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## The interplay between fluorescence and phosphorescence with luminescent gold(I) and gold(III) complexes bearing heterocyclic arylacetylide ligands<sup>†</sup>

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The photophysical properties of a series of gold(I)  $[\text{LAu}(\text{C}\equiv\text{CR})]$  ( $\text{L} = \text{PCy}_3$  (**1a–4a**), RNC (**5a**), NHC (**6a**)) and gold(III) complexes  $[\text{Au}(\text{C}^{\text{N}}\text{C})(\text{C}\equiv\text{CR})]$  (**1b–4b**) bearing heterocyclic arylacetylide ligands with narrow band-gap are compared. The luminescence of both series are derived from an intraligand transition localized on the arylacetylide ligand ( $\pi\pi^*(\text{C}\equiv\text{CR})$ ) but **1a–3a** displayed prompt fluorescence ( $\tau_{\text{PF}} = 2.7\text{--}12.0\text{ ns}$ ) while **1b–3b** showed mainly phosphorescence ( $\tau_{\text{Ph}} = 104\text{--}205\text{ }\mu\text{s}$ ). The experimentally determined intersystem crossing (ISC) rate constants ( $k_{\text{ISC}}$ ) are on the order of  $10^6$  to  $10^8\text{ s}^{-1}$  for the gold(I) series (**1a–3a**) but  $10^{10}$  to  $10^{11}\text{ s}^{-1}$  for the gold(III) analogues (**1b–3b**). DFT/TDDFT calculations have been performed to help understand the difference in the  $k_{\text{ISC}}$  between the two series of complexes. Owing to the different oxidation states of the gold ion, the Au(I) complexes have linear coordination geometry while the Au(III) complexes are square planar. It was found from DFT/TDDFT calculations that due to this difference in coordination geometries, the energy gap between the singlet and triplet excited states ( $\Delta E_{\text{ST}}$ ) with effective spin-orbit coupling (SOC) for Au(I) systems is much larger than that for the Au(III) counterparts, thus resulting in the poor ISC efficiency for the former. Time-resolved spectroscopies revealed a minor contribution (<2.9%) of a long-lived delayed fluorescence (DF) ( $\tau_{\text{DF}} = 4.6\text{--}12.5\text{ }\mu\text{s}$ ) to the total fluorescence in **1a–3a**. Attempts have been made to elucidate the mechanism for the origins of the DF: the dependence of the DF intensity with the power of excitation light reveals that triplet-triplet annihilation (TTA) is the most probable mechanism for the DF of **1a** while germinate electron-hole pair (GP) recombination accounts for the DF of **2a** in 77 K glassy solution ( $\text{MeOH}/\text{EtOH} = 4 : 1$ ). Both **4a** and **4b** contain a BODIPY moiety at the acetylide ligand and display only  $^1\text{IL}(\pi\pi^*)$  fluorescence with negligible phosphorescence being observed. Computational analyses attributed this observation to the lack of low-lying triplet excited states that could have effective SOC with the  $S_1$  excited state.

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## Introduction

Phosphorescence is a distinctive photophysical property of transition-metal complexes, which has widespread applications in diverse areas. As it is derived from the ‘forbidden’ radiative relaxation of a triplet-excited state to the singlet ground state, it

is featured by long emission lifetime (in  $\mu\text{s}$ ) and reduced emission energy compared to fluorescence commonly encountered in organic luminophores. Electronic transitions associated with a change of spin are prohibited by the spin-selection rule. However, transition-metal ions that have high atomic number and hence, large spin-orbit coupling constant ( $\xi$ ), usually lead to efficient spin-orbit coupling (SOC) that relaxes the spin selection rule. Fast intersystem crossing (ISC) in the sub-picosecond to picosecond time regime<sup>1</sup> leads to rapid depletion of a singlet excited state to a triplet excited state instead of fluorescence as fluorescence radiative lifetime is typically in the nanosecond range. Thus, in transition-metal complexes, phosphorescence normally prevails in their luminescence spectra.

However, in recent years, there are an increasing number of reports on transition-metal complexes which display slow ISC

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rate with lifetimes ranging from hundreds of ps to ns. For instance, 2,5-bis(arylethynyl)rhodacyclopentadiene complexes ( $\xi_{\text{Rh}} = 1260 \text{ cm}^{-1}$ )<sup>2</sup> were reported to display exclusively prompt fluorescence with high emission quantum yields of 0.3–0.7 and lifetimes of 1–3 ns, corresponding to ISC rate constants of  $\sim 10^8 \text{ s}^{-1}$ .<sup>3</sup> In addition, transition-metal complexes containing fused aromatic systems such as perylene, perylene diimide, pyrene and tetracene also show ligand-dominated fluorescence (see Fig. 1).<sup>4</sup> Hence, it has become clear that the presence of heavy elements does not guarantee fast ISC rate; the molecular structure and the nature of the ligands may play more critical roles in determining the ISC rate.

Luminescent Au(i) complexes are well documented to display rich photophysical properties. Although Au(i) complexes generally display phosphorescence owing to the large SOC constant of Au(i) ion ( $\xi_{\text{Au}} \sim 5100 \text{ cm}^{-1}$ ),<sup>2</sup> ligand-centered fluorescence has also been reported in a number of gold(i) complexes. For example, as revealed by the luminescence of  $[\text{TEE}(\text{AuPCy}_3)_4]$  and  $[\text{TEB}(\text{AuPCy}_3)_3]$  (TEE = tetraethynylethene; TEB = 1,3,5-triethynylbenzene), subtle changes in the electronic structure of the bridging alkynyl ligand leads to intense phosphorescence ( $\Phi_{\text{em}} = 0.46$ ,  $\tau = 285 \mu\text{s}$ ) in the latter but solely fluorescence ( $\Phi_{\text{em}} = 0.22$ ,  $\tau < 0.05 \mu\text{s}$ ) in the former.<sup>5</sup> In both cases, the luminescence originates from the ligand-centered transition mainly localized on the bridging alkynyl ligands. Che and co-workers also reported a series of Au(i)-conjugated acetylides,  $[(\text{Cy}_3\text{P})\text{Au}(\text{C}\equiv\text{C}-\text{C}_6\text{H}_4)_{n-1}(\text{C}\equiv\text{CPh})]$  ( $n \geq 2$ ), which display dual fluorescence (prompt and delayed) and phosphorescence.<sup>6</sup> Both the  $\Phi_{\text{em}}$  and ratio of fluorescence *versus* phosphorescence were found to depend on the conjugation length (number of repeating units  $n$ ) and the substitution pattern of arylacetylides ligands.

As a continuous effort to elucidate the ligand effects on the photophysics of luminescent gold complexes, heterocyclic arylacetylides ligands containing narrow band-gap moieties (benzothiadiazole (L1), coumarin (L2), naphthalimide (L3) and boron-dipyrromethene (referred to as Bodipy) (L4); Chart 1) were chosen in this study. A series of Au(i) heterocyclic arylacetylides complexes, **1a**–**4a**, were synthesized. Tricyclohexylphosphine

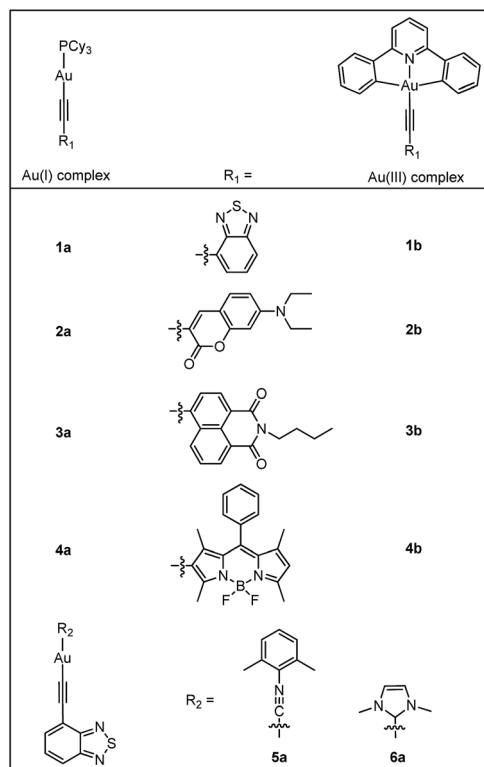


Chart 1 Au(i) and Au(III) acetylides complexes studied in this work.

( $\text{PCy}_3$ ) was used as the auxiliary ligand in these complexes because (1) it is optically transparent at wavelength  $>250 \text{ nm}$  so that it is not involved in the emissive excited states in the UV-visible spectral region and (2) its steric bulkiness would prevent the gold ions from coming into close contact that could lead to low-lying excited states originated from metal–metal and  $\pi$ – $\pi$  interactions. It is worth mentioning that several recently reported Au(i) alkynyl complexes bearing similar benzothiadiazole,<sup>7</sup> coumarin<sup>8</sup> and naphthalimide<sup>9</sup> derivatives also show similar luminescence properties as our Au(i) complexes.<sup>10,11</sup> The effects of auxiliary ligands on the photophysical properties of Au(i) complexes were also studied by comparing **1a** with two derivatives containing 2,6-dimethylphenyl isocyanide (RNC, **5a**) and 1,3-dimethylimidazol-2-ylidene (NHC, **6a**) instead of the phosphine auxiliary ligand, respectively.

The effects of the oxidation state of the metal ion on the photophysical behaviours of transition-metal complexes are relatively unexplored. Herein, an analogous series of Au(III)-acetylides supported by the cyclometalated  $[\text{C}^{\text{N}}\text{C}]$  ligand (**1b**–**4b**;  $\text{HC}^{\text{N}}\text{CH} = 2,6\text{-diphenylpyridine}$ ) were also prepared and their photophysical properties were compared with those of the Au(i) counterparts. The photophysical properties of both Au(i) and Au(III) complexes were investigated by steady-state and time-resolved spectroscopic measurements. DFT/TDDFT calculations were performed on the pairs (**1a**, **1b**) and (**4a**, **4b**) in order to understand the origin of the dramatic difference in ISC efficiencies between these Au(i) and Au(III) complexes.

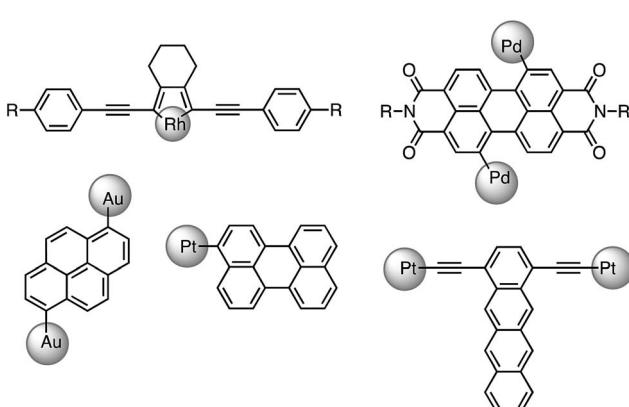


Fig. 1 Selected examples of transition-metal complexes that display dominant fluorescence instead of phosphorescence. Auxiliary ligands coordinated to the metal ions are omitted.



