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Influence of a silver salt on the nanostructure of a Au(111)/ionic liquid interface: an atomic force microscopy study and theoretical concepts†

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Ionic liquids (ILs) form a multilayered structure at the solid/electrolyte interface, and the addition of solutes can alter it. For this purpose, we have investigated the influence of the silver bis(trifluoromethylsulfonyl)amide (AgTFSA) concentration in 1-butyl-1-methylpyrrolidinium bis(trifluoromethylsulfonyl)amide ([Py_{1,4}][TFSA]) on the layering using *in situ* atomic force microscopy. AFM investigations revealed that the Au(111)/electrolyte interface indeed depends on the concentration of the salt where a typical “IL” multilayered structure is retained only at quite low concentrations of the silver salt (e.g. ≤ 200 μM). However, at 200 μM AgTFSA/[Py_{1,4}][TFSA] and above this “IL” multilayered structure is disturbed/varied. A simple double layer structure was observed at 500 μM AgTFSA in [Py_{1,4}][TFSA]. Furthermore, the widths of the innermost layers have been found to be dependent on the concentration and on the applied electrode potentials. Our AFM results show that the concentration of solutes strongly influences the structure of the electrode/electrolyte interface and can provide new insights into the electrical double layer structure of the electrode/ionic liquid interface. We also introduce a semi-continuum theory to discuss the double layer structure.

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1 Introduction

The electrodeposition of silver on Au(111) is a well-investigated example, both from surface science and from the electrochemical perspective due to quite similar lattice constants and the higher surface energy of gold.^{1,2} Furthermore, alloying,^{1,3} surface reconstruction,^{4,5} surface crystallinity,⁶ and defects¹ can play a crucial role in the initial stages of metal deposition. Moreover, the neat surfaces of gold minimize their surface energy by reconstruction and an electrode potential dependent lifting of Au has been reported.^{7,8} Specific adsorption of anions can also promote such a lifting process² and the strong influence of anions on the overlayer structure was observed using atomic force microscopy.²

Ionic liquids are interesting solvents for the electrodeposition of reactive metals such as Al and Si, which are not accessible in aqueous media.^{9,10} These liquids have wide electrochemical and

thermal windows, good ionic conductivities, usually low vapor pressures at room temperature and high solubility for a variety of compounds.^{11,12} However, the practical use of these liquids in applied electrochemistry is hindered, as there is still no thorough understanding of the species in the bulk phase and, especially, at the interface.

Several ILs have been employed as electrolytes to deposit silver, such as tetrafluoroborate,^{13,14} hexafluorophosphate,^{15,16} bis(trifluoromethylsulfonyl)amide,^{13,17} dicyanamide,¹⁸ trifluoromethylsulfonate,^{19,20} acetate,²¹ and nitrate.²¹ In most of these studies the focus was on the electrodeposition and the characterization of the deposited films.

In recent years, atomic force microscopy has become a useful tool to explore/probe the electrode/electrolyte interface (EEI) with atomic resolution.^{22–34} AFM force–distance curves can provide direct insight into the adsorption structure of ions/species in the double layer (or multilayer) on the electrode/IL interface.²²

In contrast to aqueous solutions, there are often no simple double layers in ILs, instead the ions adopt a multilayer structure.²⁸ In brief, the interfacial structure can be divided into the innermost layer, the transition zone, and the bulk liquid phase.^{29,30} The innermost layer is composed of the IL layer adsorbed next to the electrode surface. These layers experience large forces when

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the AFM tip pushes through them. At the open circuit potential (OCP) and at more negative electrode potentials the innermost layer is usually enriched with cations. At more positive potentials anions are usually dominant at the interface. The forces for penetrating these layers have been shown to be more than one order of magnitude higher than the ones for molecular liquids.^{25,29,31,35–38} Interfacial properties, including the adsorption/co-adsorption of ions at the interface, are determined by the nature of the electrode and the electrolyte. Even a small variation in the electrolyte composition, *e.g.* due to impurities³⁹ or dissolved solutes,^{25,40} can disturb the multilayer arrangement and can also influence electrochemical reactions. The Au(111)/[Py_{1,4}]TFSA interface is rather well described in the literature.^{35,41} Furthermore, the solvation behaviour of AgTFSA in [EMIm]TFSA has been investigated by spectroscopy and DFT calculations, giving a stable complex like [Ag(TFSA)₃]^{2–}.⁴² We have reported that water has an influence on the Au(111)/ionic liquid interface.⁴³ In the case of [EMIm]TfO, a multilayer structure was predominant up to 30 vol% water and above this concentration the multilayer structure was strongly disturbed. In this context the question came up as to what extent the IL layers are influenced by a solute during electrodeposition. For this purpose AgTFSA, from which Ag is easily deposited, has been taken as an example to study the influence of salt concentration on the electrode/ionic liquid interface. In the present paper we report for the first time on the influence of silver salt concentration on the electrode/electrolyte interface using atomic force microscopy.

The experimental work is accompanied by theoretical considerations. Continuum theories allow the dynamic description of macroscopic systems, *e.g.*, electrochemical cells. The nano-sized surface layers, however, are strongly affected by the particle nature of its constituents. The popular modeling approaches on this scale are density functional theory (DFT)^{44–47} and molecular dynamics (MD).⁴⁸ DFT and MD simulations resolve many of the microscopic complexities of individual molecules and intermolecular interactions, but remain limited to small length and time scales. Bazant *et al.* have proposed a phenomenological continuum theory which captures the multilayer structure of the electrochemical double layer in ionic liquids.²⁸

In this paper, we present a novel semi-continuum model which describes ionic liquids as a continuum fluid of hard spheres. It represents a unified theoretical framework for the description of both scales, bulk phase and surface layers. Our approach is derived from fundamental thermodynamical principles and basic physical assumptions, in contrast to phenomenological continuum theory.^{28,49,50} Thus, we improve the assignment between material properties and the simulated phenomena. Most importantly, we find that incompressibility leads to crowding and the particle nature of fluids leads to overscreening.

Our consistent continuum model describes both the short-range and long-range behavior of ionic liquids as well as both static and dynamic effects. This generality complements previous modeling efforts. On the one hand, DFT and MD describe electrolytes at smaller length and time scales. Thus, they resolve more detailed ionic properties than our model, *e.g.*, molecular orientation, position, and shape, and make

quantitative predictions. The fine resolution, however, comes at the cost of computational complexity. DFT and MD are limited to the nano-scale dimensions of the electrochemical double layer (EDL), whereas our model can calculate properties of ionic liquids from the nano-scale EDL to the micro-scale electrode structure. Most importantly, the simplicity and abstraction of our continuum model brings out the fundamental principles behind the interfacial behavior of ionic liquids clearer than do DFT or MD. Furthermore, our model is time dependent and able to predict the contribution of transport through the EDL structure to the apparent reaction kinetics of ionic liquids.

On the other hand, the phenomenological continuum theory of Bazant *et al.* is numerically simpler than our model, but remains restricted to binary ionic liquids and does not clearly distinguish between different electrolytes. Our model, instead, applies to more complex electrolytes and clearly connects with material properties because we derive it from non-equilibrium thermodynamics and make explicit the hardcore nature of the molecules. In this paper, for example, we can evaluate the effect of water molecules or silver ions onto the EDL structure of ionic liquids and study the effect of molecular size on electrolyte properties.

2 Methods

2.1 Experimental methods

1-Butyl-1-methylpyrrolidinium

bis(trifluoromethylsulfonyl)amide ([Py_{1,4}]TFSA) was purchased from IOLITEC, Germany, in a purity of 99.5%. Silver bis(trifluoromethylsulfonyl)amide (99.5%) was procured from Solvionics, France. The ionic liquid was dried at 100 °C under vacuum for 3 days to water values below 10 ppm. Several concentrations of AgTFSA in [Py_{1,4}]TFSA ranging from 50 μM to 1000 μM were prepared.

The electrochemical measurements were carried out using a PARSTAT 2263 potentiostat/galvanostat controlled by PowerCV and PowerStep software. A three electrode cell was used for all electrochemical experiments. Gold on glass was used as the working electrodes (WE) for cyclic voltammetry experiments. The working electrode for the AFM studies was Au(111) on mica purchased from Agilent Technologies. The electrochemical cell was made of Teflon and clamped over a Teflon-covered Viton O-ring onto the working electrode with a geometric surface area of 0.3 cm². Pt wires were used as the counter and quasi reference electrodes, respectively. The deflection–distance curves were obtained using a Molecular Imaging Pico Plus AFM in contact mode inside an argon-filled glove-box at room temperature and transformed to force–distance curves. A silicon SPM-sensor with a spring constant of 6 N m^{–1} from Nano World was employed for all AFM measurements.

