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# Electronic and photocatalytic properties of N/F co-doped anatase TiO<sub>2</sub><sup>†</sup>

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The crystal structure, formation energies, electronic structures and effective masses of charge carriers of N/F co-doped anatase TiO<sub>2</sub> were investigated by first principle calculations. We have found that the incorporation of the N atom will be easier in the presence of the F dopant. Second, the impurity band of N/F co-doped TiO<sub>2</sub> has no spin polarization and smaller effective mass, which improves the photogenerated carriers' mobility and separation. Moreover, it has a lower band edge energy, which increases the oxidation ability of the photogenerated holes. Thus, we explained the mechanism of the enhancement of photocatalytic efficiency of N/F co-doped TiO<sub>2</sub> observed from experiments.

## 1. Introduction

TiO<sub>2</sub> has attracted much interest and been widely utilized in photoelectrochemical devices, photovoltaic solar cells and the atmospheric environment, because of its low cost, nontoxicity, chemical inertness, oxidizing potential and high photocatalytic activity.<sup>1–4</sup> However, practical applications of TiO<sub>2</sub> have been restricted by its large intrinsic band gap. Thus, it can only be excited by ultra-violet (UV) light (about 5% of the solar spectrum).<sup>5–7</sup>

Doping nonmetal or metal elements is considered to be a promising method to reduce the absorption threshold of semiconductor-based photocatalysts (TiO<sub>2</sub>, ZnO *etc.*) and to move it from UV to visible light (about 45% of the solar spectrum).<sup>8–12</sup> However, it is generally found that single element doped TiO<sub>2</sub> has limited visible light photocatalytic performance.<sup>13–18</sup> Such as, although much work has suggested that N doping of TiO<sub>2</sub> extends the optical absorption edge into the visible light region, increased photocatalytic efficiency is still limited.<sup>15,19</sup> On the other hand, co-doping two elements in TiO<sub>2</sub> demonstrates remarkable enhancement of photocurrent and photocatalyst effect, such as B and N co-doping,<sup>20</sup> C and N co-doping,<sup>21,22</sup> C and F co-doping,<sup>23</sup> N and S co-doping,<sup>24,25</sup> and N and F co-doping.<sup>26–30</sup> Among them, N and F co-doping are particular interesting, because the formation of the impurity levels (ILs) effectively lowering the bulk band, and charge compensation effect between N and F atoms reducing the recombination rate of photogenerated electron-hole pairs.<sup>9</sup>

However, there is almost no relevant theoretical study performed on the influence of photogenerated carrier effective mass and ORR over the enhanced photocatalytic efficiency of the N/F co-doped TiO<sub>2</sub>. Thus, a systematically first principles study was conducted on the crystal structures, formation energy, electronic structures and effective masses of charge carriers. The results show that the Ti-rich growth condition is favorable to all doped TiO<sub>2</sub>; meanwhile, the incorporation of the N atom into the TiO<sub>2</sub> will be promoted in the presence of the F dopant. Moreover, N/F co-doped TiO<sub>2</sub> not only greatly improves the photogenerated carriers' mobility and separation, but also increases the oxidation ability of the photogenerated holes; especially N<sub>s</sub>-F<sub>s</sub> co-doped TiO<sub>2</sub>. This explains why <sup>•</sup>OH radical dominates the process of photocatalytic degradation of methylene blue.

