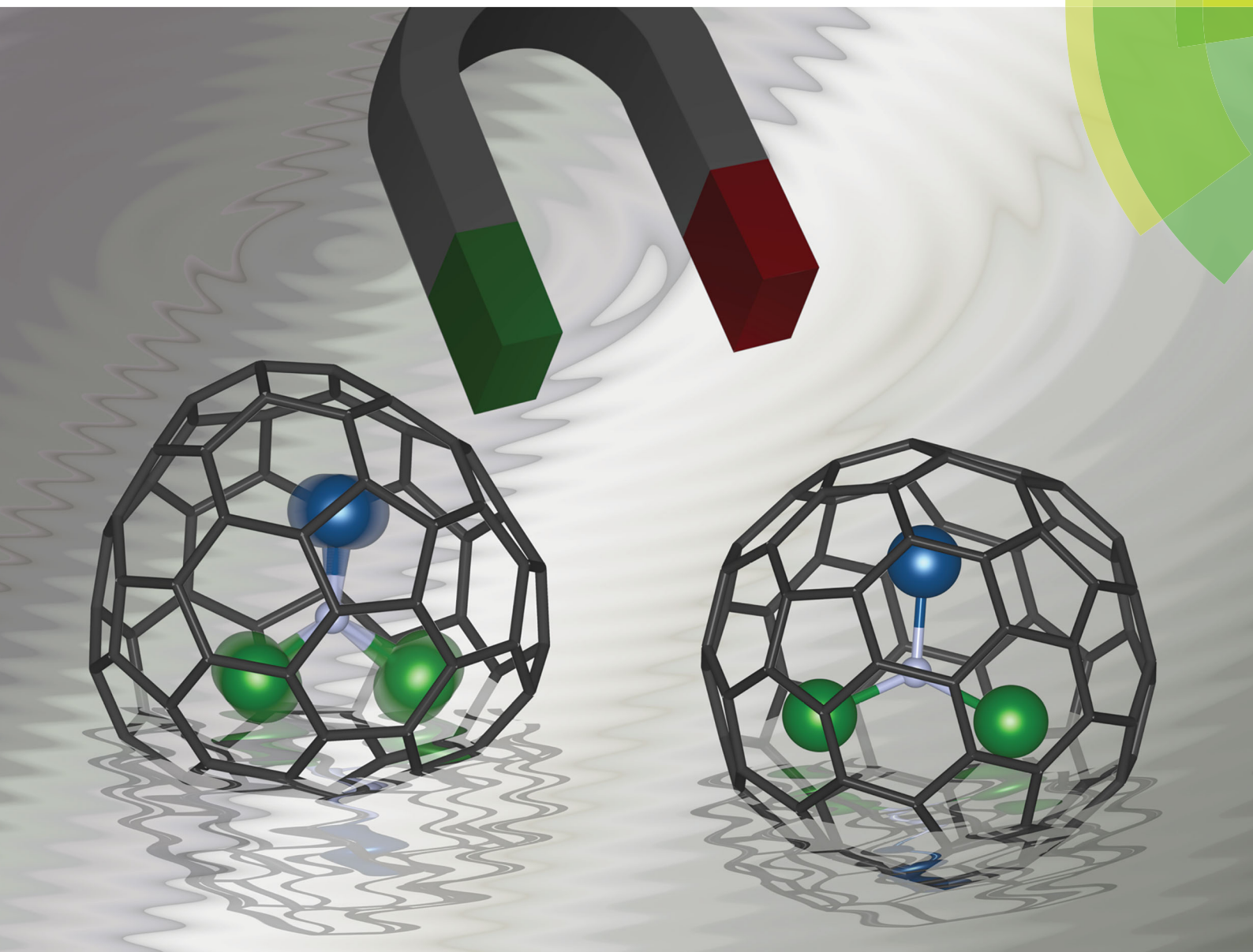


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












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Strong carbon cage influence on the single molecule magnetism in Dy–Sc nitride clusterfullerenes†

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Magnetic properties of endohedral metallofullerenes with nitride clusters DySc₂N and Dy₂ScN and different carbon cages are studied by SQUID magnetometry. All molecules behave as single molecule magnets (SMMs) and exhibit magnetic hysteresis. It is found that the blocking temperature of magnetization and relaxation times strongly depend on the fullerene cage, with the C₈₀-I_h isomer offering the best SMM properties.

