**[M_{III}^{2+}M_{III}^{3+}]^{2+}** trigonal bipyramidal cages based on diamagnetic and paramagnetic metalloigands†

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A family of five \([M_{II}^{i}M_{III}^{j}]^{2+}\) trigonal bipyramidal cages \([M_{II}^{i} = \text{Fe, Cr and Al}; M_{III}^{j} = \text{Co, Zn and Pd}; i = 0 \text{ or } 1; j = 0 \text{ or } 2]\) of formula \([\text{Fe}_2\text{Co}_2\text{L}_6\text{Cl}_6] (1), [\text{Fe}_2\text{Zn}_3\text{L}_6\text{Br}_6] (2), [\text{Cr}_2\text{Zn}_3\text{L}_6\text{Br}_6] (3), [\text{Cr}_2\text{Pd}_3\text{L}_6(dppp)_2]^{6+}(\text{OTf})_6] (4) and [\text{Al}_3\text{Pd}_3\text{L}_6(dppp)_2]^{6+}(\text{OTf})_6] (5)\) have been synthesised using the tritopic \([M_{II}^{i}L]_3\) metallocigand in combination with the salts \(\text{Co}^{ii}\text{Cl}_2\) and \(\text{Zn}^{ii}\text{Br}_2\), which both act as tetrahedral linkers. The assembly of the \(\text{cis}-\text{protected} \quad \text{Pd}^{ii}(dppp)(\text{OTf})_2\) with \([M_{III}^{j}\text{L}]_3\) afforded the anionic cages 4–5 of general formula \([M_{II}^{i}\text{Pd}^{ii}(dppp)]^{6+}(\text{OTf})_6\). The metallic skeleton of all cages describes a trigonal bipyramid with the \(M_{III}^{j}\) ions occupying the two axial sites and the \(M_{II}^{i}\) ions sitting in the three equatorial positions. Direct current (DC) magnetic susceptibility, magnetisation and heat capacity measurements on 1 reveal weak antiferromagnetic exchange between the \(\text{Fe}^{ii}\) and \(\text{Co}^{ii}\) ions. EPR spectroscopy demonstrates that the distortion imposed on the \(\text{MO}_6\) coordination sphere of \([M_{III}^{j}\text{L}]_3\) by complexation in the \([M_{III}^{j}\text{L}]_3\) supramolecules results in a small, but measurable, increase of the zero field splitting at \(M_{III}^{j}\). Complete active space self-consistent field (CASSCF) calculations on the three unique \(\text{Co}^{ii}\) sites of 1 suggest \(D_{\text{Co}} = -14 \text{ cm}^{-1}\) and \(E/D = 0.1\), consistent with the magnetothermal and spectroscopic data.

**Introduction**

Molecular magnetism relies on the ability of the synthetic chemist to make an enormous breadth of structurally diverse polymetallic cages spanning the d- and f-block of the periodic table. The structural and magnetic characterisation of such species details the magneto-structural relationship and often uncovers fascinating magnetic phenomena which, in turn, suggest feedback into the synthesis of new complexes designed to enhance and improve properties toward application. Synthetic strategies for the design of polymetallic clusters containing multiple paramagnetic metal ions span the range from serendipitous self-assembly in which coordinately flexible metal ions, that can often exist in multiple oxidation states, are combined with organic ligands capable of bridging in numerous ways to form complexes whose absolute structures are difficult to predict, through to a more ‘supramolecular’ approach whereby metal ions with defined coordination geometries are paired with rigid ligands containing donor atoms with a single, predesigned orientation preference that afford, in most cases, a predicted structure. In the field of molecular magnetism, the latter is perhaps best exemplified by cyanometalate chemistry. A similar synthetic approach is followed in the metallosupramolecular chemistry of diamagnetic cages and capsules where the combination of directional metal–ligand bonding and rigorously rigid ligands creates cages with permanent internal cavities capable of hosting guest molecules, constructed primarily for potential application in, for example, catalysis, the stabilisation of reactive molecules and photochemistry. Due to the difficulties associated with performing solution-based one- and two-dimensional NMR spectroscopy on paramagnetic species, where broad signals and a wide chemical...
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shift range are commonplace, it is perhaps not surprising that the majority of metallosupramolecular chemistry has focused on the use of diamagnetic metal centres, albeit with some notable exceptions.

