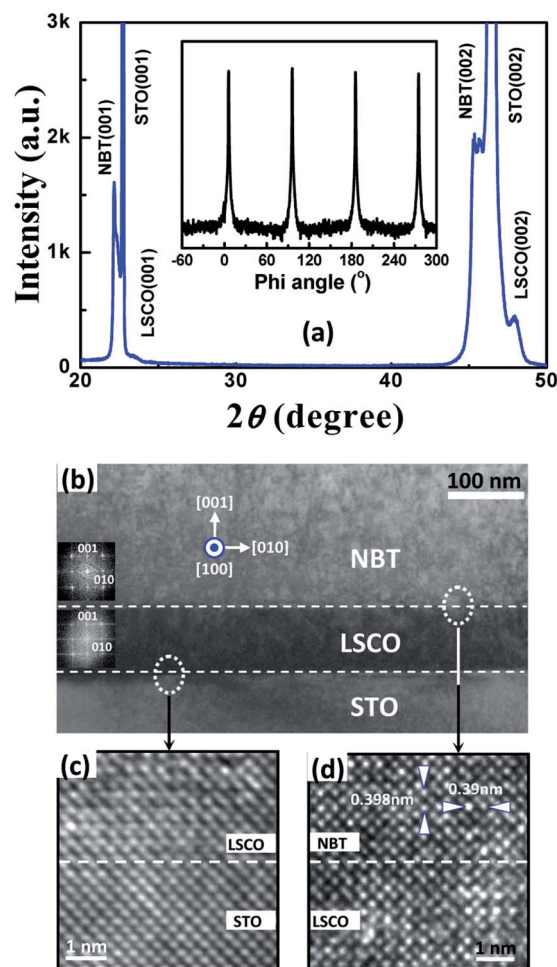


Fig. 1 XRD spectrum of NBT/LSCO/STO heterostructure, the inset is the phi-scan pattern of the NBT film (a). TEM image of the cross sectional image of the NBT/LSCO/STO heterostructure (b). The high-resolution TEM images of LSCO/STO (c) and NBT/LSCO (d).

LSCO film in the NBT/LSCO heterostructure is as follows: (001)NBT//(001)LSCO, [100]NBT//[100]LSCO, which is expected since the NBT (0.389 nm) has a good lattice match with the underlying LSCO (0.384 nm). Based on the SAED of NBT film in the Fig. 1(b), the lattice constants are 0.390 nm and 0.398 nm for the in-plane (*a* axis) and the out-plane (*c* axis), respectively. The *c* axis constant is consistent with the value, 0.397 nm, obtained from the XRD measurement. We should note that the out-plane lattice constant of NBT (0.397 nm) is larger than that of its bulk material (0.389 nm). A compressive force exists in the NBT film when NBT is epitaxially grown on LSCO surface so that out-plane lattice constant of NBT becomes larger. The polarization *P<sub>s</sub>* of a ferroelectric material is related to its tetragonality, (*P<sub>s</sub>*)<sup>2</sup> ∝ (*c/a* - 1)<sup>2</sup>, and the tetragonality (*c/a* = 1.02) of the epitaxial NBT film in the LSCO/NBT/LSCO heterostructure is larger than its bulk value (*c/a* = 1.002), which is favorable for the NBT film to have a larger remnant polarization.

A typical ferroelectric hysteresis loop of the LSCO/NBT/LSCO capacitor, measured at 250 kV cm<sup>-1</sup> bias electric field and 1 kHz frequency using capacitive method, is illustrated in Fig. 2(a).



