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Understanding three-body contributions to coarse-grained force fields†

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Coarse-graining is a systematic reduction of the number of degrees of freedom used to describe a system of interest. Coarse-graining can be thought of as a projection on the coarse-grained degrees of freedom and is therefore dependent on the number and type of basis functions used to represent the coarse-grained force field. We show that many-body extensions of the coarse-grained force field can result in substantial changes of the two-body interactions, making them much more attractive at short distances. This interplay can be alleviated by first parametrizing the two-body potential and then fitting the additional three-body contribution to the residual forces. The approach is illustrated on liquid water where three-body interactions are essential to reproduce the structural properties, and liquid methanol where two-body interactions are sufficient to reproduce the main structural features of the atomistic system. Furthermore, we demonstrate that the structural and thermodynamic accuracy of the coarse-grained models can be controlled by varying the magnitude of the three-body interactions. Our findings motivate basis set extensions which separate the many-body contributions of different order.

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1. Introduction

Coarse-graining (CG) is a systematic way to reduce the number of degrees of freedom describing a specific physical system. By combining coherently moving atoms of a molecule into a single interaction site one can increase the simulation times by ten to hundred times compared to the atomistic simulation. The computational cost is reduced in several ways: First, the coarse-grained system has less degrees of freedom. Second, smoother interaction potentials result in a smaller friction and, hence, faster dynamics. Third, one can use a bigger simulation timestep in the integration algorithm.

To coarse-grain a specific system, we need to choose coarse-grained degrees of freedom, relate them to the fine-grained description, identify a merit function which quantifies the difference between the fine- and coarse-grained representations, and optimize the coarse-grained potential energy surface (PES). Consistency between the coarse-grained and the fine-grained models can be defined in terms of consistency of the equilibrium probability densities resulting in unique expressions for the CG masses and interaction potential, *i.e.* the many-body potential of mean force (PMF).¹ Evaluating this potential is as computationally demanding as propagating the fine-grained system and, in

practice, it is approximated using a limited set of basis functions chosen to represent the interactions at the CG level of resolution.

Several methods have been developed that approximate the many-body PMF, targeting specific properties of the underlying atomistic system. For example, iterative Boltzmann inversion (IBI)² and inverse Monte Carlo (IMC)³ schemes target structural pair distribution functions. It can be shown that these methods minimize the relative entropy between the CG and atomistic ensembles.⁴ An alternative route is to match the forces of the CG system to those of the atomistic description, employing force-matching (FM) or multiscale coarse-graining (MS-CG).^{1,5,6} This approach corresponds to projecting the many-body mean force, *i.e.*, the negative gradient of the many-body PMF, into the space of force fields defined by the CG basis set. It allows to systematically increase the accuracy of the approximation of the many-body PMF by expanding the basis set. FM and structural coarse-graining can be connected *via* Yvon–Born–Green theory.^{7,8}

The structural accuracy of the force-matched CG model depends on the CG mapping scheme⁹ and on the ability of the CG basis set to capture the relevant many-body correlations of the mapped ensemble.⁸ In some cases, force-matched pair potentials reproduce the atomistic pair correlation functions, as for one-site models of liquid methanol.^{10,11} However, for liquids with a strong local orientational order, with water being a prominent example, CG pair potentials fail to fully capture the structural correlations of the underlying atomistic system.^{12–17} Even when the pair potential is able to reproduce the atomistic pair correlation functions, thermodynamic properties are usually not reproduced correctly, for example the cohesive energy can be

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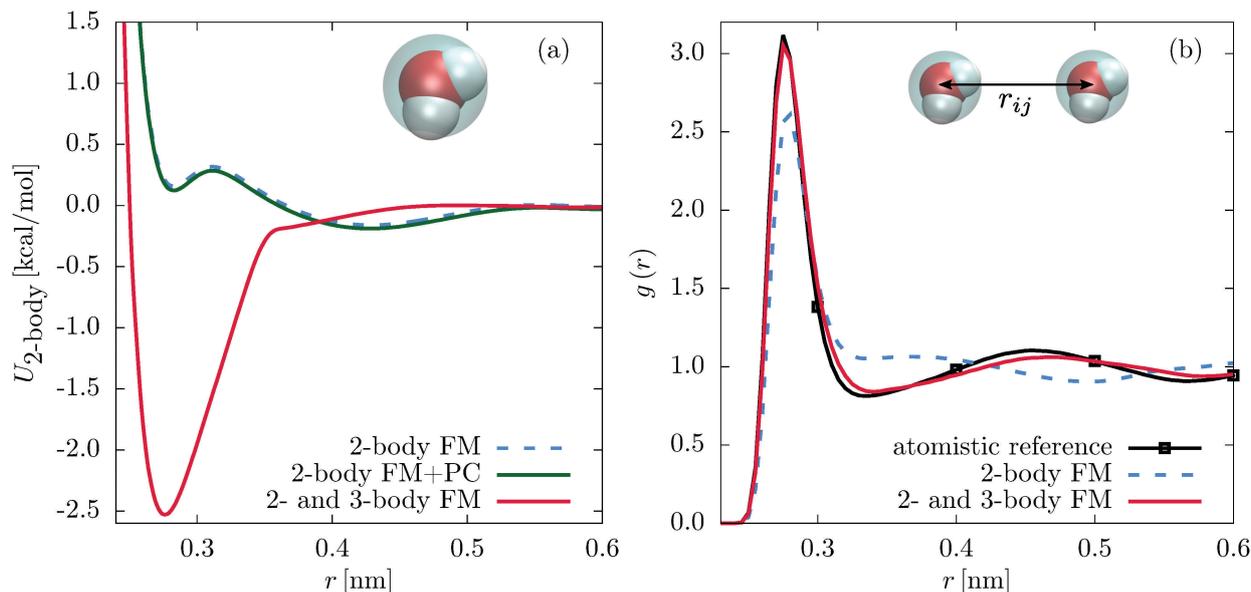


Fig. 1 The pair potentials (a) and the radial distribution functions (b) of coarse-grained SPC/E water. The following parametrizations are shown: FM with a (tabulated) pair potential only (2-body FM), and FM with a tabulated pair potential and the short-range three-body SW potential (2- and 3-body FM). For the pair potential, we also show the curve with the additional pressure correction (2-body FM + PC).

significantly underestimated.¹⁸ This implies that an extension of the basis set is necessary to describe the many-body correlations more accurately.^{19–23}

Indeed, coarse-grained models can be improved by including a local density dependent potential^{24–26} or a short-ranged non-bonded three-body potential,^{12–14,27–30} such as the Stillinger–Weber³¹ (SW) potential. These basis set extensions lead to a significant change of the two-body interactions.^{12,13,24,25} For example, an one-site CG SPC/E water model with the SW potential has a deep attractive well at 2.5 nm, as shown in Fig. 1. In this paper we demonstrate that the change of the pair potential can be tracked back to the non-orthogonality of the three-body and the two-body interactions. In addition, we propose a parametrization scheme which reduces the three-body contribution to the two-body PMF or, in other words, decouples the two- and three-body interaction terms.

