



Cite this: RSC Adv., 2017, 7, 45156

Solvothermal synthesis of three-dimensional, Fe_2O_3 NPs-embedded CNT/N-doped graphene composites with excellent microwave absorption performance[†]

Nan Zhou,^a Qingda An,^{ID *a} Zuoyi Xiao,^a Shangru Zhai,^{ID *a} and Zhan Shi^{ID b}

For high-performance microwave absorption materials, component integration, structure tailorability and light weight are of practical significance. Herein, a new three-dimensional magnetic composite, namely carbon nanotubes (CNTs)-intercalated nitrogen-doped graphene (N-GN) embedded with Fe_2O_3 NPs (referred to as Fe_2O_3 /N-GN/CNTs), were fabricated by a solvothermal method for use as synergistic microwave absorbers. Electronic microscopy investigation shows that the Fe_2O_3 NPs, which formed *in situ*, are uniformly dispersed on the N-doped graphene nano-sheets without aggregation and that the CNTs have been successfully inserted into the layers of nano-sheets. In addition, X-ray diffraction patterns (XRD), Fourier transform infrared spectra (FTIR), X-ray photoelectron spectroscopy (XPS), Raman spectrometry and a vector network analyzer (VNA) were used to investigate the structure, chemical composition, and electromagnetic parameters of the resulting Fe_2O_3 /N-GN/CNTs composite. Taking advantage of the synergistic effect between Fe_2O_3 nanocrystals with magnetic loss and N-GN/CNTs with dielectric loss, the 3D Fe_2O_3 /N-GN/CNTs composite shows excellent microwave absorption capabilities with highly efficient performance. A maximum reflection loss of -45.8 dB can be achieved at 9.32 GHz with a matching thickness of only 3 mm, and the effective absorption (below -10 dB) bandwidth reaches 14.5 GHz (3.5 to 18.0 GHz). The excellent microwave absorption of the as-prepared composites can be related to the integrated characteristics of good impedance matching and Debye relaxation, high attenuation, strong dipole polarization and interface polarization. Therefore, these results indicate that the newly designed Fe_2O_3 /N-GN/CNTs composite is an ideal candidate for use as a synergistic microwave absorber to address electromagnetic pollution.

Received 17th June 2017

Accepted 8th September 2017

DOI: 10.1039/c7ra06751h

rsc.li/rsc-advances

1. Introduction

Currently, with the rapid development of electronic devices, such as wireless communication facilities, microwave circuit devices and radar systems, electromagnetic interference (EMI) pollution has become a serious problem that not only affects the operation of electronic devices, but is also harmful to the health of human beings.¹ To resolve the issues associated with EMI pollution, it is urgently necessary to find effective ways to attenuate EM waves. Fortunately, to address this problem, microwave absorption (MA) materials can be utilized to weaken the unwanted EM energy. Over the past few decades, there has been increasing demand for a variety of materials to be used as microwave absorbers for military and commercial applications.

This has sparked extensive studies for the development of novel high-performance microwave absorbers with the advantages of light weight, low thickness, wide absorption bandwidth, and strong absorption characteristics.^{2,3} Conventional MA materials based on carbon, ceramics, and metallic magnets as well as their hybrids have been extensively used because of their good dielectric and magnetic properties, low cost, high abundance and environmental benignity.^{4–6} In recent years, ferrite has received particularly high attention in the MA field, and the research on this topic has become more popular. However, ferrite has some drawbacks, including high density, a narrow absorption band and poor absorption intensity, which severely limit its applications in practice. Therefore, there is high demand for the development of new lightweight materials/composites with excellent EM wave absorption characteristics.

In view of these issues, graphene-based hybrids have been studied as an effective material that can not only reduce the density of an absorber, but can also enhance its microwave absorbing characteristics due to their strong dielectric loss. As is well known, lightweight graphene, a new type of carbon

^aFaculty of Light Industry and Chemical Engineering, Dalian Polytechnic University, Dalian 116034, China. E-mail: anqingdachem@163.com; zhaisrchem@163.com

^bState Key Laboratory of Inorganic Synthesis and Preparative Chemistry, College of Chemistry, Jilin University, Changchun 130012, China

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c7ra06751h

material, is considered to be the most useful building block among all carbon forms. Chemical oxidation and reduction processes using graphite as a raw material can be employed for the mass production of graphene with relative cost-effectiveness. Generally, graphene produced by the method mentioned above is known as reduced graphene oxide (RGO). Owing to its excellent layered structure, larger surface area, and outstanding electronic conductivity, RGO is considered to be an excellent candidate for attaining light-weight MA properties; it can readily convert electrons within its network structure under alternating electrical fields and expend the incident microwaves by transforming them into thermal energy.^{7–9} Furthermore, the residual functional groups and defects in RGO can give rise to enhanced defect polarization relaxation and electronic dipole relaxation, which are beneficial to MA.¹⁰ Recently, researchers have reported that GN-derived composites have excellent MA properties and can be used as EM wave absorbing materials. For example, Xu *et al.* researched the MA properties of RGO and reported that RGO exhibited good MA properties.¹¹ Wu *et al.* prepared an S-PPy/RGO aerogel and found that GN had good dielectric properties.¹² Moreover, a variety of iron-based composites have been synthesized; for example, Wang and coworkers fabricated graphene/α-Fe₂O₃ nanocomposites and studied their MA properties.¹³ The maximum reflection loss of the composites reached –38 dB at 14.78 GHz, and absorption bandwidths greater than –10 dB were obtained at 5.8 GHz with a thickness of 2 mm. The MA properties of RGO/SiO₂/Fe₃O₄ composites were also reported by Xu *et al.*, who found that the composites exhibited strong magnetic loss in the low-frequency range;¹¹ the maximum reflection loss of the composites reached –56.4 dB at 8.1 GHz with a thickness of 4.5 mm. Wang *et al.* researched a facile method to synthesize GO/CNT-Fe₃O₄ composites, and the composite materials exhibited high MA properties in the low-frequency range (2 to 9 GHz).¹⁴ In addition, Mishra's research group confirmed that the presence of RGO plays an important role in enhancing the absorption of electromagnetic waves.¹⁵ However, although several types of graphene-based hybrid absorbers have been investigated in previous studies, it is still a significant challenge to disperse magnetic nano-particles in the related 3D materials in a facile, controllable manner. More unfavorably, their low-yielding and tedious fabrication processes are bottlenecking factors for the potential practical applications of magnetic nano-particles.

In view of these issues, researchers formulated a new idea to dope nitrogen atoms in graphene. Interestingly, the doping of nitrogen atoms in graphene can not only improve the conductivity of graphene, but can also introduce a large number of defects in graphene sheets which can provide nucleation sites for the growth of magnetic grains, leading to efficient control over the particle size distribution and dispersion.¹⁶ Accordingly, defect-rich nitrogen-doped graphene would be favorable for MA, which can be attributed to its particular heterogeneous interfaces. The presence of heterogeneous interfaces can enhance dielectric loss as a result of multi-interfacial polarizations. In addition, the interfaces can provide more active sites to promote the reflection and scattering of microwaves. However, it is well known that serious agglomeration usually occurs

among graphene sheets because of the strong π–π conjugation between them.¹⁷ If this problem is not solved, the large number of defects produced by the addition of nitrogen atoms on the surface of the graphene sheet would be meaningless. To address this issue, taking advantage of the excellent mechanic stability of fiber-like CNTs as an effective solution for intercalation, the effect of physical separation can prevent aggregation of the two-dimensional nitrogen-doped graphene sheets.¹⁸ CNTs are not acid-treated during this process. This is because after carboxyl functionalization, although the dispersion increases in solution, the electronic structure causes greater damage. This results in the formation of graphene/CNTs C–C interface layers with electronic structures that do not match each other, thus reducing the efficiency of electron transfer between the materials; this affects the further transmission of electromagnetic waves in the material.¹⁶

Herein, following the aforementioned discussion and in continuation of our research interest in designing high-performance MA absorbers,^{19–21} a facile one-step solvothermal method is presented for the fabrication of ternary composite materials, *i.e.* Fe₂O₃/N-GO/CNTs, which can behave as a new types of lightweight and highly efficient microwave absorber. The intercalation of CNTs in nitrogen-doped graphene sheets (N-GN/CNTs) and the formation of Fe₂O₃ nanocrystals were achieved simultaneously. The synergistic effect of the N-GN/CNTs and Fe₂O₃ nanocrystals greatly enhanced the microwave absorption properties of the material, which is an ideal candidate for next generation microwave absorbers. That is, the aim of this work is to explore a simple approach for the manufacture of 3D N-doped graphene-based composites with thoroughly dispersed magnetic Fe₂O₃ nanocrystals, which are potential microwave absorbers for practical applications.

2. Experimental

2.1 Materials

Graphite (99.95%, ≥325 mesh) was purchased from Shanghai Aladdin Reagent Co. Ltd, China. The CNTs (diameter, 20 to 40 nm; length, ~5 to 10 μm; >97%) were purchased from Shenzhen Nanoport Co. Ltd, China. Potassium permanganate, sodium nitrate, sulfuric acid (H₂SO₄, 98 wt%), hydrogen peroxide (H₂O₂, 30 wt% in water), urea, ethanol and FeSO₄·7H₂O were purchased from Tianjin Kermel Chemical Reagent Co. Ltd, China. All chemicals were of analytical grade and were used without further purification.

2.2 Synthesis of Fe₂O₃/N-GN/CNTs composites

First, graphite oxide was fabricated from natural flake graphite powder by a modified Hummers' method.²² Briefly, 2.0 g of graphite powder was dispersed in 48 ml H₂SO₄ solution in an ice bath for 1 hour. With continuous agitation, 1 g NaNO₃ and 6 g KMnO₄ were slowly added to the solution while maintaining the temperature of the solution below 25 °C. After stirring for 1 hour, 280 ml distilled water was added dropwise. After an additional 30 minutes of stirring, H₂O₂ (30%) was slowly added to the above-mentioned solution until oxygen release stopped. Graphite oxide



was washed 5 times with 5% hydrochloric acid solution and distilled water by centrifugation at 5000 rpm for 10 min, respectively. Then, the gel-like graphite oxide was freeze-dried at -50°C for 24 h to obtain dry graphite oxide for further use.

