Time-resolved SERS study of the oxygen reduction reaction in ionic liquid electrolytes for non-aqueous lithium–oxygen cells

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Superoxide (O$_2$$^{2-}$) is the key intermediate formed during oxygen reduction in non-aqueous electrolytes. One significant obstacle towards the realisation of a practical lithium–oxygen (Li–O$_2$) battery is electrolyte instability in the presence of radical oxides, principally superoxide. Here we use the Raman active bands of O$_2$$^{2-}$ as a diagnostic molecule for probing the influence of the electrolyte on reaction processes and intermediaries at the electrode surface. In situ surface enhanced Raman studies of the interface at a roughened Au electrode with controlled and dynamic surface potentials were performed in two ionic liquids with differing properties: 1-butyl-1-methylazepenium bis(trifluoromethanesulfonyl)imide (Aze14TFSI), which has a large/soft cation, and triethylsulfonium bis(trifluoromethanesulfonyl)imide (TESTFSI), which has a relatively small/hard and e$^-$ accepting cation. The counter-cation and potential were seen to significantly influence the radical nature, or Lewis basicity of O$_2$$^{2-}$. The analysis of peak intensities and Stark shifts in O$_2$$^{2-}$ related spectral bands allowed for key information on its character and electrolyte interactions to be elucidated. Time-resolved studies of dynamic surface potentials permitted real time observation of the flux and reorientation of ions at the electrode/electrolyte interface.

Introduction

Superoxide, O$_2$$^{2-}$, is the key reaction intermediary formed during the oxygen reduction reaction (ORR) at the cathode in lithium–oxygen (Li–O$_2$) cells and can cause parasitic side reactions that limit the capacity and cycle life of the cell. Therefore, understanding the interaction of O$_2$$^{2-}$ with electrolyte provides valuable information for selecting or designing a more robust electrolyte.

The solvation properties of the electrolyte have been suggested to play an important role in the cathode intermediary reaction mechanism for some time.$^4$ More recently, solvents used in the non-aqueous Li–O$_2$ electrolyte were
shown to directly affect the ORR, inducing either a surface or solution reaction mechanism. The surface mechanism produces an insulating Li2O2 film that eventually builds up and passivates the cathode surface causing premature cell death. Whereas, the solution mechanism is favoured due to its ability to bypass passivation films forming surface agglomerates on the O2-cathode instead, extending the exposed cathode surface area for charge transfer until later into discharge, allowing for larger discharge capacities. Two distinct paths for inducing the solution mechanism have since emerged: [1] high Gutmann donor number (DN) and [2] high Gutmann acceptor number (AN) mechanisms. The Gutmann donor/acceptor numbers are empirically derived methods for estimating the affinity of a solvent for donating or accepting negative charge from solvated species and thus the strength of solvent–solute interactions.