II. Methods

A. Atomistic simulations

For atomistic simulations we employ the GROMACS³² package, version 5.1. We use the SPC/E water model³³ and the OPLS-AA force field^{34,35} for methanol. Boxes of 1000 molecules are first equilibrated for 10 ns in the *NPT* ensemble at $T = 300$ K and $p = 1.0$ bar, followed by 10 ns *NVT* production runs at the average density of the preceding *NPT* simulations, $\rho = 0.998$ g cm⁻³ (water) and $\rho = 0.776$ g cm⁻³ (methanol). These values are very close to the experimental densities at normal conditions, $\rho = 0.9971$ g cm⁻³ (water)³⁶ and $\rho = 0.7872$ g cm⁻³ (methanol).³⁷ To integrate the equations of motion, we use a stochastic dynamics algorithm³⁸ in combination with a Berendsen barostat³⁹ with a time constant of 1 ps,

the compressibility parameter of water (4.5×10^{-5} bar⁻¹), and a time step of 1 fs. Electrostatic interactions are treated with a smooth particle mesh Ewald method⁴⁰ with cubic interpolation, a grid spacing of 0.12 nm and an Ewald accuracy parameter of 10^{-5} . We use a short-range cutoff of 1.2 nm and long-range dispersion corrections for energy and pressure. The OPLS-AA force field is used without constraints.

B. Coarse-graining procedure

We replace one water or methanol molecule with a CG bead positioned at its center of mass,

$$\mathbf{R}_i = \sum_{\alpha \in i} w_{\alpha} \mathbf{r}_{\alpha}, \quad (1)$$

where $\omega_{\alpha} = m_{\alpha} / \left(\sum_{\alpha \in i} m_{\alpha} \right)$ are the weights of each atom of molecule i . The mass of the CG bead is the sum of the atomistic masses: $M_i = 18.0154$ amu (water) and $M_i = 32.0374$ amu (methanol).

To parametrize the FM potentials, the reference force on each CG bead i is calculated as the sum of the atomistic forces,

$$\mathbf{f}_i^{\text{ref}} = \sum_{\alpha \in i} \mathbf{f}_{\alpha}. \quad (2)$$

The CG representation of the force is then determined by solving:

$$f_{il}^{\text{CG}}(g_1, \dots, g_M) = f_{il}^{\text{ref}}, \quad i = 1 \dots N, l = 1 \dots L. \quad (3)$$

Here, g_1, \dots, g_M are the coefficients of the CG force field basis functions, N is the number of CG beads, and L is the number of simulation snapshots per block. We use $N = 1000$ and $L = 20$. Given that the CG force field basis functions f_{il}^{CG} depend linearly on the parameters g_1, \dots, g_M , eqn (3) is a set of $N \times L$



overdetermined linear equations ($M < N \times L$) which we solve with a constrained least-squares solver.

All coarse-graining algorithms are implemented into the VOTCA-CSG package.¹⁸

C. Coarse-grained potentials

Nonbonded two-body potentials are represented by cubic splines which depend linearly on the coefficients g_1, \dots, g_M . A set of K grid points leads to $M = 2K$ spline coefficients where K of the coefficients are fixed due to constraints guaranteeing the continuity of the first derivatives.

The SW three-body potential,

$$U = \sum_{i,j \neq i, k > j} f^{(3b)}(\theta_{ijk}) \exp\left(\frac{\gamma_{ij}\sigma_{ij}}{r_{ij} - a_{ij}\sigma_{ij}} + \frac{\gamma_{ik}\sigma_{ik}}{r_{ik} - a_{ik}\sigma_{ik}}\right), \quad (4)$$

is used as a short-range three-body potential, where i is the index of the central atom, and j and k are the other two atom indices of a triplet of atoms with an angular interaction term $f^{(3b)}(\theta_{ijk})$. We do not limit ourselves to an analytic expression of $f^{(3b)}(\theta_{ijk})$ as in the original SW potential, but allow for a flexible angular dependence of $f^{(3b)}(\theta_{ijk})$.

The FM parametrization is implemented in the VOTCA-CSG package¹⁸ using a cubic spline representation of $f^{(3b)}(\theta)$ with K grid points and a linear dependence on the $M = 2K$ spline coefficients. Treating the two exponential terms, $\exp\left(\frac{\gamma\sigma}{r - a\sigma}\right)$, as prefactors yields a linear set of equations which can be solved by a constrained least-squares solver as in the case of the pair forces. To do so, the remaining parameters a_{ij} , a_{ik} , σ_{ij} , σ_{ik} , γ_{ij} , and γ_{ik} , have to be set beforehand. For one CG bead type, as in our case, the number of parameters reduces to $a_{ij} = a_{ik} = a$, $\sigma_{ij} = \sigma_{ik} = \sigma$, and $\gamma_{ij} = \gamma_{ik} = \gamma$. Here, a is the three-body short-range cutoff which we choose in a way that the three-body potential is fully switched on in the first coordination shell. Setting $\sigma = 1$, the parameter γ controls the steepness of this switching on. In the limit of $\gamma \rightarrow 0$, the SW potential is instantaneously switched on. For water, we refer to the cutoff radius of ref. 13, namely $a = 0.37$ nm. In the case of liquid methanol, we choose $a = 0.45$ nm. We systematically vary γ until the structural fit does not improve anymore. The optimal value, in both cases, is $\gamma = 0.08$ nm. This means that U^{SW} is fully switched on in the first and switched off within the second coordination shell. For more technical details on the FM implementation we refer to the ESI.†

D. CG simulations

Coarse-grained MD simulations are performed at the atomistic density using the LAMMPS package.⁴¹ Two-body CG interactions are treated as linear interpolation tables with a short-range cutoff of 1.2 nm. We extend the LAMMPS implementation of the Stillinger–Weber interaction potential to read in tabulated angular dependent potentials. We employ a time step of 1 fs and a chain of three Nose–Hoover thermostats^{42,43} with a damping parameter of 200 fs when integrating the equations of motion⁴⁴ for 10 ns in the NVT ensemble. When simulating in the NPT ensemble,

we apply an additional chain of three Nose–Hoover barostats with a damping parameter of 1000 fs.