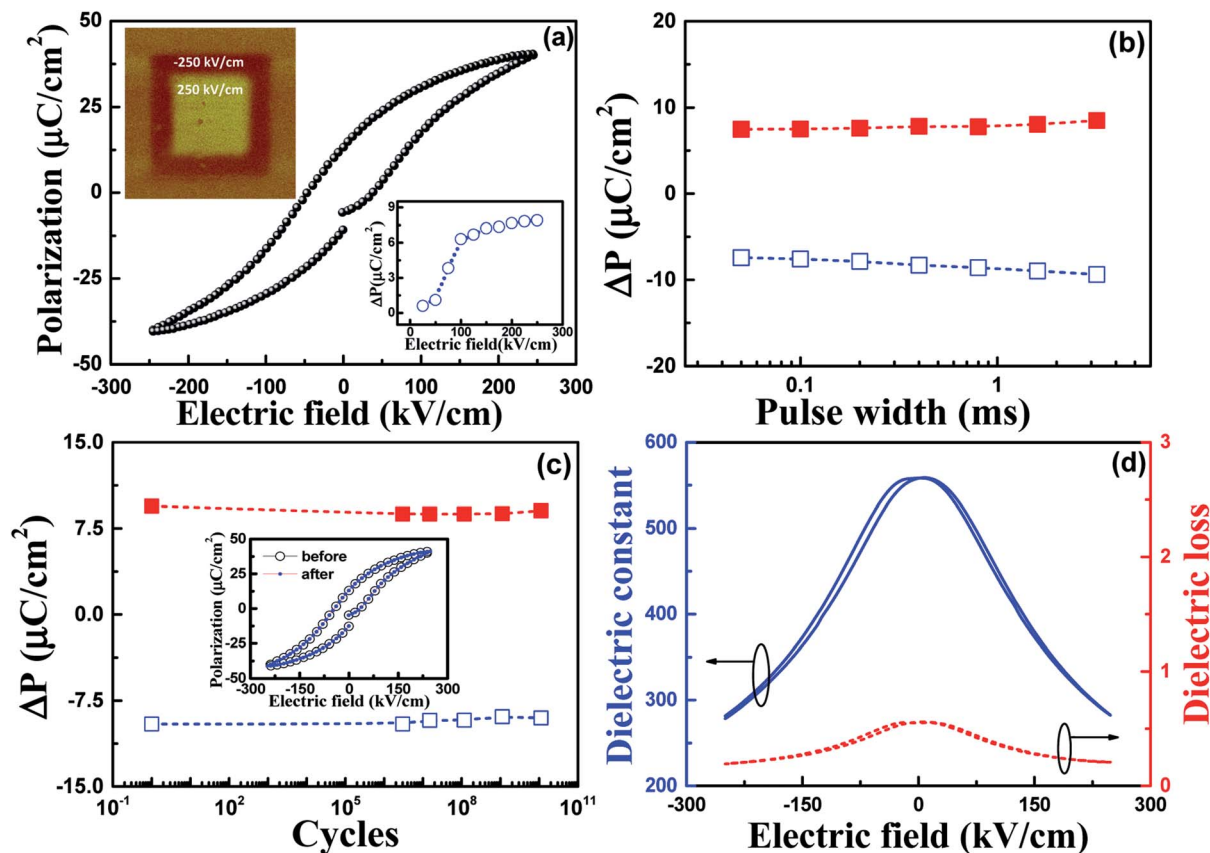


Fig. 2 Ferroelectric hysteresis loop of the LSCO/NBT/LSCO capacitor, the inset at the top left-hand side shows the phase PFM images of NBT film, the inset at the bottom right-hand side presents relation of  $\Delta P$  vs. applied field (a). Relation of  $\Delta P$  vs. pulse width of the LSCO/NBT/LSCO capacitor (b). Relation of  $\Delta P$  vs. frequency of the LSCO/NBT/LSCO capacitor, the inset is the ferroelectric hysteresis loops before and after fatigue tests (c). Relation of dielectric constant and dielectric loss vs. applied field of the LSCO/NBT/LSCO capacitor (d).

A saturated ferroelectric loop can be obtained with a remnant polarization,  $15.6 \mu\text{C cm}^{-2}$ , and a small coercive field,  $47 \text{ kV cm}^{-1}$ . The inset at the top left-hand side of Fig. 2(a) is the phase PFM images, which recorded for sample NBT after square areas of ( $7 \times 7 \mu\text{m}^2$ ) and ( $4 \times 4 \mu\text{m}^2$ ) have been polarized (black region) and reverse polarized (white region) by applying bias electric field  $-250 \text{ kV cm}^{-1}$  and  $250 \text{ kV cm}^{-1}$ , respectively. We can see the polarization switching within these domains as square areas defined by black and white contrasts, characteristic of each orientation of the polarization can be observed. This means that (001) epitaxial NBT thin film has good reversal memory properties. The switchable polarization,  $\Delta P$ , is defined as the difference between the switched ( $P^*$ ) and nonswitched ( $P^\wedge$ ) polarization. The inset at the bottom right-hand side of Fig. 2(a) demonstrates the relation of switchable polarization with applied electric field. We can see that the switchable polarization increases rapidly with the increase of the applied electric field, and it is saturated when the applied electric field is greater than  $100 \text{ kV cm}^{-1}$ . Fig. 2(b) shows the pulse width dependence of the LSCO/NBT/LSCO capacitor measured at bias electric field  $250 \text{ kV cm}^{-1}$  with a pulse delay time of 1 s. It can be seen that the switchable polarization slightly increases with the increase of the pulse width, which is favorable to the high speed ferroelectric memories. In addition, fatigue is also an important