We recently initiated a project that would enable hetero-metallic, paramagnetic coordination cages to be accessed in a modular and predictable fashion, an approach centred around the tritopic metalloligand [MIII-L] (where HL = 1-(4-pyridyl)butane-1,3-dione), which features a tris(acac) octahedral transition metal core functionalised with three p-pyridyl)butane-1,3-dione), which features a tris(acac) octahedral transition metal core functionalised with three p-pyridyl donor groups (Fig. 1). Combination of the fac-isomer of [MIII-L] with a square-planar MII connector leads to the formation of [MIII-MII] molecular cubes. Herein we show that replacement of the square planar connector with tetrahedral or cis-capped square planar metal salts leads to the formation of trigonal bipyramidal [MIII-MII] cages, where MIII = Fe, Cr, Al and MII = Co, Zn, Pd. Reports of magnetic clusters based on this skeleton are rare, the only previous examples employing cyano bridging ligands.

Experimental section

Syntheses

1-(4-Pyridyl)butane-1,3-dione (HL) and the metalloligand [CrIII-L] were prepared according to previously published procedures. All reactions were performed under aerobic conditions. Solvents and reagents were used as received from commercial suppliers. Elemental analyses were carried by Medac Ltd.

[FeIII-L]
FeCl3 (1 mmol, 0.162 g), 1-(4-pyridyl)butane-1,3-dione (3.5 mmol, 0.57 g) and NaOMe (3.5 mmol, 0.189 g) were dissolved in 100 mL of MeOH/H2O (1:1 v/v) and left to stir until a red product precipitated (∼24 h). The resultant red precipitate was filtered and washed with water. The crude product was extracted with CHCl3 and dried over anhydrous MgSO4. The CHCl3 was removed under reduced pressure to a dark-red solid product was isolated by filtration and layered with Et2O. Orange, plate-shaped X-ray quality crystals were obtained after 20 days. Yield (98 mg, 67%). Elemental analysis (%) calculated for C54H48N6O12Cl6-Fe2Co3: C 44.00, H 3.28, N 5.70. Found: C 44.12, H 3.39, N 5.77.

[AlIII-L]3
Al(NO3)3·9H2O (1 mmol, 0.375 g), 1-(4-pyridyl)butane-1,3-dione (3.5 mmol, 0.57 g) and NaOMe (3.5 mmol, 0.189 g) were dissolved in 100 mL of MeOH/H2O (1:1 v/v) and left to stir until a white product precipitated (∼1 h). The resultant white precipitate was filtered and washed with water. The crude product was extracted with CHCl3 and dried over anhydrous MgSO4. The CHCl3 was removed under reduced pressure to afford the product as a white solid. Yield (0.39 g, 76%). Elemental analysis (%) calculated (found): C 63.16 (63.06), H 4.71 (4.53), N 8.18 (8.11).

[CrIIIL-L6](OTf)6 (4)
To a solution of the metalloligand [CrIIIL-L6] (108 mg, 0.2 mmol) in 35 mL of dichloromethane was added ZnBr2 (67 mg, 0.3 mmol). The resultant mixture was stirred for 30 minutes, before being filtered and layered with Et2O. Orange, plate-shaped X-ray quality crystals were obtained after 10 days. Yield (133 mg, 76%). Elemental analysis (%) calculated for C54H48N6O12Br6Fe2Zn1: C 36.85, H 2.75, N 4.77. Found: C 36.97, H 2.87, N 4.91.

[Fe2Co3L6Cl6] (1)
To a solution of the metalloligand [FeIIIL-L6] (108 mg, 0.2 mmol) in 35 mL of acetone, was added CoCl2 (39 mg, 0.3 mmol). The resultant mixture was stirred for 30 minutes, before being evaporated to dryness. The dark-red product was re-dissolved in nitromethane, filtered and allowed to stand. Dark-red, prism-shaped X-ray quality crystals were obtained after room temperature evaporation of the mother liquor after 10 days. Yield (133 mg, 76%). Elemental analysis (%) calculated for C54H48N6O12Zn2: C 36.97, H 2.87, N 4.91.

[Cr2Zn1L6Br6] (3)
To a solution of the metalloligand [CrIIIL-L6] (108 mg, 0.2 mmol) in 35 mL of dichloromethane was added ZnBr2 (67 mg, 0.3 mmol). After 1 day of reaction a precipitate appeared. The dark-red solid product was isolated by filtration, re-dissolved in DMF and layered with MeOH. Dark-red, prism-shaped X-ray quality crystals were obtained after 10 days. Yield (142 mg, 81%). Elemental analysis (%) calculated for C54H48N6O12Br6Cr2Zn1: C 37.01, H 2.76, N 4.80. Found: C 36.92, H 2.67, N 4.67.