2.2 Mathematical model

Our mathematical model of ionic liquids is based on the methodology of rational thermodynamics and describes a thermodynamically consistent transport theory. The modeling methodology has been applied to electrolytes with a dominant



neutral solvent before, as outlined by Latz and Zausch.⁵¹ However, ionic liquids possess unique physical properties which must be captured by fundamental principles and which differ from those of uncharged solvent electrolytes. Therefore, transport theories for such media are not applicable to ionic liquids. This motivates the formulation of a novel transport theory for ionic liquids.

Furthermore, our definition of incompressibility is based on a volumetric constraint^{52,53} and differs from the one used in ref. 51. Here we show how to supplement the continuum theory with non-local interactions, modeling the particle nature of the constituents. Therefore, our approach can be used to examine both the bulk behavior, and the quasi-crystalline double-layer behavior of ionic liquids.

Thermodynamics. The equations of motion are derived following the method of Coleman and Noll supplemented by a linear Onsager Ansatz. The central part of this method is the free energy F . First, we assume that the free energy $F(\rho_\alpha, \mathbf{D})$ depends on the mass densities ρ_α and the dielectric displacement \mathbf{D} only. Based on this simple assumption and the fundamental laws of thermodynamics, the transport equations can be systematically derived. Second, by modeling its specific form, we specify all the material properties of the ionic liquid. In this way, the only freedom of modeling is the determination of a suitable free energy, the rest is a rigorous mathematical consequence of this choice.

We divide the free energy

$$F = \int_{\Omega} \rho \varphi_H dV + F^{\text{int}} \quad (1)$$

into a standard part described by a single-particle free energy density $\rho \varphi_H$ and a many-particle interaction energy F^{int} . The latter term contains energetic contributions involving more than one particle, e.g., molecular attraction. In this work we restrict F^{int} to hardcore repulsion.

The electrolyte equation of state is formulated as a constraint on the concentrations c_α and the partial molar volumes ν_α , such that $\sum_{\alpha=1}^N \nu_\alpha c_\alpha = 1$.^{52–54} If the partial molar volumes do not depend on pressure, this describes incompressibility. As consequence, we derive the convective equation of state $\nabla \mathbf{v} = -\sum_{\alpha=1}^N \nu_\alpha \nabla \mathbf{N}_\alpha$, where \mathbf{v} is the convective bulk velocity and the ionic fluxes \mathbf{N}_α measure transport in the center-of-mass frame.

The single-particle free energy density is chosen as

$$\rho \varphi_H = \frac{\mathbf{E} \cdot \mathbf{D}}{2} + \frac{\mathcal{K}}{2} \left(1 - \sum_{\alpha=1}^N \nu_\alpha c_\alpha \right)^2 + RT \sum_{\alpha=1}^N c_\alpha \ln \left(\frac{c_\alpha}{c_0} \right), \quad (2)$$

modeling the electrostatic energy, elastic energy, and entropy. The elastic energy enforces the incompressibility $\sum_{\alpha=1}^N \nu_\alpha c_\alpha = 1$ because we let the bulk modulus diverge $\mathcal{K} \rightarrow \infty$. The elastic contribution determines the force density *via* the Navier–Stokes equation

$$\rho \dot{\mathbf{v}} = \nabla \left[\frac{\mathcal{K}}{2} \left(1 - \left(\sum_{\alpha=1}^N \nu_\alpha c_\alpha \right)^2 \right) - \rho_F E \right]. \quad (3)$$

We omit the inertial forces $\rho \dot{\mathbf{v}} = 0$ in the highly viscous ionic liquids. Thus, we can express elastic forces in terms of electrostatic forces in

our transport equations below. The remnant of the single-particle free energy density is analogous to the one used in ref. 51.

Hardcore particles. We describe ionic liquids as composed of oppositely charged (valency ± 1) hard spheres. Although our formalism allows implementing different volumes for particles of different species, we assume an equal volume here. Our approach allows fast simulations of both the interfacial and bulk processes. The particle nature of the electrolyte is modeled *via* an interaction contribution to the free energy. We capture the requirement that the volume of the hard spheres is inaccessible to distinct particles by a repulsive potential between the particles. In particular, we choose

$$F^{\text{int}}[c] = \frac{1}{2} \sum_{\alpha, \beta} \int_{\Omega} c_\alpha(\mathbf{x}) \mathcal{F}_{\alpha\beta}(\mathbf{x} - \mathbf{y}) c_\beta(\mathbf{y}) d\mathbf{x} d\mathbf{y}. \quad (4)$$

This choice of the repulsive interaction $\mathcal{F}_{\alpha\beta}(\mathbf{x} - \mathbf{y})$ models the impenetrable hard spheres, *i.e.*, the repulsive interaction is non-zero in a small volume Ω determined by the particle-radii. The mathematical form of $\mathcal{F}_{\alpha\beta}$ is given in the ESI.† The modification of the chemical potentials due to this inter-particle repulsion is derived with functional derivatives

$$\mu_\alpha^{\text{int}}(\mathbf{x}) = \frac{\delta F^{\text{int}}[c]}{\delta c_\alpha} = \int_{\Omega} \mathcal{F}_{\alpha\beta}(|\mathbf{x} - \mathbf{y}|) c_\beta(\mathbf{y}) d\mathbf{y}. \quad (5)$$

Transport model. An Onsager Ansatz yields the equations of motion for ionic fluxes and the electric current for the incompressible ionic liquid ([Py_{1,4}]TFSa) in the isothermal limit. We restrict the following discussion to the ternary case with the pure ionic liquid and a single additive, e.g., water molecules or silver ions. If we know the charge density ρ_F and the additive concentration c_3 we can calculate the concentrations c_i of all the species. This follows from the electrolyte equation of state, $\nu_+ c_+ + \nu_- c_- + \nu_3 c_3 = 1$, and the expansion of the charge density, $\rho_F = F(c_+ - c_- + z_3 c_3)$. We model transport with the following equations

$$\partial_t \rho_F = -\nabla \mathbf{J} - \nabla(\rho_F \mathbf{v}), \quad (6)$$

$$\partial_t c_3 = -\nabla \mathbf{N}_3 - \nabla(c_3 \mathbf{v}), \quad (7)$$

$$\rho_F = -\varepsilon_R \varepsilon \Delta \Phi, \quad (8)$$

i.e., the continuity of charge, the continuity of the additive, and the Poisson equation, respectively. The convection velocity is given by

$$\nabla \mathbf{v} = \frac{1}{F} \frac{M_- \nu_+ - M_+ \nu_-}{M_+ z_- - M_- z_+} \cdot \nabla \mathbf{J} - \left(\frac{M_+(z_- \nu_3 - z_3 \nu_-)}{M_+ z_- - M_- z_+} + \frac{z_+(M_3 \nu_- - M_- \nu_3)}{M_+ z_- - M_- z_+} + \frac{\nu_+(M_- z_3 - M_3 z_-)}{M_+ z_- - M_- z_+} \right) \cdot \nabla \mathbf{N}_3, \quad (9)$$

where \mathbf{J} denotes the electric current in the center of the mass frame

$$\mathbf{J} = -\kappa \cdot \nabla \Phi + \frac{\kappa}{F} \cdot (1 - t_3) \cdot \frac{M_- \nabla \mu_+ - M_+ \nabla \mu_-}{M_+ z_- - M_- z_+} - \frac{\kappa}{F} \cdot t_3 \cdot \frac{M_+ \nabla \mu_3 - M_3 \nabla \mu_+}{M_+ z_3 - M_3 z_+}. \quad (10)$$



Here, the particle flux of the additive

$$\begin{aligned} N_3 = & D \cdot \frac{M_{3z-} - M_{-z3}}{M_{+z-} - M_{-z+}} \cdot \nabla \mu_+ + D \cdot \frac{M_{+z3} - M_{3z+}}{M_{+z-} - M_{-z+}} \cdot \nabla \mu_- \\ & - D \cdot \nabla \mu_3 + \frac{t_3}{F} \cdot \frac{M_+}{M_{+z3} - M_{3z+}} \cdot \mathbf{J}, \end{aligned} \quad (11)$$

depends on the electric current, on the three chemical forces and on the transport parameters of silver in the electrolyte solution, on the diffusion coefficient D , and on the transference number t_3 . The chemical forces $\nabla \mu_x$

$$\nabla \mu_x = -\nu_x \rho_F \nabla \Phi + RT \nabla \ln \left(\frac{c_x}{c_0} \right) + \nabla \mu_x^{\text{int}} \quad (12)$$

are determined by the Navier–Stokes eqn (3) which couples the elastic and electrostatic forces. The material parameters^{55,56} are listed in the ESI.† The above set of equations describes the dynamics and equilibrium of an ionic liquid with additive. In this work, we focus on their stationary solution near an electrified interface. The electrolyte structure is reflected in the electrolyte potential distribution Φ and the charge density ρ_F . In our AFM experiments, we perturb this structure and obtain information on its size and stability. In our simulations, we have access to the elastic reaction of the system when subjected to mechanical excitations. The mechanical forces in the EDL are compensated for by electrostatic forces. As a consequence, the generalized Lorentz force density $\mathbf{f} = \rho_F \nabla \Phi$ appearing in the Navier–Stokes equation resembles these mechanical forces. We obtain a force

$$\mathbf{F} = \frac{(\nu_+ + \nu_-)}{N_A} \mathbf{f} = \frac{(\nu_+ + \nu_-)}{N_A} \rho_F \nabla \Phi \quad (13)$$

by multiplying \mathbf{f} with the molecular volume. This simulated force is compared with the experimental results delivered by AFM measurements testing the stationary multilayer structure of the EDL. While the AFM tip approaches the electrode surface, it penetrates the EDL and pushes ions away. Thus, the AFM measurement is an inelastic probe of the EDL, while our simulations determine an elastic force. Nevertheless, the simulated and measured forces both represent the mechanical reaction to perturbations of the system.

