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Introduction

The metal carbide family is a large one containing several groups, including the materials based on extended 2D carbon

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Pressure induced polymerization of acetylide anions in CaC_2 and 10^7 fold enhancement of electrical conductivity[†]

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Transformation between different types of carbon–carbon bonding in carbides often results in a dramatic change of physical and chemical properties. Under external pressure, unsaturated carbon atoms form new covalent bonds regardless of the electrostatic repulsion. It was predicted that calcium acetylide (also known as calcium carbide, CaC₂) polymerizes to form calcium polyacetylide, calcium polyacenide and calcium graphenide under high pressure. In this work, the phase transitions of CaC₂ under external pressure were systematically investigated, and the amorphous phase was studied in detail for the first time. Polycarbide anions like C_6^{6-} are identified with gas chromatography-mass spectrometry and several other techniques, which evidences the pressure induced polymerization of the acetylide anions and suggests the existence of the polyacenide fragment. Additionally, the process of polymerization is accompanied with a 10^7 fold enhancement of the electrical conductivity. The polymerization of acetylide anions demonstrates that high pressure compression is a viable route to synthesize novel metal polycarbides and materials with extended carbon networks, while shedding light on the synthesis of more complicated metal organics.

structures, like the graphene-based Li-battery anode material LiC_6 , and 0D ionic compounds like the C_{60} -based superconductor K_3C_{60} , acetylide-based salts CaC_2 and Mg_2C , as well as other groups. Their numerous intriguing physical and chemical properties are related to their structural peculiarities, which benefit from the rich chemistry of carbon. Compressing carbides to tens of gigapascals will facilitate the bonding between the isolated carbon groups, which will hence change the electrical properties significantly. For example, BeC₂ and MgC₂ are predicted to contain five-membered carbon rings,¹ and Li_2C_2 and BaC₂ are predicted to have polyanionic structures,^{2,3} but only a few of these topics have been investigated experimentally.³

CaC₂ is the most important and common metal carbide used in industry. It is part of a large group of metal carbides that are composed of M^{2+} and the dumbbell-shaped C_2^{2-} anion. The polymorphs of CaC₂ have been studied for almost a century,⁴⁻⁶ and four crystalline phases have been identified at ambient pressure, named CaC₂-I (space group (SG) *I4/ mmm*); CaC₂-II (SG *C2/c*); CaC₂-III (SG *C2/m*); and CaC₂-IV (SG *Fm*3*m*). All of these phases are composed of Ca²⁺ and C₂²⁻ in various configurations. Recently, several new polymorphs of CaC₂ stable under external pressure were predicted by theoretical researchers, with the C₂²⁻ dumbbell-shaped anions connecting to one another to form carbon chains, graphene



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ribbons or sheets, which are actually Ca-graphenide, Ca-polyacenide and Ca-polyacetylide.7-9 These polymorphs are expected to be metallic, and superconductive at low temperature.7 However, up to now, experimental investigations of the high pressure transformation of CaC₂ have been subject to controversy. Two phase transitions were reported, with the first at 10-12 GPa, followed by an amorphization at a pressure above ~18 GPa.^{10,11} Very recently, another report suggested the existence of the predicted *Cmcm* phase (Ca polyacetylide) above 7 GPa, and confirmed the amorphization upon further compression again.¹² The amorphization hindered almost all further investigations. In this work, we experimentally probed both the crystallized phases before the amorphization and the amorphous phase using several cutting-edge techniques, and evidenced the existence of linear and cyclic polycarbide anions in the product of the pressure induced polymerization of CaC₂.

Results and discussion

Variation of crystal structure

To investigate the phase transitions under high pressure, multiple in situ techniques including high pressure X-ray diffraction (XRD), neutron diffraction and Raman spectroscopy were employed. The lattice parameters of CaC2-I under high pressure were determined by in situ XRD. In similarity to previous reports,^{10,11} CaC₂ starts to amorphize above 18-20 GPa, above which the XRD peaks start to broaden and become indistinguishable. The unit cell parameters of CaC2 under external pressure were determined by Rietveld refinement (Fig. 1). The reported phase transition at 10-12 GPa is too subtle to be confirmed in the plot of the pressure dependent lattice parameters, and the lattice parameter of the whole pressure range before amorphization can be fitted to a 3rd order Birch-Murnaghan (B–M) Equation of State (EOS), with $V_0 = 96.7(2) \text{ Å}^3$, $B_0 = 62(2)$ GPa, and $B_1 = 3.6(2)$.