## Results

### Synthesis and characterization

The gold(i) alkynyl complexes **1a–6a** were synthesized in 53–79% yields following the protocol of base deprotonation ( $\text{NaOMe}$ ) of terminal alkynes and substitution of chloride ion of the corresponding Au(i) precursors.<sup>5a,6a,b,12</sup> As these complexes were observed to show signs of decomposition on  $\text{SiO}_2$  column, column chromatography was not used for their purification. Analytically pure **1a–6a** were obtained by recrystallization from  $\text{CH}_2\text{Cl}_2$ /hexane mixtures. The Au(iii) complexes **1b–4b** were synthesized by copper-catalyzed Sonogashira coupling between terminal alkynes and  $[\text{Au}(\text{C}^{\wedge}\text{N}^{\wedge}\text{C})\text{Cl}]$  using deoxygenated  $\text{CH}_2\text{Cl}_2$  as the solvent, with  $\text{NEt}_3$  added to initiate the deprotonation of alkynes.<sup>13</sup> These complexes were purified by chromatography on  $\text{SiO}_2$  column using dichloromethane and hexane as eluent. The yields were 51–84%.

All complexes have been characterized by  $^1\text{H}$  and  $^{13}\text{C}$  NMR, mass spectrometry (FAB+) and elemental analyses. Ligands L1–L4 were characterized by  $^1\text{H}$  NMR and MS-EI. The complexes are stable in the solid state and in solution under ambient conditions. Complexes **1a–6a** are highly soluble in  $\text{CH}_2\text{Cl}_2$  and THF but are less soluble in alcoholic solvents such as MeOH. Complexes **1b–4b** have lower solubility compared with their Au(i) counterparts. All of these gold complexes appear as yellow or orange solids except for **4a** and **4b** that are purplish red. The  $^{31}\text{P}$  signals of **1a–4a** occur at *ca.*  $\delta$  56.3 as a singlet, characteristic of the  $^{31}\text{P}$  signals of the Au- $\text{PCy}_3$  moiety that usually appear in the range of  $\delta$  56.0–58.0.<sup>6a,b,12</sup> In the  $^{13}\text{C}$  NMR spectra, two doublets are observed at *ca.*  $\delta$  131.2–146.2 ( $^2J_{\text{CP}} \approx 130$  Hz) and 94.6–98.4 ( $^3J_{\text{CP}} \approx 24$  Hz) which can be assigned to the  $\alpha$  and  $\beta$ -acetylenic carbons.<sup>6a,b,12b</sup> In **6a**, the carbene carbon ligated to gold occurs at  $\delta$  187.7.<sup>14</sup>

### X-Ray crystallography

Crystals of **1a**, **4a** and **2b** were obtained by layering hexane over concentrated  $\text{CH}_2\text{Cl}_2$  solutions. Their crystal data and selected bond lengths and angles are given in ESI.† Fig. 2 shows the

structures of **1a** and **4a** (top panel). The  $\text{P1}-\text{Au1}-\text{C}(\text{acetylide})$  angles of **1a** and **4a** are  $175.0(13)$  and  $178.1(2)^\circ$  and  $\text{Au1}-\text{C}\equiv\text{C}$  angles are  $170.9(4)$  and  $177.6(5)^\circ$ , respectively, revealing slight deviation from linear coordination geometry. The  $\text{Au1}-\text{C}(\text{acetylide})$  distances of  $2.049(4)$  and  $2.001(5)$  Å and  $\text{C}\equiv\text{C}$  distances of  $1.146(7)$  and  $1.191(8)$  Å for **1a** and **4a**, respectively, are comparable with those of other reported gold(i) acetylide complexes.<sup>6a,b,12a,b</sup> The crystal packing diagrams of **1a** and **4a** are shown in ESI (Fig. S1†). In both cases, there are no short intermolecular contacts; the closest  $\text{Au}\cdots\text{Au}$  distances are  $5.9715(5)$  and  $5.2726(4)$  Å for **1a** and **4a**, respectively.

The crystal structure of **2b** (Fig. 2, bottom) shows a slightly distorted square-planar geometry with  $\text{C1}-\text{Au1}-\text{C17}$  angle of  $162.44(14)^\circ$ . The  $\text{Au1}-\text{C}(\text{acetylide})$  and  $\text{C}\equiv\text{C}$  distances are  $1.969(4)$  and  $1.197(5)$  Å, respectively. These parameters are similar to those found in related cyclometalated Au(iii) arylacetylide complexes.<sup>13</sup> The torsional angle between the  $\text{Au}(\text{C}^{\wedge}\text{N}^{\wedge}\text{C})$  and arylacetylide planes is approximately  $72.6^\circ$ . This non-planarity gives rise to negligible  $\pi\cdots\pi$  stacking between molecules as shown in Fig. S2 in ESI.†

### Electrochemical properties

The electrochemical properties of selected complexes, **1a–4a** and **1b–4b**, were investigated by cyclic voltammetry. The electrochemical data are summarized in Table 1. The cyclic voltammograms of the Au(i) complexes and their Au(iii) counterparts are shown in Fig. S3, ESI.† Except for **1a** and **2a**, both classes of complexes display both irreversible oxidation ( $E_{\text{pa}} = +0.6$  to  $+1.4$  V) and quasi-reversible/irreversible reduction waves ( $E_{\text{pc}} = -1.5$  to  $-1.8$  V) attributed to the redox process localized on the arylacetylide. For the pairs [**2a**, **2b**] and [**4a**, **4b**], the  $E_{\text{pa}}$  occur at relatively low potential of *ca.*  $+0.6$  and  $+0.7$  V, respectively, suggesting higher HOMO level of the conjugated coumarin and Bodipy than the other heterocyclic moieties. For the Au(iii) complexes, other than the redox reactions occurring at the arylacetylide ligands at potentials similar to the Au(i) counterparts, there are also irreversible reduction waves at *ca.*  $-1.9$  to  $-2.0$  V attributable to reduction of the  $[\text{C}^{\wedge}\text{N}^{\wedge}\text{C}]$  ligands.

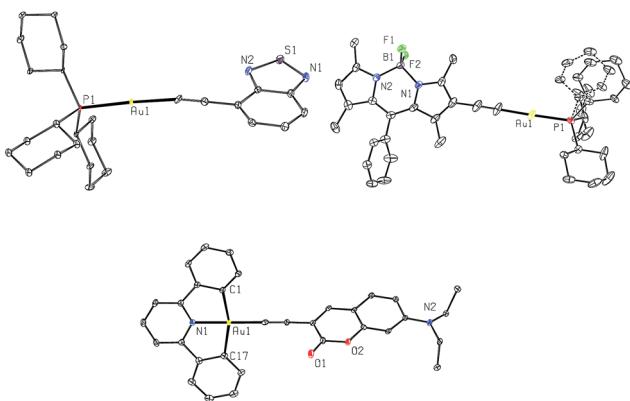


Fig. 2 Perspective drawings of the crystal structures of **1a** (top left), **4a** (top right) and **2b** (bottom) with the thermal ellipsoids shown at 30% probability level. Hydrogen atoms have been omitted for clarity.

Table 1 Electrochemical data of **1a–4a** and **1b–4b**<sup>a</sup>

Complex	$E_{\text{pc}}^b$ /V	$E_{\text{pa}}^c$ /V
<b>1a</b>	$-1.79^d$	—
<b>2a</b>	—	$0.61, 0.95$
<b>3a</b>	$-1.61^d$	$1.44$
<b>4a</b>	$-1.50^d$	$0.71, 1.14$
<b>1b</b>	$-1.82, -2.03$	$1.25$
<b>2b</b>	$-1.89$	$0.59, 1.02$
<b>3b</b>	$-1.59, -1.95$	$1.42$
<b>4b</b>	$-1.49^d, -2.02$	$0.74$

<sup>a</sup> Values determined in  $\text{CH}_2\text{Cl}_2$  ( $\text{Cp}_2\text{Fe}^{+0}$  occurs at  $E_{1/2} = +0.15\text{--}0.16$  V) at  $298$  K; values reported *versus*  $\text{Ag}/\text{AgNO}_3$  reference electrode; electrolyte:  $0.1$  M  $n\text{Bu}_4\text{NPf}_6$ ; scan rate =  $100$  mV s $^{-1}$ . <sup>b</sup> Cathodic peak potential ( $E_{\text{pc}}$ ) of irreversible wave. <sup>c</sup> Anodic peak potential of ( $E_{\text{pa}}$ ) irreversible wave. <sup>d</sup>  $E_{1/2} = E_{\text{pa}} + E_{\text{pc}}$  of quasi-reversible wave.



Table 2 Photophysical data of **1a–6a**, **1b–4b** and **L1–L4<sup>a</sup>**

	UV/Vis absorption, $\lambda_{\text{max}}/\text{nm}$ ( $10^3 \epsilon/\text{M}^{-1} \text{cm}^{-1}$ )	Emission						
		Medium	$\lambda_{\text{F}}/\text{nm}$	$\lambda_{\text{ph}}/\text{nm}$	$\tau_{\text{PF}}^b/\text{ns}$	$\tau_{\text{DF}}^c/\mu\text{s}$	$\tau_{\text{phos}}/\mu\text{s}$	
<b>1a</b>	265 (14.1), 275 (14.6), 305 (9.3), 311 (9.8), 319 (12.4), 379 (6.9)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	467	—	12.0	11.9	—	0.91
		Glassy 77 K	442	630, 688 <sup>e</sup>		n.d. <sup>f</sup>	109 <sup>e</sup>	
		Solid 298 K	504	—		5.3	—	
		Solid 77 K	492	—		11.6	—	
<b>2a</b>	271 (13.9), 316 (3.2), 331 (3.2), 410 (40.0), 423 (3.7, br)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	466	596, 652 <sup>e</sup>	2.7	4.6	n.d. <sup>f</sup>	0.70
		Glassy 77 K	460, 481	596, 653 <sup>e</sup>		11.8 <sup>g</sup>	13.6 <sup>g</sup>	
		Solid 298 K	405 (weak), 480 (sh), 515 (max)	598, 652 (sh) <sup>e</sup>		243 <sup>e</sup>	203 <sup>e</sup>	
		Solid 77 K	409 (weak), 487 (max), 517 (sh)	597, 654(sh) <sup>e</sup>		15.0 <sup>e</sup>	24.6 <sup>e</sup>	
<b>3a</b>	283 (12.2), 333 (5.6), 350 (11.3), 380 (23.8), 397 (25.5)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	441	613, 670 <sup>e</sup>	2.8	2.6, 12.5	n.d. <sup>f</sup>	0.78
		Glassy 77 K	418, 439, 462 (sh)	569, 609, 666 <sup>e</sup>		31.5 <sup>g</sup>	61.9 <sup>g</sup>	
		Solid 298 K	403 (sh), 503	610, 668 <sup>e</sup>		n.d. <sup>f</sup>	530 <sup>e</sup>	
		Solid 77 K	404 (sh), 504	627, 680 <sup>e</sup>		29.1	64.9 <sup>e</sup>	
<b>4a</b>	280 (13.5), 325 (4.5), 412 (10.1), 553 (40.2)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	593	—	0.8	—	—	0.04
<b>5a</b>	266 (18.6), 279 (15.8), 294 (12.5), 305 (11.2), 311 (10.1), 319 (12.7), 371 (7.1)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	456	—	11.0	6.8	—	0.90
<b>6a</b>	259 (18.3), 280 (17.5), 298 (7.8), 305 (10.7), 311 (11.4), 319 (14.2), 383 (8.3)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	476	—	14.0	7.4	—	0.84
<b>1b</b>	283 (20.4), 310 (24.1), 318 (26.8), 369 (13.0), 381 (13.8)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	461 <sup>h</sup>	630, 671 (sh)	13.8 <sup>h</sup>	—	104	0.003 <sup>i</sup>
<b>2b</b>	312 (13.9), 406 (35.9, br), 432 (40.7)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	473 <sup>h</sup>	610, 668		n.d. <sup>f</sup>		
				592, 642 (sh)	9.0 <sup>h</sup>	—	124	0.01 <sup>i</sup>
				585, 605, 643			1200	
<b>3b</b>	312 (12.7), 325 (11.9), 379 (26.6), 395 (29.5)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	459 <sup>h</sup>	530, 601, 660			1.4	
				603, 659 (sh)	5.2 <sup>h</sup>	—	205	0.04 <sup>i</sup>
				598, 614 (sh), 652			2200	
<b>4b</b>	312 (12.1), 320 (12.0), 366 (7.6), 384 (9.6), 401 (10.7), 515 (24.0, br), 546 (41.7)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	583	—	2.1	—	—	0.13
<b>L1</b>	303 (8.7), 309 (9.7), 316 (11.9), 341 (4.1)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	412	—	1.0	—	—	0.07
<b>L2</b>	260 (15.9), 326 (4.7), 405 (24.0, br), 417 (24.7)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	455	—	3.6	—	—	0.94
<b>L3</b>	333 (14.6), 350 (22.4), 367 (20.6)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	378, 398, 419 (sh)	—	0.5	—	—	0.11
<b>L4</b>	321 (3.9, br), 376 (5.3, br), 486 (14.8, sh), 517 (45.4)	CH <sub>2</sub> Cl <sub>2</sub> 298 K	533	—	6.4	—	—	0.83