The formalism outlined above can be generalized to electrolytes with a neutral solvent and an arbitrary number of solutes. In future work, we will extend the transport theory in this regard and include viscous media and heat transport. It is the strength of our consistent approach that the contributions of the solvent are not hidden as in standard concentrated solution theory.⁵⁷

3 Results and discussions – electrochemical and interfacial investigations of the electrolytes

3.1 Experimental results

Cyclic voltammetry has been employed to investigate the electrochemical behaviour of the electrolytes at varying concentrations of AgTFSA in $[\text{Py}_{1,4}]\text{TFSA}$ on gold at RT. Multiple-cycles have

been recorded for $\text{AgTFSA}/[\text{Py}_{1,4}]\text{TFSA}$ electrolytes to investigate changes in the interfacial processes/redox processes. All cyclic voltammograms have been recorded first in the negative direction from the Open Circuit Potential (OCP) to the chosen switching potentials at a scan rate of 10 mV s^{-1} . For the neat liquid, a rise in reduction current has been observed at $\sim -0.7 \text{ V}$, which can be due to the adsorption of IL on Au (Fig. 1a).³⁵ Quite a weak oxidation peak has been observed at -0.9 V in the backward scan of the CV. Upon addition of $50 \text{ }\mu\text{M}$ AgTFSA to $[\text{Py}_{1,4}]\text{TFSA}$, three reduction processes have been noticed in the forward scan of the CV (Fig. 1b). The first two reduction processes C1 and C2 could be due to the adsorption processes of the electrolyte species and to the underpotential deposition of Ag. The reduction process C3 can be clearly correlated with the bulk deposition of silver.

For $50 \text{ }\mu\text{M}$ AgTFSA/ $[\text{Py}_{1,4}]\text{TFSA}$, two oxidation processes have been observed in the backward scan of the CV (Fig. 1c). The oxidation process A1' is related to the reduction process C3 and the oxidation process A1 is associated with the reduction process C2. A broad reduction (C_R) wave has been noticed after the first scan, which can be related to the surface reduction process due to a change in the electrode surface. The third and fifth CVs also exhibit similar features to that of the first scan except for a broad reduction process C_R . For the CV of $100 \text{ }\mu\text{M}$ AgTFSA/ $[\text{Py}_{1,4}]\text{TFSA}$, a reduction wave is seen at C1 followed by a reduction step at C2. Here, the reduction step C1 could be due to an adsorption of the electrolyte, and C2 is due to the deposition of silver, respectively. The processes A1' and A1 are correlated with the desorption of adsorbed layers of the electrolyte from the electrode surface and with the dissolution of the deposited silver films, respectively.

The CV of $200 \text{ }\mu\text{M}$ AgTFSA/ $[\text{Py}_{1,4}]\text{TFSA}$ (Fig. 1d) consists of a clear reduction peak at C2 along with a broad reduction wave at C1 in the forward scan and three oxidation steps in the backward scan. The reduction process C2 with a peak at -1.6 V is due to the bulk deposition of Ag and the process C1 is likely to be correlated with the underpotential deposition of Ag. The oxidation process A1 can be related to the reduction process C2. Moreover, it has been observed that silver nanoparticles can form chemically under light irradiance. The third and fifth CVs exhibit similar features to those of the first scan. Furthermore, the adsorption process that was observed in the neat IL and at low concentrations of AgTFSA/ $[\text{Py}_{1,4}]\text{TFSA}$ (50 and $100 \text{ }\mu\text{M}$) has not been observed for $200 \text{ }\mu\text{M}$ AgTFSA/ $[\text{Py}_{1,4}]\text{TFSA}$, obviously an effect of the increase in concentration of AgTFSA. In the case of the zinc ions in ILs, it has been proposed that such interfacial processes can be varied by the arrangement of IL anions and cations adjacent to the electrode in the compact layer (IHP, Inner Helmholtz Plane), which can dictate the kinetics of the redox processes.⁵⁸ Based on these electrochemical results the question arises as to how the interfacial processes of ILs vary in the presence of a solute and how are they dependent on its concentration.

In order to describe the electrode/electrolyte interface (or the interfacial structure) of ILs and solutions with a varying concentration of silver salts, we have performed *in situ* Atomic Force Microscopy (AFM) investigations. To the best of the



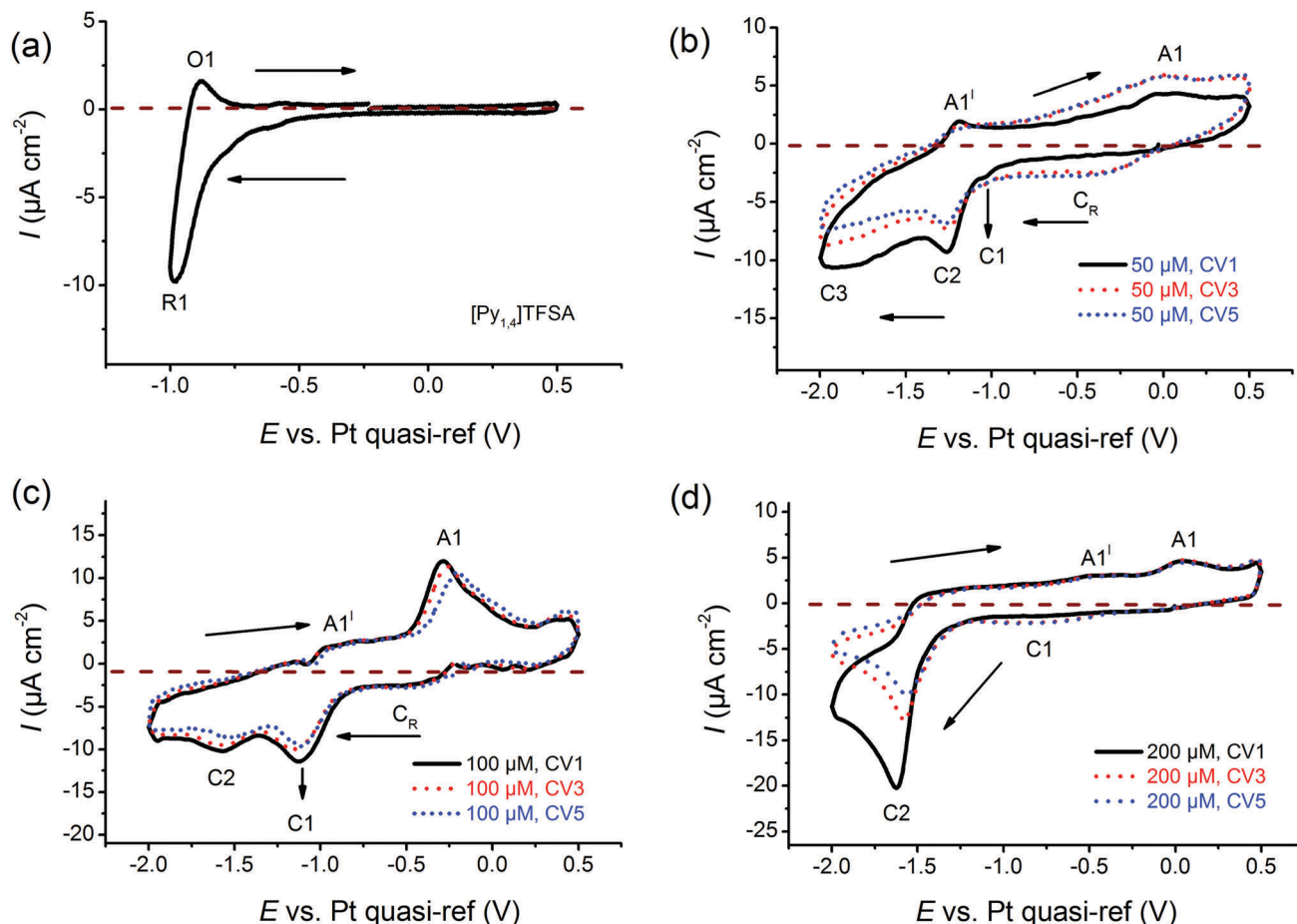


Fig. 1 Comparison of cyclic voltammograms of (a) neat $[\text{Py}_{1,4}]\text{TfSA}$, (b) $50\ \mu\text{M}$ $\text{AgTfSA}/[\text{Py}_{1,4}]\text{TfSA}$, (c) $100\ \mu\text{M}$ $\text{AgTfSA}/[\text{Py}_{1,4}]\text{TfSA}$, and (d) $200\ \mu\text{M}$ $\text{AgTfSA}/[\text{Py}_{1,4}]\text{TfSA}$. Scan rate = $10\ \text{mV s}^{-1}$.

authors' knowledge, such measurements, where a systematic variation of salt concentration has been performed, have not yet been reported so far.