An ideal electrolyte should have an ‘amphoteric-like’ nature, with both a high DN and AN component capable of effectively stabilising and solvating both charged species, reducing their propensity to react together and increasing their solvation into the bulk electrolyte away from the cathode surface (provided the metal anode is adequately protected against dissolution reactions driven by O2−(sol) crossover). Ionic liquids fit this model as they contain ions with both positive and negative formal charges. Many conventional organic solvents are unstable in the presence of lithium metal and/or reduced oxygen species. Even those electrolytes that have relatively good Li–O2 electrochemistry (such as dimethyl sulfoxide, DMSO) have been shown to be unstable after prolonged exposure to these species, as well as in the presence of lithium metal, thus ruling them out as practical battery electrolytes. One method of overcoming this may be to use ionic liquids (ILs) as additives or neat ionic electrolytes. Ionic liquids possess a number of favourable properties for open ‘air’ batteries due to their innate conductivity, universally low vapour pressures, flammability and their high thermal, chemical and electrochemical stability. Also, their properties can be tailored to application requirements (depending on constituent ions selected) making them extremely versatile solvent-electrolyte candidates. Schmeisser et al. correlated the relationship between the cation and anion of an IL with its acceptor and donor numbers. In general, increasing the mass and size or ‘charge dissociation’ across the cation/anion, decreases the AN/DN respectively of the IL. Larger, ‘bulkier’, heavier ions have slower kinetics, increased steric hindrance and are less able to interact with external charge. This can also be explained in terms of hard–soft acid–base (HSAB) theory. Larger, heavier ions are softer and interact less with hard counter ions; therefore, smaller IL cations will be expected to interact more with O2− than larger cations that have a high degree of charge dissociation. Considering the application of non-aqueous Li–O2 cells, the evaluation of the chemical nature of O2− at the electrode interface provides valuable information for the ion selection process of an IL electrolyte. To probe this effect and the cations influence on O2−, a large charge dissociated cyclic alkylammonium cation, 1-butyl-1-methyl-azepenium (Aze14+), and a relatively small cation, triethylsulfonium (TES+), based ILs were selected for this work (Fig. 1).
Electrochemical measurements were carried out within a glass, multi-necked vacuum tight cell within an inert atmosphere glovebox (<30 ppm O₂, <0.5 ppm H₂O) at ambient temperature. High purity oxygen (≥99.999%) and argon (≥99.998%), dried by passing through multiple desiccant filled drying tubes, was used to degas and oxygenate electrolytes via Young tapped gas inlet and outlets. Polycrystalline gold (Au) working disc electrodes (0.16 cm diameter) were used as working electrodes. The electrode surfaces were polished mechanically with decreasing grain sized alumina slurries (1.0 μm, 0.3 μm and 0.05 μm). These were washed in Milli-Q water (18.2 MΩ) and sonicated before being dried at 120 °C under vacuum overnight. A flame annealed platinum coiled wire was used as a counter electrode and a silver wire as a quasi-reference electrode. The quasi-reference electrode was standardised against an internal ferrocene standard which has a potential of +0.44 V vs. NHE in each individual IL to determine the corresponding Ag+/Ag potential shift.

In situ SERS measurements were setup similarly using a multi-necked gas-tight glass cell fitted with a sapphire window (TMS Vacuum Components). The cell was specially designed to minimise the volume of ionic liquid required for each test (<0.6 ml) (Fig. 2). Behind this window an electrochemically roughened Au working electrode was placed. The Au working electrode was cleaned (described previously) and roughened using an oxidation/reduction cycle (ORC) described by Tian et al.18 Spectra were recorded using a 50 × objective on a Raman spectrometer (Renishaw Invia) with a 785 nm laser calibrated against a silicon wafer. SERS measurements of the surface were conducted by taking rapid (~1.4 s) static spectra of the regions of interest. Due to the limited wavenumber range of the quick scans (ca. 400 cm⁻¹ region) spectra of νO–O and νO–surf regions were taken on consecutive cyclic-voltammetry (CV) cycles. Comparing differences in spectra of the bulk IL and the surface at OCP with spectra of the surface at lower potentials shows the flux of species at the surface.

Azē₁₄TFSI (provided by Queens University Belfast)¹⁹ and TESTFSI (Sigma, 99% purity) were dried at 100–130 °C for between 1 and 2 days inside an in-house made vacuum tube in a sand bath using a heater stirrer (Assynt) under vacuum (1 × 10⁻⁶ Pa) before being stored in an argon (≥99.998%) atmosphere glovebox. Water content was reduced to 1–6 ppm for Azē₁₄TFSI and <20 ppm for TESTFSI, measured using a coulometric Karl Fischer titrator (Mettler-Toledo). TESTFSI is more hygroscopic and required longer drying times (~3 days).
Results and discussion

As stated in the Experimental section Raman spectra were collected in rapid succession (every ~1.4 s) during potential cycling. The inherently short laser exposure times can lead to high background noise making enhancement active surfaces essential. However, with an optimum SERS electrode substrate, multiple spectra can be amassed with excellent signal-to-noise ratio allowing for detailed contour and multi-dimensional plots (Fig. 3) of bond vibrations to be generated showing the real-time flux of vibrational bonds at the surface.