<sup>a</sup> Data were obtained from steady-state measurements with degassed CH<sub>2</sub>Cl<sub>2</sub> solutions ( $2 \times 10^{-5} \text{ M}$ ) unless specified. Measurements with glassy solutions were performed in EtOH/MeOH (4 : 1) mixture at 77 K. <sup>b</sup> Emission lifetimes of prompt fluorescence ( $\tau_{\text{PF}}$ ) were determined by time-correlated single photon counting (TCSPC) measurement. <sup>c</sup> Emission lifetimes of delayed fluorescence ( $\tau_{\text{DF}}$ ) were obtained from fitting the decay of the time-resolved emission (TRE) as a mono-exponential decay in the delay time range of 0–40 ns and 1–46  $\mu\text{s}$ , respectively. Measurements were performed in degassed CH<sub>2</sub>Cl<sub>2</sub> ( $5 \times 10^{-5} \text{ M}$ ) solutions. <sup>d</sup> Emission quantum yields ( $\Phi_{\text{em}}$ ) were obtained using quinine sulfate in degassed 0.5 M H<sub>2</sub>SO<sub>4</sub> ( $\Phi = 0.546$ ) as the standard unless specified.  $\Phi_{\text{em}}$  measured in steady state is the overall emission quantum yield, *i.e.*  $\Phi_{\text{em}} = \Phi_{\text{PF}} + \Phi_{\text{DF}}$  for **1a–3a** and **5a–6a**. <sup>e</sup> Obtained from time-resolved emission spectra. <sup>f</sup> Emission lifetime was not determined (n.d.) due to weak emission signal. <sup>g</sup> Determined from time-resolved emission spectra in degassed CH<sub>2</sub>Cl<sub>2</sub> ( $1 \times 10^{-5} \text{ M}$ ) solutions. <sup>h</sup> Determined from fs-TRF spectra. <sup>i</sup> Emission quantum yields ( $\Phi_{\text{em}}$ ) were obtained using [Ru(bpy)<sub>3</sub>][PF<sub>6</sub>]<sub>2</sub> in degassed acetonitrile as the standard ( $\Phi = 0.062$ ).

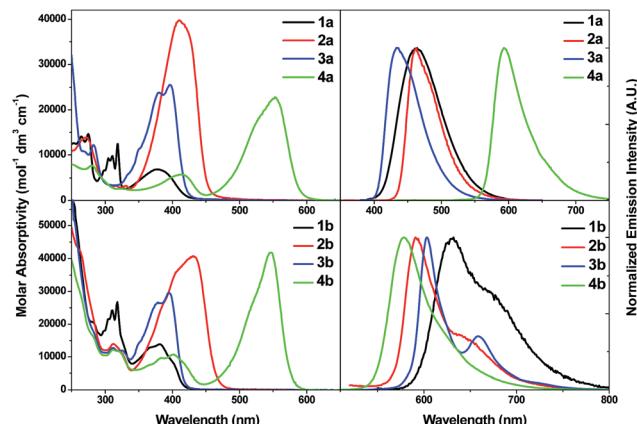


Fig. 3 UV-vis absorption spectra (left) and emission spectra (right) of **1a**–**4a** (top) and **1b**–**4b** (bottom) in  $\text{CH}_2\text{Cl}_2$  at 298 K ( $2 \times 10^{-5}$  M).

### UV-vis absorption spectroscopy

All photophysical data of the gold complexes and the free ligands L1–L4 are listed in Table 2.

**Electronic absorption of Au(i) complexes **1a**–**6a**.** Fig. 3 (top left) shows the absorption spectra of **1a**–**4a**. The lowest energy absorption bands of **1a**–**4a** are at 379, 410, 397 and 553 nm, respectively, and their molar absorptivities ( $\epsilon$ ) fall in the range  $6.9 \times 10^3$  to  $4 \times 10^4 \text{ mol}^{-1} \text{ dm}^3 \text{ cm}^{-1}$ . These lowest energy absorption bands have spectral features resembling those of L1–L4 (Fig. S4 in ESI<sup>†</sup>) and are attributable to the dipole-allowed intraligand transitions of the arylacetylides ligands ( $^1\pi\pi^*(\text{C}\equiv\text{CR})$ ) with some charge-transfer character. Similar assignments were also made for other Au(i) alkynyl complexes in the literature.<sup>6,7,8a</sup> Bathochromic shifts of  $^1\text{IL}$  transitions of arylacetylides are observed upon coordination of the arylacetylides to the Au(i) ion and are ascribed to  $\pi$ -interaction between Au(i) 5d orbitals and the ligand  $\pi$ -orbitals (see MO surfaces in Fig. 9 and 10).

Replacing the neutral auxiliary ligand  $\text{PCy}_3$  in **1a** with 2,6-dimethylphenyl isocyanide (RNC, **5a**) and 1,3-dimethylimidazol-2-ylidene (NHC, **6a**) results in a slight change in  $\lambda_{\text{max}}$  of the lowest energy absorption band ( $\lambda_{\text{max}} = 371 \text{ nm}$  (**5a**, RNC) and  $383 \text{ nm}$  (**6a**, NHC) cf.  $\lambda_{\text{max}} = 379 \text{ nm}$  (**1a**,  $\text{PCy}_3$ )) (Fig. S6 in ESI<sup>†</sup>).

**Electronic absorption spectra of Au(iii) complexes **1b**–**4b**.** Fig. 3 (bottom left) shows the UV-vis absorption spectra of the Au(iii) complexes studied in the present work. The spectral features at  $\lambda > 300 \text{ nm}$  are similar to those of the Au(i) analogues and free ligands L1–L4. Complexes **1b**–**4b** display low-energy absorption bands at  $\lambda_{\text{max}} = 381, 432, 395$  and  $546 \text{ nm}$ , respectively, and with  $\epsilon$  values in the range of  $1.38 \times 10^4$  to  $4.17 \times 10^4 \text{ mol}^{-1} \text{ dm}^3 \text{ cm}^{-1}$ . These absorption bands are, like the Au(i) complexes, attributable to  $^1\pi\pi^*(\text{C}\equiv\text{CR})$  transitions. Comparisons of the absorption spectra of **1a**–**4a** and **1b**–**4b** at  $\lambda \leq 300 \text{ nm}$  revealed that the absorption bands in this spectral region are more intense in **1b**–**4b** with  $\epsilon$  values of *ca.*  $4 \times 10^4 \text{ mol}^{-1} \text{ dm}^3 \text{ cm}^{-1}$  and these high energy absorption bands likely involve intraligand  $^1\pi\pi^*(\text{C}^{\wedge}\text{N}^{\wedge}\text{C})$  transitions (Fig. S7 in ESI<sup>†</sup>).

### Steady-state emission spectroscopy

All of the complexes are luminescent in degassed  $\text{CH}_2\text{Cl}_2$  at room temperature and in 77 K glassy solutions (EtOH : MeOH = 4 : 1) upon excitation at the corresponding lowest-energy absorption  $\lambda_{\text{max}}$ . As depicted in Table 2, there is a distinct difference between the two classes of complexes: the Au(i) complexes **1a**–**6a** display predominantly fluorescence while the Au(iii) complexes, **1b**–**3b**, exhibit exclusively weak phosphorescence. Complex **4b**, on the other hand, shows fluorescence only.

**Emission of **1a**–**4a** and **5a**–**6a**.** In dichloromethane solutions, structureless emission bands are observed at  $\lambda_{\text{max}} = 467, 466, 439$  and  $553 \text{ nm}$  for **1a**–**4a**, respectively (Fig. 3, top right). The corresponding excitation spectra of **1a**–**4a** can be found in the ESI (Fig. S8<sup>†</sup>). The emission quantum yields for **1a**–**3a** are high ( $\Phi_{\text{em}} = 0.91, 0.70$  and  $0.78$ , respectively). In the case of **4a**, its emission quantum yield is low ( $\Phi_{\text{em}} = 0.04$ ). Emission lifetimes of **1a**–**4a** are in the nanosecond time regime: 0.8–12 ns. As the emissions of these Au(i) complexes resemble those of the corresponding free ligands L1–L4, they are attributable to  $^1\pi\pi^*(\text{C}\equiv\text{CR})$  excited states, with some charge transfer character, which probably arise from mixings of metal-to-ligand charge-transfer (MLCT) character. Solvent effects on the emissions of **1a**–**3a** can be found in Fig. S10, ESI.<sup>†</sup> There is no discernible phosphorescence for **1a**–**4a** under steady-state conditions in solutions at either room temperature or 77 K (Fig. S9, ESI<sup>†</sup>).

Comparing the three Au(i) complexes bearing the benzothiadiazole moiety, the emission energies ( $\lambda_{\text{max}} = 467, 456$  and  $476 \text{ nm}$  for **1a** ( $\text{PCy}_3$ ), **5a** (RNC) and **6a** (NHC), respectively (Fig. S11 in ESI<sup>†</sup>)) and emission lifetimes ( $\tau_{\text{PF}} \sim 11$ –14 ns) are similar, indicating that the auxiliary ligand plays an insignificant role in modification of the electronic structures of the excited states.