The ability of fullerenes to stabilize unconventional species inside the protective shell of their carbon cages resulted in a plethora of different classes of endohedral metallofullerenes (EMFs).¹ The endohedral metal ions can preserve some of their physical properties (such as spin states) inside the carbon cage, whereas intramolecular interactions with other ions in the cluster as well as the electron transfer to the carbon cage can lead to the development of new properties. In 2012 it was found that DySc₂N@C₈₀-I_h exhibits slow relaxation of magnetization in a finite magnetic field of 0.2 T.² In zero field, the magnetization of DySc₂N@C₈₀-I_h relaxes fast *via* quantum tunnelling of magnetization (QTM), resulting in a characteristic butterfly shape of the magnetic hysteresis typical for many single-ion magnets (SIMs).³ Further studies revealed single molecule magnetism (SMM) in other lanthanide-based EMFs, including clusterfullerenes with nitride,⁴ sulfide,⁵ carbide,^{5,6} or cyanide units,⁷ as well as in dimetallofullerenes.⁸ The SMM properties were found to depend substantially on the choice of the metal (Dy gives the longest relaxation times) and the endohedral cluster composition. For instance, QTM at zero field is suppressed in

Dy₂ScN@C₈₀-I_h due to the ferromagnetic exchange and dipolar interactions of the two Dy ions.^{4a}

Among the various lanthanide clusterfullerenes, DySc₂N and Dy₂ScN clusterfullerenes showed the best SMM performance presumably due to the strong ligand field created by the nitride ion located at the short distance from the lanthanide.^{4b,5,9} Studies of the magnetic properties of nitride clusterfullerenes have been limited so far to the C₈₀-I_h cage because the corresponding EMFs are usually produced in much higher yields and constitute *ca.* 80% of the nitride clusterfullerenes mixtures. Yet, the arc-discharge synthesis of nitride clusterfullerenes affords a variety of other isolable carbon cages, from C₆₈ to C₈₈, as well as another cage isomer of C₈₀, with D_{5h} symmetry (the Gd–Sc and Dy–Sc systems have been studied in most detail).¹⁰ Due to the much lower production yield their properties remain largely unexplored. It is not clearly understood if the fullerene cage is merely a container for magnetic species, or if it plays an active role in the relaxation of magnetization. For Dy₂S- and TbCN-clusterfullerenes the variation of the carbon cage was found to affect the magnetic properties.^{5,7b} In this Communication, we show that the magnetic properties of several Dy–Sc nitride clusterfullerenes strongly depend on their fullerene cages and that the most symmetric C₈₀-I_h fullerene provides the best SMM properties of the endohedral clusters.

The structures chosen for this study (Fig. 1), DySc₂N@C₆₈-D₃(6140), Dy₂ScN@C₈₄-C_s(51365), DySc₂N@C₈₀-D_{5h}(6), and Dy₂ScN@C₈₀-D_{5h}(6),[†] have been obtained using arc-discharge synthesis followed by HPLC separation as described in detail earlier.^{10b,d} This set of EMFs allows to compare the influence of the carbon cage isomerism (I_h *versus* D_{5h}) as well as the size of the fullerene cage (C₆₈ *versus* C₈₀ *versus* C₈₄).

In DySc₂N@C₆₈, the D₃-symmetric fullerene cage violates the isolated pentagon rule, as it features three adjacent pentagon pairs coordinated by metal atoms of the endohedral nitride cluster. The cluster is therefore tightly fixed inside the carbon cage. Dy₂ScN@C₈₄ with the C_s(51365) fullerene cage has one pentagon adjacency and thus also violates the IPR. Computational studies

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† Electronic supplementary information (ESI) available: Mass-spectra and determination of relaxation times. See DOI: 10.1039/c8cc05029e