[Cr2Pd3L6(dppp)3](OTf)6 (4)
To a solution of the metalloligand [CrIIIL-L6] (108 mg, 0.2 mmol) in 35 mL of methanol was added [Pd(dppp)(CF3SO3)2] (245 mg, 0.3 mmol; dppp is 1,3-bis(diphenylphosphino)propane). The solution was stirred for 30 minutes, before being filtered and allowed to stand. Orange, rod-shaped X-ray quality crystals were obtained after room temperature evaporation of the mother liquor after 5 days. Yield (275 mg, 78%). Elemental analysis (%) calculated for C141H126O30N6F18P6S6Cr2Pd3: C 48.00, H 3.60, N 5.77. Found: C 47.89, H 3.47, N 2.27.

Fig. 1 The molecular structure of the fac-isomer of the metalloligand [MIII-L] (M = Fe, Cr, Al). Colour code: MIII = green, O = red, N = blue, C = black. Hydrogen atoms have been omitted for clarity.
[Al₂Pd₆(dppp)₆(OTf)₆] (5)

To a solution of the metalloligand [AlIII₃] (103 mg, 0.2 mmol) in 35 mL of acetonitrile was added [Pd(dppp)(CF₃SO₃)] (245 mg, 0.3 mmol). The solution was stirred for 15 hours at 50 °C, before being filtered and layered with diethyl ether. Colourless, rod-shaped X-ray quality crystals were obtained after 5 days. Yield (288 mg, 83%).

1H NMR (500 MHz, CD₃CN): δ 8.61 (bs, 12H, Py-H), 7.79–7.67 (m, 12H, dppp-ArH), 7.48–7.44 (m, 12H, dppp-ArH), 7.42–7.39 (m, 12H, dppp-ArH), 7.26–7.22 (m, 12H, dppp-ArH), 7.18 (d, J = 6.5 Hz, 12H, Py-H), 6.13 (s, 6H, COC), 3.28–3.11 (m, 6H, dppp-Ch₂), 3.10–2.92 (m, 6H, dppp-Ch₂), 2.60–2.36 (m, 3H, dppp-Ch₂), 2.15 (s, 18H, Ch₃), 1.93–1.75 (m, 3H, dppp-Ch₂) ppm.

13C NMR (126 MHz, CD₃CN): δ 198.79, 177.61, 147.79, 134.91–134.82 (m, 2 signals), 133.70, 133.07–132.99 (m, 3 signals), 130.60–130.51 (m, 2 signals), 130.42–130.33 (m, 2 signals), 127.58–126.88 (m, 12m, 124.37, 122.06 (q, J = 321.0 Hz), 99.58, 28.10, 22.25–21.92 (m), 18.30 ppm.

19F NMR (471 MHz, CD₃CN) δ −79.05 ppm. Diffusion coefficient (DOSY, 500 MHz, CD₃CN, 298 K) 5.99 × 10⁻¹⁰ m² s⁻¹, hydrodynamic radius 9.9 Å. ESI TOF HRMS m/z: found 1010.1238 [M – 3OTf]⁺, calculated for [C₁₃₈H₁₂₆Al₂F₉N₆O₂₁P₆Pd₃S₃]³⁺ 1010.1069. Elemental analysis found for C₁₄₃H₁₂₆O₃₆F₉N₆P₆Pd₃S₃Al₂Pd₁⁺: C 48.69, H 3.65, N 2.42. Found: C 48.42, H 3.57, N 2.35.

Crystal structure information

For compounds 1, 2 and 3 single-crystal X-ray diffraction data were collected at T = 100 K on a Rigaku AFC12 goniometer equipped with an enhanced sensitivity (HG) Saturn 724+ detector mounted at the window of an FR-E Superbright MoKα single crystal X-ray diffractometer using Cu radiation at T = 120 K. The CrystalisPro software package was used for instrument control, unit cell determination and data reduction. Unit cell parameters in all cases were refined against all data. Crystal structures were solved using the charge flipping method implemented in SUPERFLIP (1, 2, and 3), or by direct methods with ShelXS (4 and 5). All structures were refined on Fo² by full-matrix least-squares refinements using ShelXL within the OLEX2 suite. All non-hydrogen atoms were refined with anisotropic displacement parameters, and all hydrogen atoms were added at calculated positions and refined using a riding model with isotropic displacement parameters based on the equivalent isotropic displacement parameter (Ueq) of the parent atom. All five structures contain accessible voids and channels that are filled with diffuse electron density belonging to uncoordinated solvent, and CF₃SO₃⁻ anions in the case of compounds 4–5. The SQUEEZE routine of PLATON was used to remove remaining electron density corresponding to solvent and anions not reported in the calculated formula. Crystallographic summary and structure refinement details are presented in Table 1. CCDC: 1520425–1520429.