parameter to characterize ferroelectric film for ferroelectric memories. In this experiment, the fatigue characteristic of the LSCO/NBT/LSCO capacitor was tested at bias electric field  $250 \text{ kV cm}^{-1}$  and 1 MHz frequency with a pulse delay time of 1 s. As shown in Fig. 2(c), no obvious degradation can be found from the polarizations of the capacitor up to  $10^{10}$  switching cycles. The inset of Fig. 2(c) presents the hysteresis loops before and after the fatigue test, in which no obvious difference can be found from the hysteresis loops, further indicating that the epitaxial LSCO/NBT/LSCO capacitor has very good fatigue resistance characteristic. Fig. 2(d) demonstrates the relation of dielectric constant and loss tangent with applied electric field for the LSCO/NBT/LSCO capacitor at a frequency of 10 kHz. The tunability of 50.1% at the maximum applied field of  $250 \text{ kV cm}^{-1}$  is obtained, in which the tunability is defined<sup>19</sup>

$$\text{Tunability} = \frac{\varepsilon_{\text{max}} - \varepsilon_{\text{min}}}{\varepsilon_{\text{max}}} \quad (1)$$

In formula (1),  $\varepsilon_{\text{min}}$  is the minimum value of the dielectric constant at the maximum applied field,  $\varepsilon_{\text{max}}$  is the maximum value of the dielectric constant at the applied field of  $0 \text{ kV cm}^{-1}$ . The dielectric constant and dielectric loss of the (001) epitaxial NBT film measured at 10 kHz are 559 and 0.19, respectively.



To understand the local switching behavior of the NBT thin films in ferroelectrics, the piezoresponse force microscope PFM (Asylum Research, Cypher, American) was used to characterize piezoelectric properties of the NBT film. Since the effective piezoelectric coefficient  $d_{33}$  can reflect the most essential piezoelectric effect of piezoelectric film material.<sup>20</sup> Measurement of  $d_{33}$  was achieved by keeping the PFM tip fixed above the local point of NBT film and applying an electric field from  $-500$  to  $500$   $\text{kV cm}^{-1}$ , while recording the piezoresponse signal amplitude ( $A$ ) and applied field ( $E$ ). In the measurements, a bias voltage of  $800$  mV and a resonance frequency of  $320$  Hz were applied on the conductive test probe. According to the law of the converse piezoelectric effect, the relationship between amplitude ( $A$ ) and applied field ( $E$ ) can be described as follows:<sup>21</sup>

$$d_{33} = \frac{A - A_0}{d(E - E_0)} \quad (2)$$

In formula (2),  $d$  is the initial thickness of the films before deformation,  $A_0$  and  $E_0$  are the piezoelectric deformation and electric field of the intersection, respectively. The piezoelectric hysteresis loop ( $d_{33}$ - $E$ ) of NBT film is calculated from the  $A$ - $E$  curve based on eqn (1). Fig. 3(a) and (b) displays the corresponding piezoelectric phase-electric field loop and typical piezoelectric response butterfly loop. In the piezoelectric phase electric field loop, the film exhibits a significant  $180^\circ$  domain