**Emission of **1b**–**4b**.** Emission spectra of the Au(iii) complexes are depicted in Fig. 3 (bottom right). Contrary to the Au(i) analogues where the emission profiles are structureless, the emission spectra of complexes **1b**–**3b** are vibronically structured with  $\lambda_{\text{max}}$  at 630, 592 and 603 nm and quantum yields of 0.003, 0.01 and 0.04, respectively. The emission lifetimes are of hundreds of microseconds ( $\sim 100 \mu\text{s}$  for **1b** and **2b**;  $\sim 200 \mu\text{s}$  for **3b**). Taking into account the large Stokes shifts (between 6300 and  $10\,400 \text{ cm}^{-1}$ ), structured emission profiles, and long emission lifetimes, the emissions of **1b**–**3b** could be attributed to  $^3\pi\pi^*(\text{C}\equiv\text{CR})$  excited states with negligible mixings of MLCT and LLCT character (LLCT = ligand-to-ligand charge transfer). Solvent effects on the emissions of **1b** can be found in Fig. S12, ESI.<sup>†</sup> On the contrary, **4b** shows emission with a small Stokes shift of  $920 \text{ cm}^{-1}$  and emission lifetime of only 2.1 ns. Thus, the emission of **4b** is derived from fluorescence with  $^1\pi\pi^*(\text{C}\equiv\text{CBodipy})$  parentage.

### Time-resolved spectroscopies for the gold(i) complexes

Nanosecond time-resolved emission (ns-TRE) spectra of the Au(i) complexes in degassed  $\text{CH}_2\text{Cl}_2$  solutions at 298 K ( $5 \times 10^{-5}$  M) are measured at different time delays and are presented in Fig. 4 (**1a**) and ESI (**2a**–**3a**, **5a**–**6a**; Fig. S13<sup>†</sup>). There are two

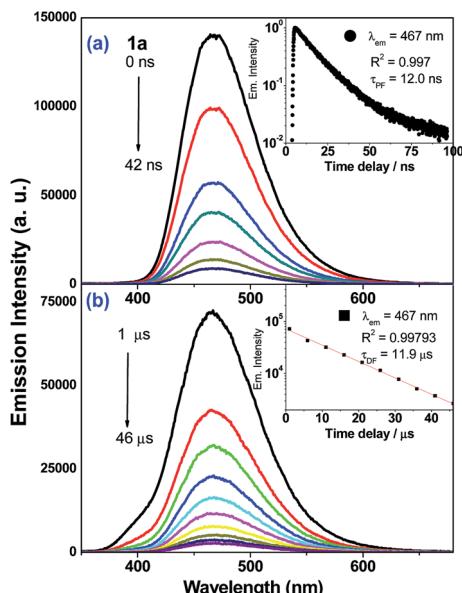


Fig. 4 ns-TRE spectra of **1a** recorded from (a) 0–42 ns and (b) after a time delay of 1  $\mu\text{s}$  in degassed  $\text{CH}_2\text{Cl}_2$  ( $5 \times 10^{-5}$  M) at 298 K. Inset shows the emission kinetic decay trace. Decay time constants were fitted as mono-exponential decay ( $\lambda_{\text{exc}} = 355 \text{ nm}$ ).

components in the emission decay: a major component which decays within nanoseconds ( $\tau_1 = 12.0$  (**1a**), 2.7 (**2a**), 2.8 (**3a**), 11.0 (**5a**) and 14.0 ns (**6a**)) and a minor component with microsecond decay lifetime ( $\tau_2 = 11.9$  (**1a**), 4.6 (**2a**), 2.6, 12.5 (**3a**), 6.8 (**5a**) and 7.4  $\mu\text{s}$  (**6a**)). For each of these Au(I) complexes, both decay components have identical emission profile and peak energy and so, the short-lived one ( $\tau_1$ ) is assigned to be prompt fluorescence (PF) while the long-lived one ( $\tau_2$ ) is delayed fluorescence (DF) of  ${}^1\pi\pi^*(\text{C}\equiv\text{CR})$  character. In the case of **4a**, only PF ( $\tau_{\text{PF}} = 0.8$  ns) is observed. The proportion of DF and PF constituting the total emission of **1a–3a** have been estimated (Table 3): the intensity of DF is minute (<3%) when compared with that of PF (>97%). It is noted that delayed fluorescence in the microsecond time regime is indicative of the emission generated from a long-lived excited state. This is further supported by nanosecond transient absorption (ns-TA) measurements that reveal the presence of long-lived absorbing species in the microsecond timescale (*vide infra*).

Table 3 Proportion of PF and DF constituting the fluorescence of **1a–3a**<sup>a</sup>

Complex	% PF	% DF
<b>1a</b>	99.2	0.81
<b>2a</b>	99.9	0.1
<b>3a</b>	97.1	2.9

<sup>a</sup> % PF and % DF are estimated by integrating the emission intensity of degassed  $\text{CH}_2\text{Cl}_2$  ( $5 \times 10^{-5}$  M) in the spectral region of  $\lambda = 350\text{--}700 \text{ nm}$  over the time range: 0–500 ns and 800 ns to 999  $\mu\text{s}$ , respectively ( $\lambda_{\text{exc}} = 355 \text{ nm}$ ).

Weak phosphorescence bands were observed for **1a–3a** under different conditions. For dilute  $\text{CH}_2\text{Cl}_2$  solutions ( $1 \times 10^{-5}$  M) at 298 K, dominant emissions were observed in the spectral region of 440–470 nm, which correspond to fluorescence (Fig. S14, left panel in ESI†). In addition, weak emission peaks at *ca.* 600 nm become discernible for **2a** and **3a** and the lifetimes measured are 13.6 and 61.9  $\mu\text{s}$ , respectively (Fig. S13, right panel in ESI†). Cooling to 77 K gives more resolved phosphorescence bands with vibrational progression spacings of 1300–1400  $\text{cm}^{-1}$  for all three complexes (Fig. 5 and 6). For **1a**, contrary to the ns-TRE spectra recorded in degassed  $\text{CH}_2\text{Cl}_2$  at room temperature (Fig. 4 ( $5 \times 10^{-5}$  M); Fig. S14, ESI† ( $1 \times 10^{-5}$  M)) where only DF could be observed over the time range 1–46  $\mu\text{s}$ , in 77 K glassy solution, phosphorescence at 630 nm is dominant and the weak DF at 467 nm vanishes after 80  $\mu\text{s}$  (Fig. 5a). The phosphorescence band decays with first-order kinetics at  $\tau_{\text{phos}} = 109 \mu\text{s}$ . Similarly, the low-temperature ns-TRE spectra of **3a** is dominated by phosphorescence at 609 nm and DF vanishes after 200  $\mu\text{s}$  (Fig. 5b). The phosphorescence band also decays mono-exponentially with  $\tau_{\text{phos}} = 530 \mu\text{s}$ . The photodynamics of **2a** at 77 K, however, is different from that of **1a** and **3a**: both DF and phosphorescence of **2a** are of comparable intensities initially ( $\sim 1 \mu\text{s}$ ) in the 77 K ns-TRE spectra (Fig. 6); in addition, DF and phosphorescence do not follow first-order kinetics but decay according to the power law ( $I \propto t^{-1}$ ) in the time interval 1  $\mu\text{s}$  to 1.2 ms (inset of Fig. 6). The thermally induced Stokes shifts ( $\Delta E_s = E_{00}$  (77 K) –  $E_{00}$  (298 K)), being  $\sim 0$  (**2a**) and  $\sim 107 \text{ cm}^{-1}$  (**3a**), are small, thus supporting that the phosphorescence bands are originated from  ${}^3\text{IL}$ .<sup>15</sup> Moreover, as the emission energies and profiles of the low-energy bands are similar to those of the steady-state emission

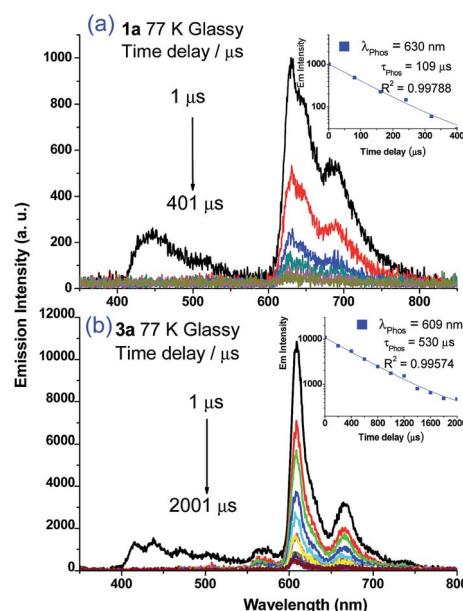


Fig. 5 ns-TRE spectra of (a) **1a** and (b) **3a** in 77 K glassy solution (EtOH/MeOH = 4 : 1) recorded at different time intervals.  $\lambda_{\text{exc}} = 355 \text{ nm}$ ; integration time: 80 and 200  $\mu\text{s}$  for **1a** and **3a**, respectively. Insets of (a) and (b) show the kinetic decay traces at the specified wavelengths with the estimated phosphorescence lifetime ( $\tau_{\text{phos}}$ ).

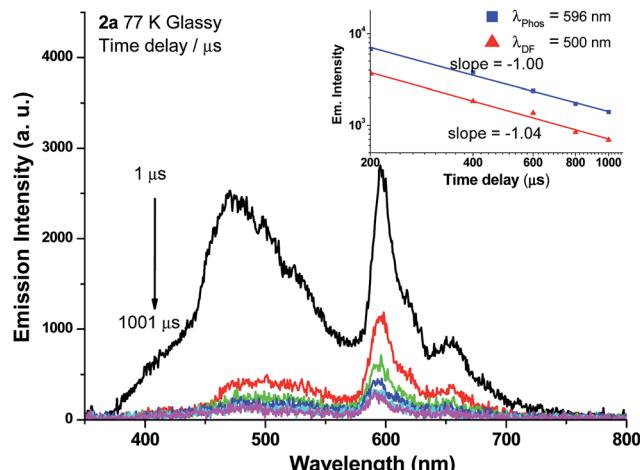


Fig. 6 ns-TRE spectra of **2a** in 77 K glassy solution (EtOH/MeOH = 4 : 1) recorded at different time intervals.  $\lambda_{\text{exc}} = 355$  nm; integration time: 200  $\mu$ s. Inset shows the log–log plot of emission intensity of DF ( $\lambda_{\text{DF}} = 500$  nm) and phosphorescence ( $\lambda_{\text{phos}} = 596$  nm) of **2a** in 77 K glassy solution against time; both decay according to a power law:  $I \propto t^{-1}$ .

spectra of the Au(III) analogues, the low-energy emission bands of **1a–3a** are assigned to be from phosphorescence decay of the  $^3\pi\pi^*(\text{C}\equiv\text{CR})$  excited state.

Nanosecond transient absorption (ns-TA) difference spectra of **1a–3a** (Fig. 7) and **5a–6a** (Fig. S15 in ESI<sup>†</sup>) have been recorded in deoxygenated  $\text{CH}_2\text{Cl}_2$  at a gate delay of 1  $\mu$ s after excitation at  $\lambda = 355$  nm. The ns-TA spectra are characterized by an intense positive signal due to excited-state absorption (ESA) within the spectral range 400–700 nm. The decay time constants of the lowest-energy ESA ( $\tau_{\text{ESA}}$ ) are 20.3 (**1a**), 39.2; 304 (**2a**), and 13.3; 70.9  $\mu$ s (**3a**) (insets of Fig. 7). Changing the auxiliary ligand from PCy<sub>3</sub> (**1a**) to RNC (**5a**) and NHC (**6a**) results in negligible changes in the ns-TA spectra and  $\tau_{\text{ESA}}$  (Fig. 7 vs. S15<sup>†</sup>), suggesting that auxiliary ligand has little effect on the photophysics of the gold(I) arylacetylide complexes.