Physical measurements

Magnetisation measurements were carried out on a Quantum Design SQUID MPMS-XXL magnetometer, operating between 1.8 and 300 K for DC applied magnetic fields ranging from 0 to 5 T. Microcrystalline samples were dispersed in eicosane in order to avoid torquing of the crystallites. Heat capacity measurements were carried out for temperatures down to ca. 0.3 K by using a Quantum Design 9T-PPMS, equipped with a 3He cryostat. The experiments were performed on thin pressed pellets (ca. 1 mg) of a polycrystalline sample, thermalised by ca. 0.2 mg of Apiezon N grease, whose contribution was subtracted by using a phenomenological expression. X- and Q-band EPR spectra were collected on powdered microcrystalline samples of [Fe₃₆] and compounds 1–4 at the UK National EPR Facility in Manchester.

Results and discussion

Solution self-assembly and structure

It could be reasonably expected that reaction of the metalloligand [AlIII₃] with a cis-protected square planar complex should yield a trigonal bipyramid. However, in the case of the archetypal 90° acceptor complex ([en]Pd(NO₃)₂)₄, it had previously been shown that instead, displacement of the bidentate ethylene diamine ligand occurs to yield the [Al₂Pd₆]₁²⁺ cube. We were thus pleased to find that when we switched to the more strongly coordinating bis(diphenylphosphino)propane (dppp), we were able to isolate the [Al₂Pd₆]₁⁹⁺ trigonal bipyramidal complex, 5, in 83% yield following reaction overnight at 50 °C between [AlIII₃] and [Pd(dppp)(OTf)] in acetonitrile. All the spectroscopic data indicate that the structure of 5, confirmed by X-ray crystallography (see below), is preserved in solution. As well as ESI-MS, which reveals the charge state corresponding to [5 – 3OTf]⁻ matching the expected isotopic distribution (see ESIF), the ¹H NMR spectrum of the product (Fig. 2b) shows just a single set of signals. The ¹H DOSY spectrum also indicates that all the resonances possess the same diffusion coefficient, which corresponds to a hydrodynamic radius of 9.9 Å, closely matching the data obtained by XRD.

It is also interesting to note that the starting metalloligand [AlIII₃] exists as a mixture of the mer and fac configurations, clearly evidenced by the multiplet for the acac CH and CH₃ signals in the ¹H NMR spectrum (Fig. 2c, resonances shown in blue and magenta), which is replaced by a singlet in the crude reaction mixture (Fig. 2d). This indicates that under the conditions of the reaction, [AlIII₃] is configurationally dynamic, and that the self-assembly process amplifies the proportion of the fac configuration through the formation of 5. While mer tris(bidentate) octahedral complexes are also known to generate discrete metallosupramolecular cages, the divergent disposition of the pendant donor groups create larger closed systems, which with a dynamic system such as this will rapidly rearrange
to give the entropically more favourable trigonal bipyramid. A comparison of the $^1$H NMR spectra of the re-dissolved crystalline sample of 5 (Fig. 2b) and the crude reaction solution, obtained by treating a slight excess of [Pd(dpdp)](OTf)$_2$ with [AlIII] in CD$_3$CN (Fig. 2d), shows that this amplification is not a solid-state packing effect, rather a solution-based effect. The single set of signals in the $^1$H NMR spectrum of the product (Fig. 2b/d) also indicates that 5 is formed with complete diastereoselectivity. This represents a second tier of self-sorting, which, unusually, involves Pd-mediated heterochiral recognition of $\Delta$ and $\Lambda$-[AlIII] enantiomers (see below).

**Solid-state structure descriptions**

The heterometallic trigonal bipyramid cages [Fe$_2$Co$_3$L$_6$Cl$_8$] (1), [Fe$_2$Zn$_3$L$_6$Br$_6$] (2), [Cr$_2$Zn$_3$L$_6$Br$_6$] (3), [Cr$_2$Pd$_3$L$_6$(dpdp)][(OTf)$_6$] (4) and [Al$_2$Pd$_3$L$_6$(dpdp)][(OTf)$_6$] (5) (Fig. 3 and 4) were all synthesised in a similar manner, by addition of either tetrahedral or cis-protected square planar M$^{II}$ compounds to the metal-ligand [M$^{III}$L$_6$] ($M^{III}$ = Fe, Cr or Al) in acetone, methanol, acetonitrile or a mixed solvent system, with crystals isolated from slow evaporation of the mother liquor, or diffusion of Et$_2$O or MeOH (see the Experimental section for full details). The metallic skeletons of the cages in 1–5 describe a trigonal bipyramidal with the M$^{III}$ ions situated on the axial positions and the M$^{II}$ ions on the equatorial sites. The approximate dimensions of the [M$^{III}$L$_6$] metalmetallic skeleton are M$^{III}$⋯M$^{II}$ (8.77–9.99 Å), M$^{III}$⋯M$^{III}$ (11.72–12.80 Å) and M$^{III}$⋯M$^{III}$ (10.75–11.20 Å).