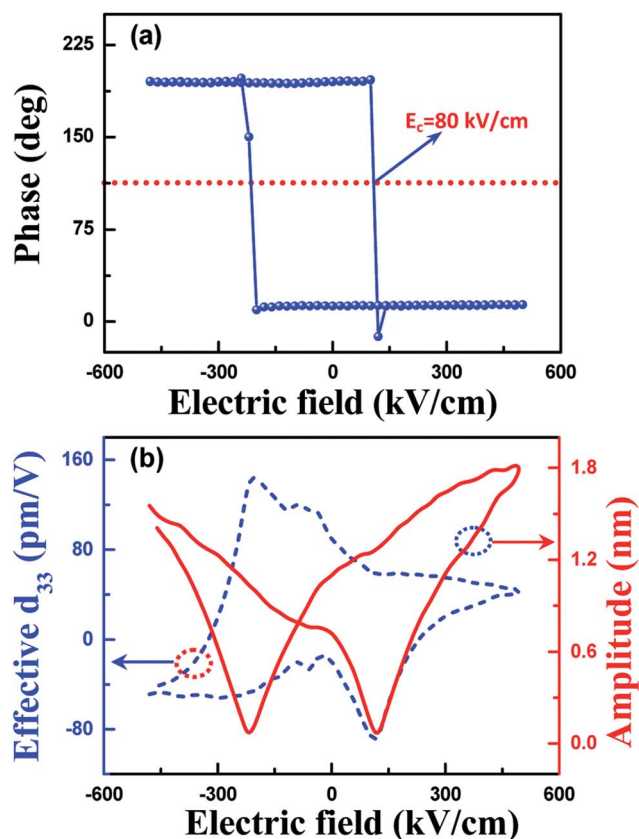


Fig. 3 Piezoelectricity of the NBT film obtained by PFM. (a) Phase loop of the NBT film. (b) Ferroelectric butterfly loop and the corresponding piezoelectric hysteresis loop of the NBT film.

flip. As can be seen, a typical well-shaped  $A$ - $E$  “butterfly” loop is obtained with a amplitude maximum of  $1.82$  nm appearing at  $475$   $\text{kV cm}^{-1}$ . This result shows a strain as high as  $\sim 0.7\%$  ( $A/d$ ) and indicates that the  $(00l)$  orientation of the epitaxial NBT film has a very good piezoelectricity. In addition, at  $-205$   $\text{kV cm}^{-1}$ , the  $d_{33}$  piezoelectric coefficients reached the  $145$   $\text{pm V}^{-1}$ , which is relative higher than the reported NBT-based films.<sup>22–24</sup>

Leakage current is an important parameter of the electrical properties of ferroelectric capacitor related to thermal effect and energy loss so that minimum leakage current density is ideal for the real capacitors used in the ferroelectric memories. Four kinds of leakage mechanisms have been proposed for ferroelectric perovskite oxides: the ohmic conduction, the bulk-limited space-charge-limited conduction (SCLC), the bulk-limited Poole-Frenkel emission and the interface-limited Schottky emission.<sup>25–28</sup> Fig. 4 shows the relationship between the leakage current density  $J$  and the electric field  $E$  for the LSCO/NBT/LSCO ferroelectric capacitor. The leakage current density at the applied electric field of  $250$   $\text{kV cm}^{-1}$  is  $3.0 \times 10^{-3}$   $\text{A cm}^{-2}$ . The leakage current curve was re-plotted as shown in the inset at the left-hand side and the right-hand side of Fig. 4 in order to further characterize the conduction mechanisms of the LSCO/NBT/LSCO capacitor in different electric field ranges. We can see two kinds of leakage behaviors account for the leakage current characteristic of the LSCO/NBT/LSCO capacitor, and  $55$   $\text{kV cm}^{-1}$  is the transitional point for the two different leakage current mechanisms. The inset at the left-hand side of Fig. 4 demonstrates the linear relation of  $\log(J)$  versus  $\log(E)$  of the LSCO/NBT/LSCO capacitor for the low electric field region, whose slope of  $1.15$  is close to  $1.0$ , implying ohmic-like conduction for the low electric field region. The relation between current density and electric field for the SCLC is given by