### Time-resolved spectroscopies for the gold(III) complexes

ns-TRE and ns-TA difference spectra of **1b–3b** have been recorded in degassed  $\text{CH}_2\text{Cl}_2$  solutions at 298 K at a gate delay of 1  $\mu$ s. The

ns-TRE spectra of **1b–3b** (Fig. S16 in ESI<sup>†</sup>) have the same emission profiles and peak positions as the corresponding steady-state phosphorescence spectra and exhibit single exponential decay lifetimes of 20.8 (**1b**), 9.7 (**2b**) and 18.3  $\mu$ s (**3b**). For the ns-TA difference spectra of **1b–3b** (Fig. 8, bottom panel), a broad positive ESA band was observed in the spectral region 450–800 nm; this ESA signal follows first-order kinetics with lifetimes determined to be 23.8 (**1b**), 13.6 (**2b**) and 25.4  $\mu$ s (**3b**), in reasonable agreement with the phosphorescence decay lifetimes determined from their respective ns-TRE spectra, thus indicating that the broad ESA is derived from  $\text{T}_1 \rightarrow \text{T}_n$  absorption.

To probe the early excited state dynamics of the gold(III) complexes, in particular the events associated with ISC, femtosecond time-resolved fluorescence (fs-TRF) and transient absorption difference spectra (fs-TA) of **1b–3b** have been recorded. Fig. 8 depicts the fs-TRF (top panel) and fs-TA spectra (middle panel) of complexes **1b–3b** in  $\text{CH}_2\text{Cl}_2$  solution at various time intervals after 400 nm excitation at 298 K. Promptly (<2 ps) after photo-excitation, an unstructured fluorescence band peaking at 461 (**1b**), 473 (**2b**) and 459 nm (**3b**) appears and decays completely within 100 ps. As the TRF emission peaks and profiles closely resemble those of their Au(I) analogues, **1a–3a**, these TRF spectra are suggested to be originated from the  $^1\pi\pi^*(\text{C}\equiv\text{CR})$  excited state. Fitting of the kinetic traces at their peaking wavelengths reveals that bi-exponential functions are required for **1b–3b** with  $\tau_1$  and  $\tau_2$  being 1.28 and 13.8 ps for **1b**, 0.95 and 9.04 ps for **2b**, and 0.74 and 5.22 ps for **3b**.

In the fs-TA of **1b–3b** (Fig. 8, middle panel), all three complexes displayed similar spectral transformations: the initially formed ( $\sim$ 1.4–2.5 ps) excited state absorption peaking at  $\sim$ 490 nm (ESA1) decays with a concomitant growth of a broad band covering a spectral region 450–800 nm (ESA2) and is fully developed within 40 ps and persists up to 2.7 ns (the longest time recorded in the fs measurements). Clear isosbestic points could be observed at  $\sim$ 500 nm (**1b**), 530 nm (**2b**) and  $\sim$ 500 and 700 nm (**3b**) during the temporal evolution. Such kind of spectral conversion points to a precursor–successor relationship between ESA1 and ESA2. Kinetic analyses at representative wavelengths of these TA spectra reveals that ESA1 of **1b** and **3b** decay bi-exponentially with  $\tau_1$  and  $\tau_2$  being 0.80 and 13.2 ps for **1b** and 0.63 and 3.49 ps for **3b**, respectively, whereas ESA1 of **2b** decays with a single exponential time constant of  $\tau_2 = 8.38$  ps.

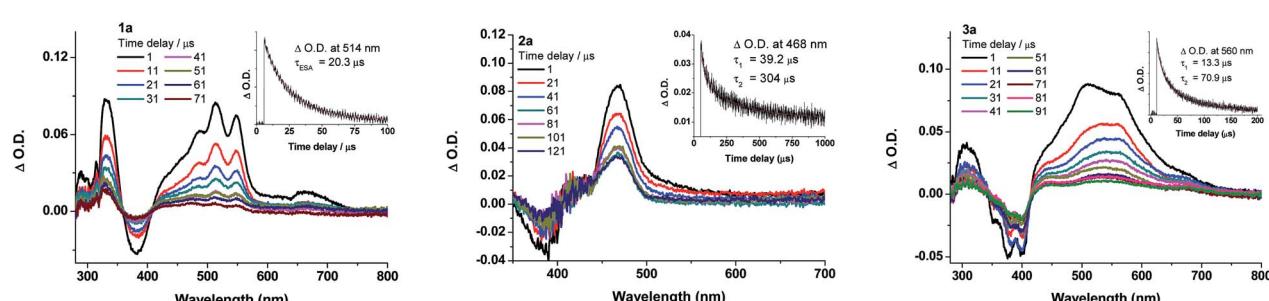


Fig. 7 Nanosecond transient absorption (ns-TA) difference spectra of **1a–3a** recorded at selected decay times in degassed  $\text{CH}_2\text{Cl}_2$  ( $5 \times 10^{-5}$  M) at 298 K. Insets show the ESA kinetic decay trace at the specified wavelengths; decay lifetimes were fitted as mono-exponential decay for **1a** and bi-exponential decays for **2a** and **3a**. ( $\lambda_{\text{exc}} = 355$  nm; integration time: 200 ns).



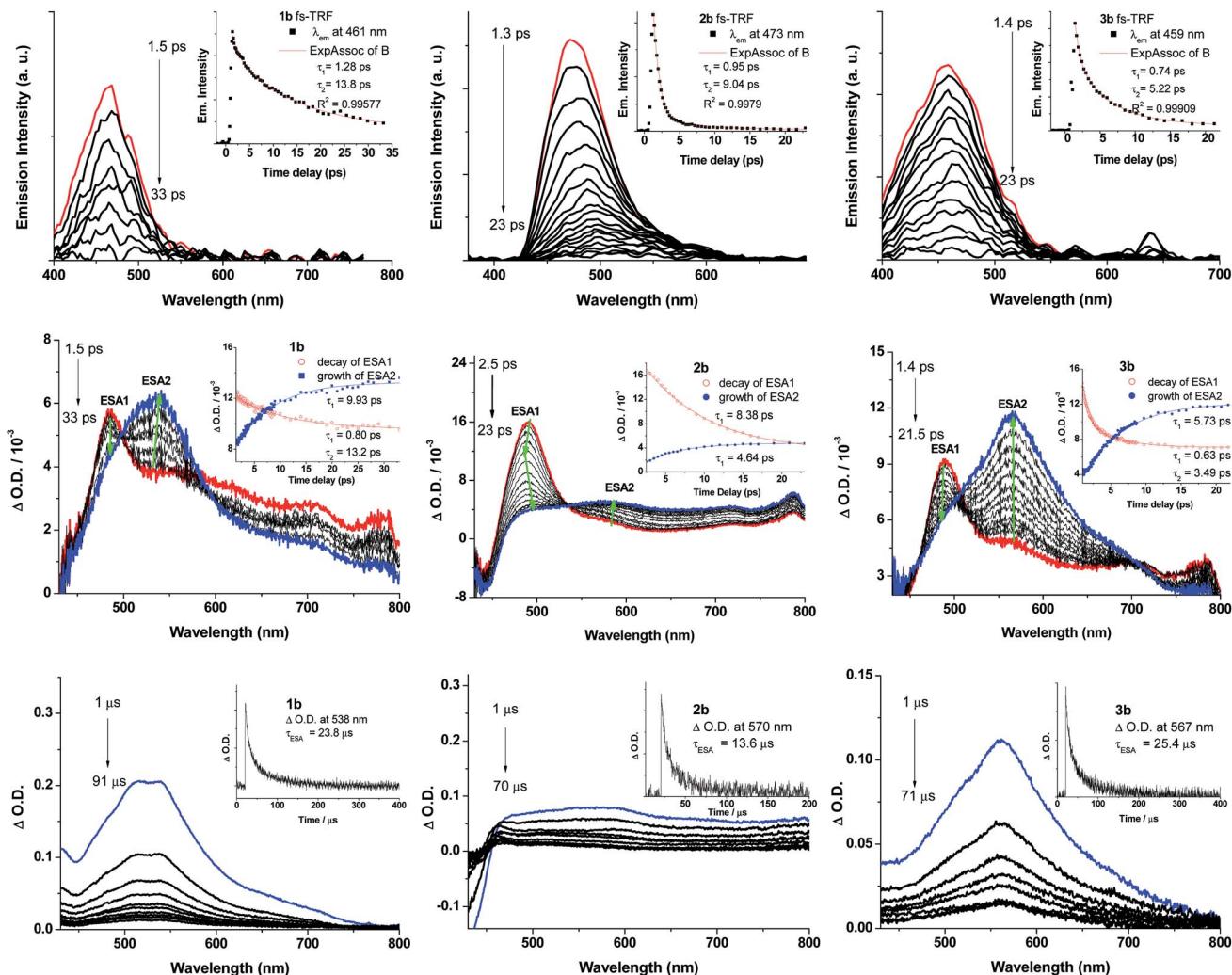


Fig. 8 (Top) fs-TRF spectra and (middle) fs-TA difference spectra of **1b**–**3b** in  $\text{CH}_2\text{Cl}_2$  ( $5 \times 10^{-5}$  M) at 298 K ( $\lambda_{\text{exc}} = 400$  nm; 120 fs fwhm). Arrows indicate the spectral evolution. (Bottom) ns-TA difference spectra of **1b**–**3b** in degassed  $\text{CH}_2\text{Cl}_2$  (laser  $\lambda_{\text{exc}}$ : 355 nm). Insets show the kinetic time profiles and the decay time constants at the specified wavelengths.

ESA2, on the other hand, grows with first-order kinetics for all three complexes **1b**–**3b** with time constants  $\tau_{\text{ESA2}} = 9.93$  (**1b**), 4.64 (**2b**) and 5.73 ps (**3b**). Given the similar decay time constants between the fs-TRF and ESA1 in fs-TA of **1b**–**3b**, the spectral dynamics for both time-resolved spectra should be originated from the same  $S_1$  excited state, namely, the  $^1\pi\pi^*(\text{C}\equiv\text{CR})$  excited state as revealed in the fs-TRF. On the other hand, comparing the ESA2 in fs-TA spectra at the longest time recorded with the corresponding ns-TA spectra for each Au(III) complex (Fig. 8, bottom panel), the two spectra are similar, indicating that ESA2 is derived from  $T_1 \rightarrow T_n$  absorption. Because there is a precursor-successor relationship between the ESA1 ( $S_1 \rightarrow S_n$  absorption) and ESA2 ( $T_1 \rightarrow T_n$  absorption),  $\tau_2$  of ESA1 is assigned to ISC from the  $S_1$  excited state to a receiving triplet excited state, which then internally converted to the  $T_1$  excited state with an ultrafast time scale. Thus,  $\tau_{\text{ISC}} = 13.2$  ps (**1b**), 8.38 ps (**2b**) and 3.49 ps (**3b**). The short  $\tau_1 = 0.80/1.28$  (**1b**), 0.95 ps (**2b**) and 0.63/0.74 ps (**3b**) of ESA1/TRF may likely correspond to the  $S_1$  vibrational relaxation.