Each of the three M$^{III}$ metal ions is coordinated by two N donors from the pyridyl groups of [M$^{III}$L$_6$]. The N–M$^{III}$–N angle of the tetrahedral Co$_2$Cl$_4$/ZnN$_2$Br$_2$ moiety for compounds 1–3 lies in the range 90.63–103.57$^\circ$; in 4–5 the equivalent N–Pd–N angle is in the range 84.40–85.39$^\circ$. Each [M$^{III}$L$_6$] corner unit consists of a six-coordinate M$^{III}$ ion with regular {M$^{III}$O$_6$} octahedral geometry. For the three different [M$^{III}$L$_6$] metallophanes used in the synthesis the M$^{III}$–O distances and angles are: Fe–O 1.98–

![Chemical Structure Image]
2.02 Å, Fe–O cis/trans angles 83.48–95.17° and 169.99–178.40°, respectively; Cr–O 1.91–1.98 Å, Cr–O cis/trans angles 86.82–94.04° and 176.35–179.85°, respectively; Al–O 1.86–1.89 Å, Al–O cis/trans angles 88.84–91.43° and 179.03–179.54°, respectively.

The CoII and ZnII ions lie in distorted tetrahedral environments with bond distances in the range 2.05–2.35 Å (Co–Cl ~ 2.23 Å, Co–N ~ 2.05 Å, Zn–Br ~ 2.35 Å and Zn–N ~ 2.06 Å) and bond angles around the metal centres ranging from 90.62° to 120.08°. In compounds 4 and 5, the PdII ion is 90° cis-blocked through the use of the dppp ligand (Pd–P bond distance ~ 2.27 Å). The coordination of Pd to [MIIIIL3] through the use of Pd–N bonds (ranging from 2.08–2.14 Å) creates a distorted square planar geometry around the Pd centre with cis/trans bond angles in the range 84.40–93.50° and 165.49–178.57°, respectively. While complexes 1–3 are neutral, charge balance is maintained in 4 and 5 through the presence of a total of six CF3SO3− anions, lying outside the cage.

While the intrametallic distances of the five trigonal bipyramids are similar, there is nonetheless a distinct diastereomeric difference between structures 1–3 and 4–5. Whereas 1–3 are all homochiral racemates in which each intact capsule features two [MIIIIL3] units that possess the same Δ or Δ chirality, in contrast structures 4 and 5 are both the achiral heterodiastereomer. While sorting of chiral octahedral metal motifs has been frequently observed in metallosupramolecular assembly reactions, for the vast majority homochiral assemblies are energetically preferred. The commonality of the [Pd(dppp)] unit in both 4 and 5 that feature different [MIIIIL3] metalloligands would suggest that either the small change in angle between pyridine donors at each MII connector and/or the interactions of the dppp protecting ligand with these donors cause the change in diastereomeric preference. Solution studies with 5 would also indicate this is not simply due to selective crystallization from a complex mixture (see above). Outwith cyanometalate
SQUID magnetometry

The dc (direct current) molar magnetic susceptibility, \( \chi \), of a polycrystalline sample of 1 was measured in an applied magnetic field, \( B \), of 0.1 T, over the 2–300 K temperature, \( T \), range. The experimental results are shown in Fig. 5 in the form of the \( \chi T \) product, where \( \chi = MB \) and \( M \) is the magnetisation of the sample. At room temperature, the \( \chi T \) product of 1 has a value of 14.4 cm\(^3\) K mol\(^{-1}\), in good agreement with the sum of Curie constants for a [Fe\(^{III}\)Co\(^{II}\)] unit (14.375 cm\(^3\) K mol\(^{-1}\), \( g_{\text{Fe}} = \frac{g_{\text{Co}}}{2} = 2.0 \)). Note that the estimation of the \( g \)-value of the Co\(^{II}\) ions here is an approximation and subject to error (e.g. lattice solvent lost upon sample drying will result in a variation of the samples diamagnetism), and a better measure comes from the EPR spectroscopy, which is consistent with \( g_{\text{Co}} = 2.3 \) (vide infra).