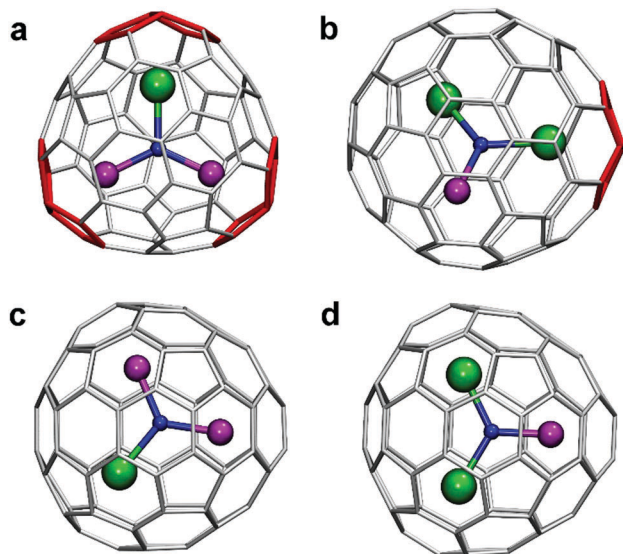


Fig. 1 Molecular structures of Dy–Sc nitride clusterfullerenes studies in this work: (a) $\text{DySc}_2\text{N}@C_{68}\text{-}D_3(6140)$, the C_3 axis of the cage is perpendicular to the paper plane; (b) $\text{Dy}_2\text{ScN}@C_{84}\text{-}C_5(51365)$, the symmetry plane of the cage is parallel to the paper plane; (c) $\text{DySc}_2\text{N}@C_{80}\text{-}D_{5h}(6)$; (d) $\text{Dy}_2\text{ScN}@C_{80}\text{-}D_{5h}(6)$. Dy atoms are green, Sc atoms are magenta, N is blue, carbon cages are grey except for adjacent pentagon pairs shown in red. In (c) and (d), the plane of the cluster is parallel to one of the σ_h planes of the cage, whereas the C_5 axis is oriented vertically and is in the plane of the paper. All structures are based on DFT calculations of Y–Sc analogues.

showed that the structure in which the pentagon pair is coordinated by the bigger metal is more stable. However, the energy of the conformer with a Sc-coordinated pentagon pair is only slightly higher, and thus both structures are likely to be present. The pentagon adjacency fixes the position of the metal cluster, which therefore cannot rotate freely in $\text{DySc}_2\text{N}@C_{68}$ and $\text{Dy}_2\text{ScN}@C_{84}$.

The $C_{80}\text{-}D_{5h}$ cage resembles the $C_{80}\text{-}I_h$ counterpart. Both isomers can be described as built from identical hemispheres, but rotated with respect to each other at different angles. The NMR studies of Lu–Sc nitride clusterfullerenes showed that the nitride cluster inside the $C_{80}\text{-}D_{5h}$ cage rotates freely at room temperature, similar to the rotation inside the $C_{80}\text{-}I_h$ cage.¹¹ Thus, disordered positions of the endohedral cluster are to be expected in $\text{DySc}_2\text{N}@C_{80}\text{-}D_{5h}$ and $\text{Dy}_2\text{ScN}@C_{80}\text{-}D_{5h}$, and Fig. 1 shows one of the lowest energy configurations determined by DFT calculations.^{10b}

As a measure of the SMM performance we use several parameters (Table 1), which can be obtained from DC SQUID magnetometry (the low yield of these minor structures precludes the isolation of amounts sufficient for AC magnetometry). First of all, magnetic hysteresis is a clear indication that the sample behaves as an SMM (Fig. 2). The width of the hysteresis and the temperature at which it is closing are characteristic parameters of a given SMM. Another important characteristic parameter of the SMM is the blocking temperature of magnetization, T_B . Here we adopt the definition of T_B based on the divergence of the magnetic susceptibility of the SMM measured during cooling the sample in a field (FC) and for the sample cooled in zero field (ZFC). T_B is then defined as the temperature of the

Table 1 Selected SMM parameters in Dy–Sc nitride clusterfullerenes

EMF	T_B^a , K	T_{B100}^b , K	H_c , T (at 2 K) ^c
$\text{DySc}_2\text{N}@C_{68}\text{-}D_3(6140)$	3.8	2.3	
$\text{DySc}_2\text{N}@C_{80}\text{-}D_{5h}(6)$	5.9	3.6	
$\text{DySc}_2\text{N}@C_{80}\text{-}I_h(7)$	7.0	4.6	
$\text{Dy}_2\text{ScN}@C_{80}\text{-}D_{5h}(6)$	5.3	2.6	0.48
$\text{Dy}_2\text{ScN}@C_{80}\text{-}I_h(7)$	8.0	5.0	0.70
$\text{Dy}_2\text{ScN}@C_{84}\text{-}C_5(51365)$	3.3	≈ 1.8	0.11

^a T_B measured in the field of 0.2 T with a sweep rate of 5 K min^{−1}.