Next, 35 mg of GO was added to deionized water (40 ml), followed by strong ultrasonic treatment (ultrasonic bath SK250H, 53 kHz) for 1 hour (maintaining the solution temperature below 25°C) to produce several layers of GO nanosheets. Then, CNTs (5 mg) were added to the abovementioned GO solution by ultrasonic treatment for another 1 h (at room temperature). Subsequently, 4.6 mmol of $\text{FeSO}_4 \cdot 7\text{H}_2\text{O}$ and 20 mmol urea were added to the abovementioned solution, which was followed by magnetic stirring for about 15 min. Finally, the solution was sealed in a 60 ml Teflon-lined stainless autoclave and then placed in an oven at 180°C for 6 h. After the autoclave cooled to room temperature naturally, the black precipitates were collected using a permanent magnet, washed with deionized water (60 ml) and ethanol (60 ml) 3 times sequentially, and dried in an oven at 60°C overnight.

Under the same experimental conditions, the $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{GN}$ composites were fabricated.

2.3 Structural characterizations

XRD patterns were obtained with a Rigaku model D/max-2700 diffractometer using $\text{Cu K}\alpha$ radiation at a scan rate (2θ) of $8^{\circ}\text{ min}^{-1}$ ranging from 10° to 70° . All measurements were performed in triplicate and repeated 3 times. The chemical structures of the prepared nanocomposite materials were evaluated using FTIR spectroscopy (Perkin-Elmer 400 spectrometer, Llantrisant, UK) over the wavenumber range of $400\text{--}4000\text{ cm}^{-1}$. All assays were performed in triplicate and repeated 3 times. The surface composition was investigated by X-ray photoelectron spectroscopy (XPS, PHI 5000). All assays were performed in triplicate and repeated 3 times. The morphologies and structures of the samples were observed by scanning electron microscopy (SEM, JEM JEOL 2100) and transmission electron microscopy (TEM, Hitachi H9000NAR). The structural deformation was recorded using Raman spectra (RM2000, Renishaw,

UK). All assays were performed in triplicate and repeated 3 times. The magnetic properties of the composites were measured with a vibrating sample magnetometer (VSM) at room temperature. All assays were performed in triplicate and repeated 3 times. The electromagnetic parameters of the samples were studied using a network analyzer (Agilent N5222A) in the range of 2 to 18 GHz. For microwave absorption measurements, the samples were pressed into toroidal shapes with an outer diameter of 7.00 mm and an inner diameter of 3.04 mm in which 40 wt% of sample and 60 wt% of paraffin were mixed uniformly. All assays were performed in triplicate and repeated 3 times.

3. Results and discussion

3.1 Characterization of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$

Scheme 1 shows the formation process of the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ hybrids synthesized by a facile solvothermal route. First, graphene oxide (GO) was synthesized by a modified Hummers' method. After oxidization, several functional groups such as $-\text{COOH}$ and $-\text{OH}$ are uniformly distributed on the surface of the GO sheets, and these groups are negatively charged after the ionization process. Second, GO, CNTs and Fe^{2+} cations were dispersed uniformly under ultrasonic treatment. The Fe^{2+} was anchored by the functional groups on the surface of the GO sheets *via* electrostatic interactions. Subsequently, urea was added to the abovementioned solution with vigorous stirring. Lastly, hydrothermal treatment was conducted at 180°C for 6 h. In this process, Fe^{2+} was oxidized to Fe^{3+} , resulting in the formation of FeO(OH) , and the GO was reduced and doped with N atoms. Under hydrothermal treatment, Fe_2O_3 nanoparticles were grown *in situ* on the surface of N-doped graphene nanosheets (N-GN) pillared with CNTs, resulting in the formation of three-dimensional $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites. The $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites exhibited high stability against strong ultrasound treatment, suggesting that all the Fe_2O_3 nanoparticles were formed *in situ* and bonded on the surface of N-GN nanosheets with favorable interactions; this can be



Scheme 1 Fabrication process of the magnetic carbonaceous absorbing material.



attributed to the stabilization effect of nitrogen atoms on the defect sites.²³

The crystalline structures of the $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites were characterized by XRD. As shown in Fig. S1,[†] the 2θ angles at 24.1° , 33.2° , 35.6° , 40.8° , 54.1° , 62.4° can be assigned to the (012), (104), (110), (113), (116), (214) planes of Fe_2O_3 (JCPDS no. 33-0664), suggesting that the material should present crystalline Fe_2O_3 components; this was further supported by the XPS results. Moreover, no obvious characteristic diffraction peaks of GO (at about 10°) or graphite (at about 26°) were detected, which demonstrates the highly disordered distribution of complexes and the low graphitization of hydrothermal reduction products.²⁴ Furthermore, comparing the XRD curves of the $\text{Fe}_2\text{O}_3/\text{GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN}$ samples, it can be clearly seen that the addition of urea did not affect the Fe_2O_3 crystal formation. The formation mechanism probably occurs due to the relatively low concentration of the precursor in the preparation conditions; moreover, Fe^{2+} was oxidized to Fe^{3+} during the high temperature hydrolysis, leading to the *in situ* formation of Fe_2O_3 by hydrolysis–dehydration reaction.²⁵ At the same time, for the $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ samples, the introduction of urea during the preparation process can hydrolyze and release OH^- under high temperature conditions, which can also promote the formation of Fe_2O_3 to a certain extent.

XPS is a very effective technique to investigate the surface states of materials; it can deliver information regarding their atomic compositions and can also identify the types of bonds between atoms. The XPS survey spectra of the $\text{Fe}_2\text{O}_3/\text{GN}$,

$\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites are shown in Fig. 1a. Clearly, four sharp peaks with binding energies at 285, 400, 530 and 711 eV can be observed, which are ascribed to the characteristic peaks of C 1s, N 1s, O 1s and Fe 2p, respectively, demonstrating the existence of C, O, Fe and N elements. When highly accurate XPS spectra were further acquired to elucidate the electronic states of the elements, the Fe 2p spectrum of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ showed two characteristic peaks at 711 and 725 eV, which were related to $\text{Fe}^{2p_{3/2}}$ and $\text{Fe}^{2p_{1/2}}$, respectively. The $\text{Fe}^{2p_{1/2}}$ and $\text{Fe}^{2p_{3/2}}$ main peaks are visibly followed by two satellite peaks on their high binding-energy side (at ~ 8 eV), as shown in Fig. 1b; these are feature peaks of Fe_2O_3 .²⁶ Moreover, the C 1s spectrum of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ shows five peaks, *e.g.* 284.7, 285.3, 285.7, 286.5 and 289.8 eV, which are attributed to C=C and C–C in aromatic rings, and C–N, C–O and C(O)O groups, respectively (Fig. 1c). These results further confirm that nitrogen atoms have been effectively doped in the graphite layers.²⁷ Furthermore, in Table S1,[†] we list the atomic compositions of the GN, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites and the content of C 1s chemical groups obtained by their spectral characteristics. The XPS results show that the total C–C content ($\text{sp}^2 + \text{sp}^3$) is higher for $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$, and the ratio between sp^2 and sp^3 C is also higher for $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ than for $\text{Fe}_2\text{O}_3/\text{N-GN}$. This further demonstrates the successful introduction of carbon nanotubes, resulting in an increase in C=C (sp^2) content. This statement is in agreement with the Raman results (see below). Moreover, on comparison, after the hydrothermal reaction, the N 1s peak was detected simultaneously in the $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ samples,

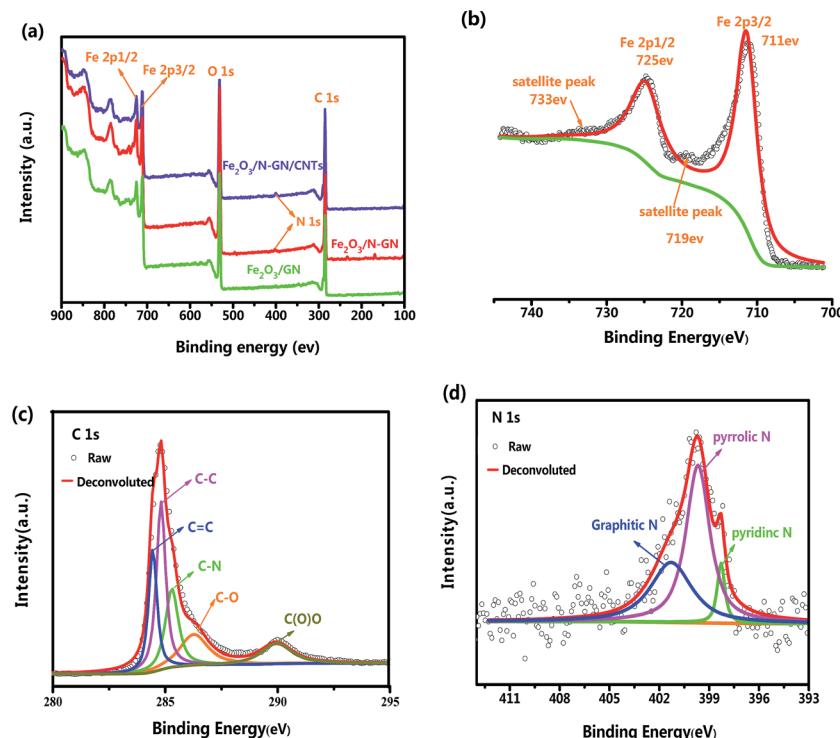


Fig. 1 (a) XPS survey spectra of the $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites. (b) Fe (2p), (c) C (1s) and (d) N (1s) core level spectra of the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites.



again indicating that N element was successfully incorporated into the composites. The N 1s spectrum of the $\text{Fe}_2\text{O}_3/\text{N-GN}/\text{CNTs}$ hybrid (Fig. 1d) can also be fitted by three deconvoluted peaks at 398.6, 400.1 and 401.7 eV, corresponding to pyridinic N bonding, pyrrolic N bonding and graphitic N bonding, respectively. It is clear that the pyrrolic N bonding configuration is the dominant N-doping type in the composites. Favorably, this form of N-doping can provide a large amount of space on the surface of the graphite layers, leading to an increased number of defect sites, which may produce polarization relaxation and electronic dipole polarization relaxation. This should be beneficial for the absorption of microwaves.²⁸ In addition, the XPS analysis results indicate that the GO was partially reduced and doped with CNTs and N atoms during the formation of Fe_2O_3 nanocrystals, leading to the formation of $\text{Fe}_2\text{O}_3/\text{N-GN}/\text{CNTs}$ composite; this was further demonstrated by the FTIR spectra.