In Fig. 3(a) and (c) the top part shows the variation of potential and current with respect to time during the CV scan. The corresponding contour plots (red colour being the most intense and purple the least) of Raman band intensities with respect to time are shown in the bottom part. The equivalent 3D plot of the data is shown in Fig. 3(b) and (d) with significant bands labelled on the plot. For both TESTFSI and Aze14TFSI a number of potential dependent bands are observed to emerge and vanish as the potential is swept. As can be seen in the Fig. 3 plots, two spectral bands for O2•− appear during ORR and disappear on the reverse scan: (1) the asymmetric dioxygen stretch, νO−O, (ca. 1100 cm⁻¹) and (2) the oxygen–surface bond, νO–surf (ca. 470 cm⁻¹). The νO–O band position is heavily dependent on the immediate coordinating environment and counter-cation size.20 However, acidity/basicity of aqueous electrolytes,21,22 neighbouring anions in a doped structures23 and H-bonding impurities24 have also been shown to significantly influence the O2•− bond force constant and νO–O spectral band position by directly interacting with O2•− charge. Additionally the counter-cation coordination strength, lattice parameters in superoxide salts or increasing mutual repulsion between electrolytic anions can also indirectly influence the O2•− bond force constant causing shifts in νO–O.
Fig. 3  Dynamic surface potential in situ SERS multi-dimensional plots. (a) and (c): Top: potential (black line)/current (blue line) vs. time plots of ORR/OER in oxygenated TESTFSI analogous with CV in Fig. 4(b) and (d) and contour plots of SERS data in $\nu_{O-O}$ and $\nu_{O-surf}$ regions respectively at corresponding times with the electrochemistry (analogous with SERS spectra in Fig. 4(a) and (c)). Regions where key surface species and $O_2^{2-}$ vibrational intensities are present and strongest/weakest highlighted in red/green respectively (top). Regions of low and high $\nu_{O-O}$ and $\nu_{O-surf}$ band intensities highlighted in green and red in the electrochemical plots. (b) and (d): 3D visual plot of spectral changes with time. Changes in intensity of $\delta$CF$_3$ bond are more clearly visible than in the contour plot. These are associated with reduced anion present at the surface at lower potentials that recovers when potentials are returned to positive values and anions repopulate the double layer. Decreasing $O_2^{2-}$ and electrolyte band (peak 2–4) intensities at low potentials are also clearly visible.

Fig. 4(a) and (c) show the corresponding selected stacked spectra of the data shown in Fig. 3. From this plot band shifts of $\nu_{O-O}$ and $\nu_{O-surf}$ become more apparent. The top half of Fig. 4(b) and (d) show ORR/OER cyclic voltammograms (CVs) of oxygen saturated TESTFSI (relatively hard cation) and Aze$_{14}$TFSI (soft cation) ILs respectively during standard (lower potential limit of 2.0 V) and extended (lower potential limit down to either 1.6 or 0.6 V) potential cycling. Normal reduction cycles are shown by dashed lines, however, in order to improve
**Fig. 4** Dynamic surface potential *in situ* SERS. (a) and (c) Selection of stacked SERS spectra of the rAu working electrode surface collected in rapid succession in oxygen saturated TESTFSI and Aze14TFSI respectively. Left and right are scans of $\nu_{O-\text{surf}}$ and $\nu_{O-O}$ wave-number regions respectively. (b) Top, CVs of oxygenated TESTFSI at 100 mV s$^{-1}$ with normal ORR/OER cycle (dashed line) and extended reduction cycle (blue line). Top right inset shows potential region where O$_2$ bands first appear with small current response due to an initial adsorbed dioxygen surface layer prior to ORR. Bottom, $\nu_{O-O}$ and $\nu_{O-\text{surf}}$ wavenumber position and intensity plots vs. potential. There is a visible blue-Stark shift. (d) Top, CVs of oxygenated Aze14TFSI with normal ORR/OER cycle (black dashed line) and extended reduction cycle (red line). Bottom, $\nu_{O-O}$ and $\nu_{O-\text{surf}}$ wavenumber position and intensity plots versus potential. There is a visible red-Stark shift. Both (b) and (d): starting data points when O$_2$ bands appear shown in black, where possible coloured arrows are included as visual aids indicating reduction (red) and oxidation (blue) directions of potential scanning.
visibility and investigate the influence of potential on the O$_2^-$ bond vibrations, spectra were collected during reduction cycles where the potential was taken to the second reduction maxima generally associated with O$_2^{2-}$ production.\textsuperscript{25} It should be noted within this study no bands pertaining to O$_2^{2-}$ were observed. During full cycles, rapid spectra of the roughened gold (rAu) working electrode surface were continuously collected. The appearance and disappearance of characteristic $v_{O-O}$ and $v_{O-surf}$ spectral bands (highlighted in purple and orange respectively) were apparent in both ILs as the potential was swept. However, the signals for O$_2^-$ bands were almost an order of magnitude weaker in the Aze$_{14}$TFSI than in TESTFSI, and were only detected at potentials after the reduction maxima and slower scan rates were required (20 mV s$^{-1}$ in Aze$_{14}$TFSI compared with 100 mV s$^{-1}$ in TESTFSI). Similarly weak O$_2^-$ bands have been reported previously with alkylammonium salts dissolved in MeCN where the $v_{O-O}$ signal was much stronger in smaller tetraethylammonium than tetrabutylammonium based electrolytes.\textsuperscript{26} Both bands were observed to blue or red Stark shift with decreasing potential in TESTFSI and Aze$_{14}$TFSI respectively, and a plot of their peak fitted band positions and intensities are shown in the bottom half of Fig. 4(b) and (d). Electrolyte peaks labelled 2–4, were also observed to be potential dependent in both TESTFSI and Aze$_{14}$TFSI. Identity of key potential dependent peaks of bothionic liquids, with observations are summarised within Table 1.