^b T_{B100} is estimated from the temperature dependence of relaxation times (in zero field for Dy_2ScN -EMFs, and in the field of 0.2 T for DySc_2N -EMFs). ^c Coercive field H_c is measured at 2 K with an average sweep rate of 2.9 mT s^{−1}.

maximum of the peak observed in χ_{ZFC} (measured in a field of 0.2 T with a sweep rate of 5 K min^{−1} in this work; note that the use of a different DC field or a temperature sweep rate may give somewhat different T_B values^{3b}). In these measurement conditions, the magnetic relaxation time at T_B is *ca.* 10 seconds. At higher temperature, the relaxation becomes fast on the time scale of the experiment, and the two curves coincide. In this definition, T_B is close to the temperature at which magnetic hysteresis is closing, but the latter also depends on the sweep rate. The temperature at which the relaxation time is equal 100 s is another common parameter denoted here as T_{B100} . Finally, magnetic relaxation times and their evolution with temperature provide information on the SMM performance. Here relaxation times are measured by first magnetizing the sample to saturation, followed by sweeping the field as fast as possible to the desired value (zero field or 0.2 T). Then the decay of the magnetization is recorded, and the resulting decay curves are fitted with a stretched exponential. Due to the finite sweep rates and the time necessary for the stabilization of the field, this procedure gives reliable values for relaxation times longer than 100 s, whereas shorter relaxation times may be somewhat overestimated. A detailed discussion of the relaxation time determination is given in ref. 4d.

Fig. 2 shows the magnetization curves of the samples measured at low temperature. All samples exhibit magnetic hysteresis, showing that SMM properties are intrinsic to Dy–Sc nitride clusterfullerenes. However, the width of the magnetic hysteresis and the temperature at which it is closing varies strongly within the series. Both $\text{DySc}_2\text{N}@C_{68}$ and $\text{DySc}_2\text{N}@C_{80}\text{-}D_{5h}$ show a butterfly hysteresis shape with a fast relaxation near zero magnetic field, but the width of the hysteresis is much narrower for the $\text{DySc}_2\text{N}@C_{68}$. Likewise, the T_B value of $\text{DySc}_2\text{N}@C_{68}$ is only 3.8 K, which is much lower than the T_B value of 5.9 K in $\text{DySc}_2\text{N}@C_{80}\text{-}D_{5h}$. For comparison, the blocking temperature of the magnetization of $\text{DySc}_2\text{N}@C_{80}\text{-}I_h$ measured in the same conditions is 7 K. Among the studied DySc_2N -clusterfullerenes, $\text{DySc}_2\text{N}@C_{80}\text{-}I_h$ also shows the broadest hysteresis at 2 K (see Fig. 2c for a comparison).

A similar behavior is found in the series of Dy_2ScN -clusterfullerenes. $\text{Dy}_2\text{ScN}@C_{84}$ has the lowest blocking temperature, $T_B = 3.3$ K, and the narrowest magnetic hysteresis. In $\text{Dy}_2\text{ScN}@C_{80}\text{-}D_{5h}$, the blocking temperature is increased to