Fig. 1 Lattice parameters of CaC₂ under high pressure. The red and blue markers including error bars are for the a and c-axis respectively, and the lines are a guide for the eyes. The black dots are the experimental data of the unit cell volume, and the black solid line is the fitting of the BM-EOS.

To understand the structural details of CaC2 under high pressure, in situ neutron diffraction was carried out (Fig. 2a), which is more sensitive to the positions of carbon atoms. The diffraction patterns below 10 GPa can be well fitted using the CaC₂-I structural model. The refined results are shown in Table S1[†] and a selected plot of Rietveld refinement is shown in Fig. S1.[†] However, above 10–12 GPa, the *hkl* peaks with $l \neq 0$ are significantly broadened, the C=C bond length decreases significantly to unreasonable values and the nearest C···C intergroup distances increase correspondingly (Fig. 2b). This indicates that the CaC2-I model is not suitable for the data above 10-12 GPa, though the averaged structure does not deviate from CaC2-I dramatically. The reason for peak broadening comes from strain or disordering, most likely in the substructure of C instead of Ca, because the broadening is more pronounced in neutron diffraction.

The anomaly at 10–12 GPa was also found in the impedance spectrum, which will be discussed later. This can probably be attributed to the instability and interruption of the CaC2-I lattice, because CaC₂-VI is more stable at this pressure, as predicted by theoretical investigations.7,8,10 Our Density Functional Theory (DFT) calculations also show that CaC₂-VI is more stable than CaC₂-I above 9 GPa (Table S2[†]). As shown in Fig. 3, CaC₂-VI (Fig. 3b) is a distortion of CaC₂-I (Fig. 3a), with β deviating from 90° and its space group I2/m being a subgroup of I4/mmm. Because no characteristic peak of CaC2-VI can be identified under current experimental conditions, it seems the I-VI phase transition did not really go through, and the sample still stays in an intermediate state, like that shown in Fig. 3c. In this state, CaC₂-I is unstable; the phase transition starts locally (local disordering), but the new phase is not crystallized and cannot be identified.

Spectroscopic investigation

To probe the functional groups, Raman spectroscopy was employed to detect the transition of the carbon-carbon bonds. In similarity to the reported results,¹⁰ both the C-C stretching and C₂ libration modes show minor alterations at 10-12 GPa (Fig. 4a), which corresponds to the transition uncovered by the neutron diffraction experiment. Above 18-20 GPa, the Raman peaks start to degrade and disappear into the background, corresponding to the amorphization (Fig. S2[†]). After being maintained at \sim 30 GPa for 2 days, a minor peak at 1839 cm⁻¹ was observed, which is significantly lower than the $C \equiv C$ stretching of CaC_2 before its disappearance (1916 cm⁻¹, Fig. 4b). When decompressed, the peak becomes more obvious and splits to two peaks at 1690 and 1850 cm^{-1} (3.9 GPa) respectively. Such a peak should be attributed to an intermediate (or conjugated) state between a triple bond and a double bond, because its Raman shift is located between typical values for those bonds. The peak at 1850 cm⁻¹ is still around the region of triple bonds, while the peak at 1690 cm^{-1} is approaching the region of double bond. This splitting provides further evidence that the peak at 1839 cm^{-1} (before splitting) is from polymerized (oligomerized) carbon anions created by high



Fig. 2 In situ neutron diffraction patterns of CaC_2 under different values of pressure. (a) Neutron diffraction patterns of CaC_2 under high pressure. The peaks of diamond anvils and CaO are marked by d and c respectively. (b) C–C bond length and the nearest C···C distance between neighbouring C_2^{2-} units under different values of external pressure.

pressure. Its decomposition at low pressure results in isolated triple bonds and other complex containing double bonds.