In addition, FTIR spectroscopy can provide information on the bonding of organic functional groups. The FTIR adsorption spectra of the GO, $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN}/\text{CNTs}$ composites are displayed in Fig. 2. It can be observed from the FTIR spectrum of GO that a band appearing at about 3440 cm^{-1} was associated with (O-H) stretching vibrations, indicating that the material still contains a small amount of adsorbed water molecules. In addition, the peaks that appeared at around 2930 and 2921 were ascribed to the stretching frequency of C-H. The band at about 1620 cm^{-1} was related to stretching vibrations from (C=C) and assigned to the skeletal vibrations of unoxidized graphitic domains or contributions from the stretching deformation vibration of intercalated water.²⁹ The bands located at about 1429 and 1057 cm^{-1} were ascribed to (O-H) symmetric deformation vibrations and (C-O) (epoxy or alkoxy) stretching vibrations, respectively. Moreover, the band for (C=O) in carboxylic acid and carbonyl moieties (carbonyl) was observed at around 1743 cm^{-1} .³⁰ By comparing the FTIR spectra of $\text{Fe}_2\text{O}_3/\text{GN}$ and GO, it can be found that most of the oxygen-containing groups, such as (C=O) bonds, (C-OH)

bonds and (C-O) bonds, were partly removed during the hydrothermal process, indicating that GO could be reduced to a certain extent under high temperature. However, it can also be observed from the FTIR spectrum of $\text{Fe}_2\text{O}_3/\text{GN}$ that a band appeared at around 576 cm^{-1} ; this was attributed to Fe-O stretching vibrations. Moreover, on comparing the FTIR spectra of $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN}/\text{CNTs}$, it is found that the characteristic peaks appearing in the spectra, such as those at ~ 3440 , ~ 1620 and $\sim 576\text{ cm}^{-1}$, are similar. However, there is a small difference in the 1100 to 1600 cm^{-1} region. The differences in the feature peaks at ~ 1545 and $\sim 1150\text{ cm}^{-1}$ were ascribed to C=N stretching vibrations and C-N stretching vibrations, respectively.³¹ Therefore, these results show that the addition of urea could simultaneously enable hydrothermal reduction of GO and introduce nitrogen atoms into the final sample.

Raman spectroscopy, an efficient and non-multilative technique, can be used to identify the bonding and microstructures of C-species materials. In order to understand the chemical bonds in the carbon skeleton of in the hybrid system and any structural changes, Raman spectroscopy was adopted. Fig. 3 shows two representative Raman spectra of the GO, $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN}/\text{CNTs}$ composites. As shown in Fig. 3, the D band at $\sim 1340\text{ cm}^{-1}$ and the G band at $\sim 1590\text{ cm}^{-1}$ can be clearly detected in the three composites. The G band is a characteristic feature of graphitic layers, while the D band corresponds to disordered carbon.³² The intensity ratio of the D and G bands (I_D/I_G) is related to the disordered structure of carbon.³³ Compared with GO and $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ reveals an obviously enhanced value of I_D/I_G , indicating that more defects are introduced into the $\text{Fe}_2\text{O}_3/\text{N-GN}$ composite during the hydrothermal reduction process. This probably originates from the incorporation of N atoms, which may disrupt the sp^2 structure of the hexagonal lattice domain, resulting in the formation of sp^3 defects and disordered sites; this is helpful to improve the dielectric loss of microwave

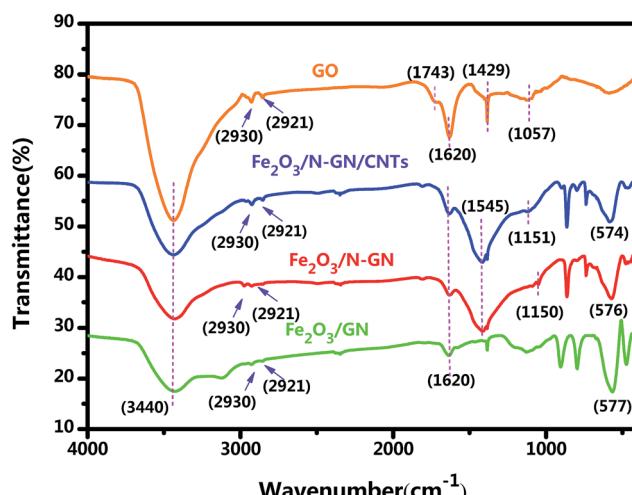


Fig. 2 FTIR spectra of the GO, $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN}/\text{CNTs}$ composites.

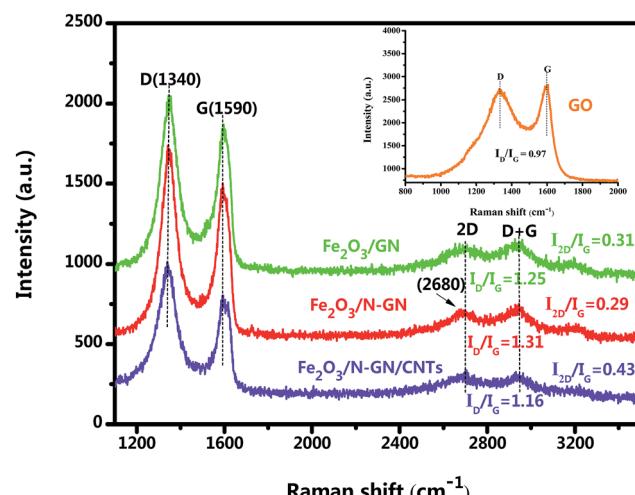


Fig. 3 Raman spectra of the GO, $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN}/\text{CNTs}$ composites.



absorbing materials.³⁴ However, in contrast to $\text{Fe}_2\text{O}_3/\text{N-GN}$, the $I_{\text{D}}/I_{\text{G}}$ ratio of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ has a certain degree of decline, which can be associated with the insertion of more structured carbon CNTs in the N-doped graphite layers.³⁵ In addition, we analyzed the variation of the 2D peak at about 2680 cm^{-1} . The number of layers of graphene was calculated using the $I_{\text{2D}}/I_{\text{G}}$ ratio; the layers increased with increasing $I_{\text{2D}}/I_{\text{G}}$ ratio. Clearly, the intensity ratio of $I_{\text{2D}}/I_{\text{G}}$ of the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ hybrid is greater than that of $\text{Fe}_2\text{O}_3/\text{GN}$ and the $\text{Fe}_2\text{O}_3/\text{N-GN}$ hybrid (Fig. 3). This result indicates that the insertion of CNTs may play an important role in the separation of graphene sheets, leading to the construction of the 3D composite structure.

The morphologies and microstructures of the synthesized $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites were investigated by TEM, and the results are presented in Fig. 4. In

Fig. 4a, it can be clearly seen that the GO consists of transparent silky nanosheets with some wrinkles on their surfaces. In addition, the GO nanosheets are almost transparent, indicating that they are very thin. Fig. 4b shows the TEM micrograph of $\text{Fe}_2\text{O}_3/\text{GN}$, which clearly reveals that their shapes are irregular spheres with diameters from 300 to 500 nm, and the distribution is uneven. As shown in Fig. 4c and d, the GO nanosheets were decorated with large Fe_2O_3 nanoparticles *via* a hydrothermal process. Most of the Fe_2O_3 nanoparticles supported on the N-doped graphene sheets (N-GN) are about 40 to 50 nm, which is much smaller than the Fe_2O_3 nanoparticles on the graphene sheets (GN), and the fine Fe_2O_3 nanocrystals are uniformly decorated on the surfaces of the transparent and wrinkled graphene sheets without aggregation. All the Fe_2O_3 nanoparticles remained firmly attached to the N-GN sheets even