E. Calculation of the pair PMF

To separate the two-body and three-body contributions of the CG interaction potential, we evaluate the pair PMF between two coarse-grained sites:

$$U_{\text{PMF}}(r) = -\int_0^r F_{\text{rad}}(r') dr'. \quad (5)$$

Here, $F_{\text{rad}}(r) = \langle \mathbf{F}_j \cdot \hat{\mathbf{r}}_{ij} \rangle \equiv \langle -\mathbf{F}_i \cdot \hat{\mathbf{r}}_{ij} \rangle$ is the projection of the total CG force acting on a CG bead i or j onto the unit distance vector $\hat{\mathbf{r}}_{ij}$ connecting this pair of beads, averaged over all CG bead pairs with distance r . Note that the total pair $U_{\text{PMF}}(r)$ is equivalent to the radial distribution function, since $g(r) \sim \exp(-U_{\text{PMF}}(r)/k_{\text{B}}T)$. When calculating $U_{\text{PMF}}(r)$ for the atomistic simulations, \mathbf{F}_i and \mathbf{F}_j are evaluated as a weighted sum of all atomistic forces acting on each CG bead i and j , according to eqn (2).

$F_{\text{rad}}(r)$ and, thus, $U_{\text{PMF}}(r)$ can be evaluated separately for the two-body and three-body interactions by a rerun of the CG trajectory and switching on and off the individual force contributions. That gives the contributions of the two-body and three-body interactions to the total pair PMF, including the effects of correlated forces from the environment. This is comparable to calculating the decomposition of the pair PMF in terms of the MS-CG G-matrix.⁴⁵

F. Inner product of force field functions

To quantify the degree of correlation between two-body and three-body force field functions, we define an inner product as⁷

$$\mathbf{F}^{(2)} \cdot \mathbf{F}^{(3)} = \int d\mathbf{R} p_{\mathbf{R}}(\mathbf{R}) \sum_{i=1}^N \mathbf{F}_i^{(2)}(\mathbf{R}) \cdot \mathbf{F}_i^{(3)}(\mathbf{R}). \quad (6)$$

Here, $\mathbf{F}_i^{(\xi)}$ is the total force on bead i , $\xi = 2, 3$ corresponds to the two- and three-body contributions, and $p_{\mathbf{R}}(\mathbf{R})$ is the CG configuration distribution. We evaluate a normalized inner product of the two-body and three-body contributions to the total force on each CG bead as:

$$p = \frac{\mathbf{F}^{(2)} \cdot \mathbf{F}^{(3)}}{|\mathbf{F}^{(2)}| |\mathbf{F}^{(3)}|}. \quad (7)$$

The normalized inner product corresponds to the ensemble average of the inner product of the two-body and three-body contribution on each CG bead and can be converted into an angle:

$$\Phi = \arccos(p) = \arccos\left(\frac{\mathbf{F}^{(2)} \cdot \mathbf{F}^{(3)}}{|\mathbf{F}^{(2)}| |\mathbf{F}^{(3)}|}\right). \quad (8)$$

III. Results and discussion

We choose one-site models of liquid water and methanol as test systems where two-body basis functions fail (water) and succeed (methanol) to accurately reproduce the atomistic pair correlation



functions. We employ force-matching (FM) to parametrize the tabulated two-body in combination with the three-body Stillinger-Weber (SW) interactions with a flexible angular dependence.

A. Structural properties

First, we compare the structural correlation functions of the different models. As mentioned above and shown in Fig. 1(b), the CG model of SPC/E water with the three-body SW potential (2- and 3-body FM) reproduces the radial distribution function (RDF) of the atomistic reference better than the model with only pairwise interactions (2-body FM). This also holds for the angular distribution function, as shown in Fig. 5(b). For liquid methanol, the two-body FM model already reproduces the RDF nearly perfectly. Therefore, there is no visible change upon including the three-body SW term, as shown in Fig. S6 of the ESI.† At the same time, the angular distribution function significantly improves, as shown in Fig. 2(b).

Next, we focus on the interplay of the two-body and three-body interactions. In Fig. 3, we show the total pair PMF (red solid lines), the two-body (grey dashed lines) and three-body (grey dotted lines) contributions to this potential of mean force for water and methanol. In addition, we show the pair PMF generated by the two-body FM parametrizations as blue crosses. The evaluation of the pair PMF and the decomposition are described in Section IIE. To some extent, the attractive forces of the two-body potential within the first coordination shell are compensated by correlation effects due to two-body interactions with surrounding molecules. This can be seen in the reduced magnitude of the two-body contribution to the pair PMF compared to the direct pair potential (compare Fig. 1(a) and 3(a) for water and Fig. 2(a) and 3(b) for methanol).

However, most of the compensation of the two-body attractions comes from short-range repulsive forces of the three-body interactions.

The strong coupling of the two-body and three-body interactions motivates a procedure similar to the subtraction of the Coulomb force from the atomistic reference.⁴⁶ This can be viewed as a Gram-Schmidt orthogonalization scheme which reduces the coupling between different terms in the CG force field. First, we obtain the CG pair potential by two-body FM. In a second step, we determine the residual force Δf_{FM} acting on each CG bead by subtracting the two-body force from the total reference force, f^{ref} , given by eqn (2):

$$\Delta f_{\text{FM}} = f^{\text{ref}} - f^{\text{FM}}. \quad (9)$$

We now parametrize the three-body SW potential using the residual instead of the reference force. The main effect of the residual force parametrization is a vertical shift of the angular part of the SW potential $f^{(3b)}(\theta)$ to smaller values (see Fig. S2(b) and S5(b) of the ESI†). The decompositions of the pair PMFs for the new parametrizations (see Fig. 4, orange solid and dashed lines with dots) clearly show significantly reduced three-body contributions and less attractive two-body contributions to the pair PMFs compared to the unconstrained parametrizations (red solid and dashed lines).