## 2. Computational methods

First-principle density functional theory (DFT) calculations are performed using the CASTEP code. Norm-conserving pseudopotentials and Perdew–Burke–Ernzerhof for Solids (PBESOL) function of the generalized gradient approximation (GGA) are used for the electron-ion interactions and exchange-correlation potential, respectively.<sup>31,32</sup> All simulations are carried out for a 2 × 2 × 1 supercell with 16 Ti atoms and 32 O atoms. In the N-doped system, we only consider the N atom substitutes the O atom (N<sub>s</sub>), which is consistent with the previous experimental results.<sup>15,26</sup> For co-doped TiO<sub>2</sub> model, N atom and F atom are simultaneously introduced into the supercell of TiO<sub>2</sub>, they substitutes the O atom (N<sub>s</sub> or F<sub>s</sub>) or occupy the interstitial position (N<sub>i</sub> or F<sub>i</sub>). This is resulted in four different kinds of N/F co-doped TiO<sub>2</sub>, namely, N<sub>s</sub>-F<sub>s</sub>-, N<sub>s</sub>-F<sub>i</sub>-, N<sub>i</sub>-F<sub>s</sub>- and N<sub>i</sub>-F<sub>i</sub>-co-doped TiO<sub>2</sub>. In addition, we have considered the physical doping position of N and F atoms, and calculated the total energies of all possible configurations (based on the doped distance between N and F), as shown in ESI Fig. S1.<sup>†</sup> Only the most stable

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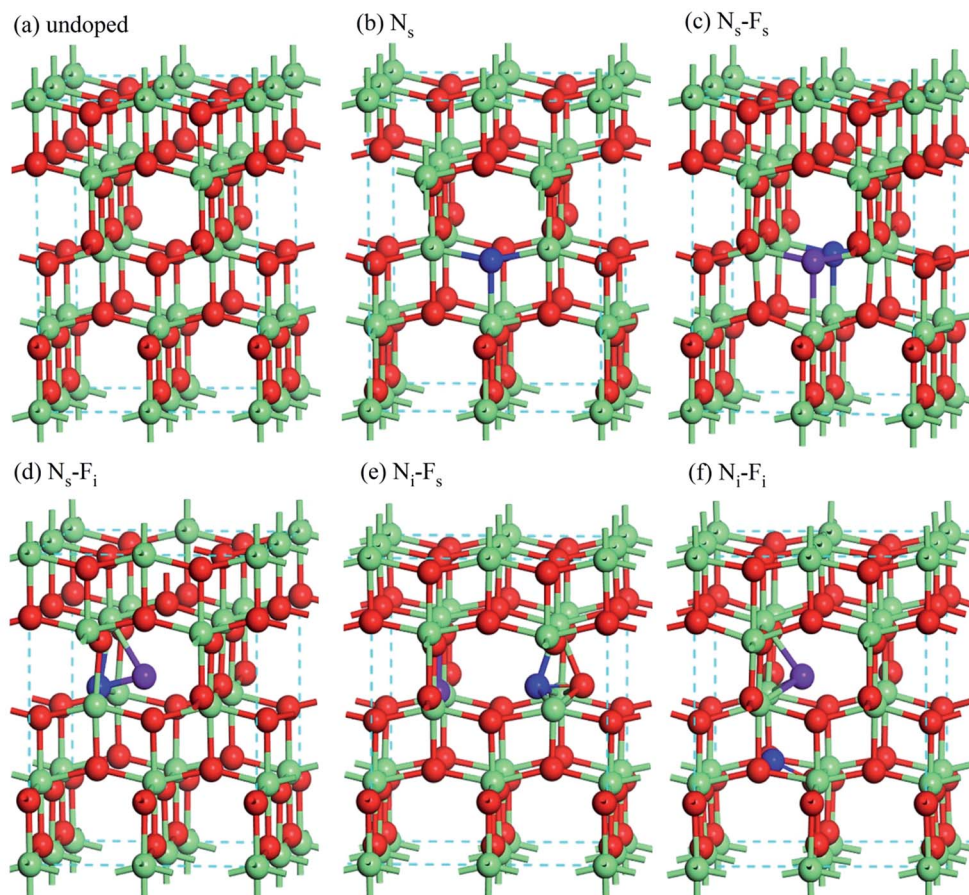


Fig. 1 The crystal structures of undoped and doped  $\text{TiO}_2$ , where green, red, blue and purple balls are the Ti, O, N and F atoms, respectively.

configurations for each case of doped  $\text{TiO}_2$  are shown in Fig. 1. The doping concentration (atom ratio) of the N or F ranges is chosen as around 2.08%, similar to the concentration used in the experiments.<sup>26,29,33,34</sup>