Fig. 2 shows the force vs. separation curves for an AFM tip approaching the Au(111) surface in the neat IL and in mixtures with two different concentrations of AgTfSA and at three different electrode potentials (OCP, $+0.2\ \text{V}$ and $-0.2\ \text{V}$). The appearance of the AFM force curves differ with varying concentrations of AgTfSA (addition of solute) from the force curves of the Au(111)/IL interface.^{25,40,41} For the neat IL, the interfacial structure consists of at least four discrete steps at the measured electrode potentials, which is in agreement with the literature. This behavior reveals a strong ion-ion interaction between the solvation layers, as in this case by displacing the innermost solvation layers, the AFM tip pulls the other layers as well. For the Au(111) surface the first layer is observed at $\sim 0.6\ \text{nm}$ followed by the 1.5, 2.4, and 3.3 nm steps, which correspond to widths of ~ 0.6 and $0.9\ \text{nm}$, respectively. The steps correspond to the rupturing of successive adjacent surface layers as the AFM tip approaches the gold surface. Furthermore, the force required to rupture decreases for the layers farther away from the electrode surface to the bulk, indicating a weaker ordering of the layered structure away from the electrode surface.

Upon changing the electrode potentials to either $+0.2\ \text{V}$ or to $-0.2\ \text{V}$, a similar multi-layered structure with a variation in the separation of the layers can be detected (Fig. 2). For the neat IL, the force required to rupture the layer adjacent to the electrode was found to vary with the applied electrode potential. The rupturing forces were found to be $10\ \text{nN}$, $6\ \text{nN}$, and $\sim 3\ \text{nN}$ for the neat IL at $-0.2\ \text{V}$, at OCP, and at $+0.2\ \text{V}$, respectively, revealing that a relatively strong near surface structure is present at $-0.2\ \text{V}$. This indicates an increase in the strength of adsorption at more negative electrode potentials. For the neat IL at $-0.2\ \text{V}$, the widths of the innermost and successive layers were found to be 0.5 and $\sim 0.9\ \text{nm}$, revealing a strong near surface interaction of the ions due to the applied electrode potential. This observation is in agreement with literature data.

For $50\ \mu\text{M}$ $\text{AgTfSA}/[\text{Py}_{1,4}]\text{TfSA}$, a few changes have been recorded in the force-distance curves at all measured electrode potentials. At OCP, the width of the innermost layer is $\sim 0.8\ \text{nm}$, which can either be due to a rearrangement of the ions or to a change in the composition of the electrolyte (or species like AgTfSA complexes are present). The widths of the subsequent layers have been slightly increased for $50\ \mu\text{M}$ AgTfSA in comparison to the neat IL, with an average width of $1\ \text{nm}$ at all electrode potentials. However, the innermost



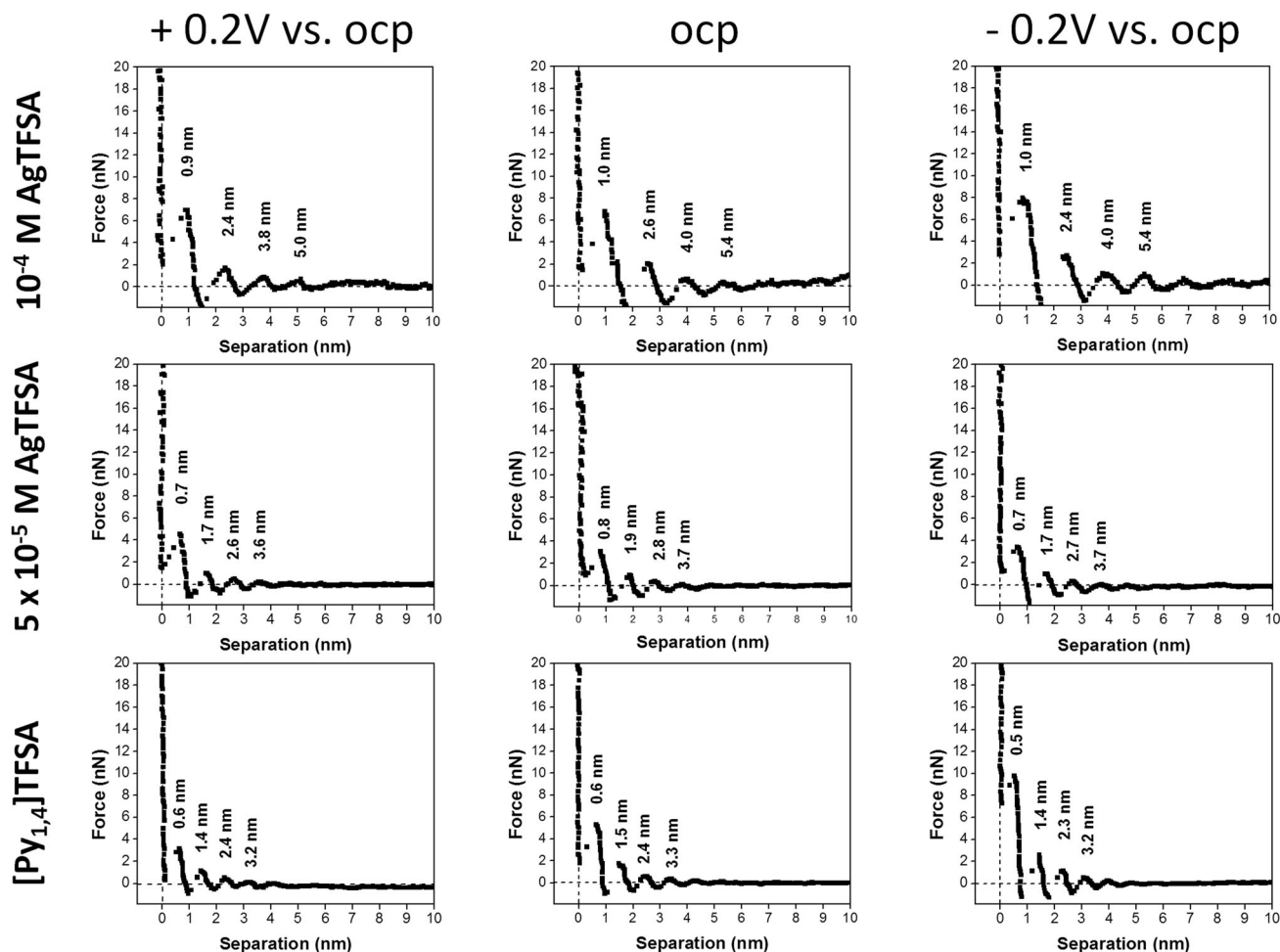


Fig. 2 A comparison of the typical force–separation curves of neat $[\text{Py}_{1,4}]\text{TFSA}$ and of two different concentrations of AgTFSA in $[\text{Py}_{1,4}]\text{TFSA}$. The curves have been obtained at +0.2 V vs. OCP (left column), at OCP (middle column), and at –0.2 V vs. OCP (right column). $[\text{AgTFSA}] = 50 \mu\text{M}$ and $100 \mu\text{M}$.

layer widths were found to be 0.7 nm at –0.2 V and at +0.2 V. The increase in the width of the layers for $50 \mu\text{M}$ AgTFSA can be attributed either to the presence of AgTFSA complexes or to their changes in orientation (or conformers) of the species. Furthermore, there is a marginal change in the rupturing force in the force–distance curves. The rupturing forces were found to be 3 nN for all electrode potentials, indicating a rather weak near surface interaction of the ions with the electrode.