### Intersystem crossing rate

The spectroscopically determined intersystem crossing rate constants ( $k_{\text{ISC}}$ ) and the corresponding time constants ( $\tau_{\text{ISC}}$ ) for both Au(I) and Au(III) complexes studied herein are tabulated in Table 4. For **1a**–**3a** and **5a**–**6a**, assuming that the major non-

Table 4  $k_{\text{ISC}}$  and  $\tau_{\text{ISC}}$  of the Au(I) and Au(III) complexes

Complex	$k_{\text{ISC}}^a/10^7 \text{ s}^{-1}$	$\tau_{\text{ISC}}^b/\text{ns}$	Complex	$k_{\text{ISC}}^b/10^{10} \text{ s}^{-1}$	$\tau_{\text{ISC}}^c/\text{ps}$
<b>1a</b>	<0.75	133	<b>1b</b>	7.57	13.2
<b>2a</b>	<11.1	9.0	<b>2b</b>	11.9	8.38
<b>3a</b>	<7.9	12.7	<b>3b</b>	28.7	3.49
<b>5a</b>	<0.91	110			
<b>6a</b>	<1.14	87.5			

<sup>a</sup>  $k_{\text{ISC}}$  for **1a**–**3a**, **5a**–**6a** are calculated according to eqn (1). <sup>b</sup>  $\tau_{\text{ISC}} = 1/k_{\text{ISC}}$ .

<sup>c</sup>  $\tau_{\text{ISC}}$  of **1b**–**3b** is estimated from  $\tau_2$  obtained from the fs-TA spectra of ESA1.

radiative decay of the  $S_1$  excited state is ISC, *i.e.*  $k_{nr} \approx k_{ISC}$ , an upper bound approximation of the  $k_{ISC}$  values could be obtained by eqn (1):

$$\Phi_{em} = \frac{k_r}{k_r + k_{ISC}} \quad (1)$$

The estimated  $k_{ISC}$  for the gold(I) complexes are  $7.5 \times 10^6$  to  $1.1 \times 10^8 \text{ s}^{-1}$  and the intersystem crossing time constants ( $\tau_{ISC}$ ) are 9.0–133 ns. These  $\tau_{ISC}$  are much larger than those of many phosphorescent transition-metal complexes ( $\tau_{ISC}$  in the femtosecond to picosecond timescale). For **1b**–**3b**, the  $\tau_{ISC}$  values are more than three orders of magnitude faster than their gold(I) analogues; these ISC rates, nevertheless, are comparable to other transition-metal complexes where  $S_1 \rightarrow T_1$  ISC is mediated by a higher-lying  $T_n$  triplet excited state.<sup>16,17</sup>

## Computational study

The different luminescence behaviors between the Au(I) and Au(III) systems were investigated by DFT/TDDFT calculations. The pair (**1a**, **1b**) was chosen as a representative example to examine why the Au(I) complexes studied herein display only fluorescence while the Au(III) counterparts exhibit exclusively phosphorescence. As the Bodipy-functionalized complexes give fluorescence for both Au(I) and Au(III) complexes, DFT/TDDFT calculations were also performed on the pair, (**4a**, **4b**). To save computational time, the cyclohexyl groups of the phosphine ligands in **1a** and **4a** were replaced by methyl groups.

### Calculations on **1a** and **1b**

The frontier MO diagrams of **1a** and **1b** are shown in Fig. 9. The HOMO and LUMO for both complexes **1a** and **1b** are predominantly localized on the arylacetylide ligand. For **1b**, a considerable contribution (17%) from the  $C^N^C$  moiety to the LUMO is also noted. The energy gap between HOMO and H-1 in **1a** is approximately 0.88 eV. For **1a**, the H-1 is comprised of the anti-bonding combinations of  $Au(d_{xy})$  and  $\pi(C\equiv C)$  orbitals with little

involvement of the phosphine ligand. For **1b**, H-1 is composed of  $Au(d_{xz})$  and the  $\pi(C^N^C)$  orbitals; the energy gap between HOMO and H-1 in **1b** is only 0.2 eV. The H-2 of **1b** is made up of an antibonding combination of the  $Au(d_{xy})$ ,  $\pi(C\equiv C)$  and  $\sigma(C^N^C)$  orbitals, with a HOMO/H-2 orbital energy gap of only  $\sim 0.4$  eV. Clearly, the cyclometalated  $[C^N^C]$  ligand has a role in destabilizing the  $Au(d)$  orbitals. Therefore, the HOMO and H-1/H-2 energy gaps in **1b** are much smaller than that in **1a**.

The energies of the singlet and triplet excited states and the associated nature and composition for **1a** and **1b** at their respective optimized singlet ground state geometries are obtained by TDDFT and are shown in Table S5 and S6 in ESI.† For **1a**, there is only one triplet excited state ( $T_1$ ) which is more than  $10\,000 \text{ cm}^{-1}$  below  $S_1$ . In addition, both  $S_1$  and  $T_1$  excited states are of the same parentage and are derived from  $\text{HOMO} \rightarrow \text{LUMO}$  transition ( $\sim 90\%$ ) and thus, there would be no effective SOC between them. The triplet excited states above  $S_1$  were also considered; the closest lying  $T_m$  excited state with efficient SOC is when  $m = 4$ , which is derived from a H-1 to LUMO transition (90%  $H-1 \rightarrow L$ ). However, the energy separation  $\Delta E(S_1-T_4)$  is  $-3180 \text{ cm}^{-1}$ , which is too large to be overcome by thermal activation.

On the other hand, for **1b**, there are four triplet excited states which are lower-lying than  $S_1$ , of which the closest-lying  $T_4$  excited state is only  $\sim 70 \text{ cm}^{-1}$  below the  $S_1$  excited state. Thus, thermal energy at room temperature assists facile ISC, even though SOC is small between the  $S_1$  and  $T_4$  excited states ( $|\langle S_1 | H_{SOC} | T_4 \rangle|^2 \sim 1.5 \text{ cm}^{-2}$ ). In addition, among the triplet excited states above  $S_1$ , there is a close-lying  $T_5$  excited state derived from the  $H-2 \rightarrow \text{LUMO}$  transition (79%) which lies only  $390 \text{ cm}^{-1}$  above the  $S_1$  excited state and ISC from  $S_1$  to  $T_5$  could be thermally activated. Besides, owing to the different orientations of the d-orbitals in HOMO and H-2, the  $S_1$  and  $T_5$  excited states could have effective SOC ( $|\langle S_1 | H_{SOC} | T_5 \rangle|^2 \sim 4.1 \times 10^3 \text{ cm}^{-2}$ ).

### Calculations on **4a** and **4b**

The Frontier MOs for **4a** and **4b** are shown in Fig. 10. Relative to **1a** and **1b**, the HOMO is destabilized and the LUMO is stabilized

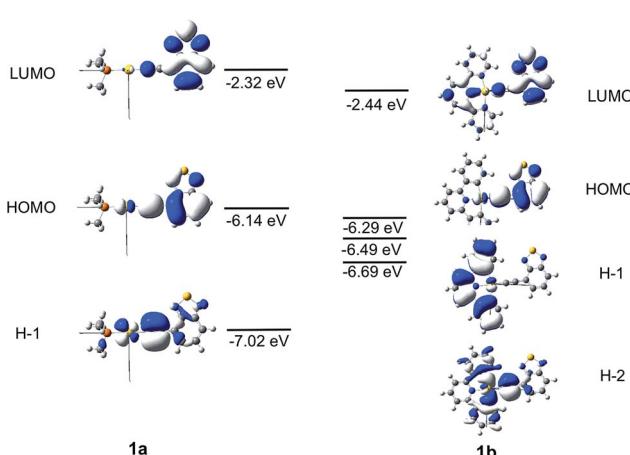


Fig. 9 Frontier MOs of **1a** and **1b** at the optimized  $S_0$  geometries. Orbital energies are also given in eV.

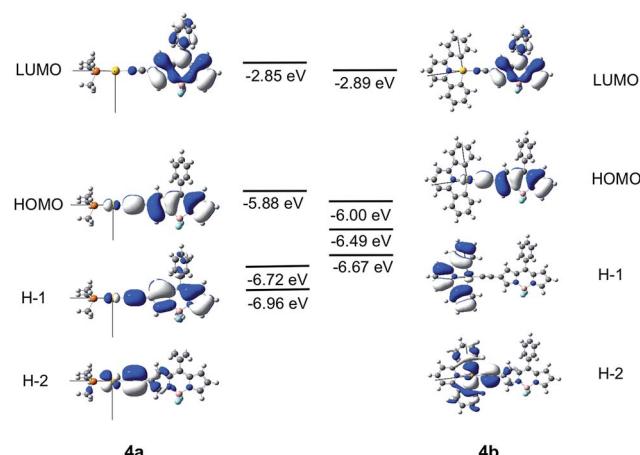


Fig. 10 Frontier MOs of **4a** and **4b** at their optimized  $S_0$  geometries. Orbital energies are also given in eV.



for **4a** and **4b**. Even for **4b**, which contains a cyclometalated  $[\text{C}^{\text{N}}\text{C}]$  ligand, the LUMO is predominantly localized on the Bodipy-functionalised arylacetylide ligand. The HOMO is composed of an antibonding combination of the  $\text{Au}(\text{d}_{yz})$  and  $\pi(\text{C}\equiv\text{CBodipy})$  orbitals. The energies and compositions of the singlet and triplet excited states of **4a** and **4b** at their respective optimized singlet ground state geometries were obtained by TDDFT and are collected in Tables S7 and S8 in ESI.† For this pair, (**4a**, **4b**), the two triplet excited states,  $\text{T}_1$  and  $\text{T}_2$ , are more than  $2000\text{ cm}^{-1}$  below the  $\text{S}_1$  excited state and are all composed of  $\text{Au}(\text{d}_{yz})$  orbitals. As SOC between the coupling singlet and triplet excited states would be ineffective with d-orbitals of the same orientation, ISC from  $\text{S}_1$  to  $\text{T}_2$  (or  $\text{T}_1$ ) for **4a** and **4b** would be sluggish.

For **4a**, the  $\text{T}_3$  excited state is the closest-lying triplet excited state that could have effective SOC with the  $\text{S}_1$  excited state due to a minor contribution of the  $\text{H}-5 \rightarrow \text{LUMO}$  transition to the  $\text{T}_3$  excited state ( $\text{H}-5$  is composed of the  $\text{Au}(\text{d}_z^2)$  orbital); however,  $\Delta E(\text{S}_1-\text{T}_3) = -2955\text{ cm}^{-1}$  (negative sign indicates that  $\text{T}_3$  lies above  $\text{S}_1$ ) which is much larger than the thermal energy. For **4b**, the  $\text{T}_3$  excited state is also the closest-lying triplet excited state that could have effective SOC with the  $\text{S}_1$  excited state due to a minor contribution of  $\text{H}-1 \rightarrow \text{LUMO}$  transition in the  $\text{T}_3$  excited state (the d-orbitals of the  $\text{Au}(\text{m})$  ion at the HOMO and  $\text{H}-1$  of **4b** are of different orientations, Fig. 10). However, the singlet-triplet gap,  $\Delta E(\text{S}_1-\text{T}_3) = -1192\text{ cm}^{-1}$ , is also much larger than the thermal energy. Thus, the pair (**4a**, **4b**) is expected to have slow ISC rates, when taking into consideration both the singlet-triplet energy gaps and SOC.