The cutoff energy with 750 eV and  $4 \times 4 \times 3$   $k$ -point sampling set are sufficiently large for the systems considered. The convergence tolerance of energy, maximum displacement, and maximum force are  $5.0 \times 10^{-6}$  eV per atom,  $5.0 \times 10^{-4}$  Å, and  $0.01$  eV Å<sup>-1</sup>, respectively. On the other hand, DFT calculations underestimate the band gap, and DFT+ $U$ , as the improvement of standard DFT can overcome this situation. In this work,  $U$  values was selected 7.0 eV for Ti elements based on the calculated band-gap for anatase  $\text{TiO}_2$  as a function of the  $U$  parameter in our previous work.<sup>35</sup> And, the energy for the band gap ( $E_g$ ) of undoped anatase  $\text{TiO}_2$  is 3.05 eV, which is consisted with the experimental value (3.20 eV) and other DFT+ $U$  and HSE06 calculated values.<sup>36,37</sup>

We calculated the  $E_f$  to assess the stability of all the doped systems:

$$E_s = E_{\text{doped}} - E_{\text{undoped}} - m\mu_N - n\mu_F + x\mu_O \quad (1)$$

where  $E_{\text{doped}}$  and  $E_{\text{undoped}}$  are the total energies of doped and undoped  $\text{TiO}_2$ . The  $m$  and  $n$  denotes the number of N and F atoms doped into the supercell of  $\text{TiO}_2$ , respectively. The  $x$  represents the numbers of O atoms substituted by dopants

atoms. In view of the growth of  $\text{TiO}_2$  being a dynamic process, the formation energy is unfixed and depends on the growth condition (changing from O-rich condition to Ti-rich condition). Relations between  $\mu_O$  and  $\mu_N$ ,  $\mu_F$ ,  $\mu_{\text{Ti}}$  are satisfying the equations:

$$\mu_{\text{Ti}} + 2\mu_O = \mu(\text{TiO}_2) \quad (2)$$

$$\mu_N + 2\mu_O = \mu(\text{NO}_2) \quad (3)$$

$$3\mu_F + \mu_O = \mu(\text{OF}_3) \quad (4)$$

Under O-rich condition,  $\mu_O$  is the chemical potential of the ground-state of the  $\text{O}_2$  molecule;  $\mu_{\text{Ti}}$ ,  $\mu_N$  and  $\mu_F$  are obtained by the eqn (2), (3) and (4), respectively. Under Ti-rich condition, the  $\mu_{\text{Ti}}$  is the chemical potential of Ti bulk, while  $\mu_O$ ,  $\mu_N$  and  $\mu_F$  are calculated by the eqn (2), (3) and (4), respectively. The simulation bond lengths of the  $\text{O}_2$ ,  $\text{NO}_2$  and  $\text{OF}_3$  molecules are close to the experimental value.<sup>38</sup>

### 3. Result and discussion

#### 3.1 Crystal structure and formation energies

Table 1 lists the calculated lattice constants of the undoped and doped  $\text{TiO}_2$ . The lattice constants of undoped  $\text{TiO}_2$  is consistent with the previous experimental results.<sup>39</sup> Overall, the effect of



**Table 1** The equilibrium supercell lattice constants  $a$ ,  $b$  and  $c$  with unit Å.  $V$  is the volume of all systems with unit Å<sup>3</sup>.  $E_f$  (Ti-rich and O-rich)  $E_f$ ,  $E_g$ ,  $E_{\max}$ ,  $E_{\text{CBM}}$  and  $E_{\text{VBM}}$  are the calculated formation energy, Fermi level, band gap, the maximum energy gap in the band gap, the energy of CBM and VBM, respectively, with unit eV.  $m^*$  is the carrier effective mass compared with electronic mass