Variances in spectral band intensities give qualitative information on surface coverage, molecular orientation and the concentration of species at the electrode/electrolyte interface during potential cycling. Differences (or shifts) in $v_{O-O}$ band positions between electrolytes from lower (red) to higher (blue) wavenumbers generally indicate more ionic or covalent, respectively, interactions between O$_2^-$ and its coordinating environment. A more ionic O$_2^{*-}$ species is a less coordinated, more radical, harder Lewis base whilst a more covalent O$_2^{*-}$ species is a more coordinated, softer Lewis base. Likewise, blue or red Stark shifts in $v_{O-O}$ bands during potential cycling indicate surface potential induced fluctuations in the ionic/covalent character, respectively, of generated O$_2^-$\textsuperscript{.} The Stark effect, being consistent with Guoy–Chapman theory, is heavily dependent on the distance of the probe molecule (O$_2^-$) from the surface and therefore indicates observation of species bonded, adsorbed or adjacent to the surfaces.\textsuperscript{26} For $v_{O-surf}$, differences in band positions between electrolytes and red/blue Stark shifts indicate longer (weaker) and shorter (stronger) adsorption bonds, respectively. Both $v_{O-O}$ and $v_{O-surf}$ bands are closely related appearing on ORR and disappearing on OER; though there were significant differences between the two IL electrolytes which will be discussed. As such, the focus on the effect of the cation on electrochemically generated surface O$_2^-$ \textsuperscript{25} by varying the substituent cation in two distinct ILs with either a relatively small, TES, or large charge dissociated cyclic alkylammonium, Aze$_{14}$, as the cation is given thorough consideration.

In TESTFSI the O$_2^-$ bond vibrations first appear (weakly) at 2.75 V vs. Li$^+/\text{Li}$ in TESTFSI and grow in intensities to a maximum just after the reduction potential current maxima. The top right inset of Fig. 4(b) shows this to be the precise onset point of an exponential decrease in current and beginning of bulk O$_2$ reduction. This is below the thermodynamic reduction potential of oxygen reduction in lithiated systems (2.97 V vs. Li$^+/\text{Li}$)\textsuperscript{27} but significantly higher than the $E_{1/2}$ of the ORR/OER reaction in TESTFSI.
Comparing the $\text{O}_2^{\cdot-}$ bands between the two ILs the $\nu_{\text{O-O}}$ band is much higher wavenumber for TESTFSI (1120 cm$^{-1}$) than in the Aze$\sub{14}$TFSI (1105 cm$^{-1}$). This agrees with the trend in coordinating cation size observed in our previous work on ORR within tetraalkylammonium based electrolytes, indicating a more Lewis basic $\text{O}_2^{\cdot-}$ is generated at the surface in Aze$\sub{14}$TFSI than TESTFSI, due to being less strongly coordinated by the Aze$\sub{14}$+ counter cation in the electrolyte. There is a similar difference in the $\nu_{\text{O-surf}}$ bands, with lower wavenumber values in TESTFSI (466 cm$^{-1}$) than Aze$\sub{14}$TFSI (488 cm$^{-1}$). Lower $\nu_{\text{O-surf}}$ wavenumbers indicate longer, weaker oxygen-to-surface adsorption bonds that suggests that some $\nu_e$ density is shared with the electrode surface.