Upon cooling, the \( \chi T \) product of 1 remains essentially constant down to approximately 100 K, wherefrom it decreases upon further cooling to 9.5 cm\(^3\) K mol\(^{-1}\) at 2 K. Given that the anisotropy of Fe\(^{III}\) is negligible, this behaviour is consistent with a relatively large single-ion magnetic anisotropy for the Co\(^{II}\) centres and/or an antiferromagnetic exchange interaction between the Fe\(^{III}\) and Co\(^{II}\) centres. To better define the low-temperature magnetic properties of 1, low temperature variable-temperature-and-variable-field (VTVB) magnetisation data were measured in the temperature and magnetic field ranges \( T = 2–12 \) K and \( B = 0–5 \) T (Fig. 5). At the highest investigated field (5 T) and the lowest investigated temperature (2 K), the magnetisation of 1 is of 13.7 \( \mu_B \) (\( \mu_B \) is the Bohr magneton). Furthermore, when the VTVB data of 1 are plotted against the reduced quantity \( \mu_B B/kT \), little nesting of the VTVB data is observed. This observation indicates that the part of the energy spectrum of 1 probed under these experimental conditions does not present significant anisotropy splitting with respect to the temperature of measurement at zero magnetic field.

For the quantitative interpretation of the magnetisation data, we used spin-Hamiltonian (1)

\[
H = \mu_B B \sum_i g_i \hat{S}_i - 2 \sum_{i<j} J_{ij} \hat{S}_i \cdot \hat{S}_j + \sum_i D_i \left( \hat{S}_{iz}^2 - S_i (S_i + 1)/3 \right)
\]

where the summation indexes \( i,j \) run through the constitutive metal centres, \( g_i \) is the \( g \)-factor of the \( i \)th centre, \( \hat{S} \) is a spin operator, \( J \) is the isotropic exchange interaction parameter, \( D \) is the uniaxial anisotropy parameter and \( S \) is the total spin.

In our spin-Hamiltonian model, we assume for simplicity that all \( g \)-factors are equal to 2, \( S_{\text{Fe}} = 5/2, S_{\text{Co}} = 3/2 \), we only consider exchange interactions between Co\(^{II}\) and Fe\(^{III}\) centres, and neglect the single-ion anisotropy of Fe\(^{III}\). Furthermore, we fix the uniaxial anisotropy of Co\(^{II}\) to \( D_{\text{Co}} = -14 \) cm\(^{-1}\), as extracted from the modelling of the EPR data and theoretical calculations, which are discussed further in the following sections. Thus, at this point our model contains only one free parameter, namely, the isotropic exchange between Fe\(^{III}\) and Co\(^{II}\), \( J_{\text{Fe-Co}} \). The \( \chi T \) product of 1 was fitted to spin-Hamiltonian (1) by full matrix numerical diagonalisation of the spin-Hamiltonian of the full system of dimension 2304 by 2304, through use of the Levenberg–Marquardt algorithm.\(^{53}\) This resulted in the best-fit parameter \( J_{\text{Fe-Co}} = -0.04 \) cm\(^{-1}\). In order to verify the validity of our model, \( J_{\text{Fe-Co}} \) was fixed to the determined best-fit value, \( J_{\text{Fe-Co}} = -0.04 \) cm\(^{-1}\), and \( D_{\text{Co}} \) was maintained fixed at \(-14 \) cm\(^{-1}\). At this point our model contains no free parameters. Thereafter, the VTVB data of 1 were simulated by use of spin-Hamiltonian (1). The simulated curves are shown as solid red lines in Fig. 5. With these parameters, the energy spectrum of 1 consists of four groups of densely packed states, each separated by approximately 2D\(_{\text{Co}}\) (Fig. 6). It is interesting to note that multiple ground level crossings simultaneously occur at approximately 0.47 T when the magnetic field is applied parallel to the quantisation axis.

Heat capacity

Fig. 7 shows the collected heat capacity data, normalised to the gas constant, \( c_p/R \) of 1 as a function of temperature (between ca. 0.3 K and 30 K) for zero-applied magnetic field. As is typical for molecular magnetic materials,\(^{54}\) lattice vibrations contribute predominantly to \( c_p \) as a rapid increase above liquid-helium temperature. The lattice contribution can be described by the Debye model (dotted line in Fig. 7), which simplifies to \( c_p = aT^3 \) dependence at the lowest temperatures, where \( a = 7.6 \times 10^{-3} \) K\(^{-3}\) for 1.