In situ infrared spectra were also measured to detect the vibration of carbon species (Fig. 5). The C–C stretching of the acetylide anion is symmetrical and is hence infrared inactive. This is why no absorption peak was observed at low pressure (up to ~21 GPa). Above 22 GPa, absorption peaks around 1750 cm⁻¹ and 1200 cm⁻¹ are observed. These peaks are in the range of the stretching modes of C=C double bonds and C–C single bonds. The bonds have to be asymmetric as they are infrared active, which means the carbon atom has at least two neighbours covalently bonded, and hence evidences the bonding between $C_2^{2^-}$.

Identification of polycarbide anions in the recovered sample

From the crystallographic and spectroscopic analysis, we conclude that the polymerization of acetylide anions happens in the amorphous state. The polycarbide anions would hence hide in the amorphous phase. Besides the spectroscopic investigations, neutron pair distribution function (PDF) analysis is employed to investigate the local structure of the amorphous phase, which gives the likelihood of finding an atom at a given radius from another atom at an arbitrary origin.¹³ Obtaining good *in situ* PDF data above 20 GPa for CaC₂ is very challenging, if not impossible, therefore samples recovered from external pressure were measured (Fig. S3a†). No other periodic phase was identified from the recovered samples

except CaC₂ (phase I, starting material). However, a distinct difference was found at the C=C triple bond distance ($d \sim 1.24$ Å). The peak of the sample recovered from ~28 GPa is shifted to a greater radial distance and shows evidence for enhanced density between 1.3 Å and 1.5 Å when compared to that recovered from below 24 GPa (Fig. 6, more patterns in Fig. S3†). This verifies the presence of longer carbon–carbon bonds (C=C bonds) in the sample recovered from the highest pressure, which would be the product of the addition reaction between C₂²⁻ ions.

More solid evidence comes from the Gas Chromatography-Mass Spectrometry (GC-MS) analysis. The hydrolysis of the polymerized C_m^{x-} (here, *x* is supposed to be equal to *m*) anions in the recovered sample will produce $C_{2n}H_{2n}$ ($Ca_nC_{2n} + 2nH_2O = nCa(OH)_2 + C_{2n}H_{2n}\uparrow$). In the hydrolyzed products, several tens of hydrocarbons were identified by GC-MS, which were not detected in the raw material (Fig. 7). It is worthy to point out that even with sub-microgram amounts of sample synthesized using a diamond anvil cell (DAC), similar results were obtained to those with a milligram amount of sample synthesized using a Paris-Edinburgh (PE) cell. This approach demonstrates that GC-MS has a significant application in the characterization of chemical reactions under extremely high pressure.

With the data obtained from the high-resolution quadrupole-time- of -flight-mass spectrometer (QTOF-MS), the molecular formulas of most peaks in the total ion chromatograms (TIC) can be determined unambiguously, as listed in Table S3.† These include C_3H_4 , C_5H_6 , C_5H_4 , C_6H_8 , C_6H_6 , C_6H_4 , and C_8H_7 in



Fig. 3 Structures of CaC_2 -I (a) and VI (b) viewed along [010]; (c) a schematic model of possible local disordering of CaC_2 above 12 GPa. The black spheres are for carbon atoms and the blue spheres are for calcium atoms.



Fig. 4 (a) Pressure dependence of Raman shifts of CaC_2 and (b) selected Raman spectra of CaC_2 upon compression and decompression.



Fig. 5 Infrared absorption spectra of CaC₂ under external pressure.

the gas phase, with various numbers of isomers. More complicated molecules like $C_{12}H_{12}$, $C_{12}H_{10}$ and $C_{12}H_{14}$ can be identified in the liquid phase (Fig. S4 and Table S4†), and their diversities will be discussed in a following paper. All the identified molecules have a C : H ratio around 1 : 1, in consistency with the valence of carbon in CaC₂. The slight deviations from 1 : 1 probably result from the heterolytic cleavage of some unstable molecules or anions after or during the hydrolysis. The molecules with odd numbers of carbon atoms like C_3H_4 and C_5H_4 also likely result from cleavage, because the combination of C_2^{2-} anions can only result in molecules with even numbers of carbon atoms. Hydrocarbons like C_2H_6 , C_2H_4 , C_4H_8 and those with a C : H molar ratio severely deviating from 1 : 1 are not detected, which excludes the existence of C_2^{6-} , C_2^{4-} , C_4^{8-} and the corresponding calcium carbide. This indicates that the non-ox/red polymerization (or oligomerization) instead of disproportionation (or other ox/red reactions) dominates the reaction process, and the Ca : C molar ratio does not really change during the reaction.