Dioxygen ligand-to-metal bonds in organometallic chemistry are analogous to surface adsorption bonds and can help rationalise the bond structures of the observed adsorption bond vibrations. The three most common dioxygen–metal bond structures are side-on (flat), kinked and bent, shown in Fig. 5. The longest/weakest adsorption bond for the side-on, or flat, dioxygen–metal structure is often associated with the peroxo-complex though side-on superoxo-complexes have also been reported. The energy of the dioxygen $\pi^*$ levels lies below that of metal d-orbitals therefore the flat structure is mainly composed of d-to-$\pi^*$

Table 1 Summary of key ionic liquid bands. C$^+$ refers to cation, either TES$^+$ or Aze$\sub{14}$+ (note peak 1 is the designation of the flat conformer of adsorbed superoxide at ca. 430 cm$^{-1}$)

<table>
<thead>
<tr>
<th>TESTFSI</th>
<th>Aze$\sub{14}$TFSI</th>
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<tr>
<td>$\delta\text{CF}_3$ (TFSI$^-$)</td>
<td>- Intensity decreases/increases with potential, anion bands are still present throughout and changes indicate that the anion reorientates at lower potentials from flat to side-on, balancing surface charge</td>
</tr>
<tr>
<td>$\omega\text{SO}_2$ (TFSI$^-$)</td>
<td>- Intensity decreases/increases with potential</td>
</tr>
<tr>
<td>Peaks 2,3 (C$^+$)</td>
<td>- Supports anion reorientation at low potentials</td>
</tr>
<tr>
<td></td>
<td>- 590–610 (doublet), 650 cm$^{-1}$</td>
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<tr>
<td></td>
<td>- Present at all potentials</td>
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<tr>
<td></td>
<td>- Assigned to cation (match Spartan DFT simulated spectra for TES$^+$)</td>
</tr>
<tr>
<td></td>
<td>- Intensity changes mirror $\text{O}_2^{\cdot-}$ bands</td>
</tr>
<tr>
<td></td>
<td>- Cation–$\text{O}_2^{\cdot-}$ coupling causes intensity changes</td>
</tr>
<tr>
<td></td>
<td>- Increased cation concentration at surface</td>
</tr>
<tr>
<td></td>
<td>- Higher energy cation vibration, stronger cation and cation–$\text{O}_2^{\cdot-}$ bond</td>
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<tr>
<td>Peak 4 (C$^+$)</td>
<td>- 1000 cm$^{-1}$</td>
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<tr>
<td></td>
<td>- Only present when $\text{O}_2^{\cdot-}$ bands present</td>
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<tr>
<td></td>
<td>- Intensity changes mirror $\text{O}_2^{\cdot-}$ bands</td>
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<tr>
<td></td>
<td>- S–O, cation–$\text{O}_2^{\cdot-}$ coupling bond</td>
</tr>
<tr>
<td></td>
<td>- Changes in coupling intensity correspond with changes in $\text{O}_2^{\cdot-}$ concentration at surface</td>
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bonding with equal bonding between both oxygen atoms and the metal centre,\textsuperscript{28} provided that the metal Fermi level is in a state that permits it. The superoxo-ligand however is known to be preferentially bent-bonded to metal sites, through ligand-to-metal $\sigma$ donation. Back donation from the metal into the antibonding orbitals can occur, if the surface energetics permit, which weakens the O-to-surface bond, relative to the bent-bonded structure, lowering the bond force constant producing a longer kinked adsorption bond.\textsuperscript{28} A kinked bond has characteristics between the flat and bent bonded species with both oxygen atoms interacting to differing degrees with the metal centre and at an angle with the surface, the extent of back bonding and coordination dictates the angle.

\textit{Aze}_{14}\textit{TFSI} is comprised of soft poorly accepting and donating ions. The weak accepting nature of the electrolyte increases the propensity of adsorbed $\text{O}_2^{-}$ to share charge with the surface rather than the electrolyte. Giving rise to higher $\nu_{\text{O-surf}}$ wavenumbers indicative of more energetic shorter bond vibrations in \textit{Aze}_{14}\textit{TFSI} relative to \textit{TESTFSI}. This $\nu_{\text{O-surf}}$ vibration can be attributed to a bent bonded structure that has little or no surface back bonding.