For \( T < \text{ca.} \, 3 \) K, the zero-field \( c_p \) shows a wide bump-like anomaly, which we attribute to the splitting of the spin levels...
by zero-field splitting and magnetic interactions. At such low temperatures, the magnetic measurements are very sensitive to the applied magnetic field, as seen in the experimental behaviour for fields of 3 T and higher (inset of Fig. 7). Such large intensities of the applied magnetic field are sufficient for promoting full decoupling between the individual spin centres (we recall that the exchange interaction is as small as \(J_{\text{Fe-Co}} = -0.04 \text{ cm}^{-1}\) on the basis of the fit of the magnetometry data). Therefore, the temperature and field dependence of the \(c_p\) data in Fig. 7 (inset), collected for \(B \geq 3\) T, are particularly suitable for probing the influence of crystal fields on \(1\), down to temperatures significantly lower than the ones obtained in the magnetisation measurements.

The solid lines in Fig. 7 are the curves calculated for Hamiltonian (1), using the best-fit parameters from the magnetothermal and spectroscopic data and theoretical calculations, i.e., \(D_{\text{Co}} = -14 \text{ cm}^{-1}\) and the here-negligible \(J_{\text{Fe-Co}} = -0.04 \text{ cm}^{-1}\). The agreement with the experimental data is good, though not outstanding. Anticipating the discussion on the EPR spectra (vide infra), we have checked that adding a zero-field splitting (ZFS) of \(D_{\text{Fe}} = -0.2 \text{ cm}^{-1}\) at the \(\text{Fe}^{III}\) sites does not improve the fit. The discrepancy is most evident below \(ca. 1\) K, where the experimental data have lower values than the calculated ones. This behaviour can be explained by a wider broadening of the low-lying energy spectrum, likely induced by higher-order anisotropy terms, which are not taken into account in Hamiltonian (1).

EPR spectroscopy

We previously reported EPR spectra of \([\text{CrL}_3]\), giving the ZFS of the \(\text{Cr}^{III}\) site as \(D = -0.55 \text{ cm}^{-1}\) with a small rhombicity of \(|E/D| = 0.045\)\(^{29}\). Q-Band spectra of 3 and 4 are similar to that of \([\text{Cr}^{III}\text{L}_3]\), and give \(D = -0.64\) and \(-0.61 \text{ cm}^{-1}\), respectively (Fig. S11; \(|E/D| = 0.03-0.04\)\(^{26}\). Hence, the distortion imposed on the \{CrO\}_6 coordination sphere of \([\text{Cr}^{III}\text{L}_3]\) by complexation in the \{Cr\}_2\{M\}_3 supramolecules results in a small, but measurable, increase of the ZFS at \(\text{Cr}^{III}\). The \{CrO\}_6 metric parameters do not appear to be very different.

Such an increase in \(D\) is also found for the \(\text{Fe}^{III}\) \((s = 5/2)\) systems. X- and Q-band EPR spectra of \([\text{Fe}^{III}\text{L}_3]\) reveal a rather small ZFS of \(D = 0.08 \text{ cm}^{-1}\) with \(|E/D| = 1/3\) (Fig. 8 and S12; note the sign of \(D\) has no significance with a fully rhombic D-tensor). These values are similar to those reported for \([\text{Fe}(\text{acac})_3]\) \((|D| = 0.16 \text{ cm}^{-1}, E/D = 0.3)\)\(^{30}\) and \([\text{Fe}(\text{dpm})_3]\) \((\text{dpm = dipivaloylmethane}; D = -0.20 \text{ cm}^{-1}, |E/D| = 0.25)\).\(^{37}\) On incorporation into the \{Fe\}_2\{Zn\}_3 complex 2, a much richer spectrum is observed (Fig. 8 and S12), giving \(D = 0.20 \text{ cm}^{-1}\) \((E/D = 1/3)\). Angular overlap model studies on \([\text{Fe}(\text{acac})_3]\) and \([\text{Fe}(\text{dpm})_3]\) show \(D\) to be very sensitive to the trigonal distortion at \(\text{Fe}^{III}\)\(^{37}\), and there is a more significant structural difference in the \{Fe\}_2\{O\}_6 coordination spheres when bound in \{Fe\}_2\{Zn\}_3, with longer Fe–O bonds and wider O–Fe–O angles in the \{py\}_3 face, than in the equivalent \(\text{Cr}^{III}\) systems.