System		TiO <sub>2</sub>	N <sub>s</sub>	N <sub>s</sub> -F <sub>s</sub>	N <sub>s</sub> -F <sub>i</sub>	N <sub>i</sub> -F <sub>s</sub>	N <sub>i</sub> -F <sub>i</sub>
$a$		7.551	7.587	7.564	7.573	7.565	7.602
$b$		7.551	7.543	7.607	7.554	7.595	7.550
$c$		9.627	9.615	9.589	9.696	9.655	9.695
$V$		548.91	550.25	551.74	554.67	554.74	556.45
$E_f$	Ti-rich	—	−1.79	−6.92	−3.73	−4.77	−0.01
	O-rich	—	5.77	4.01	4.68	3.63	5.89
$E_F$		6.41	4.08	4.09	3.65	4.10	5.19
$E_{\text{CBM}}$		9.46	7.06	6.47	6.70	6.72	6.79
$E_{\text{VBM}}$		6.41	4.08	3.53	3.65	3.48	3.74
$E_g$		3.05	2.98	2.94	3.05	3.24	3.05
$E_{\max}$		3.05	1.88	2.39	3.05	2.61	1.49
$m^*$	CBM	0.35	0.38	0.40	0.42	0.46	0.39
	IL	—	5.43	1.21	—	2.73, 1.08	1.63, 2.10, 2.27
	VBM	0.92	1.28	0.95	1.30	1.25	1.21

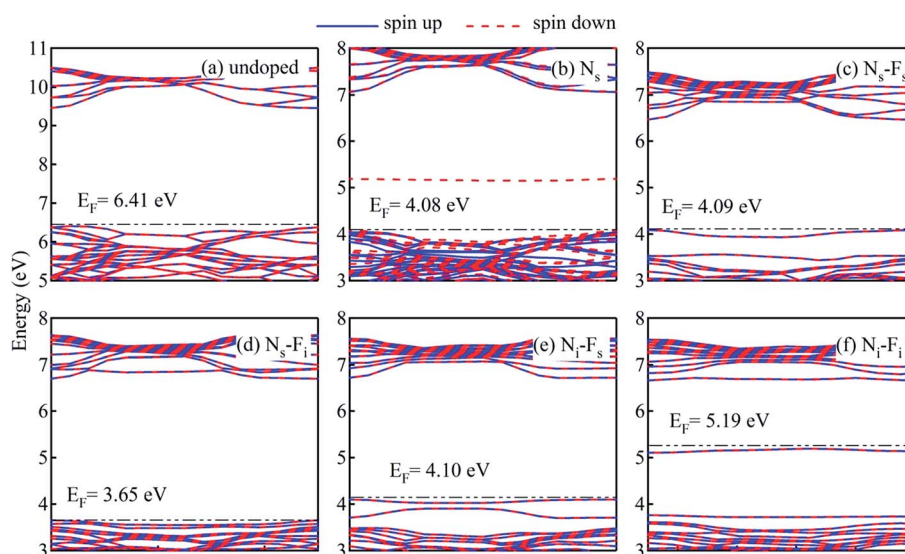
dopants and various doped models on the lattice constants and volume is very minor, within 0.74% and 1.37%, respectively. The reason may be that the atomic radius of the dopant is similar to that of the O atom. However, further observation from Fig. 1 found that dopants significantly affect the local structure and thus the electronic properties dramatically.

We have calculated the  $E_f$  to assess the stability of dopants into TiO<sub>2</sub>, as shown in Table 1. All of them are positive (or negative) under O-rich (or Ti-rich) conditions. Negative  $E_f$  suggests that dopant tends to enter the TiO<sub>2</sub> lattice. Thus, this confirms that the N(F) atoms prefer to substitute the O atom under Ti-rich conditions, since more oxygen vacancies are present in this case.<sup>29,34</sup> Meanwhile, the  $E_f$  of these doped systems have the relative relationships:  $E_f(\text{N}_s\text{-F}_s) < E_f(\text{N}_i\text{-F}_s) < E_f(\text{N}_s\text{-F}_i) < E_f(\text{N}_s) < E_f(\text{N}_i\text{-F}_i)$  under Ti-rich conditions. Thus, N<sub>s</sub>-F<sub>s</sub> co-doped TiO<sub>2</sub> is the most stable system, due to its lowest

$E_f$  (−6.92 eV). Moreover, the  $E_f$  of all co-doped TiO<sub>2</sub> (except for N<sub>i</sub>-F<sub>i</sub> co-doped TiO<sub>2</sub>) is smaller than N doped TiO<sub>2</sub> indicates that the incorporation of the N atoms into the TiO<sub>2</sub> will be easier to achieve in the presence of the F dopant. Further, considering the relationship between  $E_f$  and  $\mu_{\text{O}}$ , a variety of N/F co-doped TiO<sub>2</sub> can be synthesized by controlling the flow of O<sub>2</sub> and supplying enough energy.