By checking the corresponding mass spectrum in the National Institute of Standards and Technology (NIST) library, the peaks in the TIC are recognized. The list of possible molecules includes both linear and cyclic molecules, which result from the hydrolysis of linear and cyclic polycarbide anions respectively. By comparing the retention times, benzene can be identified unambiguously (Fig. S5 and S6[†]) among the peaks, which indicates the formation of C_6^{6-} cyclic structures in the recovered CaC2. It was predicted that carbon atoms tend to polymerize following the sequence of chain, belt and sheet, with increasing pressure. Both of the latter two structures (Ca polyacenide and Ca graphenide) contain six-membered rings.7 Our experiment evidenced that under the current experimental conditions it is possible for C_2^{2-} to polymerize to polyacenide (Immm phase in ref. 7) or its fragment, which was predicted to be stabilized above \sim 15 GPa. The chain structures can also exist due to incomplete reactions.

Calibrated by the working curve, the molar ratio of benzene to acetylene in the hydrolysis product is ~0.005 : 1 (Fig. S7†). Even if supposing other produced molecules respond at the same efficiency as benzene in MS (usually lower, thus the actual concentration is higher), the molar ratio of product to acetylene is 0.11 : 1. This indicates that a significant amount of C_2^{2-} is reacted under the current experimental conditions. The details of the quantitative analysis are discussed in the ESI.†

Meta-dynamic simulation

To build up a model to understand what happens in CaC_2 under external pressure, meta-dynamic simulations were performed at 30 GPa and two models were obtained (referred as to chain and ribbon models, shown in Fig. 8). In the chain structure, the acetylide anions are connected to form polyacetylide chains. Both cis- and trans- connections are formed, and branched chains are also obtained. Starting from this structure, a meta-dynamic simulation was conducted again and the ribbon structure was obtained, which is constructed by edge-sharing six-membered rings. These two models show the structural features of the sample under pressure and are consistent with the theoretical predictions in ref. 7, though with some defects. These two models are closely related. In the chain model, six-membered rings are identified and in the ring model, chains still exist. The co-existence of ring and chains is most likely the real case, as evidenced by the GC-MS experiment. It is worthy to note that the simulation is carried out using a box containing 64 Ca and 128 C atoms, constrained by



Fig. 6 Neutron pair distribution functions (G(r)) of CaC₂ samples recovered from external pressures. (a) Selected normalized first peak after background subtraction. This peak corresponds to the carbon–carbon bond. (b) Differences between the normalized G(r) patterns of the samples recovered from high pressure (24 GPa, 28 GPa) and that of the raw material (0 GPa).

periodic boundary conditions. It can be anticipated that in the experiment the local structure will be more diverse, and it would be difficult for the produced covalent carbon network to crystallize.

This diversity uncovered by the simulation and experiment actually traced the whole process of polymerization, from small species to big ones, from chains to rings and ribbons. This is most likely what happened in the amorphous phase. The monomers polymerize with various degrees of polymerization and even in various dimensions. If the selectivity of the reaction is enhanced by controlling reaction conditions, more complex and kinetically stable metal carbides on milligram scales can be obtained besides the predicted phases.⁷ Because the calcium polycarbide obtained is nucleophilic and highly reactive, more unexpected compounds can be synthesized through reaction with acid or *via* other nucleophilic reactions. The neutral molecules can even be separated by chromatography and purified.



Fig. 7 Total ion chromatograms (TIC) of the product of CaC_2 recovered from 26 GPa and CaC_2 raw material reacting with water. (a) Sample synthesized by PE cell. (b) Sample synthesized by DAC. (c) CaC_2 raw material (before compression).