The \{Fe\}_2\{Co\}_3 complex 1 gives Q-band EPR spectra with very broad features at \(ca. 5, 9\) and \(12\) kG that line up with the main features of the spectrum of the \{Fe\}_2\{Zn\}_3 complex 2. Hence, the ZFS at \(\text{Fe}^{III}\) must be similar. The very large ZFS at \(\text{Co}^{II}\) means that only transitions within the ground Kramers doublet of this ion are observed (the microwave energy, \(hν \ll |D|\)), and there must be a significant rhombicity in order for these transitions to
fall within the observed features. The spectra also show that the $J_{\text{FeCo}}$ exchange interaction must be very weak, resulting only in severe broadening of the peaks. Test calculations on a simple $\{\text{Fe}^{III}\text{Co}^{II}\}$ model, with fixed ZFS at the $s = 5/2$ and $3/2$ spins (the latter taking $D = -14$ cm$^{-1}$ with $E/D = 0.1$; averaging the results of CASSCF calculations − see below) suggest that if $|J_{\text{FeCo}}| > ca. 0.02$ cm$^{-1}$ then additional features would be observed in the Q-band EPR spectrum. Note that the limit for the full, five-spin system would be different.

The $D_{\text{Fe}}$ values obtained from EPR would have a negligible effect on the calculated $\chi(T)$ and $c_p(T,B)$ curves for 1, and a negligible effect on the global level structure in Fig. 6a, because both $|D_{\text{Fe}}|$ and $|J_{\text{FeCo}}|$ are $\ll |D_{\text{Co}}|$. However, it would affect the detail of the states within each of the densely packed multiplets of Fig. 6a, because $|D_{\text{Fe}}|$ and $|J_{\text{FeCo}}|$ are of similar magnitude.

**Theoretical studies**

In order to independently verify the large ZFS of Co$^{II}$ we have performed complete active space self-consistent field (CASSCF) calculations on the three unique Co$^{II}$ sites of 1, see the SI for details. The results suggest $D_{\text{Co}} = -14$ cm$^{-1}$, $E/D = 0.1$ (Table S1) which is entirely consistent with the magnetometry and heat capacity data. The calculations also suggest that the principal axes of the local ZFS tensors are oriented roughly perpendicular to the Fe$^{III}$-Fe$^{III}$ axis and canted approximately $120^\circ$ with respect to one another in the plane (Fig. 9). Accounting for the non-collinearity in spin-Hamiltonian (1) did not improve the quality of the fits to the magnetometry or heat capacity data.

**Conclusions**

Complexes 1–5 represent a novel, and unusual family of trigonal bipyramidal cage complexes, built with the tritopic $[\text{ML}_3]$ metallogligand, featuring a tris(acac) octahedral transition metal core functionalised with three $p$-pyridyl donor groups, and a series of transition metal salts. Outwith cyanometalate chemistry, compound 1 represents the first example of such a cage containing paramagnetic metal ions. Complementary studies investigating the diamagnetic variants using $^1$H NMR spectroscopy reveal some interesting features about the solution self-assembly process. Firstly, the $[\text{M}^{III}\text{L}_3]$ metallogligand is a highly dynamic tritopic building block as evidenced by $fac$ configurational isomer being amplified at the expense of the $mer$ during the course of cage formation. The self-assembly process also occurs with high and unusual stereoselectivity wherein the trigonal bipyramids are formed exclusively from twisted pyramidal components of opposite Λ/Λ-handedness. Solution stability of the cage is also confirmed via mass spectrometry. SQUID magnetometry and heat capacity measurements on 1 reveal weak antiferromagnetic exchange between the Fe$^{III}$ and Co$^{II}$ ions, with $|D_{\text{Co}}| = 14$ cm$^{-1}$. EPR spectroscopy reveals that the distortion imposed on the $[\text{MO}_3]$ coordination sphere of $[\text{M}^{III}\text{L}_3]$ by complexation in the $[\text{M}_2^{III}\text{M}_1^{II}]$ supramolecules results in a small, but measurable, increase of the zero field splitting at M$^{III}$. CASSCF calculations on the three unique Co$^{II}$ sites of 1 suggest that the principal axes of the local ZFS tensors are oriented perpendicular to the Fe$^{III}$-Fe$^{III}$ axis, but canted $\sim 120^\circ$ with respect to each other.

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**Notes and references**

Furthermore, in situ monitoring of the self-assembly reaction in CD$_3$CN with excess Pd salt using $^1$H NMR spectroscopy (e.g. Fig. 2d), showed this not to be in fast exchange.

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