### 3.2 Electronic properties

To demonstrate how doped atoms modify the electronic properties of TiO<sub>2</sub>, the band structures and partial density of states (PDOS) of individual Ti, O, N and F atoms of all the systems are shown in Fig. 2 and 3, respectively. And the  $E_g$ , the Fermi level ( $E_F$ ), the maximum energy gap in the band gap ( $E_{\max}$ ), the energy of conduction band minimum (CBM) ( $E_{\text{CBM}}$ ) and valence band maximum (VBM) ( $E_{\text{VBM}}$ ) are listed in Table 1.



**Fig. 2** The band structures of undoped and doped TiO<sub>2</sub>, where the blue full (red dash) lines represent the spin up (spin down) states and the horizontal dashed lines denote the Fermi level ( $E_F$ ).





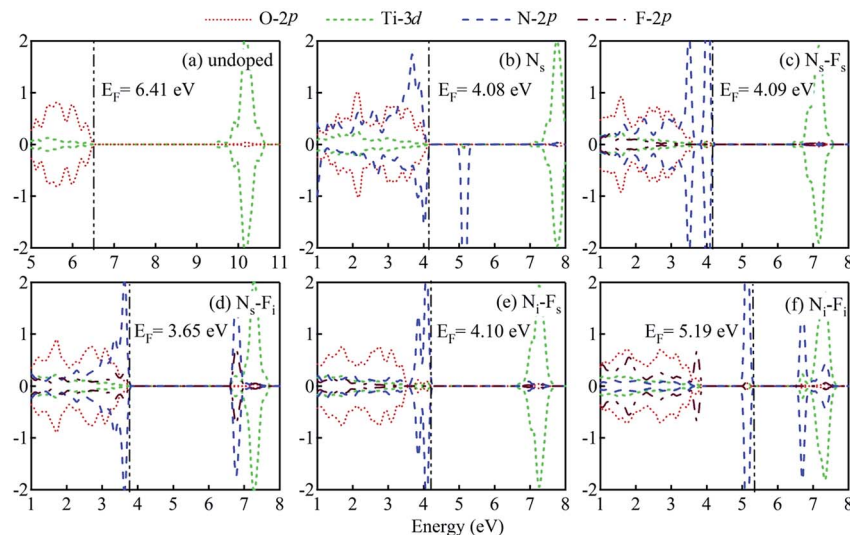


Fig. 3 PDOS of undoped and doped  $\text{TiO}_2$ , the vertical dashed lines denote the Fermi level ( $E_F$ ).

The band gap of  $\text{TiO}_2$  is mostly controlled by O-2p and Ti-3d states, but the ILs of the dopants can also modify it. Thus the dopants can fine tune the band gap of the doped system. The details of the band gap tuning can be described as the followings.