For $100 \mu\text{M}$ $\text{AgTFSA}/[\text{Py}_{1,4}]\text{TFSA}$, a further increase in the widths of the layers has been observed at all measured electrode potentials. Typically, the width of the innermost layer was found to be 0.9 ± 0.1 nm at +0.2 V, OCP and at –0.2 V. The average widths of the successive layers were found to be 1.4 nm at OCP and at –0.2 V. The force required to rupture the layers increased to 8 nN at all the measured potentials. However, a clear change was noticed at the $200 \mu\text{M}$ concentration wherein the number of layers varied upon varying the electrode potential. Four clear and distinct layers with step widths of about 1.0, 1.4, 1.5, and 1.4 nm have been observed at OCP and 0.9, 1.2, 1.4, 1.4 nm have been noticed at –0.2 V. However, only two steps with widths of about 0.9 and 1.5 nm have been recorded at +0.2 V. The force required to rupture the innermost layers is found to be large at

OCP (10 nN) and at –0.2 V (14 nN) compared to the one at +0.2 V, which indicates a difference in the strength of the adsorption of the innermost layer with the substrate. Furthermore, the number of steps has been found to be the same at –0.2 V while a decrease in the step width has been noticed (*ca.* 0.9 nm). Quite an interesting result has been obtained at and above the $500 \mu\text{M}$ AgTFSA concentration, wherein only a single step was noticed at all measured electrode potentials. Furthermore, a change in the widths of the layers by varying the electrode potential has been found for $500 \mu\text{M}$ AgTFSA (Fig. 3). Here in this case, one clear step with a layer width of 1 nm has been observed at OCP, whereas the step widths were found to decrease upon varying the electrode potential (*e.g.* 0.9 nm at +0.2 V and 0.8 nm at –0.2 V). Moreover, the forces required to rupture these layers have been nearly the same (*ca.* 2–3 nN). However, no clear and discrete steps have been observed at the OCP and at 0.2 V for 1 mM AgTFSA in the force–distance profiles, rather gradients have been noticed. These gradient widths were found to be ~ 0.3 and ~ 0.2 nm at the OCP and at +0.2 V, respectively. The forces needed to break the layers were around 6–7 nN in both cases. Furthermore, no clear signal was seen at –0.2 V, which indicates that the applied electrode



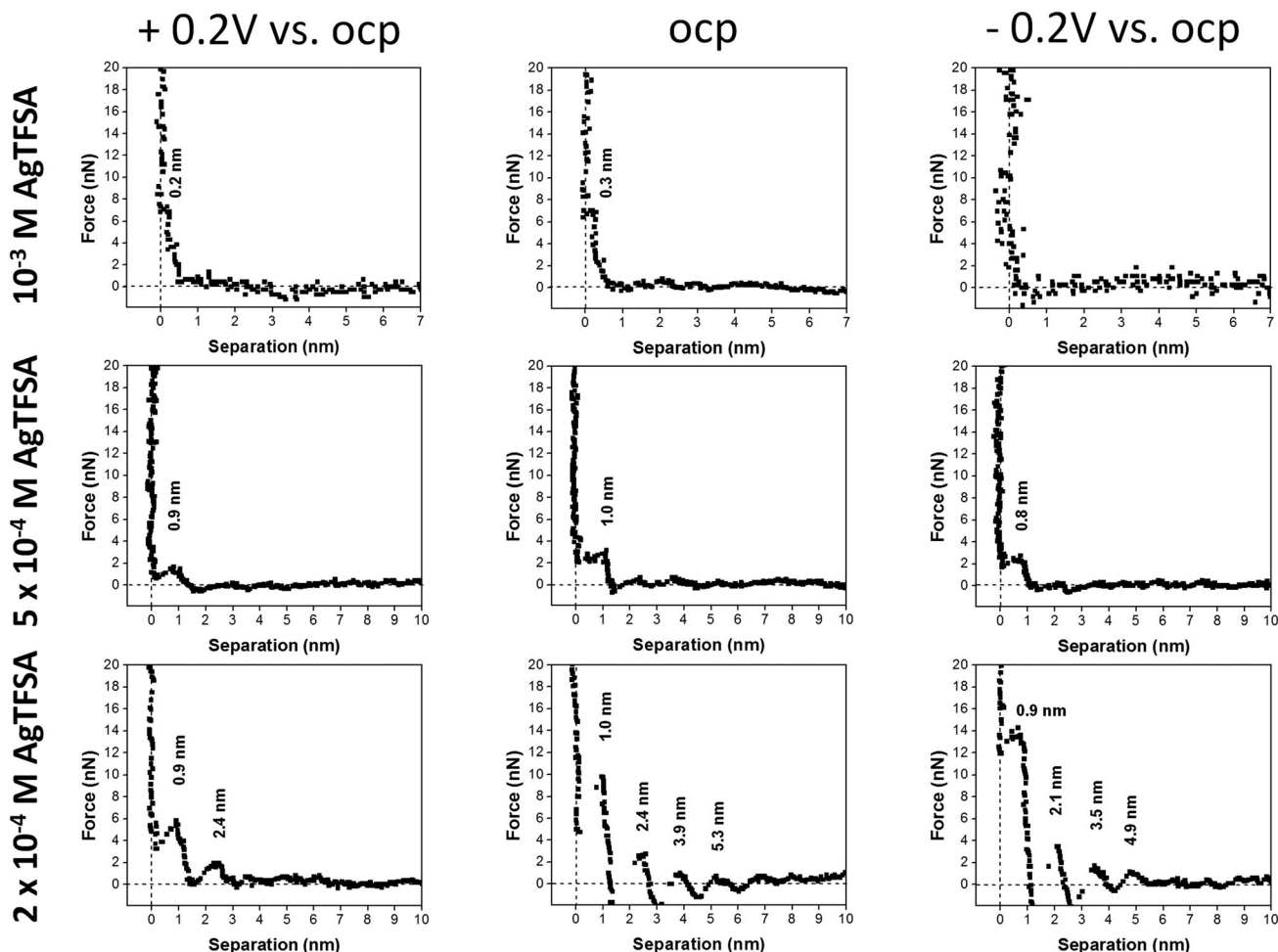


Fig. 3 A comparison of the typical force–separation curves of three different concentrations of AgTFSA in $[\text{Py}_{1,4}]\text{TFSA}$. The curves have been obtained at +0.2 V vs. OCP (left column), at OCP (middle column), and at –0.2 V vs. OCP (right column). $[\text{AgTFSA}] = 200 \mu\text{M}$, $500 \mu\text{M}$, and 1 mM .

potential affects the interfacial structure. The analysis of force–separation profiles at various concentrations of AgTFSA in $[\text{Py}_{1,4}]\text{TFSA}$ has shown that the number of solvation layers decreases from four layers to a single layer with an increase in concentration from 50 to $200 \mu\text{M}$ AgTFSA in $[\text{Py}_{1,4}]\text{TFSA}$. A significant change in the widths of the innermost layers can also be observed at 1 mM AgTFSA.

In order to interpret the present findings of the $\text{Au}(111)/[\text{Py}_{1,4}]\text{TFSA}$ interfacial nanostructure at various concentrations of AgTFSA, we compare the results with the reported ones in the literature at varying concentrations of alkali metal salts in ILs^{41,59} (e.g. $\text{LiTFSA}/[\text{Py}_{1,4}]\text{TFSA}$ ⁴¹ and $\text{NaFSA}/[\text{Py}_{1,4}]\text{TFSA}$ ⁵⁹). At OCP, five solvation layers were reported for 0.1 M $\text{LiTFSA}/[\text{Py}_{1,4}]\text{TFSA}$ on $\text{Au}(111)$. The innermost layer width was found to be 0.16 nm for 0.1 M $\text{LiTFSA}/[\text{Py}_{1,4}]\text{TFSA}$ while the innermost layer width was found to be 0.65 nm for $[\text{Py}_{1,4}]\text{TFSA}$. In the case of the 0.1 M $\text{LiTFSA}/[\text{Py}_{1,4}]\text{TFSA}$, the widths of the successive layers were found to be same ($\sim 0.7 \text{ nm}$) and these steps could be related to the presence of the ion-pairs of $[\text{Py}_{1,4}]\text{TFSA}$. A similar behaviour was also seen for $[\text{Py}_{1,4}]\text{TFSA}$, where a multi-layered structure can be observed. The widths of the innermost layers rose by increasing the concentration from 0.1 M to