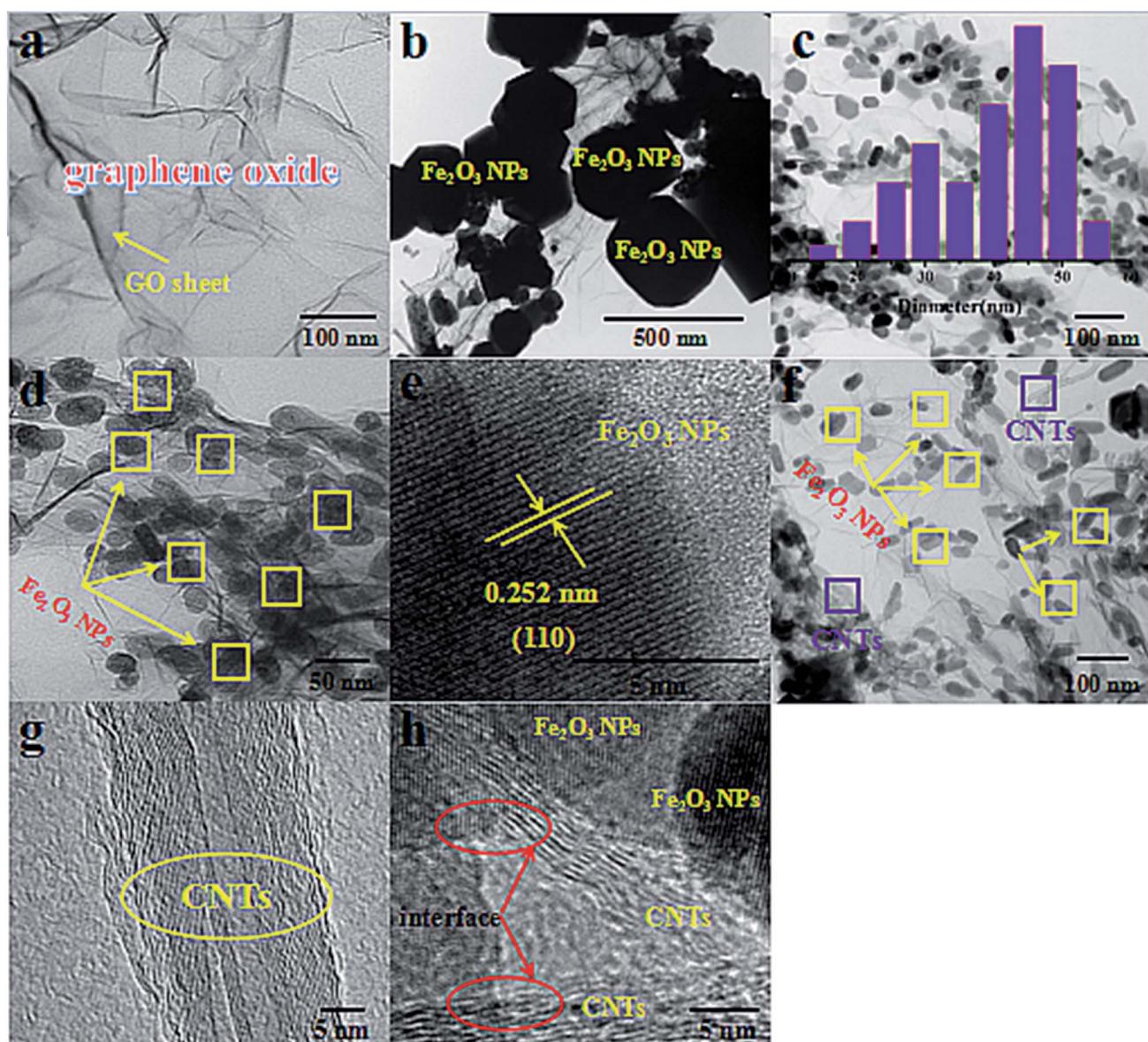


Fig. 4 TEM images of (a) GO (150000 \times), (b) $\text{Fe}_2\text{O}_3/\text{GN}$ (50000 \times) and (c) (100000 \times), (d) $\text{Fe}_2\text{O}_3/\text{N-GN}$ hybrid (150000 \times); (e) HRTEM image of the Fe_2O_3 lattice area of $\text{Fe}_2\text{O}_3/\text{N-GN}$ (1500000 \times); TEM and HRTEM images of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$: (f) (100000 \times), (g) (1000000 \times), (h) (1200000 \times).



after sonication was applied during the preparation of TEM specimens, indicating that strong adhesion between the Fe_2O_3 nanoparticles and N-GN sheets was obtained. The presence of oxygen-containing functional groups such as carboxyl groups ($-\text{COOH}$) and hydroxyl groups ($-\text{OH}$) on the surface of the GO allows iron ions to be adsorbed onto the surface by electrostatic attraction.³⁶ Moreover, GO interacted with the nitrogen material from urea; thus, N could be doped on the surface of GO, resulting in the formation of defects on the surface of the graphene.³⁷ These doped-N sites can potentially coordinate with iron ions in an Fe-N complex arrangement. This can improve the dispersion of ions and reduce agglomeration, which can indirectly control the size of magnetic particles. Thus, we can conclude that the doping of nitrogen atoms can lead to the formation of defects in the surfaces of the graphene sheets. These defects provide a large number of nucleation sites for the formation of Fe_2O_3 nanoparticles, thus suppressing the particle size of the Fe_2O_3 nanoparticles and enabling them to distribute more uniformly on the *in situ*-assembled N-doped graphene sheets.³⁸ In addition, Fig. 4e shows a lattice image of Fe_2O_3 ; an ordered array of crystal lattices with an interplanar spacing of 0.252 nm, belonging to the (110) plane, can be clearly seen, which is further supported by the XRD results given above. Furthermore, in contrast to $\text{Fe}_2\text{O}_3/\text{N-GN}$, upon the incorporation of CNTs, the internal structure clearly became more complex. Fig. 4f and g displays the TEM image of the as-prepared $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites. It is worth noting that CNTs with diameters of 20–40 nm are inserted into the N-doped graphene sheets and form a new type of three-dimensional conductive grid. The insertion of CNTs into two-dimensional N-GN sheets not only increased the spacing of the N-GN sheets, but also exposed their surface defects. More significantly, these defects can readily form collapses, unevenness and terraces, which are more conducive to the loss of electromagnetic waves.³⁹

In order to further analyze the intercalation effect of CNTs on N-GN lamellae and the distribution of Fe_2O_3 nanoparticles in the three-dimensional structure, the $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composite materials were studied by SEM (Fig. 5). As can be seen from Fig. 5a₁ and a₂, in the absence of doped nitrogen atoms, the Fe_2O_3 nanoparticles were larger and severely agglomerated. Compared with $\text{Fe}_2\text{O}_3/\text{GN}$, after doping with nitrogen atoms, the particle size of Fe_2O_3 obviously decreased and the distribution became uniform (see insets of Fig. 5b₁ and b₂). However, as can be seen from Fig. 5b₁ and b₂, without the addition of CNTs, although the Fe_2O_3 nanoparticles are evenly distributed on the nitrogen-doped graphene sheet layer, the doping of nitrogen atoms can result in defects and twisting of the two-dimensional graphene plane skeleton. That is, nitrogen atom doping can cause the edges of the nitrogen-doped graphene sheets to curl and overlap with each other, which is not conducive to the formation of electromagnetic wave channels. By comparison, it can also be found that the addition of CNTs can create a large amount of voids on the aggregates, effectively suppressing the overlap between the nitrogen-doped graphene layers (see insets of Fig. 5c₁ and c₂). As is proposed, when the microwaves pass through the surface of

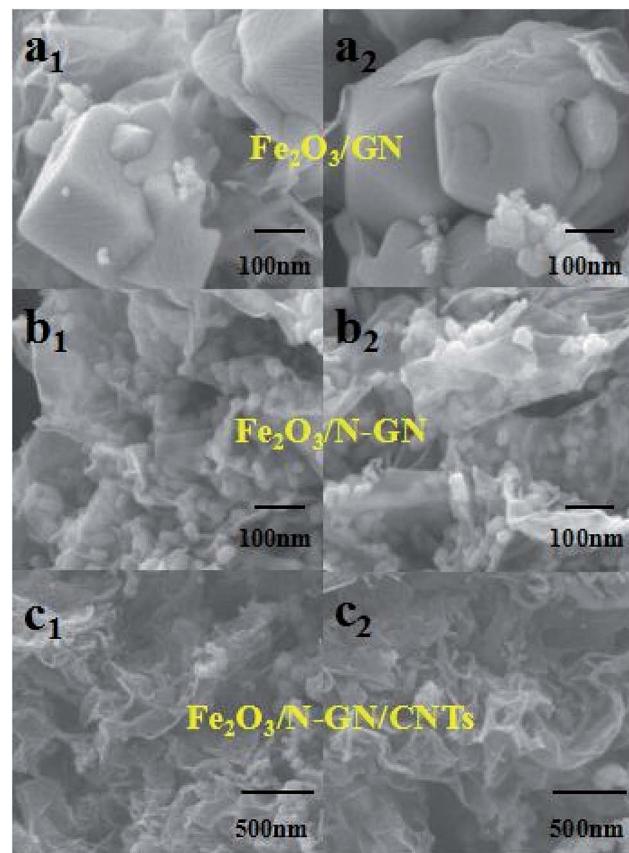


Fig. 5 SEM images of $\text{Fe}_2\text{O}_3/\text{GN}$ (a_{1,2}) (50000 \times), the $\text{Fe}_2\text{O}_3/\text{N-GN}$ composite (b_{1,2}) (50 000 \times) and the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composite (c_{1,2}) (20 000 \times) with different magnifications.

the material, this special heterogeneous interface may cause many electromagnetic waves to diffuse into the material, by which the electromagnetic waves would be effectively weakened.

Finally, the specific surface areas of the $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites were examined by nitrogen adsorption-desorption isotherms, as displayed in Fig. S2.† It can be seen that the introduction of CNTs causes the specific surface area and pore volume of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ to increase from $23.2 \text{ m}^2 \text{ g}^{-1}$ and $0.17 \text{ cm}^3 \text{ g}^{-1}$ to $45.5 \text{ m}^2 \text{ g}^{-1}$ and $0.28 \text{ cm}^3 \text{ g}^{-1}$, respectively. This indicates that the insertion of CNTs can effectively prevent overlap of the flexible nitrogen-doped graphene sheets. More importantly, a relatively high surface area is favorable for the formation of more dipoles, which may cause dipole polarization and improve microwave absorption performance.⁴⁰

3.2 Magnetic properties

All the above characterization results demonstrated that the 3D $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composite was successfully fabricated by a one-step solvothermal process. The 3D layered structure of the prepared $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites is expected to possess excellent electromagnetic wave absorption properties. To start, the magnetic properties of the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ and $\text{Fe}_2\text{O}_3/\text{N-GN}$ composites were studied at room temperature using

a vibrating sample magnetometer (VSM). Fig. S3† shows the typical magnetic hysteresis loops of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ and $\text{Fe}_2\text{O}_3/\text{N-GN}$, which are derived from the presence of magnetic iron oxide particles. The saturation magnetization (M_s), coercivity (H_c), and remnant magnetization (M_r) are 16.467 emu g⁻¹, 146.41 Oe, and 0.56 emu g⁻¹ for $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ and 19.437 emu g⁻¹, 175.7 Oe, and 3.54 emu g⁻¹ for $\text{Fe}_2\text{O}_3/\text{N-GN}$, respectively. It can be clearly seen that the samples show S-like shapes with strong coercivities (H_c) and negligible remnant magnetization (M_r). In addition, the difference in the M_s values of these composites may be attributed to the addition of non-magnetic CNTs.⁴¹

As is known, for an EM wave absorber, the electromagnetic wave absorption properties and electromagnetic parameters are closely related, *i.e.* complex permittivity ($\epsilon_r = \epsilon' - j\epsilon''$) and permeability ($\mu_r = \mu' - j\mu''$). The real parts of the complex permittivity (ϵ') and complex permeability (μ') represent the storage capabilities of the electric and magnetic energies, respectively. In addition, the imaginary parts (ϵ'') and (μ'') are associated with the dissipation energies of electromagnetic energy, respectively.⁴² Fig. 6 presents the electromagnetic

parameters of the composites containing 40% $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$, and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$; the EM absorption properties were measured at room temperature in the frequency range of 2 to 18 GHz. As shown in Fig. 6a, for these composites, the ϵ' values exhibit a rapid decline in the 2 to 18 GHz frequency range with slight fluctuations. This is due to the higher frequency of the delayed response of the dipole polarization electric field changes.⁴³ The ϵ' values decrease gradually from 10.53 at 2 GHz to 8.7 at 18 GHz. Moreover, $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ reveals the highest ϵ' values among the three samples. The higher ϵ' values for the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composite are associated with the addition of CNTs. The CNTs may enhance the electric conductivity and electric polarization because ϵ_r is a polarizability expression of the absorber. This is also beneficial to improve EM wave absorbing performance.