For \textit{Aze}_{14}\textit{TFSI}, both $\text{O}_2^{-}$ bands were seen to red-Stark shift with decreasing surface potentials. Shifting the electrode potential negative increases the Fermi level of surface molecules, inducing coulombic repulsion between the surface and adsorbed anions and molecules with negative dipoles. When the electrolyte has a poor accepting nature that is less effective at balancing excess negative surface charge, as in \textit{Aze}_{14}\textit{TFSI}, then surface $e^{-}$ density will be transferred and concentrated in any free molecular orbitals of surface bonded or adsorbed species closest in energy. In the case of $\text{O}_2^{-}$ this will be the $\pi^*$ orbital, therefore back bonding increases with decreasing potentials and the bond becomes longer and increasingly kinked. This can also be explained as the increasingly negative surface charge contributes more $e^{-}$ density to surface adsorbed species reducing the effective nuclear charge experienced by valence bond electrons, causing both the $\nu_{\text{O-surf}}$ and dioxigen bonds to weaken and lengthen. Hence a concomitant red-Stark shift in the $\nu_{\text{O-O}}$ bond vibrations of $\text{O}_2^{-}$ close to the surface. This charge localisation on $\text{O}_2^{-}$ at low potentials is indicative of a more Lewis basic species being generated at the surface with a weaker adsorption bond.

For \textit{TESTFSI}, which can be considered as an $\text{O}_2^{-}$ accepting electrolyte, $\nu_{\text{O-surf}}$ is significantly lower than \textit{Aze}_{14}\textit{TFSI} ($\sim$16–27 cm$^{-1}$), see Table 2, indicating a propensity for valence $e^{-}$ density on $\text{O}_2^{-}$ to interact more with the electrolyte.

![Fig. 5](image-url) Metal-to-dioxygen surface bond structures left to right: flat, kinked and bent. Visible bond vibrations based on the surface selection rule are shown above each structure.\textsuperscript{28,29,32} The tick and cross indicate which bands will be observable by Raman.
than the surface indicating a weak adsorption bond with a kinked structure and some surface back bonding. An additional shoulder at ∼430 cm$^{-1}$, labelled peak 1 (area highlighted in green in Fig. 4(a)), is also present at ORR/OER surface potential above 1.8 V vs. Li$^+/\text{Li}$. Peak 1 has been characterised previously as a conformationally flat O$_2$–$\pi$ surface species on catalytically active Pt and Pd surfaces using SHINERS.\textsuperscript{33} Raman selection rules dictate that bond vibrations that are perpendicular to the surface are not visible. $\nu_{\text{O-O}}$ would therefore be absent or weak when a conformationally flat species is prominent at the surface.

This was the case when holding the potential at onset reduction potentials (Fig. 6), however at the same potential during cycling only a mixed (predominantly kinked) species is present. Therefore, the kinked structure appears to be kinetically more stable during potential cycling when the double layer is in flux; yet when the potential is held and the double layer equilibrates, the O$_2$–$\pi$ relaxes to a conformationally flat species. This was visible at onset potentials in TESTFSI electrolyte, but not in Aze$_{14}$TFSI. However, this effect was also present and far more pronounced in other ionic liquid electrolytes evaluated, not shown here. At extremely low potentials, below 1.8 V vs. Li$^+/\text{Li}$, peak 1 disappears completely with only the $\nu_{\text{O-O}}$ (kinked) species visible. Blue-Stark shifts of 4 and 7 cm$^{-1}$ were present with decreasing potentials. These observations suggest that the flat surface bond becomes less kinked and back bonding with the surface decreases with decreasing potential, the opposite of what was observed in Aze$_{14}$TFSI. Counterintuitively, this would suggest that the $\pi^*$ orbital is being depleted of charge as the surface becomes increasingly more negative. This can be accounted for by considering changes within the double layer. When the surface potential is decreased the ionic liquid double layer becomes increasingly populated with positively charged species\textsuperscript{34} where the anions initially reorientate, followed by their expulsion\textsuperscript{35} in order to balance surface charge.