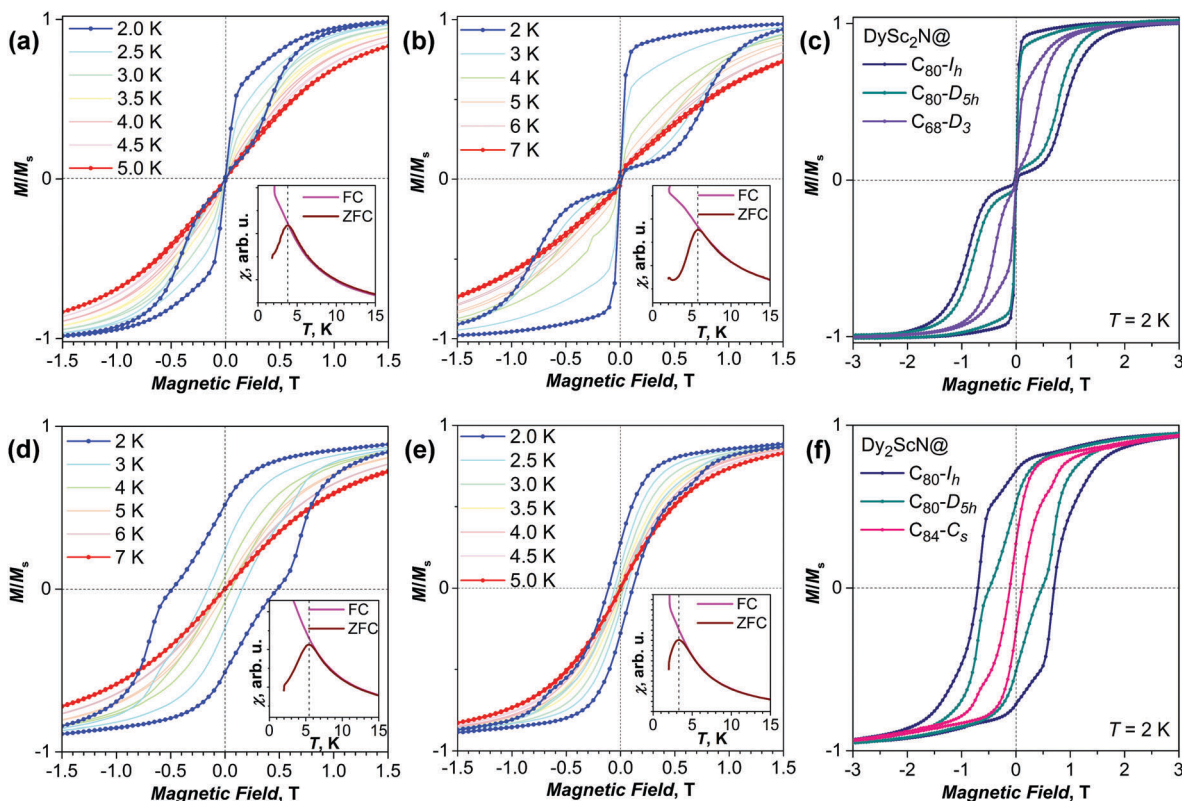


Fig. 2 Low temperature magnetization of Dy–Sc nitride clusterfullerenes: (a) DySc₂N@C₆₈-D₃; (b) DySc₂N@C₈₀-D_{5h}; (c) comparison of magnetic hysteresis curves of DySc₂N@C₆₈, DySc₂N@C₈₀-D_{5h}, and DySc₂N@C₈₀-I_h measured at 2 K in identical conditions; (d) DySc₂N@C₈₀-D_{5h}; (e) DySc₂N@C₈₄-C_s; (f) comparison of magnetic hysteresis curves of DySc₂N@C₈₄, DySc₂N@C₈₀-D_{5h}, and DySc₂N@C₈₀-I_h measured at 2 K in identical conditions. The insets in (a), (b), (d), and (e) show determination of the blocking temperature of magnetization T_B from temperature dependence of magnetic susceptibility (FC – field cooled, ZFC – zero-field cooled). Magnetic field sweep rate is 2.9 mT s⁻¹, temperature sweep rate is 5 K min⁻¹.

5.3 K, and the magnetic hysteresis is considerably broader. But the SMM parameters of the D_{5h} isomer are inferior to those of Dy₂ScN@C₈₀-I_h, which exhibits the T_B value of 8 K.^{4b}

Thus, both for the DySc₂N and Dy₂ScN cluster series, we conclude that the C₈₀-I_h cage gives the best SMM properties, the C₈₀-D_{5h} isomer has somewhat lower T_B values and narrower hysteresis, and the non-IPR C₆₈-D₃ and C₈₄-C_s cages give the softest SMMs with T_B values below 4 K and a narrow hysteresis.

Relaxation times of magnetization (τ) and their temperature dependence provide another tool to follow the SMM behavior at different temperatures (Fig. 3). For DySc₂N@C₈₀-I_h our earlier measurements^{4d} revealed a linear dependence of the τ values in Arrhenius coordinates between 1.8 and 5 K (Fig. 3a). Such a dependence corresponds to the Orbach relaxation mechanism:

$$\tau^{-1} = \tau_0^{-1} \exp(-U^{\text{eff}}/T) \quad (1)$$

where U^{eff} is the effective barrier and τ_0 is the attempt time, which for DySc₂N@C₈₀-I_h equals 23.6 ± 1 K and 0.6 ± 0.2 s, respectively.^{4d} For DySc₂N@C₈₀-D_{5h}, a clear deviation from the linear dependence is found in the same temperature range (Fig. 3a). The data can be equally well described either by a combination of an Orbach relaxation ($U^{\text{eff}} = 17.7 \pm 0.4$ K and an