Fig. 8 Simulated structure of CaC_2 at 30 GPa by meta-dynamics. (a) Chain model. (b) Ribbon model. The C–C bond limitation is set at 1.6 Å. Ca ions are omitted for clarity.

Decreased resistivity

One of the most attractive properties of the metal carbide materials is their electrical property, which is closely related to its structure.7 In this work, the pressure dependent resistivity of CaC₂ was measured. From \sim 2 GPa to \sim 22 GPa, the resistivity decreases by seven orders of magnitude (Fig. 9), with a turning point at the minor transition around 10-12 GPa. From ~ 10 GPa to 22 GPa, the resistivity decreases by 1 order of magnitude, while below 10 GPa, it decreases by six orders. This significant decrease of resistivity (R) should be attributed to the decrease of both $R_{\rm g}$ and $R_{\rm gb}$, where $R_{\rm g}$ and $R_{\rm gb}$ stand for the resistance of grains and grain boundaries respectively. Because $R = R_g + R_{gb}$ and R_g and R_{gb} are usually in the same order of magnitude, for a 10⁷ fold decrease, both $R_{\rm g}$ and $R_{\rm gb}$ should decrease. The decrease of $R_{\rm g}$ is most likely attributed to the narrowing of the band gap. The band gap of CaC₂-I decreases when compressed (Fig. S8[†]) and reaches 0 eV at ~18 GPa as indicated by the theoretical calculations (Fig. S9† by interpolation). The interaction between C_2^{2-} units is enhanced when the C…C distance is decreased, which broadens the energy bands and reduces the band gap. Quantitatively, from the formula $\sigma = ne(\mu_{-} + \mu_{+}) = 2\sqrt{N_{-}N_{+}}e^{-\frac{E_g}{2kT}} \cdot e(\mu_{-} + \mu_{+}),^{14}$ we can derive $\lg \rho = -\lg(2\sqrt{N_-N_+}\cdot e(\mu_-+\mu_+)) + \frac{E_{\rm g}}{2kT}\cdot 0.4343$ $= -\lg(2\sqrt{N_-N_+} \cdot e(\mu_- + \mu_+)) + 8.46E_g \text{ (}E_g \text{ in eV), where } \rho \text{ and}$



Fig. 9 Pressure dependent electrical resistivity of CaC₂.

σ are the resistivity and conductivity, *n* is the concentration of electrons and holes, *N*_− and *N*₊ are the density of states at the bottom of the conduction band and the top of the valence band, $µ_−$ and $µ_+$ are the carrier mobilities, E_g is the band gap, *k* is the Boltzmann constant, *T* is the absolute temperature, and *e* is the charge of an electron. From 2 GPa to 6 GPa, $\Delta E_g = 0.23$ eV, and the variation of the second term $\Delta(8.46E_g) = 2$, which is partly responsible for the decrease of lg ρ (~4) (Table S2†). The other remainder of the decrease should be attributed to the enhancing of *N*_−, *N*₊, $µ_+$, and/or $µ_−$ in the first term, $-lg(2\sqrt{N_-N_+} \cdot e(µ_- + µ_+))$, or other extrinsic reasons. In the view of real (direct) space, the decrease of *R* can simply be attributed to the increasing overlap of the π and $π^*$ -orbitals, which enhances the mobilities ($µ_-$ and $µ_+$). This is similar to fused-ring molecules under high pressure.¹⁵

The irreversibility of the resistance decrease is evidenced by the two compression loops. This irreversibility is most likely attributed to the covalent bonding between the acetylide anions. The transitions accompanied with the formation of covalent bonds are often (kinetically) stable, just like the transition from graphite to diamond. In this work we show it is also true for metal carbides in the experimental time scale. The high conductivity is also recovered, which suggests that this pressure induced polymerization can be used to prepare conductive metal polycarbide materials from insulating monomers.