This was apparent in TESTFSI (Fig. 3(a) and (b)) by a continuous decrease in the intensity of the characteristic $\delta\text{CF}_3$ vibration of the TFSI anion (740 cm$^{-1}$)\textsuperscript{36} with lowering of potential. The $\delta\text{CF}_3$ intensity increased again when the direction of scanning was reversed due to the anion repopulating the double layer. Furthermore, as the potential decreases a broad Au–S adsorption band from the S in TES (200–300 cm$^{-1}$) appears with ORR, indicative of increased cation population in the double layer (Fig. 6). Bands in the range of 600–700 cm$^{-1}$ are seen to change in shape and intensity with respect to potential, which may arise from conformational changes of TES at the electrode surface, as seen previously for tetraalklyammonium cations.\textsuperscript{20} In addition, the intensity of electrolyte related peaks 2–4 of TESTFSI mirrored the intensity changes in the $\nu_{\text{O-surf}}$ bond vibration

<table>
<thead>
<tr>
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<th>$\nu_{\text{O-o}}$</th>
<th>$\Delta\nu_{\pi}$/cm$^{-1}$</th>
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<tbody>
<tr>
<td>TESTFSI</td>
<td>1109–1120</td>
<td>11</td>
</tr>
<tr>
<td></td>
<td>459–466</td>
<td>7</td>
</tr>
<tr>
<td></td>
<td>475–488</td>
<td>13</td>
</tr>
<tr>
<td></td>
<td>1097–1105</td>
<td>8</td>
</tr>
<tr>
<td>Aze$_{14}$TFSI</td>
<td>1097–1105</td>
<td>8</td>
</tr>
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</table>
(Fig. 3(a) and (b)). Peak 5 (1001 cm\(^{-1}\)) similarly mirrored the \(\nu_{\text{O-O}}\) bond vibration, however it was dependant on the ORR/OER potentials appearing and disappearing in tandem with \(\nu_{\text{O-O}}\). Thereby showing a strong correlation of population of interfacial species from the electrolyte with respect to \(\text{O}_2\). Similar electrolyte conformational changes were present, but much more difficult to observe in Aze14TFSI as shown in the contour plots in Fig. 3(c) and (d) due to the greater signal-to-noise ratio.

**Conclusions**

The cation effect on the radical character of \(\text{O}_2^-\) in non-aqueous ionic liquid media was investigated with surface enhanced Raman spectroscopy. Two ionic liquids with differing properties: 1-butyl-1-methyl-azepenium bis(trifluoromethanesulfonyl)imide (Aze\(_{14}\)TFSI), which has a large/soft cation, and triethylsulfonium bis(trifluoromethanesulfonyl)imide (TESTFSI), which has a relatively small/hard and \(\text{e}^-\) accepting cation were investigated.

Substantial surface enhancements to the Raman signal permitted the accumulation of a Raman spectrum within ca. 1.4 seconds with sufficient signal-to-noise ratio. This allowed the observation of the effect of dynamic potential sweeping on both the appearance/disappearance and influence of \(\text{O}_2^-\) upon the anion and cations of the ionic liquids to be evaluated.

The experimental trends provide insight into the nature of \(\text{O}_2^-\) and its character in non-aqueous purely ionic media. By contrasting the surface adsorbed \(\text{O}_2^-\) bands between the two ILs the \(\nu_{\text{O-O}}\) bands were found to occur at a much higher wavenumber for TESTFSI (1120 cm\(^{-1}\)) than in the Aze\(_{14}\)TFSI (1105 cm\(^{-1}\)) and indicates that a more Lewis basic superoxide is generated at the gold electrode surface in Aze\(_{14}\)TFSI than TESTFSI due to being less strongly coordinated by the Aze\(_{14}\) counter cation in the electrolyte. There was a similar difference in the \(\nu_{\text{O-surf}}\) bands with lower values in TESTFSI than Aze\(_{14}\)TFSI (466 cm\(^{-1}\) vs. 488

**Fig. 6** Extended spectra acquired with surface potential control at specified held potentials of ORR in oxygen saturated TESTFSI.
cm\(^{-1}\)). Lower \(v_{\text{O-surf}}\) wavenumbers indicate a longer, weaker oxygen-to-surface adsorption bond.

The differing effect upon the radical nature (Lewis basicity) of \(O_2^{-}\) at the interface depending on choice of ionic liquid opens up a promising avenue of research, given the chemical flexibility of ionic liquids to be able to ‘design’ an optimum electrolyte system that provides a stable ionic medium to permit long-life and reversible cycling of Li–O\(_2\) cells. If such an electrolyte can then also be tailored to provide acceptable levels of oxygen solubility and diffusivity, then progress towards the realisation of practical Li–O\(_2\) cells becomes closer. Undoubtedly our observations on the cation effect on the nature of \(O_2^{-}\) will need to be verified with further SERS studies on additional ionic liquids of differing cations and anions, and this will be the subject of future investigations.

Acknowledgements

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