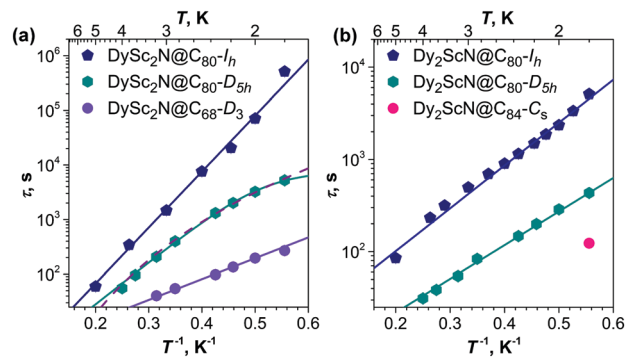


Fig. 3 Temperature dependence of the relaxation times of magnetization measured by DC magnetometry for (a) DySc₂N-EMFs in the field of 0.2 T to avoid QTM, and (b) for Dy₂ScN-EMFs in zero field. The lines are fits of the experimental data with the model discussed in the text.

attempt time of $\tau_0 = 0.8 \pm 0.1$ s) with a temperature-independent quantum tunneling (QTM, $\tau_{\text{QTM}} = 0.8 \pm 0.1$ s) denoted by a solid line in Fig. 3a, or by a power function of temperature, $\tau^{-1} = AT^n$, with $n = 5.6 \pm 0.2$ and $A = 6.7 \pm 0.8 \times 10^{-6} \text{ s}^{-1} \text{ K}^{-n}$. The power function with such parameters corresponds to the Raman relaxation mechanism with a contribution of optical phonons.¹² Relaxation times of DySc₂N@C₆₈ between 1.8 and 3 K can be described

by the Orbach mechanism with parameters $U^{\text{eff}} = 7.6 \pm 0.4$ K and $\tau_0 = 4.1 \pm 0.9$ s. Thus, for all three DySc₂N-EMFs the Orbach relaxation with a small barrier appears to be the main relaxation mechanism at low temperatures. Since the barrier is much smaller than expected crystal-field splitting of the Dy ⁶H_{15/2} manifold,^{9a} the observed barriers have to be assigned to low-frequency molecular vibrations such as frustrated rotations of the endohedral cluster.^{4d,5,8}

The low-temperature relaxation of dinuclear Dy-EMFs follows the Orbach mechanism *via* the exchange/dipolar excited state. That is, exchange and dipolar interactions favour the ferromagnetic (FM) coupling of two Dy spins in the magnetic ground state, whereas the first excited state is ascribed to the antiferromagnetic (AFM) coupling of the moments.^{4a,5} For Dy₂ScN@C₈₀-I_h, the energy difference between the FM and AFM states has been estimated as $U^{\text{eff}} = 10.7 \pm 0.3$ K with an attempt time of $\tau_0 = 11.9 \pm 1.5$ s.^{4b} For Dy₂ScN@C₈₀-D_{5h}, this study revealed a somewhat lower barrier of 8.4 ± 0.2 K and an attempt time of 4.1 ± 0.3 s. Thus, while the structure of the Dy–Sc nitride cluster is hardly affected, the D_{5h} isomeric cage reduces both the barrier and the attempt time of the Orbach relaxation *via* the exchange/dipolar excited state. Even shorter relaxation times of Dy₂ScN@C₈₄ preclude the detailed exploration of the temperature dependence. At 1.8 K, the relaxation time of this EMF is 123 s, which is considerably shorter than for both C₈₀ cages.

To conclude, the studies of the low-temperature magnetic properties of a series of Dy–Sc nitride clusterfullerenes showed that the SMM properties of these molecules strongly depend on the fullerene cage. The EMFs with the C₈₀-I_h cage show the best SMM properties, the non-IPR C₆₈-D₃ and C₈₄-C_s fullerenes give the lowest blocking temperature and narrow magnetic hysteresis curves, whereas the EMFs with C₈₀-D_{5h} cages exhibit an intermediate behaviour. The size of the fullerene cage does not seem to be as important as its shape. The magnetic relaxation rate can be correlated with the freedom of motion of the endohedral cluster inside the fullerene cage. In the molecules, which exhibit free rotation of the cluster at room temperature, the spin-phonon coupling is presumably weaker than in the molecules with the fixed position of the cluster, which leads to the faster relaxation of magnetization in the latter.

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Conflicts of interest

There are no conflicts to declare.

Notes and references

‡ Fullerene cages are denoted following the spiral algorithm.¹³ For the cages with adjacent pentagon pairs (non-IPR), the full numbering system is used. For IPR cages, more common abridged system counting only IPR isomer is employed.

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