It is also notable that the d-orbital of Ca^{2+} also contributes to the valence band and the conduction band, as indicated by the theoretical calculations (shown in Fig. S8†). In compounds with the same structure, such as UC₂, U donates 6d electrons to C₂, forming a C₂⁴⁻ anion, and UC₂ is metallic.¹⁶ As such, the d- π interaction acts as another important factor that may affect the conductivity. Additionally, it also suggests the possibility of introducing a transition metal for doping, which is always a promising method to improve the conductivity, as in polyacetylene.¹⁷

Conclusions

In summary, the phase transitions of CaC₂ under external pressure were systematically investigated and the polymerization of C₂²⁻ was confirmed using *in situ* XRD, neutron diffraction, Raman, IR and impedance spectra, as well as neutron PDF and GC-MS. The reaction in the amorphous CaC₂ was firstly probed. The polymerized product can be recovered to ambient pressure. Benzene and several other hydrocarbon molecules are identified in the hydrolyzed product of the recovered sample, which clearly evidences the formation of C₆⁶⁻ and other polycarbide anions. Accompanied with the structural variations, the conductivity of CaC_2 was enhanced by more than 10^7 fold, which demonstrates the dramatic effects of high pressure on the electrical properties of materials. The polymerization helps to stabilize the high conductivity at ambient pressure, and suggests the possibility to stabilize the predicted phases including calcium graphenide and calcium polyacetylide. This will have a great impact on the current research of functional metal-carbon materials, and more novel metal polycarbide materials can be expected from the high-pressure synthesis. In addition, since these highly charged species like C_6^{6-} are easy to be functionalized, this method can be further applied to synthesize more complicated Ca organics and other metallic organics.

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References

- P. Srepusharawoot, A. Blomqvist, C. Moysés Araújo, R. H. Scheicher and R. Ahuja, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2010, 82, 125439.
- 2 X. Q. Chen, C. L. Fu and C. Franchini, J. Phys.: Condens. Matter, 2010, 22, 292201.
- 3 I. Efthimiopoulos, K. Kunc, G. V. Vazhenim, E. Stavrou, K. Syassen, M. Hanfland, St. Liebig and U. Ruschewitz, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, 85, 054105.
- 4 N.-G. Vannerberg, Acta Chem. Scand., 1962, 16, 1212–1220.
- 5 N.-G. Vannerberg, Acta Chem. Scand., 1961, 15, 769-774.
- 6 M. Knapp and U. Ruschewitz, *Chem.-Eur. J.*, 2001, 7, 874-880.
- 7 Y. L. Li, W. Luo, Z. Zeng, H.-Q. Lin, H.-k. Mao and R. Ahuja, *Proc. Natl. Acad. Sci. U. S. A.*, 2013, **110**, 9289–9294.
- 8 A. Kulkarni, K. Doll, J. C. Schon and M. Jansen, *J. Phys. Chem. B*, 2010, **114**, 15573–15581.
- 9 D. Benson, Y. Li, W. Luo, R. Ahuja, G. Svensson and U. Häussermann, *Inorg. Chem.*, 2013, **52**, 6402–6406.
- 10 J. Nylen, S. Konar, P. Lazor, D. Benson and U. Haussermann, J. Chem. Phys., 2012, 137, 224507.
- 11 I. Efthimiopoulos, G. V. Vajenine, E. Stavrou, K. Syassen, St. Liebig, U. Ruschewitz and M. Hanfland, *Acta Crystallogr.*, 2010, **66**, s197.
- 12 L. Wang, X. Huang, D. Li, Y. Huang, K. Bao, F. Li, G. Wu, B. Liu and T. Cui, *J. Chem. Phys.*, 2016, **144**, 194506.
- 13 S. J. L. Billinge and M. G. Kanatzidis, *Chem. Commun.*, 2004, 749–760.
- 14 K. Huang and R. Han, *Solid State Physics*, Higher Education Press, Beijing, 1st edn, 1988.
- 15 G. A. Samara and H. G. Drickamer, *J. Chem. Phys.*, 1962, 37, 474–479.
- 16 J. Li and R. Hoffmann, Chem. Mater., 1989, 1, 83-101.
- 17 H. Shirakawa, E. J. Louis, A. G. Macdiarmid, C. K. Chiang and A. J. Heeger, J. Chem. Soc., Chem. Commun., 1977, 16, 578–580.