For the ϵ'' values (Fig. 6b), all of the samples showed similar trends and nonlinear phenomena. Namely, multiple peaks can be seen in the ϵ'' curves; it is worth noting that the ϵ'' curves of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ exhibit two remarkable resonance peaks in the range of 8 to 11 GHz. This can be ascribed to dipole polarizations arising from dielectric CNTs and interfacial

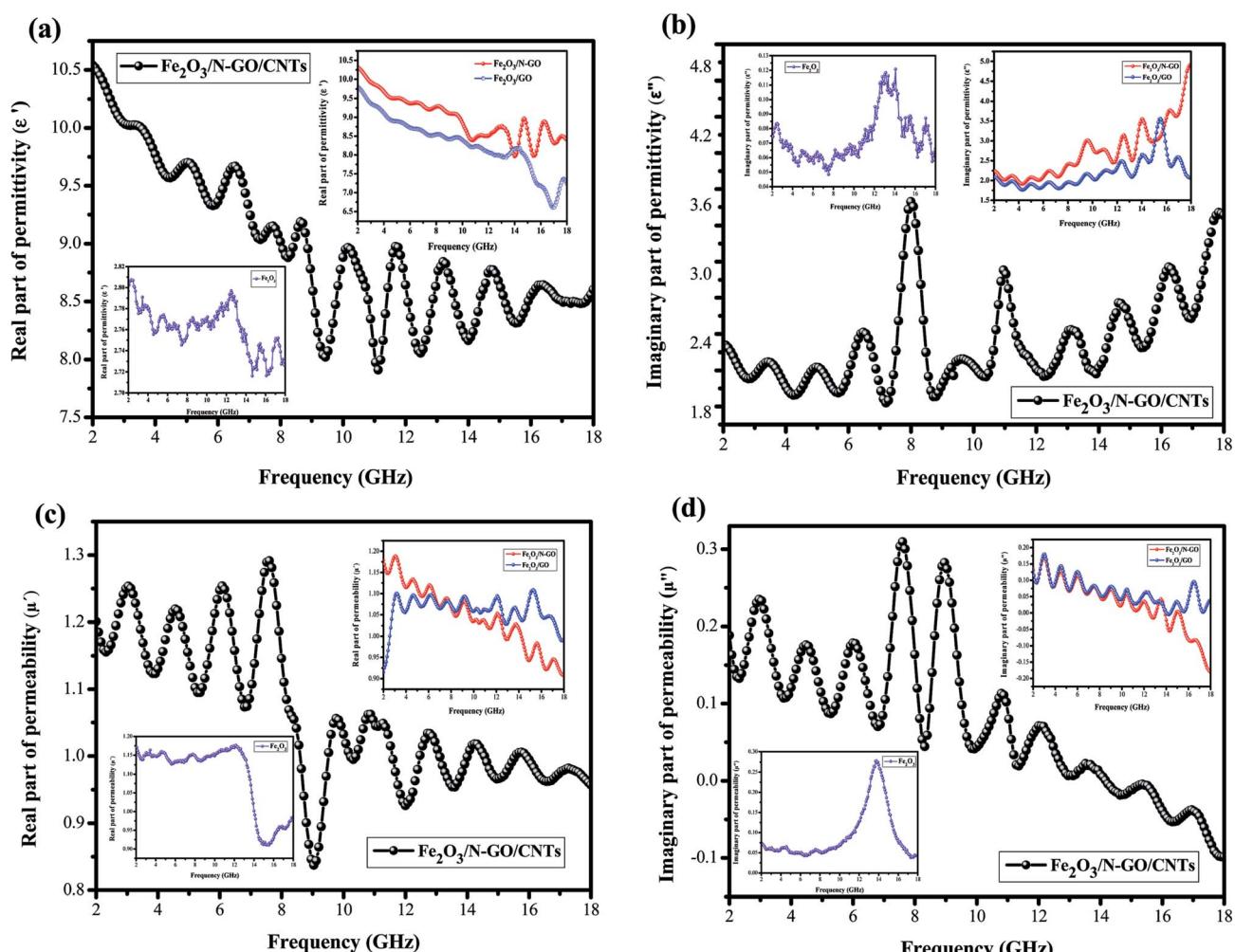


Fig. 6 Frequency dependence of the (a) real and (b) imaginary parts of the complex permittivity ((ϵ') and (ϵ'')) and the (c) real and (d) imaginary parts of the complex permeability ((μ') and (μ'')) for the samples.

polarization originating from the interface between nitrogen-doped graphite (N-GN) and the CNTs. This suggests that $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ have a strong dielectric loss against EM waves. In addition, the ϵ'' values of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ are higher than those of the $\text{Fe}_2\text{O}_3/\text{GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN}$ samples at frequencies under 8.2 GHz. According to free electron theory, $\epsilon'' = \sigma/2\pi\epsilon_0 f$, where σ is the conductivity; it can be found that the conductivity of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ is higher at frequencies under 8.2 GHz, resulting in electron loss under electromagnetic radiation.

Debye dipolar relaxation is an important mechanism for dielectric loss. According to the Debye equation, relative complex permittivity ϵ_r can be expressed as:⁴⁴

$$\epsilon_r = \epsilon' + i\epsilon'' = \epsilon_\infty + (\epsilon_s - \epsilon_\infty)/(1 + i\omega\tau_0) \quad (1)$$

where ϵ_∞ represents the relative dielectric permittivity at the high-frequency limit, ϵ_s is the static permittivity, τ is the relaxation time, and ω is the frequency. According to eqn (1), ϵ' and ϵ'' can be described by

$$\epsilon' = \epsilon_\infty + (\epsilon_s - \epsilon_\infty)/[1 + (\omega\tau_0)^2] \quad (2)$$

$$\epsilon'' = \omega\tau_0(\epsilon_s - \epsilon_\infty)/[1 + (\omega\tau_0)^2] \quad (3)$$

According to eqn (2) and (3), the relationship between ϵ' and ϵ'' can be deduced as

$$[\epsilon' - (\epsilon_s + \epsilon_\infty)/2]^2 + (\epsilon'')^2 = [(\epsilon_s - \epsilon_\infty)/2]^2 \quad (4)$$

$$\epsilon' = \epsilon''/(2\pi f\tau) + \epsilon_\infty \quad (5)$$

We can see that a plot of ϵ' versus ϵ'' should be a single semicircle, which is defined as a Cole–Cole semicircle. Each semicircle is associated with one Debye relaxation process.⁴⁵ Fig. 7a–c shows the $\epsilon'-\epsilon''$ patterns of the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$, and $\text{Fe}_2\text{O}_3/\text{GN}$ samples. It can be clearly seen from Fig. 7 that the plots of ϵ' versus ϵ'' for the composites shift to higher values with the addition of nitrogen and the intercalation of the CNTs. These findings suggest an enhanced contribution of Debye relaxation to dielectric loss.⁴⁶ In addition, as shown in the curve, the Cole–Cole semicircles are contorted, indicating that except for the loss effect that resulted from the Debye relaxation, the cooperative mechanism of Maxwell–Wagner relaxation, conductance loss, dipolar polarization and electron polarization may occur over the resulting composites with inserted CNTs and embedded Fe_2O_3 NPs.⁴⁷ Furthermore, the space charge polarization, known as interfacial polarization, often occurs at the heterogeneous interface between two components with different permittivities.⁴⁸ For the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites, the multiple phases led to the formation of abundant heterointerfaces between the N-GN, CNTs and Fe_2O_3 phases, producing strong interfacial polarization relaxation and accounting for the large dielectric loss.

Fig. 6c and d show the complex permeabilities of the three samples. The complex permeabilities of all the samples reveal decreasing trends with increasing measured frequency in the μ' curves; this can be attributed to Snoek's limitation in the GHz

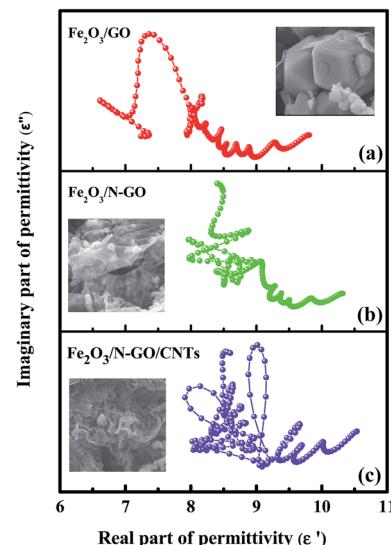


Fig. 7 $\epsilon'-\epsilon''$ curves of the $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites.

frequency range,⁴⁹ particularly for magnetic materials with relatively high conductivities. Therefore, it is not difficult to understand that the μ' values of $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ rapidly decrease above 8.1 GHz (Fig. 6c). For μ'' (Fig. 6d), two obvious resonance peaks can be observed at about 7 to 10 GHz; these are related to natural resonance.⁵⁰ In general, the magnetic loss mainly arises from hysteresis loss, domain wall resonance, natural resonance, exchange resonance and eddy current effects. The hysteresis loss in the weak field is negligible, and the domain wall resonance is usually located at a much lower range of the MHz frequency; thus, the magnetic loss caused by domain wall resonance and hysteresis loss can be eliminated. If the eddy current dominates the magnetic loss mechanism, the values of $C_o = \mu''(\mu')^{-2}f^{-1}$ should be constant when the frequency varies.⁵¹ As shown in Fig. 8, the value of C_o decreases rapidly at the lower range of 2 to 10 GHz; this is mainly caused by the natural resonance and exchange resonance. However, when the frequency exceeds 10 GHz, the values of $C_o = \mu''(\mu')^{-2}f^{-1}$ are nearly constant, suggesting that the composites have a significant eddy current effect at the higher range of 10 to 18 GHz. Thus, we can conclude that the contributions to magnetic loss result from natural resonance and that both exchange resonance and eddy current effects are crucial aspects.