$\geq 0.5 \text{ M}$ $\text{LiTFSA}/[\text{Py}_{1,4}]\text{TFSA}$. A decrease in the number of layers can also be observed with an increase in the concentration of LiTFSA . Moreover, the widths of the successive layers (e.g. 2nd, 3rd and 4th layers for 0.5 M LiTFSA in $[\text{Py}_{1,4}]\text{TFSA}$ and 2nd and 3rd layers for 1 M $\text{LiTFSA}/[\text{Py}_{1,4}]\text{TFSA}$) were found to be between ~ 0.85 and $\sim 0.95 \text{ nm}$. This indicates that the presence of $[\text{Py}_{1,4}]\text{TFSA}$ ion-pairs in the successive layers was not influenced by the presence of Li^+ at the interface for varying concentrations of Li^+ between 0.1 M and 0.5 M . However, the interfacial nanostructure was disturbed at 1 M $\text{LiTFSA}/[\text{Py}_{1,4}]\text{TFSA}$ upon varying the electrode potential, which was attributed to the formation of upd layers on the electrode. This study clearly showed that the addition of Li^+ can significantly change the interfacial structure where $[\text{Py}_{1,4}]^+$ ions are replaced by Li^+ , which was observed as a decrease in the innermost layer widths on $\text{Au}(111)$. Furthermore, the widths of the innermost layers were increased on increasing the concentration of LiTFSA , which was attributed to the change in configuration between Li^+ and TFSA^- .⁴¹

For NaFSA in $[\text{Py}_{1,4}]\text{FSA}$, the interfacial nanostructure did not vary upon the addition of 0.05 M NaFSA to $[\text{Py}_{1,4}]\text{FSA}$ ⁵⁹ in which the innermost layer width was found to be nearly the



same (~ 0.65 nm) as that of the innermost layer width for $[\text{Py}_{1,4}]\text{TFSA}$. However, the interfacial nanostructure was changed upon increasing the concentration of NaFSA from 0.05 to ≥ 0.25 M in $[\text{Py}_{1,4}]\text{FSA}$, wherein the width of the innermost layer was found to be ~ 0.35 nm for 0.25 M and 0.5 M NaFSA in $[\text{Py}_{1,4}]\text{FSA}$. In the case of NaFSA/ $[\text{Py}_{1,4}]\text{FSA}$ at OCP, the widths of the innermost layers were found to be invariant upon increasing the concentration of NaFSA. Based on these two studies we can conclude that the structure of the Au(111)/IL interface at a higher concentrations of metal salts (e.g. 1 M LiTFSA $[\text{Py}_{1,4}]\text{TFSA}$ or 0.5 M NaFSA/ $[\text{Py}_{1,4}]\text{FSA}$) differed with that of the interfacial structure at lower concentrations (e.g. 0.1 M LiTFSA in $[\text{Py}_{1,4}]\text{TFSA}$ or 0.05 M NaFSA in $[\text{Py}_{1,4}]\text{FSA}$). However, in contrast to the present studies with AgTFSA rather high concentrations of LiTFSA and NaFSA were needed to disturb the interfacial layers.

The following salient features can be drawn from the above discussion. An increase in the width of the innermost layer has been noticed upon increasing the concentration of LiTFSA from 0.1 to 0.5 M in $[\text{Py}_{1,4}]\text{TFSA}$ (0.16 nm for 0.1 M vs. 0.25 nm for 0.5 M), which was attributed to changes in the complexation (i.e. a change in the coordination number, C.N.) and to configurational changes (i.e. *cis/trans* and/or monodentate/bidentate).⁴¹ For the nanostructure of the Au(111)/NaFSA- $[\text{Py}_{1,4}]\text{FSA}$ interface, the authors have observed no significant increase in the push-through forces for 0.05 M NaFSA in $[\text{Py}_{1,4}]\text{FSA}$ at various electrode potentials. These push-through forces were almost comparable with those of $[\text{Py}_{1,4}]\text{FSA}$ under similar conditions. A significant change in the nanostructure of Au(111)/NaFSA- $[\text{Py}_{1,4}]\text{FSA}$ has been observed for 0.25 M NaFSA in $[\text{Py}_{1,4}]\text{FSA}$, where the innermost layer width significantly decreased (0.37 nm for 0.25 M FSA/ $[\text{Py}_{1,4}]\text{FSA}$ vs. 0.68 nm for 0.05 M

NaFSA/ $[\text{Py}_{1,4}]\text{FSA}$). A stable interfacial structure has been found upon further increasing the concentration from 0.25 M to 0.5 M NaFSA.⁵⁹ Here, in the aforementioned two cases, the concentrations of metal salts (LiTFSA/NaFSA) have been varied in the regime of mM to M. However, in the case of AgTFSA reported in the present paper, the interfacial structure was studied at quite low concentrations of the salt from 50 μM to 1 mM of AgTFSA. Moreover, we observe a similar solvation structure for salt concentrations between 50 and 200 μM AgTFSA. The solvation layers formed due to the presence of AgTFSA complexes from 50 to 200 μM should be the same. Therefore, the widths of the innermost layers do not differ largely in this regime of concentration. We indeed observed a similar solvation structure at the interface of the Au(111)/electrolyte in this regime of AgTFSA concentrations, where the innermost layer widths were between 0.8 and 1 nm. However, not only the orientation (*cis/trans* configurations) of the solvated species on the electrode surface change but also the solvation layers change upon further increasing the concentration to 500 μM and above. The above *in situ* AFM results reveal that the Au(111)/AgTFSA/ $[\text{Py}_{1,4}]\text{TFSA}$ interfacial structure is rather complex and further thorough investigations are needed to comprehensively describe the near surface ionic arrangement at the electrode/electrolyte interface.

In order to investigate the dynamics of the interfacial processes that occur on Au(111), we have recorded the force–distance curves of 50 μM AgTFSA/ $[\text{Py}_{1,4}]\text{TFSA}$ by varying the electrode potentials to more negative values like -0.8 V vs. OCP and then back to a more positive value, e.g. -0.3 V. The corresponding force–distance curves are shown in Fig. 4. Four distinct steps have been observed in the force–distance profiles upon varying the electrode potentials from OCP to -0.8 V. The average widths of the innermost and the

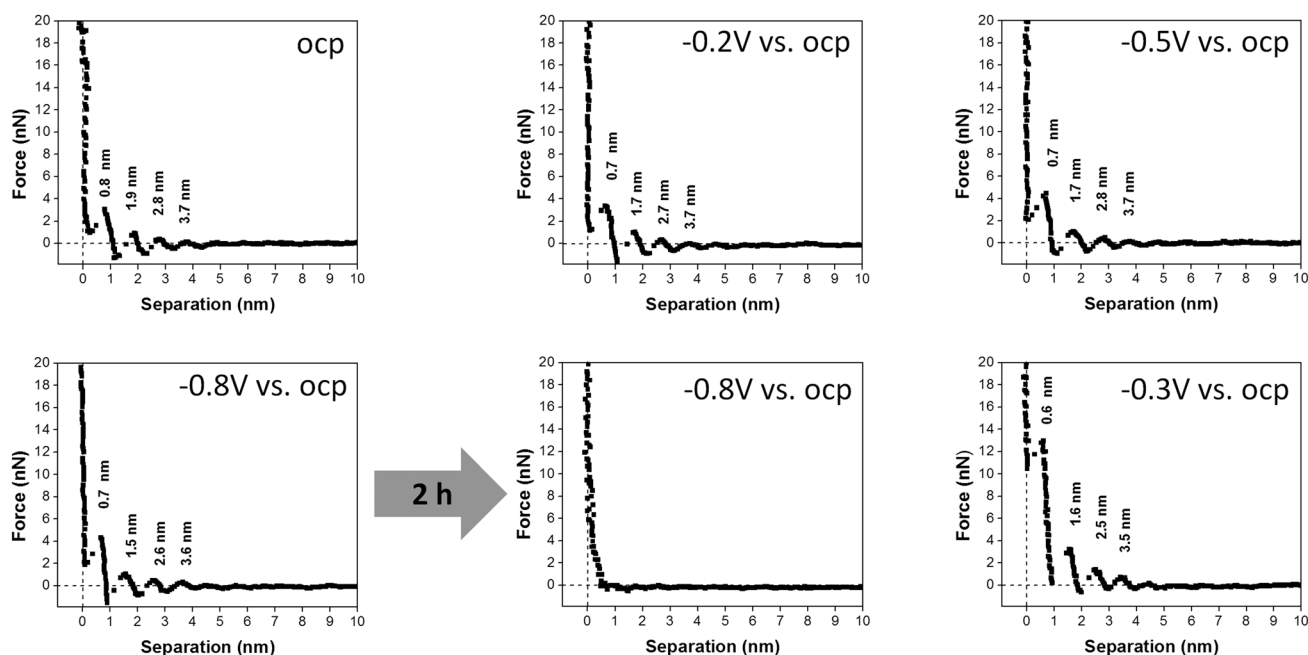


Fig. 4 A comparison of typical force–separation profiles of 50 μM AgTFSA in $[\text{Py}_{1,4}]\text{TFSA}$. The curves have been obtained at various electrode potentials from OCP to -0.8 V, the potential was held at -1 V for 2 h, and switched back to -0.3 V.



next successive layers were found to be 0.7 and 1 nm, respectively. However, by applying -0.8 V for two hours under the conditions of Ag deposition the force–distance profiles show no steps indicating that the interfacial structure is totally disturbed during the deposition of silver on gold. The deposition process consumes silver ions that are present at the interface resulting in a depletion of silver complexes at the interface. A multilayered structure is resumed after changing the potential back to -0.3 V where the deposited silver is slowly oxidized. The reappearance of this multilayer is typical for quite a low concentration of AgTFSA/[Py_{1,4}]TFSA and for the neat IL, which is due to the presence of either AgTFSA complexes or to the ion-pair of the IL at the interface.