It is essential to check how the residual force parametrizations perform in terms of the liquid structure. In Fig. 5, we compare the RDFs and angular distribution functions of CG beads within the first coordination shell for liquid water. The angular distribution functions of the two different parametrization schemes are very similar. The constrained parametrization (2-b FM and 3-b Δf_{FM})

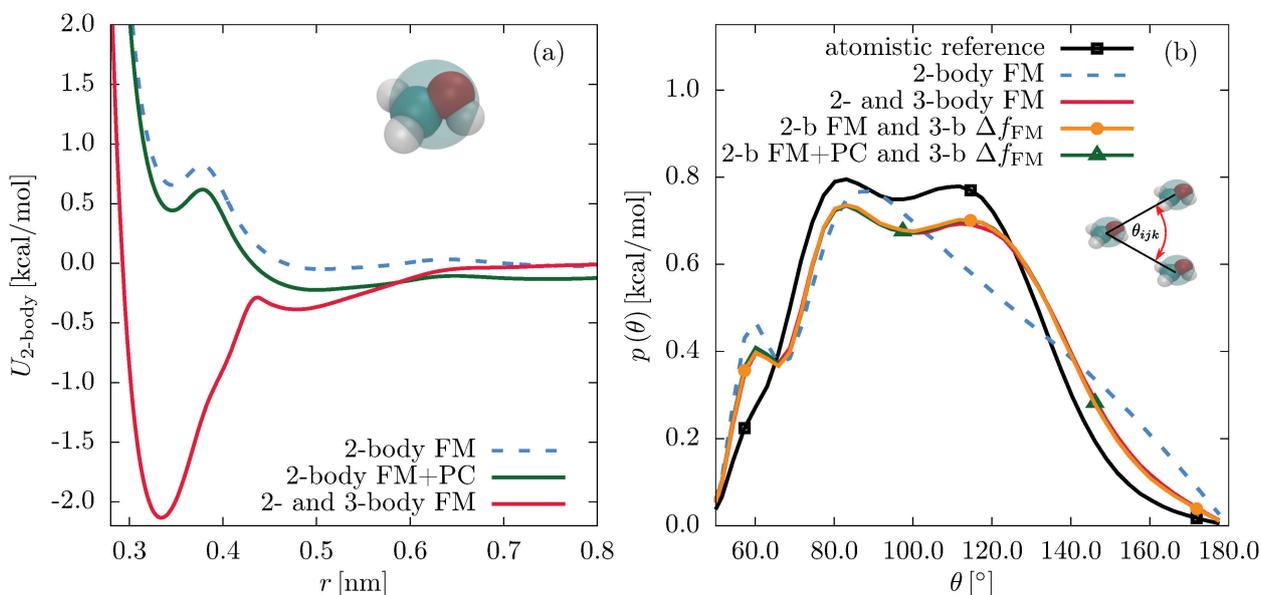


Fig. 2 The pair potentials (a) and the angular distribution functions (b) of coarse-grained liquid methanol. Four parametrizations are shown: FM with a (tabulated) pair potential only (2-body FM), FM with a tabulated pair potential and the short-range three-body SW potential (2- and 3-body FM), three-body FM using the residual force of the two-body FM potential (2-b FM and 3-b Δf_{FM}), as well as, three-body FM using the residual force of the two-body FM potential with additional pressure correction (2-b FM + PC and 3-b Δf_{FM}). The second pair potential (2-body FM + PC) refers to the pair potential with the additional pressure correction. The angular distribution function is calculated using a cutoff of 0.38 nm.



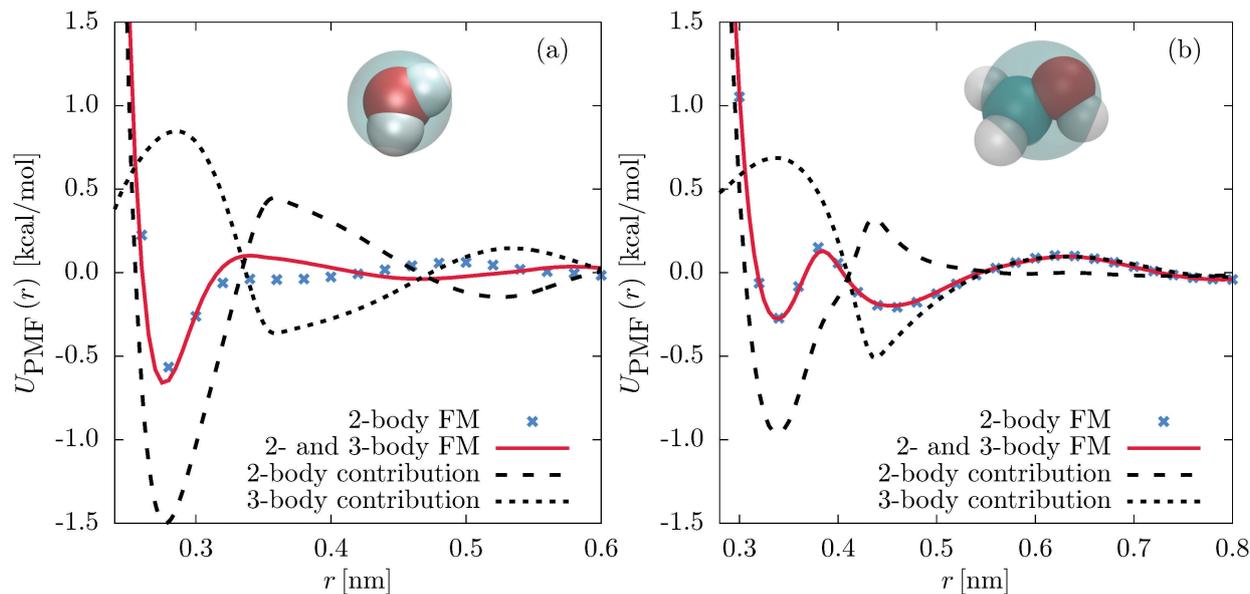


Fig. 3 Decomposition of the CG two-body PMF for (a) SPC/E water and (b) liquid methanol. The labels 2-body FM, 2- and 3-body FM refer to the two different CG parametrizations: FM with a (tabulated) pair potential only, and concurrent FM with tabulated pair and short-range three-body SW potential. The 2-body and 3-body contributions refer to the decomposition of the concurrent two- and three-body FM parametrization.