As is known, the two possible contributions to electromagnetic wave absorption are dielectric loss and magnetic loss.⁵² In order to further study which contribution is dominant for the resulting composites, the dielectric loss tangents ($\tan \delta_e = \epsilon''/\epsilon'$) and magnetic loss tangents ($\tan \delta_m = \mu''/\mu'$) of the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$, and $\text{Fe}_2\text{O}_3/\text{GN}$ samples were measured (see Fig. S4†). Throughout the frequency range, the dielectric loss tangent ($\tan \delta_e = \epsilon''/\epsilon'$) is greater than the magnetic loss tangent ($\tan \delta_m = \mu''/\mu'$), which indicates that dielectric loss plays an important role in the EM absorption of these samples. Moreover, assuming that the magnetic loss is nearly equal to the



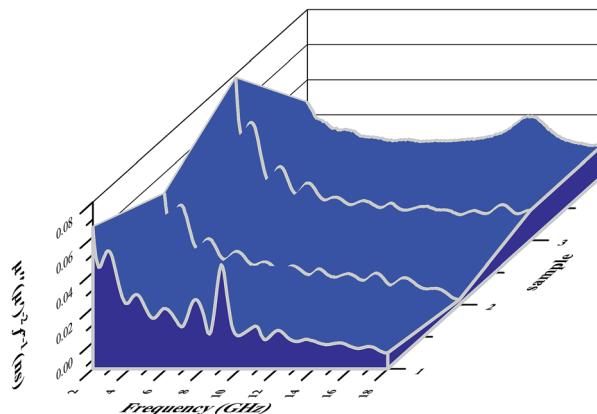


Fig. 8 The values of $\mu''(\mu')^{-2}f^{-1}$ of the $\text{Fe}_2\text{O}_3/\text{GN}$, $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites.

dielectric loss, the absorber will show great microwave absorption.⁵³ From Fig. S4,† it is worth noting that $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ has similar dielectric and magnetic losses, indicating that these composites may have excellent microwave absorption properties. In order to clarify the impedance matching phenomenon, we used the impedance matching ratio ($Z_r = Z_{\text{in}}/Z_0$) to explain the degree of impedance matching of the composites.⁵⁴

$$Z_r = |Z_{\text{in}}/Z_0| = (\mu_r/\epsilon_r)^{1/2} \tanh[j(2\pi f d)/c(\mu_r \epsilon_r)^{1/2}] \quad (6)$$

when Z_r is close to 1, most of the electromagnetic waves can easily enter the absorber and be converted into heat or dissipated by interference, enhancing the absorption performance. The frequency dependence of Z_r for the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composite was obtained, as shown in Fig. 9a. As can be seen, at a frequency of 9.4 GHz, the Z_r value is close to 1 and greater than those of the other two samples, indicating that $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ has the best impedance matching properties. This can cause more electromagnetic waves to enter the absorber, resulting in a reduction of reflection. In addition, the multiple reflections and scatterings between N-GN, CNTs and Fe_2O_3 are helpful for the attenuation of electromagnetic waves.

To obtain superior absorbers, in addition to considering good impedance matching, it is also necessary to enable electromagnetic waves to penetrate into the material, increasing the microwave attenuation. The dissipation characteristics of an absorbent material are usually expressed by the attenuation constant, α . Accordingly, we calculated the attenuation constants α of the three composites as follows:⁵⁵

$$\alpha = (\sqrt{2\pi f/c}) \times \sqrt{[(\mu''\epsilon'' - \mu'\epsilon')^2 + (\mu''\epsilon'' + \mu'\epsilon')^2]} \quad (7)$$

where f is the frequency of the electromagnetic waves and c is the velocity of the electromagnetic waves in free space. As shown in Fig. 9b, the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composite material has strong electromagnetic wave attenuation capability in the low frequency region. The α values of the composite materials in the whole frequency range increase in the sequence: $\alpha(\text{Fe}_2\text{O}_3/\text{N-GN}) > \alpha(\text{Fe}_2\text{O}_3/\text{GN}) > \alpha(\text{Fe}_2\text{O}_3/\text{N-GN/CNTs})$.

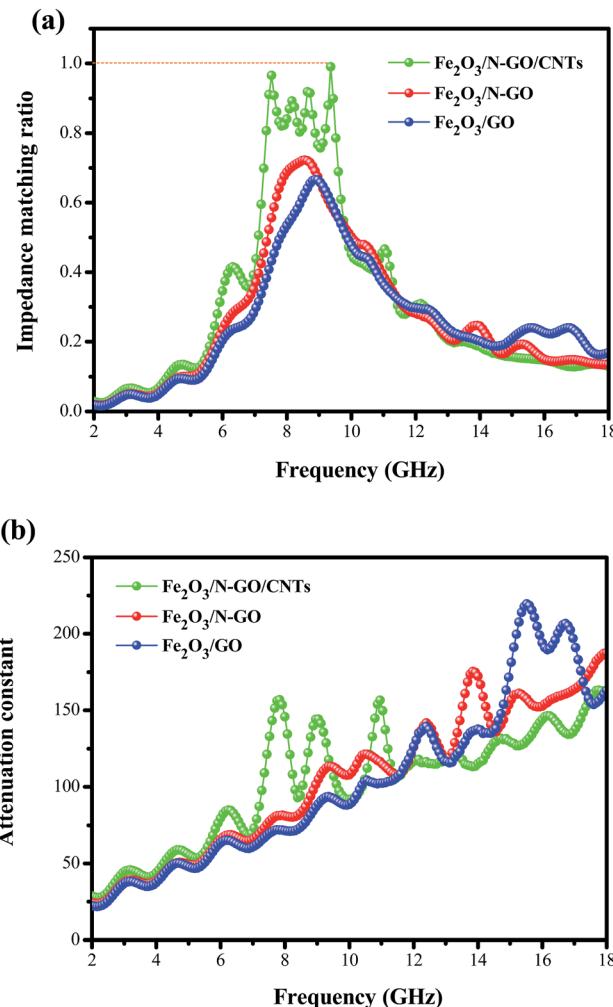


Fig. 9 Frequency dependence of the normalized input impedance (a) and attenuation constant (b).

$> \alpha(\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}) > \alpha(\text{Fe}_2\text{O}_3/\text{GN})$. Thus, taking into account its superior impedance matching and greater attenuation capability, it can be inferred that $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ may exhibit excellent electromagnetic wave absorption performance compared to the $\text{Fe}_2\text{O}_3/\text{N-GN}$ and $\text{Fe}_2\text{O}_3/\text{GN}$ samples.

To evaluate the electromagnetic wave absorption properties of the as-synthesized samples, the reflection loss (RL) values can be calculated from the measured relative complex permittivity (ϵ_r) and permeability (μ_r) dates for the given frequency range (2 to 18 GHz) and different absorber thicknesses; according to transmission line theory, the reflection loss RL (dB) can be calculated according to the following equations:⁵⁶

$$Z_{\text{in}} = Z_0(\mu_r/\epsilon_r)^{1/2} \tanh[j(2\pi f d)/c(\mu_r \epsilon_r)^{1/2}] \quad (8)$$

$$\text{RL} = 20 \log |(Z_{\text{in}} - Z_0)/(Z_{\text{in}} + Z_0)| \quad (9)$$

where Z_{in} stands for the input impedance of the absorber, Z_0 is the impedance of free space, ϵ_r and μ_r are the relative permittivity and permeability, d stands for the thickness of the absorber, f is the frequency of microwaves, and c is the velocity

of light. For most applications, MA materials need RL values less than -10 dB; this means that the material absorbs 90% of the incident microwaves. Fig. 10 shows the variation of RL values *versus* frequency for Fe_2O_3 (Fig. 10a), $\text{Fe}_2\text{O}_3/\text{GN}$ (Fig. 10b), $\text{Fe}_2\text{O}_3/\text{N-GN}$ (Fig. 10c), and $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ (Fig. 10d) at different thicknesses. Obviously, the RL values for pure Fe_2O_3 cannot reach -10 dB within a thickness range of 1.0 to 5.0 mm. The maximum RL value for the pure Fe_2O_3 nanocrystals is only -8.3 dB. The poor RL values indicate that Fe_2O_3 cannot be used for practical applications. In comparison with pure Fe_2O_3 , the $\text{Fe}_2\text{O}_3/\text{GN}$ paraffin composite showed better microwave absorption performance. The maximum reflection loss is -16.3 dB at 4.3 GHz with a thickness of 5.5 mm. More interestingly, when the Fe_2O_3 nanocrystals were decorated on the N-doped graphene sheets, forming the $\text{Fe}_2\text{O}_3/\text{N-GN}$ composite structure, the microwave absorption performance was significantly improved. For $\text{Fe}_2\text{O}_3/\text{N-GN}$ (Fig. 10c), RL values exceeding -10 can be obtained from 3.5 to 17.8 GHz when the absorber thickness is 1.0 to 5.5 mm, and the maximum reflection loss is -27.933 dB at 13.92 GHz with a thickness of only 2 mm. For $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ (Fig. 10d), an optimal RL of -45.8 dB is achieved at 9.32 GHz with a matching thickness of 3 mm. The effective absorption (below -10 dB) bandwidth reaches

14.5 GHz (3.5 to 18.0 GHz) for the absorber with a low thickness of 1.0 to 5.5 mm. Remarkably, it can be seen that when the CNTs were intercalated between the N-doped nano-sheets, the microwave absorption performance was further enhanced. Moreover, as shown in Fig. S5,[†] an interesting phenomenon was observed where the RL values are highly dependent on the sample thickness. For the $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$, it can be seen that the absorption peaks shift to the lower frequency range with increasing thickness. According to the quarter-wavelength match model, the relationship between the absorber thickness (t_m) and corresponding matching frequency (f_m) can be described by the following equation:⁵⁷

$$t_m = n\lambda/4 = nc/(4f_m\sqrt{\mu_r|\epsilon_r|}) \quad (n = 1, 3, 5\dots) \quad (10)$$

where λ is the wavelength of the microwaves, and c is the velocity of light in a vacuum. When t_m and f_m meet the above equation, the EM wave will be reflected from various interfaces with opposite phases. In this case, extinction of the overall EMW occurs on the absorber interface, and the maximum RL can be obtained. Moreover, to further optimize the microwave performance, we used the same method while changing the amounts of iron salt, urea and CNTs to analyze their impacts on