In summary, the force–distance profiles obtained for [Py_{1,4}]TFSA and for AgTFSA/[Py_{1,4}]TFSA are different. The results obtained in this investigation indicate that the multilayered structure is retained in IL solutions with quite low concentrations of the silver salt, and slight changes have been observed for 50 and 100 μ M AgTFSA at all measured electrode potentials. The innermost layer thickness for the neat IL varies upon changing the electrode potentials (*e.g.* ± 0.2 V vs. OCP). There is a slight increase in the width of the innermost layer for AgTFSA/[Py_{1,4}]TFSA at the measured potentials, indicating an increase in the dimensions of the ions, which could be due to the presence of the ion-pairs of the IL in a different orientation along with the Ag⁺ complexes in the interface. Furthermore, the presence of quite low concentrations of AgTFSA complexes may not influence the interfacial structure strongly and the innermost layer is still dominated by IL ion pairs. A transition in the interfacial structure has been observed at 200 μ M AgTFSA in [Py_{1,4}]TFSA, wherein the step width is increased with a further increase in the average widths of subsequent layers (*ca.* 1.5 nm). The increase in width of the innermost layer and the layers extending to the bulk (*i.e.* either in the transition zone or in the diffusion zone) could be due to the presence of AgTFSA complexes as these species are large in their sizes compared to

the ion-pair sizes of [Py_{1,4}]TFSA. A comparably large force is required to rupture the innermost layer at this concentration compared to the one for the neat IL and for lower concentrations of AgTFSA/[Py_{1,4}]TFSA, which is due to an increased strength of adsorption of the AgTFSA complexes on the electrode surface.⁴¹ Furthermore, the number of steps has been found to decrease to 2 at $+0.2$ V. The decrease in the number of layers indicates a change in the ordering of the layers (layers in the diffusion zone/transition zone) away from the electrode surface, which is also dependent on the electrode potential due to changes in the electrostatic attraction.^{60,61}

Rather a simple double layer structure is observed at 500 μ M AgTFSA, where only a single step is observed in the force–distance profiles at all electrode potentials. Furthermore, the tip experiences quite a low rupturing force of 2 nN indicating that the adsorbed layer can be ruptured easily with minimum force. Such behaviour indicates that the adsorbed layer weakly interacts with the surface. An increase in the concentration of AgTFSA to 1 mM changes the interfacial structure completely as we could not observe any steps in the force–distance measurements. This indicates that the cantilever cannot find any adsorbed layers until it reaches the electrode surface. The AFM results signify that the concentration has a considerable influence on the interfacial structure for Au(111)/AgTFSA in [Py_{1,4}]TFSA. This behaviour is rather complex and it is obviously different for aqueous electrolytes, where a classical EDL structure can be observed at very dilute solutions of aqueous electrolytes.

3.2 Simulation results

For comparison with experiments, we simulate the static equilibrium EDL at 200 mV overpotential. First, we simulate the behavior of the binary ionic liquid [Py_{1,4}]TFSA at a metallic surface. Second, we discuss the effect of neutral or charged additives. The details of our computational method can be found in the ESI.†

In Fig. 5, we neglect the hardcore constraint to reach a better understanding of the model ($F^{\text{int}} = 0$). The simulations show the

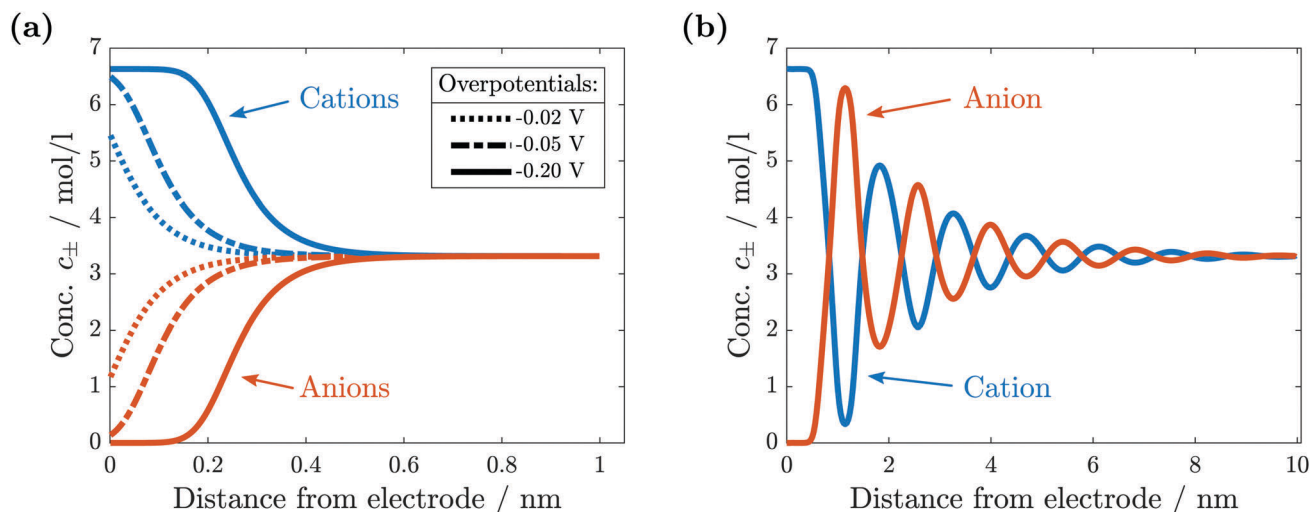


Fig. 5 Comparison of ionic concentrations in the vicinity of a negatively charged electrode surface. (a) Without hardcore interactions we observe the depletions of anions and the crowding of cations. (b) With hardcore interactions a quasi-crystalline structure of alternating layers forms.



crowding of cations at the negative electrode (see Fig. 5(a)). The shape of the ionic layers depends on the applied potential difference. In particular, at small electrode potentials (-0.02 V in Fig. 5(a)), the concentration of cations exhibits a distribution of the Boltzmann-type. The cation concentration at the surface increases with electrode potential until, at sufficiently large polarization (-0.05 V in Fig. 5(a)), it becomes saturated due to incompressibility because limited space is available for the ions. This implies that the concentration of cations at the electrode surface is bound from above. Thereby, the concentration profile changes to the Fermi-Dirac-type. During further increase of electrode polarization (up to -0.2 V in Fig. 5(a)), the cation layer broadens to screen the electric field of the electrode. The decay length of the ionic layers corresponds to the Debye-length expected for a dilute electrolyte, $\lambda_D = \sqrt{2\epsilon_R \epsilon_0 RT \nu / F^2} \approx 0.14$ nm.

In conclusion, without the hardcore constraint we observe the standard behavior for electrolytes at electrode interfaces. The electric potential jump is screened by the accumulation of ionic charges.

The simulation results change significantly if the hardcore constraint is taken into account (see Fig. 5(b)). We observe the formation of various alternating layers extending over 5 nm. The first layer adjacent to the negatively charged electrode consists entirely of cations, whereas the anions are depleted completely. The successive layers alternate in the dominating species and decay in amplitude, *i.e.*, the total depletion of the minor species does not appear any more. Finally, the concentrations of the two species converge in the bulk electrolyte. The electrolyte potential is in accordance with the damped oscillations of the ionic concentrations (see Fig. 6). These oscillations show that the ionic diameter is incommensurate with the charge required to screen the electric potential jump.