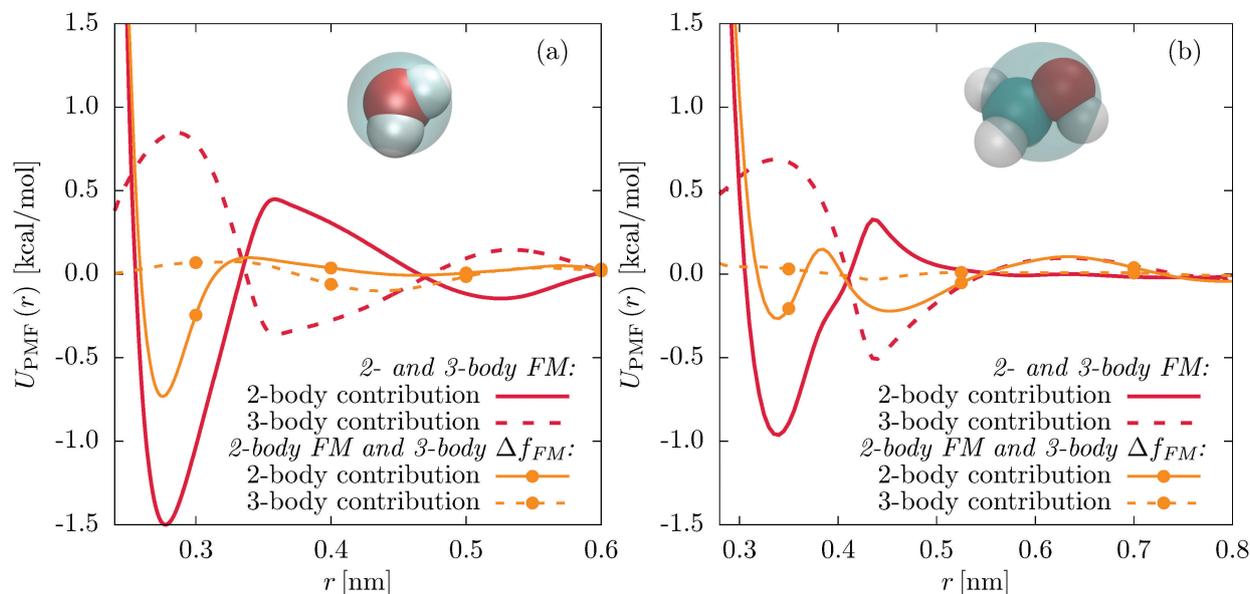


Fig. 4 Decomposition of the CG two-body PMF for different CG parametrization schemes for (a) SPC/E water and (b) methanol. The results shown are for the concurrent two-body and three-body FM parametrization (2- and 3-body FM), and the three-body FM using the residual force of the two-body FM potential (2-body FM and 3-body Δf_{FM}).

also improves the RDF compared to the two-body FM result. For liquid methanol, the same is observed for the angular distribution functions (see Fig. 2(b)). Regarding the RDF, the first peak is slightly suppressed in case of the residual force parametrization (2-body FM and 3-body Δf_{FM} , Fig. S6 of the ESI†). A similar effect has been observed when adding a many-body local density-dependent potential to a two-body pair potential for liquid methanol.²⁶

As a complementary information to the decomposition of the pair PMF into the two- and three-body interaction terms, we evaluate

the inner product p and average angle Φ between the two-body and three-body forces on each CG bead (see eqn (7) and (8)). We find that the constrained parametrization scheme indeed reduces the inner products from $p \approx -0.5$ ($\Phi \approx 120^\circ$) to $p \approx -0.03$ ($\Phi \approx 92^\circ$) (water) and $p \approx -0.6$ ($\Phi \approx 125^\circ$) to $p \approx 0.05$ ($\Phi \approx 87^\circ$) (methanol).

B. Thermodynamics

We now assess the performance of the CG models in terms of thermodynamic quantities. In particular, we evaluate the virial