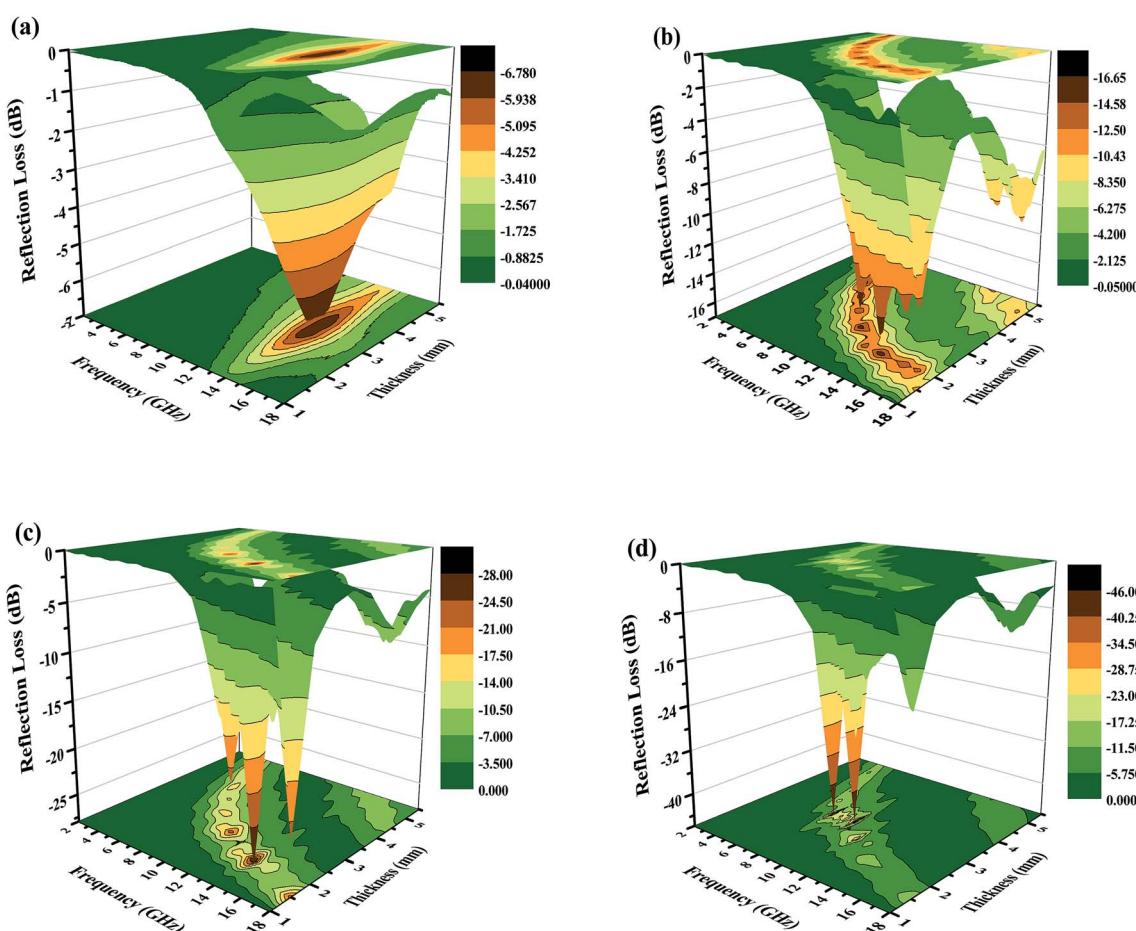


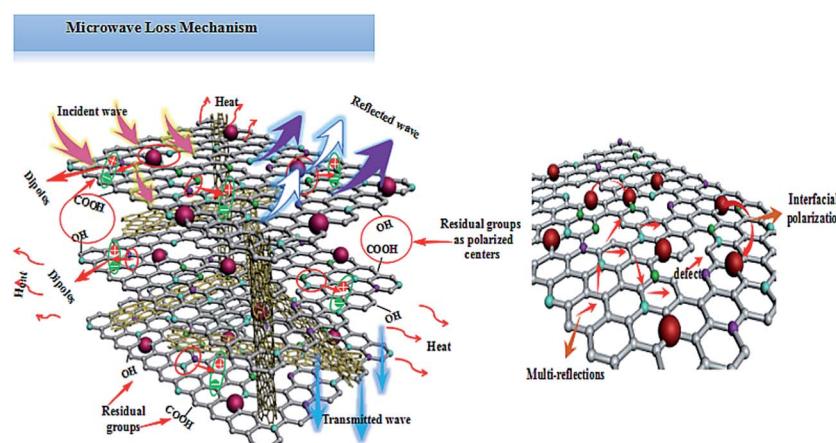
Fig. 10 Three-dimensional images of the calculated RL values of the (a) Fe_2O_3 , (b) $\text{Fe}_2\text{O}_3/\text{GN}$, (c) $\text{Fe}_2\text{O}_3/\text{N-GN}$ and (d) $\text{Fe}_2\text{O}_3/\text{N-GN/CNTs}$ composites.

the microwave performance. Fig. S6[†] shows the variation of the RL values *versus* frequency for 1.5a-Fe₂O₃/N-GN/CNTs (Fig. S6(A)[†]) (*a* refers to the initial amount of iron salt added), Fe₂O₃/N-GN/CNTs-1.5b (Fig. S6(B)[†]) (*b* refers to the initial amount of CNTs added), Fe₂O₃/1.5c-N-GN/CNTs (Fig. S6(C)[†]) (*c* refers to the initial amount of urea added) at different thicknesses. As can be seen from Fig. S6(A) and (B),[†] changing the amounts of iron salt and CNTs did not demonstrate large contributions to the absorbing properties. In addition, as shown in Fig. S6(C),[†] the values of RL for the Fe₂O₃/1.5c-N-GN/CNTs hybrids were obviously enhanced in comparison with those of the 1.5a-Fe₂O₃/N-GN/CNTs and Fe₂O₃/N-GN/CNTs-1.5b hybrids in the whole frequency band, and the values become larger with increasing N-GN content. However, compared with the Fe₂O₃/N-GN/CNTs hybrids, the value of RL for Fe₂O₃/1.5c-N-GN/CNTs was significantly lower than that for Fe₂O₃/N-GN/CNTs. Therefore, we can conclude that the appropriate ratio can exhibit excellent electromagnetic wave absorption performance. All the results suggest that the Fe₂O₃/N-GN/CNTs hybrids show highly improved microwave absorption performance in a wide frequency range. The enhancement of the microwave absorbing properties is mainly derived from the synergistic effect between the N-doped graphene nanosheets (N-GN), CNTs and magnetic Fe₂O₃ nanocrystals.

Based on the above analysis results, one may safely draw the conclusion that good impedance matching and strong attenuation capability are responsible for the superior microwave absorbing properties of the samples. Furthermore, the unique layered structure can enhance the electromagnetic wave absorption performance. Scheme 2 is a schematic of the microwave absorbing mechanism of the Fe₂O₃/N-GN/CNTs composite. First, the excellent microwave absorption properties are related to the unique 3D layered structure. Nitrogen-doped graphene (N-GN) sheets can act as a microwave receiver. When the electromagnetic waves infiltrate the N-GN sheets, they can be trapped by the inner void space, and the relatively large surface area and high pore volume can result in multiple reflections and diffuse scattering of the incident electromagnetic waves for long periods, converting electromagnetic

energy into heat energy or other forms of energy.⁵⁸ Second, the introduction of magnetic Fe₂O₃ particles in nitrogen-doped graphene (N-GN) sheets and CNTs can improve the characteristic impedance matching of the composite. Therefore, more EM waves can enter the absorbers and dissipate by the following mechanisms.⁵⁹ Third, according to the XPS and FTIR spectra, residual oxygen functional groups, such as hydroxyl and carbonyl groups, remain after the hydrothermal reaction; moreover, disordered carbon and defects exist in the nitrogen-doped graphene surfaces. These residual groups and defects act as polarized centers that can introduce defect polarization relaxation and relaxation processes, which all support microwave absorption.⁶⁰ Therefore, it is proposed that a novel absorption mechanism can be illustrated based on the layered structure; that is, the hierarchical structure provides a container which attenuates electromagnetic energy by extension through the travel path while gradually reducing energy.

Generally, these high-performance absorbers are required to have strong absorption, broad effective frequency (*f_E*), thinness and light weight. To evaluate the absorption properties of the Fe₂O₃/N-GN/CNTs composite, we list the reflection-loss properties of various GN-derived absorbers reported in previous references in Table 1. It was found that both the transition metal materials and the hybrids with GN had strong reflection loss (less than -30 dB), including RGO/SiO₂/Fe₃O₄,¹¹ RGO/ α -Fe₂O₃ (ref. 13) and GO/CNT-Fe₃O₄.¹⁴ Furthermore, the optimal reflection loss and effective bandwidth of Fe₂O₃/N-GN/CNTs in this work are very competitive with those of previously reported GN-derived composites. For example, compared with RGO/ α -Fe₂O₃ and GO/CNT-Fe₃O₄, Fe₂O₃/N-GN/CNTs shows strong microwave absorption ability and broader effective bandwidth, which are attributable to the introduction of nitrogen atoms and CNTs. However, compared with RGO/SiO₂/Fe₃O₄ and S-PPy/RGO, although the electromagnetic wave loss performance of Fe₂O₃/N-GN/CNTs is reduced, the effective bandwidth and the matching thickness of Fe₂O₃/N-GN/CNTs are even better than those of RGO/SiO₂/Fe₃O₄ and S-PPy/RGO. In addition, further comparisons revealed that the optimized GN-derived absorber obtained in this experiment has very strong absorption



Scheme 2 Possible mechanism for the microwave absorption of the Fe₂O₃/N-GN/CNTs composites.



Table 1 Electromagnetic wave absorption properties of reported and studied GN-derived composites

Sample	RL (dB)	Thickness (mm)	Frequency range (GHz)	Ref.
RGO/SiO ₂ /Fe ₃ O ₄	−56.4	4.5	3.5	11
S-PPy/RGO	−54.4	5.0	3.32	12
RGO/α-Fe ₂ O ₃	−38	2.0	5.8	13
GO/CNT-Fe ₃ O ₄	−37	5.0	2.1	14
CNT/RGO/BaFe ₁₂ O ₁₉	−19.03	2.5	3.8	61
3D GNS/INAW	−12.81	2.0	2.1	62
Ni-graphene	−13	2.0	2.6	63
Fe ₃ O ₄ -graphene	−26.4	4.0	2.0	64
GN-polyaniline	−36.9	3.5	5.3	65
Fe ₂ O ₃ /N-GN/CNTs	−45.8	3.0	3.7	This work

characteristics, very low thickness and wide absorption frequency ranges. In view of these advantages, the newly fabricated Fe₂O₃/N-GN/CNTs composite material can be considered to be a new high-performance electromagnetic wave absorber with great potential.