## Discussion

### General remarks on the photophysical properties

The emissions of the  $\text{Au}(\text{i})$  complexes, **1a–6a** are attributable to  $^1\text{IL} \pi\pi^*(\text{C}\equiv\text{CR})$  excited states. Pronounced red shifts in emission  $\lambda_{\text{max}}$  of arylacetylides can be observed upon their coordination to  $\text{Au}(\text{i})$  ion (e.g.  $\lambda_{\text{max}} = 412\text{ nm}$  (**L1**) vs.  $467\text{ nm}$  (**1a**)). Considering the complexes **1a**, **5a** and **6a**, which have different neutral auxiliary ligands (phosphine ( $\text{PCy}_3$ , **1a**), isocyanide ( $\text{RNC}$ , **5a**) and N-heterocyclic carbene ( $\text{NHC}$ , **6a**)) but the same acetylide ligand with a benzothiadiazole moiety, the lowest energy emission  $\lambda_{\text{em}}$  red shifts with the auxiliary ligand from  $456\text{ nm}$  ( $\text{RNC}$ ) to  $467\text{ nm}$  ( $\text{PCy}_3$ ) to  $476\text{ nm}$  ( $\text{NHC}$ ). A rationalization would be that the  $\text{NHC}$ , being the strongest  $\sigma$ -donor ligand among the auxiliary ligands in the three complexes, destabilizes the  $\text{Au}(\text{d})$  orbital to the greatest extent. From DFT calculations, the HOMO is comprised of a  $\text{Au}(\text{d})$  orbital and  $\pi(\text{C}\equiv\text{CR})$  (Fig. 9). Thus, the more electron-donating the auxiliary ligand, the more destabilized the HOMO, and hence, the smaller the HOMO-LUMO gap and the  $^1\pi\pi^*(\text{C}\equiv\text{CR})$  energy. A similar trend in the lowest energy absorption  $\lambda_{\text{abs}}$  can also be observed on changing the auxiliary ligand from  $\text{RNC}$  ( $371\text{ nm}$ ) to  $\text{PCy}_3$  ( $379\text{ nm}$ ) and  $\text{NHC}$  ( $383\text{ nm}$ ).

Most of the reported luminescent cyclometalated  $\text{Au}(\text{m})$  complexes display phosphorescence that comes from the  $^3\pi\pi^*$  IL excited state localized on the cyclometalated ligands.<sup>13</sup> For the  $\text{Au}(\text{m})$  complexes studied herein, **1b–3b**, the lowest-

energy triplet excited states are of  $^3\pi\pi^*(\text{C}\equiv\text{CR})$  in nature, with  $\lambda_{\text{max}} = 630$ ,  $592$  and  $603\text{ nm}$  respectively. These complexes, however, exhibit rather weak phosphorescence, with  $\Phi_{\text{em}}$  values in the range of  $0.003$ – $0.04$ . It is noted that in  $77\text{ K}$  glassy solutions, the phosphorescence lifetimes are significantly increased compared with those obtained in degassed  $\text{CH}_2\text{Cl}_2$  at  $\text{RT}$  (e.g.  $205\text{ }\mu\text{s}$  at  $\text{RT}$  to  $2.2\text{ ms}$  at  $77\text{ K}$  for **3b**). Since low temperature and rigid glassy matrix can impede structural distortion, the lifetimes obtained at  $77\text{ K}$  could reflect the intrinsic radiative lifetime of the complexes. The especially long emission lifetimes can reflect the predominant localization of the emitting  $\text{T}_1$  excited state on the arylacetylide ligand, *i.e.*  $^3\pi\pi^*(\text{C}\equiv\text{CR})$  with little participation of the metal ion. This is also corroborated by the small thermally induced Stokes shifts ( $\Delta E_{\text{S}} = E_{00}(77\text{ K}) - E_{00}(298\text{ K})$ ) of less than  $600\text{ cm}^{-1}$  (Table 2).

### Intersystem crossing in $\text{gold}(\text{i})$ and $\text{gold}(\text{m})$ complexes

ISC is usually fast in transition-metal complexes with time constants ( $\tau_{\text{ISC}}$ ) in the fs to ps time regimes. In the literature, there are numerous examples which show ultrafast ISC,<sup>1,18–25</sup> *e.g.*  $[\text{M}(\text{bpy})_3]^{2+}$  ( $\text{M} = \text{Ru}$  or  $\text{Fe}$ ,  $\tau_{\text{ISC}} = 30\text{ fs}$ ),<sup>19a,b</sup>  $[\text{Re}(\text{L})(\text{CO})_3(\text{bpy})]$  ( $\tau_{\text{ISC}} = 100$ – $140\text{ fs}$ ),<sup>20</sup>  $[\text{Ir}(\text{piq})_3]$  (piq: 1-phenylisoquinoline;  $\tau_{\text{ISC}} = 70\text{ fs}$ ),<sup>21b</sup>  $[\text{Pt}(\text{PBu}_3)_2(\text{C}\equiv\text{CPh})_2]$  ( $\tau_{\text{ISC}} = 70\text{ fs}$ ),<sup>22</sup> and  $[\text{Cy}_3\text{PAu}(2\text{-naphthyl})]$  ( $\tau_{\text{ISC}} = 230\text{ fs}$ )<sup>23</sup> *etc.* These  $\tau_{\text{ISC}}$  correspond to rates of intersystem crossing ( $k_{\text{ISC}}$ ) in the range of  $10^{12}$  to  $10^{13}\text{ s}^{-1}$ . The fast  $k_{\text{ISC}}$  in transition-metal complexes is traditionally attributed to a large spin-orbit coupling (SOC) constant inherited from the heavy metal atom. However, there are increasing number of reports revealing slow ISC rate ( $k_{\text{ISC}} \sim 10^8\text{ s}^{-1}$ ) in spite of the presence of heavy transition metal, such as the cases of  $\text{Rh}(\text{i})$ - and  $\text{Ir}(\text{m})$ -bis(arylethynyl)cyclopentadiene,<sup>3</sup>  $\text{Au}(\text{i})$ -pyrene,<sup>4b</sup>  $\text{Pt}(\text{u})$ -perylene/<sup>4c</sup> tetracene,<sup>4e</sup> and  $\text{Pd}(\text{u})$ -perylene diimide;<sup>4f</sup> all these complexes contain highly conjugated ligand systems and display ligand-dominated  $^1\pi\pi^*$  fluorescence. There are also cases where comparable  $k_{\text{ISC}}$  and  $k_{\text{r}}$  of  $\text{S}_1 \rightarrow \text{S}_0$  leads to the observation of dual fluorescence–phosphorescence under steady-state condition, *e.g.*,  $[\text{Pt}(\text{L})(\text{acac})]$  and  $[\text{Ir}(\text{L})(\text{acac})]$  ( $\text{L} = 2\text{-}( \text{oligothienyl})\text{pyridine}$ ),<sup>26</sup>  $[\text{Os}(\text{L})(\text{CO})_3\text{X}]$  ( $\text{L} = 8\text{-quinolino} \text{-} \text{olate}$ <sup>27</sup> or isoquinoline-triazole),<sup>28</sup> and  $[\text{Bu}_4\text{N}]_4[\text{Pt}_2(\mu\text{-P}_2\text{O}_5(\text{BF}_2)_2)]$ ,<sup>29</sup> *etc.*

In this work, the luminescence behaviour of the  $\text{Au}(\text{i})$  and  $\text{Au}(\text{m})$  complexes are drastically different, even though they have the same metal and arylacetylide ligands. Ligand-dominated fluorescence has been observed with the  $\text{Au}(\text{i})$  complexes, **1a–3a** and **5a–6a**, with  $k_{\text{ISC}}$  estimated to range from  $7.5 \times 10^6$  to  $1.1 \times 10^8\text{ s}^{-1}$ . The  $\text{Au}(\text{m})$  complexes **1b–3b**, on the other hand, display phosphorescence, with  $k_{\text{ISC}}$  estimated to be larger than  $10^{10}\text{ s}^{-1}$ . The major difference between the two series of gold complexes is the oxidation state of  $\text{Au}$  ion, that dictates the coordination geometry, *i.e.* a linear geometry for the  $\text{Au}(\text{i})$  complexes, **1a–6a**, and a square-planar geometry for  $\text{Au}(\text{m})$  complexes, **1b–4b**. The coordination geometry has a significant impact on the relative energies of the frontier orbitals (specifically, the d-orbital energies) and hence the relative energies of the singlet and triplet excited states, which subsequently affect the  $k_{\text{ISC}}$ .



The two factors that determine the  $k_{ISC}$  are (1) the SOC matrix element  $\langle S_n | H_{SOC} | T_m \rangle$ , and (2) the energy gap ( $\Delta E_{ST}$ ) between the coupling singlet ( $S_n$ ) and triplet ( $T_m$ ) excited states. The larger the  $H_{SOC}$  and the smaller the energy gap ( $\Delta E_{ST}$ ), the faster will be  $k_{ISC}$ . For effective SOC, this requires the metal d-orbitals of the coupling singlet and triplet excited states to have different orientations. For example, if  $S_n$  is derived from a metal-to-ligand charge transfer (MLCT) excited state where  $\text{Au}(\text{d}_{xz})$  orbital is involved,  $H_{SOC}$  would be zero if the triplet excited state is also an MLCT state that involves  $\text{Au}(\text{d}_{xz})$  orbital because of symmetry reasons.

The pair (**1a**, **1b**) has been chosen as a representative example to illustrate the different photophysical properties exhibited by the  $\text{Au}(\text{i})$  and  $\text{Au}(\text{m})$  arylacetylides studied in this work. From the DFT/TDDFT calculations, it is revealed that owing to the inherent linear coordination geometry of the  $\text{Au}(\text{i})$  complex, the d-orbitals of the gold(i) ion is mainly destabilized by the arylacetylides ligand (Fig. 9 and Table S9 in ESI<sup>†</sup>). On the other hand, as  $\text{Au}(\text{m})$  complexes are assumed to have a square-planar four-coordinated geometry, thus, in addition to the antibonding interactions with the arylacetylides ligand, the d-orbitals of gold(m) ion could also be destabilized by the cyclometalated [ $\text{C}^{\text{N}}\text{C}$ ] ligand (both  $\pi$ -type, *e.g.* H-1, and  $\sigma$ -type, *e.g.* H-2 in **1b**; Fig. 9 and Table S9 in ESI<sup>†</sup>); these latter interactions result in smaller d-orbital splittings in the  $\text{Au}(\text{m})$  series than the  $\text{Au}(\text{i})$  series. In effect,  $S_1$  and  $S_2$  excited states are  $\sim 4200 \text{ cm}^{-1}$  apart for **1a** while the analogous splitting (between  $S_1$  and  $S_3$  excited states) is only  $\sim 1300 \text{ cm}^{-1}$  for **1b**. As  $S_2$  of **1a** is derived from  ${}^1[\text{Au}(\text{d}_{xy}) \rightarrow \pi^*(\text{C} \equiv \text{CR})]/[{}^1[\pi(\text{C} \equiv \text{C}) \rightarrow \pi^*(\text{C} \equiv \text{CR})]]$  ( ${}^1\text{MLCT}/{}^1\text{ILCT}$ ) and  $S_3$  of **1b** from  ${}^1[\text{Au}(\text{d}_{xy}) \rightarrow \pi^*(\text{C} \equiv \text{CR})]/[{}^1[\pi(\text{C} \equiv \text{C}) \rightarrow \pi^*(\text{C} \equiv \text{CR})]]/{}^1[\pi(\text{C}^{\text{N}}\text{C}) \rightarrow \pi^*(\text{C} \equiv \text{CR})]$  ( ${}^1\text{MLCT}/{}^1\text{ILCT}/{}^1\text{LLCT}$ ), *i.e.*, both are of charge-transfer type excited states, the singlet-triplet energy gaps for this type of transitions are small ( $\Delta E(S_2 - T_4) \sim 1000 \text{ cm}^{-1}$  for **1a** and  $\Delta E(S_3 - T_5) \sim 900 \text{ cm}^{-1}$  for **1b**) ( $T_4$  (**1a**) and  $T_5$  (**1b**) are the triplet counterpart of  $S_2$  (**1a**) and  $S_3$  (**1b**) respectively). As depicted in Fig. 11, the  $S_1/T_5$  energy gap for **1b** is small but the  $S_1/T_4$  energy gap for **1a** is large. In other words, the oxidation state of the gold ion affects the coordination geometry of the complex, which in turn change the interactions between the metal d-orbitals and ligand orbitals, giving rise to different d-orbital splitting and subsequently the singlet-triplet splitting ( $\Delta E_{ST}$ ) of the two coupling excited states in the gold complexes.