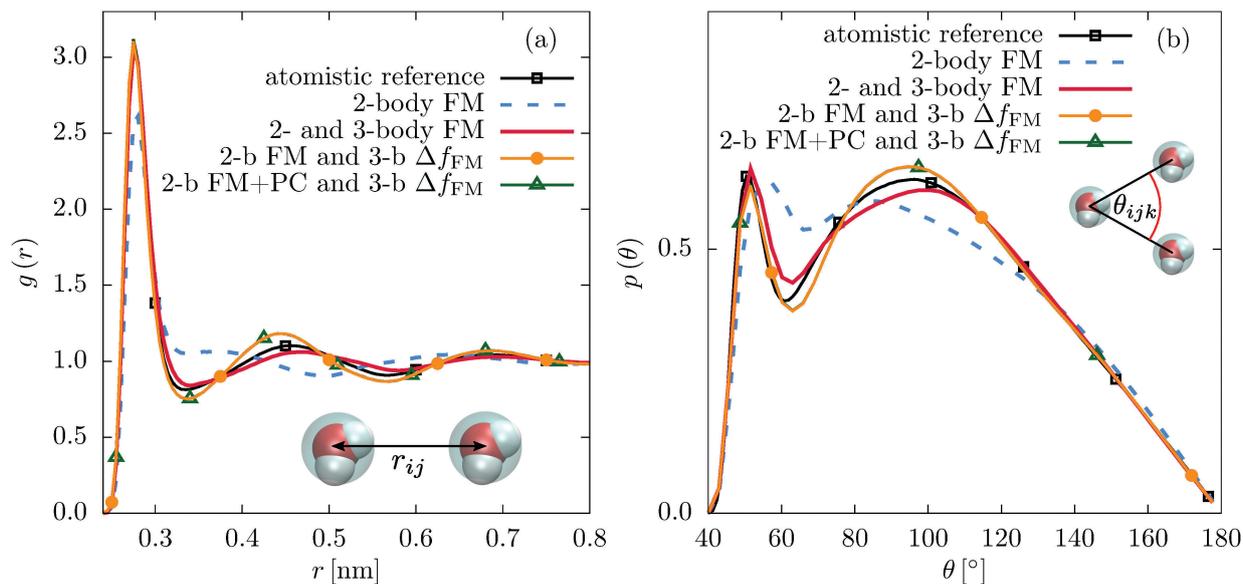


Fig. 5 (a) Radial and (b) angular atomistic reference distribution functions of SPC/E water compared to the ones of different CG parametrizations: FM with a (tabulated) pair potential only (2-body FM), concurrent two-body and three-body FM parametrization (2- and 3-body FM), the three-body FM parametrization using the residual force of the two-body FM potential (2-b FM and 3-b Δf_{FM}), and the three-body FM parametrization using the residual force of the two-body FM potential with additional pressure correction (2-b FM + PC and 3-b Δf_{FM}). The angular distribution function is calculated for the first coordination shell using a cutoff of 0.37 nm.

pressure P which should be equal to the ambient pressure when simulated at the density of the atomistic model, as well as the enthalpy of vaporization $\Delta H = H_{\text{gas}} - H_{\text{liq}}$. Note that we do not explicitly conduct CG gas phase simulations and assume an ideal CG gas.^{28,37} In addition, we take into account the actual average pressure of the liquid phase CG simulations. Calculation details can be found in the ESI.†

So far, all structural properties referred to a SW interaction parameter of $\gamma = 0.08$ nm (see eqn (4)), which is the optimal value regarding the structural accuracy of the models. Fig. 6 shows how P changes with the SW interaction parameter γ for all different parametrization schemes. For the unconstrained models, P decreases with increasing γ until reaching its minimum value at $\gamma = 0.08$ nm, namely $P = 0.12$ kbar for water and $P = 0.83$ kbar for

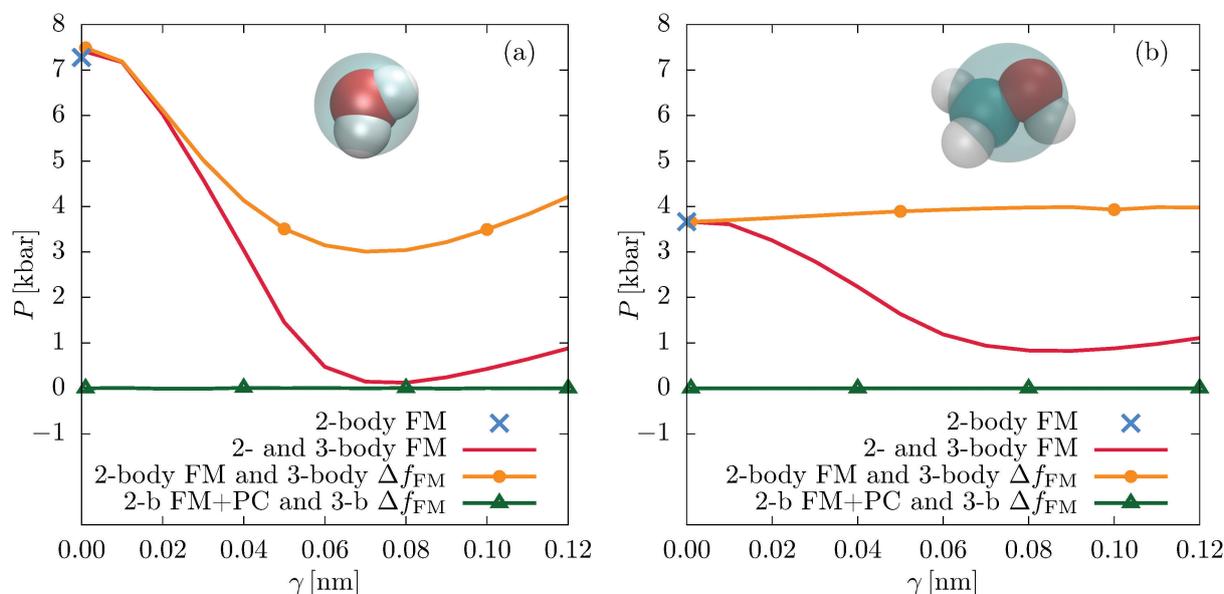


Fig. 6 Average pressures P for different CG parametrization schemes for (a) SPC/E water and (b) methanol. The results shown are for the two-body FM (2-body FM), the concurrent two-body and three-body FM parametrization (2- and 3-body FM), the three-body FM using the residual force of the two-body FM potential (2-body FM and 3-body Δf_{FM}), and the three-body FM parametrization using the residual force of the two-body FM potential with additional pressure correction (2-b FM + PC and 3-b Δf_{FM}).