We calculate the mechanical force as outlined in Section 2.2. The simulated force (see Fig. 7) exhibits the characteristic peaks observed in the AFM experiments. The spacing between the positive peaks agrees with the molecular radius of 0.67 nm (see the ESI†) within numerical accuracy. As shown in Fig. 5(b),

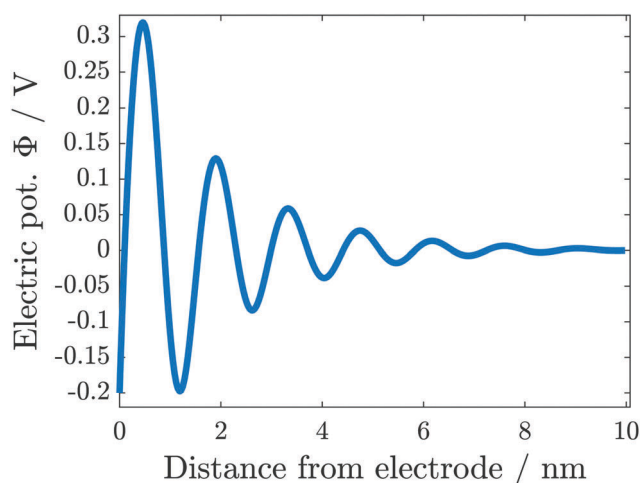


Fig. 6 Electric potential in the electrolyte in the vicinity of a negatively charged electrode surface.

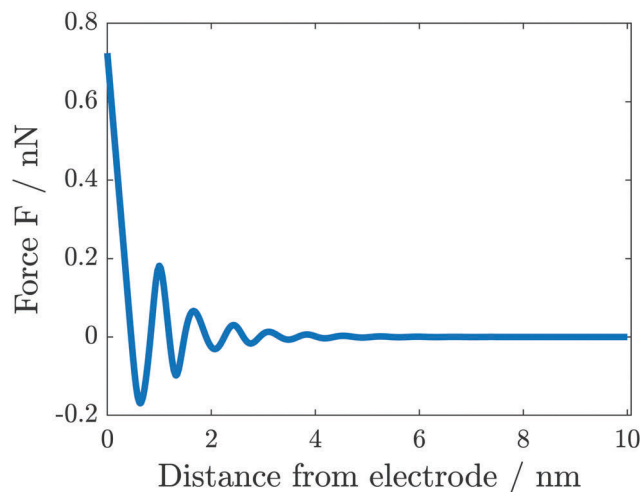


Fig. 7 Simulated force on AFM tip in the vicinity of a negatively charged electrode surface.

the innermost layers are dominated by a single species, in alternating sequence. As these layer widths are qualitatively in good agreement with the experimentally measured layer widths for the pure IL, our model supports the interpretation of AFM measurements to correlate the force peaks with molecular mono-layers. The deviations between the modeled and measured layer widths are not surprising, as our modeling approach neglects microscopic properties, *e.g.*, non-spherical molecular shapes and orientation-dependent interactions.

The deviation between the simulated and experimental forces lies between one to two orders of magnitude. This agreement is very good if we take into account the differing set-ups. The atomic force microscope measures the penetration into a many-particle crystalline structure, including the replacement of particles. In contrast, we simulate elastic perturbations of an equilibrium state.

We discuss the impact of neutral and charged additives, *e.g.*, water H_2O and ionic silver Ag^+ onto the EDL. The dynamics of this third species in the electrolyte is determined by eqn (11). The small concentration of the additive species allows us to neglect the corresponding contributions to convection and to the electric current, setting $M_3 \approx 0$, $\nu_3 \ll \nu_{\pm}$, and $t_3 \ll 1$ in this principal discussion of the EDL structure. In the stationary case ($\mathbf{J} = 0$ and $\mathbf{N}_3 = 0$) we find from eqn (10) and (11)

$$\nabla \mu_3 = -\frac{Fz_3}{RT} \nabla \Phi. \quad (14)$$

This equation determines the chemical potential and the concentration of the additive from the electric potential. In particular, using eqn (12), we get

$$\nabla \left(\ln \frac{c_3}{c_0} \right) = -\frac{Fz_3}{RT} \nabla \Phi. \quad (15)$$

Note that this relation connects the entropic and electric contributions of the free energy functional.

For the neutral additive water, our theoretical analysis in eqn (15) predicts a uniform distribution throughout the EDL, $c_3 = c_{H_2O} = \text{const}$. This is a consequence of the entropy forcing



neutral additives into the EDL. Therefore, large bulk concentrations of neutral additives are required to reach significant concentrations in the EDL and to affect its multilayer structure. Indeed, previous experiments⁴³ show that the addition of water perturbs the EDL and destroys the multilayer structure only if it makes up more than half of the solution volume.

Electric forces are to be discussed for charged additives. If we consider silver ($z_{\text{Ag}} = 1$) as a third species, eqn (15) shows that its concentration obeys a Boltzmann distribution,

$$c_{\text{Ag}}(x) = c_0 \cdot \exp\left(-\frac{F\Delta\Phi(x)}{RT}\right), \quad (16)$$

where $\Delta\Phi(x)$ is the electric potential difference between the bulk electrolyte and point x in the electrolyte. If the electrode is polarized negatively, the Boltzmann distribution describes an exponential growth of silver concentration towards the surface. Even if the electrode is polarized positively, the electric potential in the electrolyte reaches large negative values due to its oscillations shown in Fig. 6. In this case, silver accumulates behind the layer of anions adsorbed at the electrode surface. In the case of the potential difference $\Delta\Phi = -0.2$ V, the silver concentration at the electrode increases to $c_{\text{Ag}} = 2.4 \times 10^3 c_{\text{Ag},0}$. We expect that the EDL structure is affected by silver ions once their concentration at the electrode becomes comparable with the bulk concentration of the ionic liquid $c_{\text{Ag}} \sim 1$ M. This corresponds to the bulk silver concentration $c_{\text{Ag},0} = 0.4$ mM. Indeed, as shown experimentally in Fig. 3, the EDL structure is changing at these concentrations.

The screening of charges and electric fields cause the electrochemical surface layer. The quasi-crystalline structure emerges because the screening length does not fit the molecular size of the ionic liquids. Metal ions in the surface layer can efficiently shield electric forces and introduce an additional very small molecular length scale. Therefore, we expect that the silver ions very effectively destroy the EDL multi-layer structure.

This argument, however, is complicated by the presence of chemical interactions between the molecules. The chemical interactions could, for example, lead to strong pairing forces. In this case metal ions would mainly be present as a neutral complex and might behave as neutral solutes. Alternatively, neutral solutes could have preferred and strongly polarized bonds with solvent molecules that destroy the quasi-crystalline structure at low concentrations. It remains for further studies to evaluate and classify the effect of the chemical interactions on the EDL.

4 Conclusions

The electrode/electrolyte interface of Au(111)/[Py_{1,4}]TFSA and of Au(111)/AgTFSA/[Py_{1,4}]TFSA has been investigated using atomic force microscopy. The interfacial structure (or the arrangement of ions) remarkably varies with the concentration of AgTFSA until the layer structure is completely disrupted. Furthermore, changing the electrode potential has some influence on the interfacial structure. The force-separation profiles indicate that a multi-layer structure is found for the neat IL and this multi-layer structure is retained at quite low concentrations of the silver salt (50 to 100 μM AgTFSA/[Py_{1,4}]TFSA). However, a

clear change in the electrical double layer structure has been noticed at a 200 μM concentration of AgTFSA in [Py_{1,4}]TFSA wherein the interfacial layers were found to be different on changing the applied electrode potentials. Furthermore, the widths of the corresponding innermost layers at 200 μM concentration are different from those ones observed at lower concentrations of AgTFSA. A simple double layer structure has been observed for 500 μM AgTFSA, which does not vary with the applied electrode potential. A remarkable difference is observed at this concentration and above, in which no layer structure has been noticed at 1 mM concentration. The observed changes in the widths of the innermost layer from AFM measurements could be attributed to the orientation of the Ag species (or to the conformational changes of the anions that are coordinated to the Ag⁺).

Furthermore, we present a model analysis supporting these measurements of the EDL structure. We simulate force profiles that show the multilayer structure and agree with the experimental observations. The distances between force peaks correspond to ionic monolayers. We can show that in contrast to neutral additives ionic additives like AgTFSA are attracted by an electrified interface. Therefore, we predict in good agreement with experiments that roughly 400 μM concentrations are sufficient to reach large concentrations in the EDL and disrupt its structure. By neglecting the microscopic details of chemistry, our model highlights the fundamental relations behind the multilayer EDL structure of ionic liquids.

The AFM results on the dynamics of the interfacial processes reveal that the multilayer structure is disrupted completely during the deposition of silver upon applying -0.8 V for 2 h and the multilayer structure is retained after changing the electrode potential back to -0.3 V. This study signifies that the EDL structure of Au(111)/IL is indeed influenced by the presence of varying concentrations of an added salt to the IL. The general question arises to what extent different solutes alter IL interfacial layers and whether the dynamic processes alter the deposit morphology.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

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