For  $\text{N}_s$  doped  $\text{TiO}_2$  (Fig. 2b and 3b), the valence band now has the half-filled  $\text{N}_s$ -2p state which is just above the original VBM, while the conduction band (CB) remain unchanged. Thus, the valence band (VB) lifts toward to the CB and the  $E_g$  is reduced by 0.07 eV. Similar to this, the filled IL of the  $\text{N}_s$ - $\text{F}_s$  co-doped  $\text{TiO}_2$  (Fig. 2c and 3c) is also above the VBM. Thus, the  $E_g$  is lowered to 2.94 eV compared to undoped  $\text{TiO}_2$ . For the case of  $\text{N}_s$ - $\text{F}_i$  co-doped  $\text{TiO}_2$  (Fig. 2d and 3d) and  $\text{N}_i$ - $\text{F}_i$  co-doped  $\text{TiO}_2$  (Fig. 2f and 3f), the ILs overlap the original VB and CB, thus the band gap remains the same. For the case of  $\text{N}_i$ - $\text{F}_s$  co-doped  $\text{TiO}_2$  (Fig. 2e and 3e), the filled ILs locate in the middle of the band gap. And the VB is lowered due to missing one O atom. Thus the band gap increases compared to the undoped  $\text{TiO}_2$ . This verifies that N and F atoms can modulate the electronic structure of the  $\text{TiO}_2$ , by changing the position and amount of ILs. Furthermore, the spin polarization states present in the  $\text{N}_s$  doped  $\text{TiO}_2$  is not observed in the co-doped system due to the charge compensation effect between N and F atom.

Overall, the doped system only reduces the band gap below 0.11 eV. In other words, the maximum red-shifted optical absorption edge is only 15.3 nm. However, the ILs in the band gap not only reduce the  $E_g$ , but also separate and promote photo-excited electrons pumping from the VB into the CB. These characteristics are likely to improve visible light photocatalytic activity. More details will be discussed later.

### 3.3 Effective mass of photogenerated carriers

We all know that the limited number of density of states and the low charge carrier mobility in the flat ILs limits its role as a stepping stones for electronic transition, and the electrons are easily annihilated by recombination with holes. The carrier mobility has been calculated according to the equation  $\mu = e\tau/m^*$

where  $\tau$  is the mean free time,  $m^*$  is the effective mass (both related to the curvature of the band level).<sup>40,41</sup> Smaller efficient masses mean higher carrier mobility. The results show that there is a flat IL ( $m^* = 5.43$ ) in the  $\text{N}_s$  doped  $\text{TiO}_2$  that limits the separation of photogenerated carriers. In contrast, all co-doped systems have smaller  $m^*$ , especially  $\text{N}_s$ - $\text{F}_s$  co-doped  $\text{TiO}_2$  ( $m^* = 1.21$ ). Thus, this IL suppresses the recombination rate. We believe that  $\text{N}_s$ - $\text{F}_s$  co-doped  $\text{TiO}_2$  will have the highest photocatalytic efficiency.

### 3.4 Band edge energy and photocatalytic properties

The band edges energy ( $E_{\text{VBM}}$  and  $E_{\text{CBM}}$ ) determines the redox potential of the semiconductor photocatalyst. Fig. 4 and Table 1 indicate that the  $E_{\text{VBM}}$  and  $E_{\text{CBM}}$  of all doped  $\text{TiO}_2$  move down to the low energy region (more than 2.33 eV), compared with the undoped  $\text{TiO}_2$ . Such as, the oxidizability of photogenerated holes at the VBM is enhanced 2.33–2.93 eV, while the reducibility of photogenerated electrons at the CBM is reduced 2.40–

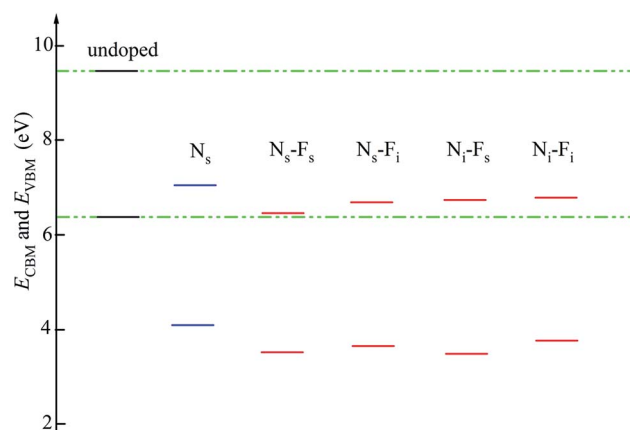


Fig. 4 The  $E_{\text{CBM}}$  and  $E_{\text{VBM}}$  of undoped (green dotted line) and doped (blue and red line)  $\text{TiO}_2$ .