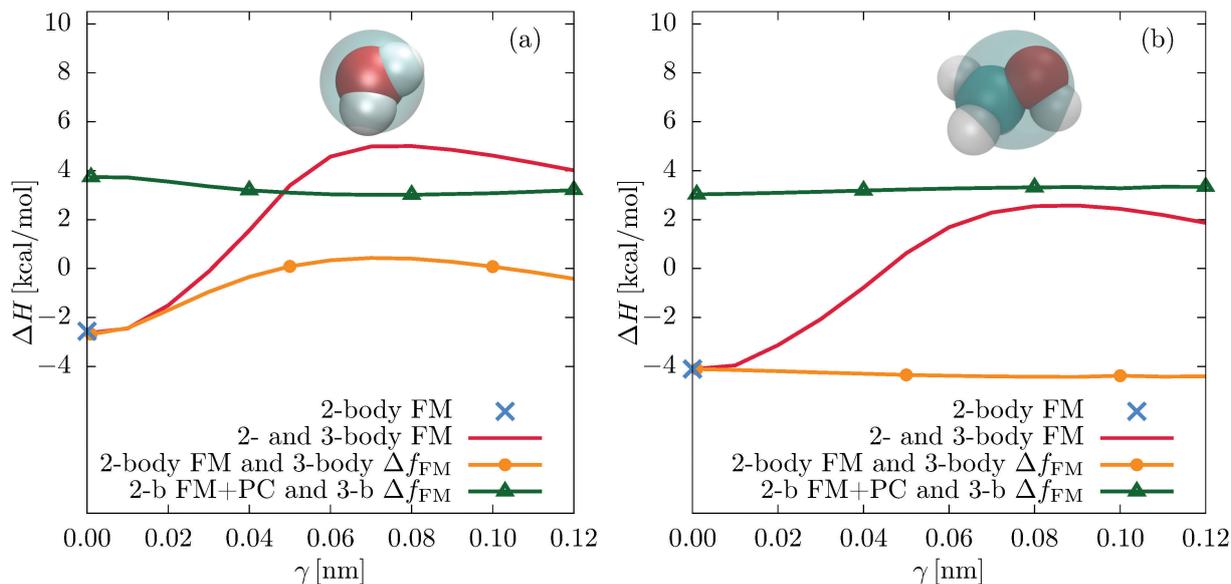


Fig. 7 Average enthalpies of vaporization ΔH for different CG parametrization schemes for (a) SPC/E water and (b) methanol. The results shown are for the two-body FM (2-body FM), the concurrent two-body and three-body FM parametrization (2- and 3-body FM), the three-body FM using the residual force of the two-body FM potential (2-body FM and 3-body Δf_{FM}), and the three-body FM parametrization using the residual force of the two-body FM potential with additional pressure correction (2-b FM + PC and 3-b Δf_{FM}).

methanol. This corresponds to system densities of $\rho = 0.990 \text{ g cm}^{-3}$ (water) and $\rho = 0.613 \text{ g cm}^{-3}$ (methanol) for *NPT* simulations at 1 atm. These values are close to the atomistic values of $\rho = 0.998 \text{ g cm}^{-3}$ (water) and $\rho = 0.776 \text{ g cm}^{-3}$ (methanol), and the experimental reference values at normal conditions: $\rho = 0.9971 \text{ g cm}^{-3}$ (water³⁶) and $\rho = 0.7872 \text{ g cm}^{-3}$ (methanol³⁷), even though these thermodynamic quantities have not been explicitly targeted in our parametrizations as, for example, with an explicit pressure matching scheme.⁴⁷

The constrained parametrizations (2-body FM and 3-body Δf_{FM}) yield significantly larger values of P . This is a known issue of the FM method which can be alleviated by adding a small linear perturbation to the two-body FM potential (pressure correction):^{2,18}

$$\Delta U_{\text{PC}} = -A \left(1 - \frac{r}{r_{\text{cut}}} \right), \quad r_{\text{cut}} = 1.2 \text{ nm}, \quad (10)$$

where A is a positive constant. The pressure correction adds a small constant attractive force to each particle pair within the cutoff radius r_{cut} . The value of A can be adjusted for each γ (see Table S1 of the ESI†) to shift the pressure to zero (see Fig. 6). The pressure corrected two-body FM potentials are shown in Fig. 1(a) (water) and Fig. 2(a) (methanol). For the full potential range until $r_{\text{cut}} = 1.2 \text{ nm}$ we refer to Fig. S2(a) and S5(a) of the ESI.† *NPT* simulations at 1 atm yield system densities of $\rho = 0.996 \text{ g cm}^{-3}$ (water) and $\rho = 0.777 \text{ g cm}^{-3}$ (methanol) which are identical to the atomistic values.

The pressure correction only slightly perturbs the structure of the system. Indeed, the pair and the angular distribution functions of the corrected and uncorrected constrained parametrizations are practically identical (see Fig. 2(b) and 5, and Fig. S6 of the ESI†). The same holds for the decompositions of

the two-body PMFs, implying that the pressure correction does not change the inner product of the two- and three-body interactions (see Fig. S3, S4, S7, and S8 of the ESI†).

Fig. 7 shows the enthalpies of vaporization ΔH for all different models. Again, the two-body FM value is indicated with a blue cross. The unconstrained parametrizations (2- and 3-body FM) result in a maximum of ΔH at $\gamma = 0.08 \text{ nm}$, where $\Delta H = 5.01 \text{ kcal mol}^{-1}$ (water) and $\Delta H = 2.55 \text{ kcal mol}^{-1}$ (methanol). The corresponding atomistic (experimental) values are $11.76 \text{ kcal mol}^{-1}$ ($10.52 \text{ kcal mol}^{-1}$) for water²⁸ and $8.94 \text{ kcal mol}^{-1}$ ($8.95 \text{ kcal mol}^{-1}$) for methanol.³⁷

The pressure correction significantly improves the vaporization enthalpies of the constrained parametrizations by increasing the cohesive energy. For example, for $\gamma = 0.08 \text{ nm}$, ΔH of water changes from $0.40 \text{ kcal mol}^{-1}$ to $3.01 \text{ kcal mol}^{-1}$ and the vaporization enthalpy of methanol changes from $-4.42 \text{ kcal mol}^{-1}$ to $3.32 \text{ kcal mol}^{-1}$.

IV. Conclusions

We have shown that adding a Stillinger–Weber three-body potential to the coarse-grained force field of liquid water and methanol leads to significantly more attractive pair potentials. This short-range attraction is then compensated by the three-body term. Parametrization of the three-body term using residual forces substantially reduces the three-body contribution to the two-body potential of mean force. At the same time it worsens the thermodynamic properties of the coarse-grained model, quantified by pressure and enthalpy of vaporization. The thermodynamic properties can be corrected by adding a weak long-range attractive potential (pressure correction)



without changing the structure of the system and the orthogonality of the three- and two-body contributions. Our findings motivate basis set extensions which separate the many-body contributions of different order.

Conflicts of interest

There are no conflicts to declare.