$$J = \frac{9\mu\epsilon}{8d} E^2 \quad (3)$$

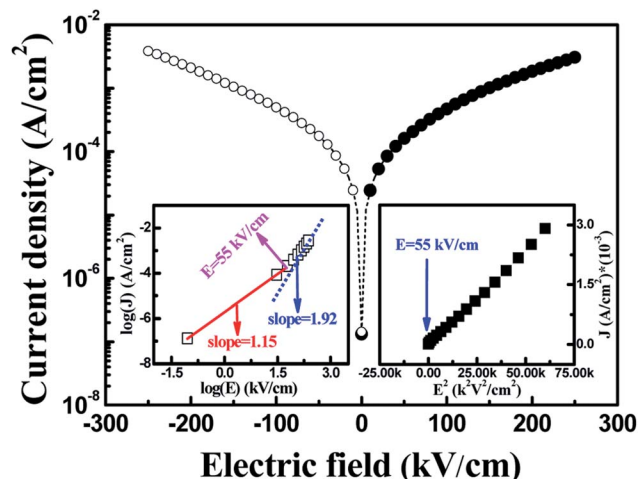


Fig. 4 Leakage current density vs. applied electric fields of the LSCO/NBT/LSCO capacitor. The inset at the left-hand side presents relation of  $\log(J)$  vs.  $\log(E)$ . The inset at the right-hand side presents relation of  $J$  vs.  $E^2$ .



In formula (3),  $J$  is the current density,  $E$  is electric field,  $\mu$  is the magnetic permeability,  $\epsilon$  is the permittivity,  $d$  is film thickness. Regarding high field region, a straight line can be yielded using the relation of  $J$  versus  $E^2$ , as shown in the inset at the right-hand side of Fig. 4, indicating SCLC conduction accounts for the conduction of the capacitor when the applied electric field is higher than  $55 \text{ kV cm}^{-1}$ . Big differences can be found from the leakage current behavior between LSCO/NBT/LSCO capacitor and Pt/NBT/Pt capacitors.<sup>17</sup> The current density of the LSCO/NBT/LSCO capacitor is  $3.0 \times 10^{-3} \text{ A cm}^{-2}$ , while it is  $1.4 \times 10^{-2} \text{ A cm}^{-2}$  for the Pt/NBT/Pt capacitor measured at the same electric field of  $250 \text{ kV cm}^{-1}$ . The  $J$ - $E$  curves of LSCO/NBT/LSCO capacitor for both positive and negative applied fields are very symmetric, however, very different leakage conduction behaviors can be found for the Pt/NBT/Pt capacitor. Moreover, the leakage current behavior of the Pt/NBT/Pt capacitor is mainly attributed to the interface-limited Schottky emission for the high electric field range, which is different from the leakage conduction behavior (SCLC) for the LSCO/NBT/LSCO capacitor. This difference can be ascribed to the fact that Pt and LSCO have different intrinsic property so that different interfaces are formed at Pt/NBT and LSCO/NBT interfaces, which further impact current conduction mechanism of the resulting capacitors.<sup>29</sup>

## 4. Conclusions

Compared to the polycrystalline films, the epitaxial LSCO/NBT/LSCO capacitor possesses larger remnant polarization, high dielectric constant and effective piezoelectric coefficient as well as good fatigue resistance and small pulse width dependence. Moreover, the capacitor satisfies ohmic conduction behavior at electric fields lower than  $55 \text{ kV cm}^{-1}$  and SCLC behavior above  $55 \text{ kV cm}^{-1}$ . Our results pave a way for the research and development of lead-free sodium bismuth titanate ferroelectric memories.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

This work is supported by the National Natural Science Foundation of China (11374086, 11474174), the Natural Science Foundation of Hebei Province (E2014201188 and E2014201063), and the Foundation of Agriculture University of Hebei (20160614).