Moreover, DFT/TDDFT calculations also revealed that there is a triplet excited state ( $T_4$ ) almost isoenergetic with the  $S_1$  excited state ( $< 70 \text{ cm}^{-1}$  below the  $S_1$  excited state) in **1b** such that even though the SOC between  $S_1$  and  $T_4$  is small due to the similar d-orbital orientations involved in both excited states, thermal energy could promote facile ISC. With **1a**, the closest triplet excited state ( $T_2$ ) to the  $S_1$  excited state is more than  $500 \text{ cm}^{-1}$  above the  $S_1$  excited state, which is more than twice the thermal energy at room temperature and SOC is also small between these two excited states as the d-orbitals involved are also of the same orientations. Thus, taken together both the SOC and  $\Delta E_{ST}$ , **1b** should have a much faster  $k_{ISC}$  than **1a**.

On the other hand, for the Bodipy-functionalized complexes, **4a** and **4b**, only  ${}^1\pi\pi^*(\text{C} \equiv \text{C}\text{Bodipy})$  fluorescence

with no long-lived species are observed under ns-TRE and ns-TA measurements. The photophysical behavior of the Bodipy-functionalized complexes can be attributed to the intrinsically small band-gap of the Bodipy moiety. Due to the highly conjugated structure of Bodipy, the HOMO is much destabilized and there is a wide orbital energy gap between the HOMO and other occupied MOs, even in the case of **4b** which contains a [ $\text{C}^{\text{N}}\text{C}$ ] ligand. As a result, the HOMO/H- $x$  orbital energy gap is the largest among the four arylacetylides studied herein (H- $x$  is the other occupied orbitals lower in energy than the HOMO;  $x = 1, 2, \dots$ ). In effect, the closest  $T_m$  excited state that could have effective SOC with  $S_1$  is more than  $1000 \text{ cm}^{-1}$  above the  $S_1$  excited state. With such a large  $\Delta E(S_1 - T_m)$ , thermal energy would be insufficient to promote ISC. Therefore, similar to the scenario in the case of **1a** (Fig. 11, left), ISC is sluggish for  $\text{Au}(\text{i})$  and  $\text{Au}(\text{m})$  arylacetylides bearing Bodipy.

### Mechanism for the generation of delayed fluorescence

From the ns-TRE measurements of **1a**–**3a**, DF contributes to the total fluorescence, though only a minute proportion (<3%, Table 3). In general, the mechanism of DF could be inferred from the dependence of the DF intensity ( $I_{DF}$ ) with the power of excitation light.<sup>30</sup> According to Bässler, a quadratic dependence of excitation power with the DF intensity indicates that the mechanism of the DF is TTA with dominant phosphorescence.<sup>30b</sup> On the other hand, a linear dependence of DF intensity with excitation power could be due to three possible mechanisms: TTA with dominant delayed-fluorescence, TADF, and GP-recombination. As depicted in Fig. 12, the plot of  $I_{DF}$  against excitation intensity in double-logarithm scale gave a slope of  $1.71 \approx 2$  for **1a**; this nearly quadratic dependence is most consistent with the TTA mechanism with dominant phosphorescence. However, for **2a** and **3a** (slope = 0.916 and 1.10), both display nearly linear dependence between  $I_{DF}$  and excitation intensity. Therefore, it is not possible to confirm the mechanism for DF in the case of **2a** and **3a** by solely considering the excitation power dependence.

Time-dependence of  $I_{DF}$  and phosphorescence intensity ( $I_P$ ) could also give hints to the DF mechanism.<sup>30,31</sup> For TTA with dominant DF, phosphorescence intensity decays with a power law,  $I_P \propto t^{-1}$  while  $I_{DF}$  is approximately constant at short time and  $I_{DF} \propto t^{-2}$  at longer time. For the GP-recombination mechanism, both DF and phosphorescence decay in accordance with the power law,  $I_{DF} \propto t^{-1}$ , at both short and long times.<sup>30a</sup> In the case of **2a** in 77 K glassy solution, both DF and phosphorescence decayed according to the power law:  $I \propto t^{-1}$  over the time intervals investigated (1  $\mu\text{s}$  to 1.2 ms) (Fig. 6, inset), suggesting that the DF mechanism under this condition is most likely the GP-recombination mechanism. As for **3a**, there is no power law decay relation with both DF and phosphorescence and so it seems unlikely that GP-recombination is the mechanism for the generation of DF in **3a**. There is still not enough information to conclude on the DF mechanism for **3a**.



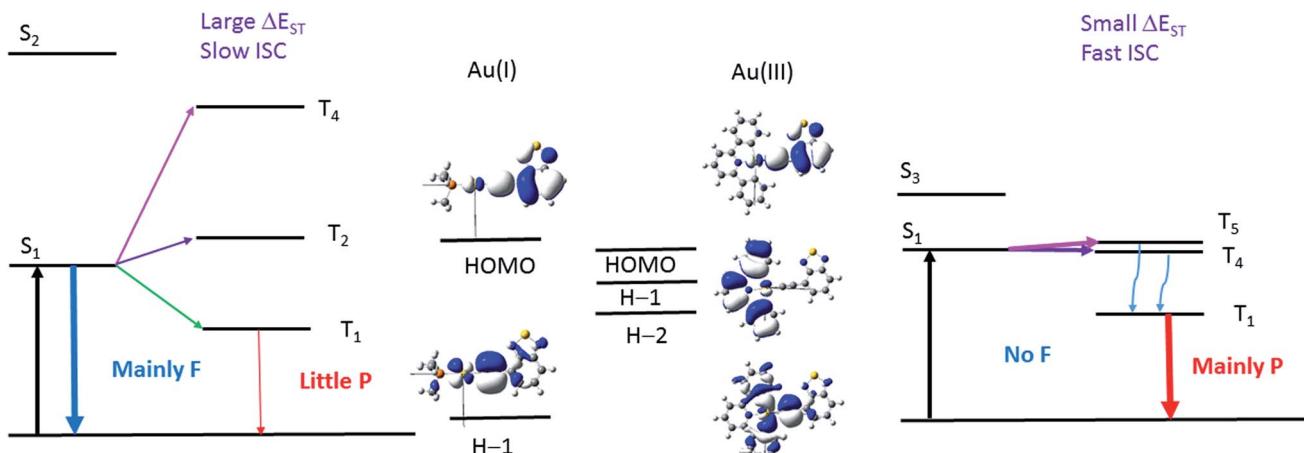


Fig. 11 Illustration of the low-lying singlet and triplet excited states of Au(I) (left) and Au(III) (right) complexes that accounts for the different photophysical behaviour of the Au(I) and Au(III) complexes investigated in this work.  $S_1$  and  $T_1$  for both complexes are derived from HOMO  $\rightarrow$  LUMO transitions;  $S_2$  and  $T_4$  excited states of **1a** are derived from  $^{1,3}[\text{H}-1 \rightarrow \text{LUMO}]$  transitions while  $S_3$  and  $T_5$  excited states of **1b** are derived from  $^{1,3}[\text{H}-2 \rightarrow \text{LUMO}]$  transitions. The d-orbitals involved in the  $T_2$  of **1a** and  $T_4$  of **1b** have the same orientations as their respective  $S_1$  excited state (see Tables S5, S6 and S9 in ESI†). The wavy blue arrows indicate internal conversion (IC) from the  $T_5$  to  $T_1$  excited state. F = fluorescence and P = phosphorescence.

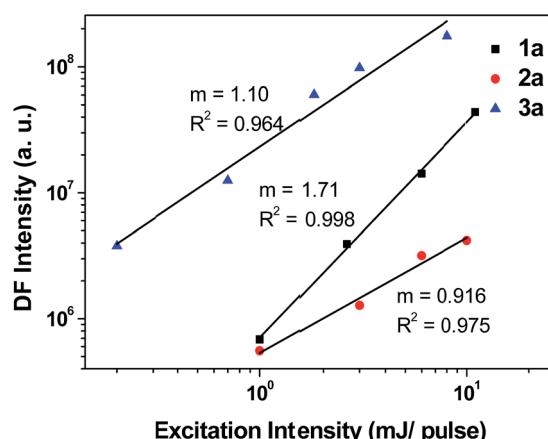


Fig. 12 Dependence of delayed fluorescence of **1a**–**3a** in  $5 \times 10^{-5}$  M degassed  $\text{CH}_2\text{Cl}_2$  with excitation power intensity. Emission intensity measured after a time delay of 1  $\mu\text{s}$ . (Laser  $\lambda_{\text{exc}} = 355$  nm; 0.2–11 mJ per pulse; diameter = 8 mm, integration time: 800  $\mu\text{s}$ .)

## Conclusion

A series of gold complexes with different oxidation states, gold(I) complexes  $[\text{LAu}(\text{C}\equiv\text{CR})]$  and gold(III) complexes  $[\text{Au}(\text{C}^{\text{N}}\text{N}^{\text{C}})(\text{C}\equiv\text{CR})]$  bearing the same heterocyclic arylacetylide with narrow bandgap were synthesized and characterized. The photophysical behaviors with the gold ion in different oxidation states are strikingly different: fluorescence dominates the luminescence of the Au(I) complexes while phosphorescence takes over in the Au(III) complexes. Detailed computational studies by DFT/TDDFT have accounted for these phenomena as a result of different coordination environments inherited from the gold ion in a particular oxidation state: a linear coordination geometry for Au(I) and a square-planar coordination geometry for Au(III). This difference in

coordination geometry subtly affects the energy separation between the coupling singlet and triplet excited states, leading to smaller  $\Delta E_{\text{ST}}$  of the Au(III) complexes than the Au(I) complexes and hence, larger  $k_{\text{ISC}}$  in the Au(III) complexes than the Au(I) complexes. For the complexes bearing Bodipy-functionalized acetylide ligand, they only display prompt fluorescence. Computational analyses revealed that, due to the especially narrow bandgap of Bodipy, the  $\Delta E_{\text{ST}}$  is still large even in the Au(III) complex so that  $k_{\text{ISC}}$  could not compete with fluorescence radiative decay. Additionally, the mechanisms for the generation of DF in Au(I) complexes have been explored. To the best of our knowledge, this is the first report which systematically studies the effects of the metal ion oxidation state on the photophysical behaviours of transition-metal complexes.

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