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References

- W. G. Noid, J.-W. Chu, G. S. Ayton, V. Krishna, S. Izvekov, G. A. Voth, A. Das and H. C. Andersen, *J. Chem. Phys.*, 2008, **128**, 244114.
- D. Reith, M. Pütz and F. Müller-Plathe, *J. Comput. Chem.*, 2003, **24**, 1624.
- A. P. Lyubartsev and A. Laaksonen, *Phys. Rev. E: Stat. Phys., Plasmas, Fluids, Relat. Interdiscip. Top.*, 1995, **52**, 3730.
- M. S. Shell, *J. Chem. Phys.*, 2008, **129**, 144108.
- S. Izvekov, M. Parrinello, C. J. Burnham and G. A. Voth, *J. Chem. Phys.*, 2004, **120**, 10896.
- G. Tóth, *J. Phys.: Condens. Matter*, 2007, **19**, 335222.
- J. W. Mullinax and W. G. Noid, *J. Phys. Chem. C*, 2010, **114**, 5661.
- J. F. Rudzinski and W. G. Noid, *Eur. Phys. J.: Spec. Top.*, 2015, **224**, 2193.
- J. F. Rudzinski and W. G. Noid, *J. Phys. Chem. B*, 2014, **118**, 8295.
- S. Izvekov and G. A. Voth, *J. Chem. Phys.*, 2005, **123**, 134105.
- J. W. Mullinax and W. G. Noid, *J. Chem. Phys.*, 2009, **131**, 104110.
- V. Molinero and E. B. Moore, *J. Phys. Chem. B*, 2009, **113**, 4008.
- L. Larini, L. Lu and G. A. Voth, *J. Chem. Phys.*, 2010, **132**, 164107.
- A. Das and H. C. Andersen, *J. Chem. Phys.*, 2012, **136**, 194114.
- A. Das and H. C. Andersen, *J. Chem. Phys.*, 2009, **131**, 034102.
- M. Jochum, D. Andrienko, K. Kremer and C. Peter, *J. Chem. Phys.*, 2012, **137**, 064102.
- C. Scherer and D. Andrienko, *Eur. Phys. J.: Spec. Top.*, 2016, **225**, 1441.
- V. Rühle, C. Junghans, A. Lukyanov, K. Kremer and D. Andrienko, *J. Chem. Theory Comput.*, 2009, **5**, 3211.
- A. P. Lyubartsev and A. Laaksonen, *Phys. Rev. E: Stat. Phys., Plasmas, Fluids, Relat. Interdiscip. Top.*, 1997, **55**, 5689.
- A. P. Bartók and G. Csányi, *Int. J. Quantum Chem.*, 2015, **115**, 1051.
- A. Tsourtis, V. Harmandaris and D. Tsagkarogiannis, *Entropy*, 2017, **19**, 395.
- M. Dijkstra, J. M. Brader and R. Evans, *J. Phys.: Condens. Matter*, 1999, **11**, 10079.
- P. G. Bolhuis, A. A. Louis and J. P. Hansen, *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.*, 2001, **64**, 021801.
- T. Sanyal and M. S. Shell, *J. Chem. Phys.*, 2016, **145**, 034109.
- J. W. Wagner, T. Dannenhoffer-Lafage, J. Jin and G. A. Voth, *J. Chem. Phys.*, 2017, **147**, 044113.
- M. R. DeLyser and W. G. Noid, *J. Chem. Phys.*, 2017, **147**, 134111.
- F. Zipoli, T. Laino, S. Stolz, E. Martin, C. Winkelmann and A. Curioni, *J. Chem. Phys.*, 2013, **139**, 094501.
- J. Lu, Y. Qiu, R. Baron and V. Molinero, *J. Chem. Theory Comput.*, 2014, **10**, 4104.
- G. A. Cisneros, K. T. Wikfeldt, L. Ojamäe, J. Lu, Y. Xu, H. Torabifard, A. P. Bartók, G. Csányi, V. Molinero and F. Paesani, *Chem. Rev.*, 2016, **116**, 7501.
- A. Das and H. C. Andersen, *J. Chem. Phys.*, 2012, **136**, 194113.
- F. H. Stillinger and T. A. Weber, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1985, **31**, 5262.
- M. J. Abraham, T. Murtola, R. Schulz, S. Páll, J. C. Smith, B. Hess and E. Lindahl, *SoftwareX*, 2015, **1–2**, 19.
- H. J. C. Berendsen, J. R. Grigera and T. P. Straatsma, *J. Phys. Chem.*, 1987, **91**, 6269.
- W. L. Jorgensen and J. Tirado-Rives, *J. Am. Chem. Soc.*, 1988, **110**, 1657.
- W. L. Jorgensen, D. S. Maxwell and J. Tirado-Rives, *J. Am. Chem. Soc.*, 1996, **118**, 11225.
- J. L. F. Abascal and C. Vega, *J. Chem. Phys.*, 2005, **123**, 234505.
- C. Caleman, P. J. van Maaren, M. Hong, J. S. Hub, L. T. Costa and D. van der Spoel, *J. Chem. Theory Comput.*, 2012, **8**, 61.
- W. F. Van Gunsteren and H. J. C. Berendsen, *Mol. Simul.*, 1988, **1**, 173.
- H. J. Berendsen, J. v. Postma, W. F. van Gunsteren, A. DiNola and J. R. Haak, *J. Chem. Phys.*, 1984, **81**, 3684.
- U. Essmann, L. Perera, M. L. Berkowitz, T. Darden, H. Lee and L. G. Pedersen, *J. Chem. Phys.*, 1995, **103**, 8577.
- S. Plimton, *J. Comput. Phys.*, 1995, **117**, 1.
- S. Nosé, *Mol. Phys.*, 1984, **52**, 255.
- W. G. Hoover, *Phys. Rev. A: At., Mol., Opt. Phys.*, 1985, **31**, 1695.
- W. Shinoda, M. Shiga and M. Mikami, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2004, **69**, 134103.
- J. F. Rudzinski and W. G. Noid, *J. Phys. Chem. B*, 2012, **116**, 8621.
- W. G. Noid, P. Liu, Y. Wang, J.-W. Chu, G. S. Ayton, S. Izvekov, H. C. Andersen and G. A. Voth, *J. Chem. Phys.*, 2008, **128**, 244115.
- N. J. H. Dunn and W. G. Noid, *J. Chem. Phys.*, 2015, **143**, 243148.