## References

- J. F. Scott, *Science*, 2007, **315**, 954–959.
- H. Pan, Y. Zeng, Y. Shen, Y. H. Lin, J. Ma, L. Li and C. W. Nan, *J. Mater. Chem. A*, 2017, **5**, 5920–5926.
- M. Zhu, Z. Du, Q. Liu, B. Chen, S. H. Tsang and E. H. T. Teo, *Appl. Phys. Lett.*, 2016, **108**, 233502.
- B. T. Liu, K. Maki, Y. So, V. Nagarajan, R. Ramesh, J. Lettieri, J. H. Haeni, D. G. Schlom, W. Tian, X. Q. Pan, F. J. Walker and R. A. McKee, *Appl. Phys. Lett.*, 2002, **80**, 4801.
- Y. J. Fu, G. S. Fu, M. Li, D. M. Jia, Y. L. Jia and B. T. Liu, *Appl. Phys. Lett.*, 2014, **104**, 223505.
- T. Zhang, W. Li, W. Cao, Y. Hou, Y. Yu and W. Fei, *Appl. Phys. Lett.*, 2016, **108**, 162902.
- I. Katsouras, K. Asadi, W. A. Groen, P. W. M. Blom and D. M. de Leeuw, *Appl. Phys. Lett.*, 2016, **108**, 232907.
- D. Gobeljic, V. V. Shvartsman, A. Belianinov, B. Okatan, S. Jesse, S. V. Kalinin, C. Groh, J. Rödel and D. C. Lupascu, *Nanoscale*, 2016, **8**, 2168–2176.
- G. A. Smolenskii, V. A. Isupov, A. I. Agranovskaya and N. N. Krainik, *Sov. Phys. Solid State*, 1961, **2**, 2651–2654.
- W. Cao, W. Li, Y. Feng, T. Bai, Y. Qiao, Y. Hou, T. Zhang, Y. Yu and W. Fei, *Appl. Phys. Lett.*, 2016, **108**, 202902.
- L. Luo, P. Du, W. Li, W. Tao and H. Chen, *J. Appl. Phys.*, 2013, **114**, 124104.
- P. Du, L. Luo, W. Li and Q. Yue, *J. Appl. Phys.*, 2014, **116**, 014102.
- A. M. Balakt, C. P. Shaw and Q. Zhang, *J. Eur. Ceram. Soc.*, 2017, **37**, 1459–1466.
- Z. H. Zhou, J. M. Xue, W. Z. Li, J. Wang, H. Zhu and J. M. Miao, *Appl. Phys. Lett.*, 2004, **85**, 804.
- J. Xu, Y. Liu, R. L. Withers, F. Brink, H. Yang and M. Wang, *J. Appl. Phys.*, 2008, **104**, 116101.
- X. G. Tang, J. Wang, X. X. Wang and L. W. Chan, *Chem. Mater.*, 2004, **16**, 5293–5296.
- M. Bousquet, J. R. Duclère, C. Champeaux, A. Boule, P. Marchet, A. Catherinot, A. Wu, P. M. Vilarinho, S. Députier, M. Guilloux-Viry, A. Crunteanu, B. Gautier, D. Albertini and C. Bachelet, *J. Appl. Phys.*, 2010, **107**, 034102.
- M. Bousquet, J. R. Duclère, B. Gautier, A. Boule, A. Wu, S. Députier, D. Fasquelle, F. Rémondrière, D. Albertini, C. Champeaux, P. Marchet, M. Guilloux-Viry and P. Vilarinho, *J. Appl. Phys.*, 2012, **111**, 104106.
- Y. Wang, B. Liu, F. Wei, Z. Yang and J. Du, *Appl. Phys. Lett.*, 2007, **90**, 042905.
- M. Stewart, S. Lepadatu, L. N. McCartney, M. G. Cain, L. Wright, J. Crain, D. M. Newns and G. J. Martyna, *APL Mater.*, 2015, **3**, 026103.
- Y. C. Yang, C. Song, X. H. Wang, F. Zeng and F. Pan, *Appl. Phys. Lett.*, 2008, **92**, 012907.
- D. Y. Wang, N. Y. Chan, S. Li, S. H. Choy, H. Y. Tian and H. L. W. Chan, *Appl. Phys. Lett.*, 2010, **97**, 212901.
- S. K. Acharya, S. K. Lee, J. H. Hyung, Y. H. Yang, B. H. Kim and B. G. Ahn, *J. Alloys Compd.*, 2012, **540**, 204–209.
- H. Foronda, M. Deluca, E. Aksel, J. S. Forrester and J. L. Jones, *Mater. Lett.*, 2014, **115**, 132–135.
- M. A. Khan, T. P. Comyn and A. J. Bell, *Appl. Phys. Lett.*, 2008, **92**, 072908.
- S. Y. Wang, X. Qiu, J. Gao, Y. Feng, W. N. Su, J. X. Zheng, D. S. Yu and D. J. Li, *Appl. Phys. Lett.*, 2011, **98**, 152902.
- W. Sun, Z. Zhou, J. Luo, K. Wang and J. F. Li, *J. Appl. Phys.*, 2017, **121**, 064101.
- Y. Li, G. I. Ng, S. Arulkumaran, G. Ye, Z. H. Liu, K. Ranjan and K. S. Ang, *J. Appl. Phys.*, 2017, **121**, 044504.
- B. Nagaraj, S. Aggarwal and R. Ramesh, *J. Appl. Phys.*, 2001, **90**, 375.