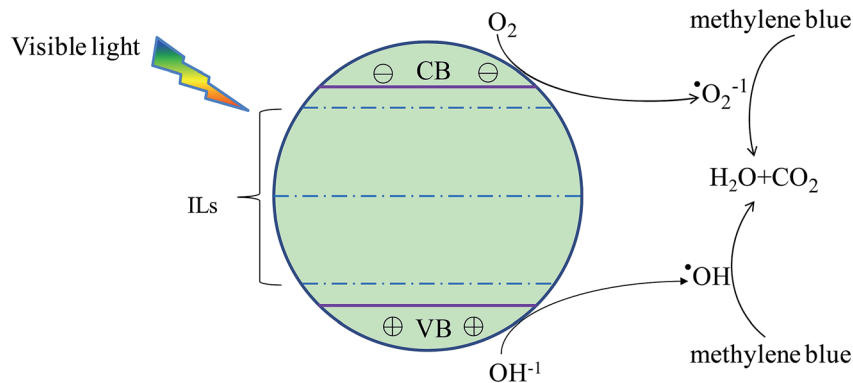


Fig. 5 The mechanism for photocatalytic degradation of MB by doped  $\text{TiO}_2$  under visible light irradiation.

2.99 eV. Moreover, the band edges energy of N/F co-doped  $\text{TiO}_2$  move more to the lower energy region compared with  $\text{N}_s$  doped  $\text{TiO}_2$ . So, they have greater ability to modify the redox potentials; especially the  $\text{N}_s\text{-F}_s$  co-doped  $\text{TiO}_2$ .

On the other hand,  $\text{O}_2^-$  ion and  $\cdot\text{OH}$  radical all are important oxidants in the process of photocatalytic degradation of methylene blue. The degradation mechanism is shown in Fig. 5. On the surface of  $\text{TiO}_2$ ,  $\text{O}_2^-$  ions are reduced by the photogenerated electrons from the  $\text{O}_2$ , while  $\cdot\text{OH}$  radicals are oxidized by the photogenerated holes from the  $\text{OH}^-$ . Thus, all doped systems tend to produce more  $\cdot\text{OH}$  than  $\text{O}_2^-$  during the photocatalytic process. For that, as discussed above, we conclude that  $\text{N}_s\text{-F}_s$  co-doped model can explain the mechanism of enhanced photocatalytic activity in the recent experiments that the  $\cdot\text{OH}$  radical played a leading role during the visible light photocatalytic process.<sup>33,42–45</sup> Meanwhile, volumes were 550.25 and 551.74  $\text{\AA}^3$  for  $\text{N}_s$  doped and  $\text{N}_s\text{-F}_s$  co-doped  $\text{TiO}_2$ . This expanded surface area also promotes the photocatalytic ability of  $\text{N}_s\text{-F}_s$  co-doped  $\text{TiO}_2$ .

## 4. Conclusions

In summary, the crystal structure, formation energy, electronic structures and effective masses of charge carriers of N doped and N/F co-doped  $\text{TiO}_2$  were investigated by first-principles calculations. We have found that the Ti-rich growth condition is beneficial to all doped  $\text{TiO}_2$ . Meanwhile, the incorporation of the N atom into the  $\text{TiO}_2$  will be easier to achieve in the presence of the F dopant. N/F co-doped  $\text{TiO}_2$  has no spin polarization, smaller effective mass and lower band edge energy. They, especially  $\text{N}_s\text{-F}_s$  co-doped  $\text{TiO}_2$ , not only greatly improve the photogenerated carriers' mobility and separation, but also increase the oxidation ability of the photogenerated holes. Thus, we have explained the mechanism of the enhancement of photocatalytic efficiency of N/F co-doped  $\text{TiO}_2$  observed by experiments. Moreover, this work also provides new insights into synthesis and design of the various-doped  $\text{TiO}_2$  by controlling the flow of  $\text{O}_2$  during sample preparation.

## Conflicts of interest

There are no conflicts to declare.

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