4. Conclusions

In summary, a new 3D Fe₂O₃/N-GN/CNTs composite with significantly enhanced comprehensive MA properties has been successfully synthesized *via* a facile *in situ* solvothermal method. Due to nitrogen doping into graphene nano-sheets, the Fe₂O₃ particles can be well dispersed within the composite structure without distinct aggregation; moreover, CNTs were successfully inserted into the two-dimensional lamellae of N-doped graphene to form a continuous conductive network and fully expose the lattice defects on the N-GN surface, which provides a large number of nucleation sites for Fe₂O₃ crystal growth *in situ*. When evaluated as a MA material, Fe₂O₃/N-GN/CNTs exhibited excellent MA properties in terms of maximum RL value, broad bandwidth, light weight and low thickness. The maximum RL of −45.8 dB can be achieved at 9.32 GHz with a matching thickness of only 3 mm. The effective absorption (below −10 dB) bandwidth reaches 14.5 GHz (3.5 to 18.0 GHz). These excellent MA properties can be mainly associated with the good impedance matching, strong loss characteristics, interfacial polarization and polarization relaxation of the material. Overall, it is believed that this newly designed Fe₂O₃/N-GN/CNTs composite with enriched interfaces and a controllable hierarchical structure will be a high-performance MA absorbing candidate for practical applications.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

Financial support from the National Natural Science Foundation of China (21546008, 21676039) and the Opening Foundation of State Key Laboratory of Inorganic Synthesis and

Preparative Chemistry of Jilin University (2016-04) are kindly acknowledged.

References

- V. K. Tyagi and S. L. Lo, *Renewable Sustainable Energy Rev.*, 2013, **18**, 288–305.
- R. Kumaran, K. S. Dinesh, N. Balasubramanian, M. Alagar, V. Subramanian and K. Dinakaran, *J. Phys. Chem. C*, 2016, **120**, 13771–13778.
- P. Saini, M. Arora, G. Gupta, B. K. Gupta, V. N. Singh and V. Choudhary, *Nanoscale*, 2013, **5**, 4330–4336.
- K. Jia, R. Zhao, J. C. Zhong and X. B. Liu, *J. Magn. Magn. Mater.*, 2010, **322**, 2167–2171.
- S. B. Ni, X. L. Sun, X. H. Wang, G. Zhou, F. Yang, J. M. Wang and D. Y. He, *Mater. Chem. Phys.*, 2010, **124**, 353–358.
- F. L. Wang, J. R. Liu, J. Kong, Z. J. Zhang, X. Z. Wang, M. Itoh and K. Machida, *J. Mater. Chem.*, 2011, **21**, 4314–4320.
- Y. Zhang, Y. Huang, T. F. Zhang, H. C. Chang, P. S. Xiao and Y. S. Chen, *Adv. Mater.*, 2015, **27**, 2049–2053.
- Y. Zhang, Y. Huang, H. H. Chen, Z. Y. Huang, Y. Yang and Y. S. Chen, *Carbon*, 2016, **105**, 438–447.
- J. J. Liang, Y. Wang, Y. Huang, Y. F. Ma, Z. F. Liu and J. M. Cai, *Carbon*, 2009, **47**, 922–925.
- K. Singh, A. Ohlan, V. H. Pham, S. Varshney, J. Jang and J. S. Chung, *Nanoscale*, 2013, **5**, 2411–2420.
- Y. J. Xu, Q. Wang, Y. F. Cao, X. F. Wei and B. Q. Huang, *RSC Adv.*, 2017, **7**, 18172–18177.
- F. Wu, A. M. Xie, M. X. Sun, Y. Wang and M. Y. Wang, *J. Mater. Chem. A*, 2015, **3**, 14358–14369.
- T. H. Wang, Y. F. Li, L. Wang, C. Liu, S. Geng, X. L. Jia, F. Yang, L. Q. Zhang, L. P. Liu, B. You, X. Ren and H. T. Yang, *RSC Adv.*, 2015, **5**, 60114–60120.
- L. N. Wang, X. L. Jia, Y. F. Li, F. Yang, L. Q. Zhang, L. P. Liu, X. Ren and H. T. Yang, *J. Mater. Chem. A*, 2014, **2**, 14940–14946.
- M. Mishra, A. P. Singh, B. P. Singh, V. N. Singh and S. K. Dhawan, *J. Mater. Chem. A*, 2014, **2**, 13159–13168.
- Z. S. Wu, W. C. Ren and L. Xu, *ACS Nano*, 2011, **5**, 5463–5471.
- L. F. Shen, X. G. Zhang and S. H. Li, *J. Phys. Chem. Lett.*, 2011, **2**, 3096–3101.
- M. Beidaghi and C. L. Wang, *Adv. Funct. Mater.*, 2012, **22**, 4501–4510.
- J. L. Lv, S. R. Zhai, C. Gao, N. Zhou, Q. D. An and B. Zhai, *Chem. Eng. J.*, 2016, **289**, 261–269.
- N. Zhou, Q. D. An, W. Zheng, Z. Y. Xiao and S. R. Zhai, *RSC Adv.*, 2016, **6**, 98128–98140.
- N. Zhou, Q. D. An, Z. Y. Xiao, S. R. Zhai and Z. Zhan, *ACS Sustainable Chem. Eng.*, 2017, **5**, 5394–5407.
- W. S. Hummers and R. E. Offeman, *J. Am. Chem. Soc.*, 1958, **80**, 1339.
- Z. X. Li, X. H. Li, Y. Zong, G. G. Tan, Y. Sun, Y. Y. Lan, M. He, Z. Y. Ren and X. L. Zheng, *Carbon*, 2017, **115**, 493–502.
- D. W. Wang, Y. Q. Li and Q. H. Wang, *J. Solid State Electrochem.*, 2012, **16**, 2095–2102.



25 B. Wang, J. S. Chen and X. W. Lou, *J. Am. Chem. Soc.*, 2011, **133**, 4940–4947.

26 A. Nasibulin, S. Rackauskas, H. Jiang, Y. Tian, P. Mudimela, S. Shandakov, L. Nasibulina, S. Jani and E. Kauppinen, *Nano Res.*, 2009, **2**, 373–379.

27 P. Su, H. L. Guo and S. Peng, *Acta Phys.–Chim. Sin.*, 2012, **28**, 1–9.

28 J. Feng, F. Z. Pu, Z. X. Li, X. H. Li, X. Y. Hu and J. T. Bai, *Carbon*, 2016, **104**, 214–225.

29 H. K. Jeong, Y. P. Lee, R. J. Lahaye, M. H. Park, K. H. An, I. J. Kim, C. W. Yang, C. Y. Park, R. S. Ruoff and Y. H. Lee, *J. Am. Chem. Soc.*, 2008, **130**, 1362–1366.

30 H. L. Guo, X. F. Wang, Q. Y. Qian, F. B. Wang and X. H. Xia, *ACS Nano*, 2009, **9**, 2653–2659.

31 X. L. Zheng, J. Feng, Y. Zong, H. Miao, X. Y. Hu, J. T. Bai and X. H. Li, *J. Mater. Chem. C*, 2015, **3**, 4452–4463.

32 J. J. Tang, J. Yang, X. Y. Zhou, H. M. Yao and L. M. Zhou, *J. Mater. Chem. A*, 2015, **3**, 23844–23851.

33 M. Chen, X. W. Yin, M. Li, L. Q. Chen, L. F. Cheng and L. T. Zhang, *Ceram. Int.*, 2015, **41**, 2467–2475.

34 X. H. Li, J. Feng, H. Zhu, C. H. Qu, J. T. Bai and X. L. Zheng, *RSC Adv.*, 2014, **4**, 33619–33625.

35 Y. Bo, L. L. Wang and L. Yao, *ChemComm*, 2013, **49**, 5016–5018.

36 Y. Nabae, S. Moriya, K. Matsubayashi, S. M. Lyth, M. Malon, L. B. Wu, N. M. Islam and Y. Koshigoe, *Carbon*, 2010, **48**, 2613–2624.

37 Z. Y. Lin, G. Waller, Y. Liu, M. L. Liu and C. P. Wong, *Adv. Energy Mater.*, 2012, **2**, 884–888.

38 M. Du, C. H. Xu and J. Sun, *Electrochim. Acta*, 2012, **80**, 302–307.

39 X. F. Zhnag, P. F. Guan and X. L. Dong, *Appl. Phys. Lett.*, 2010, **96**, 223111.

40 T. S. Wang, Z. L. Liu, M. M. Lu, B. Wen, Q. Y. Ouyang and M. S. Cao, *J. Appl. Phys.*, 2013, **113**, 024314–024321.

41 X. H. Li, J. Feng, H. Zhu, C. H. Qu, J. T. Bai and X. L. Zheng, *RSC Adv.*, 2014, **4**, 33619.

42 R. Qiang, Y. C. Du, H. T. Zhao, Y. Wang, C. H. Tian, Z. G. Li, X. J. Han and P. Xu, *J. Mater. Chem. A*, 2015, **3**, 13426–13434.

43 H. Wang, Y. Y. Dai, D. Y. Geng, S. Ma, D. Li, J. An, J. He, W. Liu and Z. D. Zhang, *Nanoscale*, 2015, **7**, 17312–17319.

44 S. He, G.-S. Wang, C. Lu, J. Liu, B. Wen, H. Liu, L. Guo and M.-S. Cao, *J. Mater. Chem. A*, 2013, **1**, 4685–4692.

45 P. H. Fang, *J. Chem. Phys.*, 1965, **42**, 3411–3413.

46 T. Wu, Y. Liu, X. Zeng, T. T. Cui, Y. T. Zhao, Y. N. Li and G. X. Tong, *ACS Appl. Mater. Interfaces*, 2016, **8**, 7370–7380.

47 X.-M. Meng, X.-J. Zhang, C. Lu, Y.-F. Pan and G.-S. Wang, *J. Mater. Chem. A*, 2014, **2**, 18725–18730.

48 D. Chuai, X. F. Liu, R. H. Yu, J. R. Ye and Y. Q. Shi, *Composites, Part A*, 2016, **89**, 33–39.

49 J. Snoek, *Physica*, 1948, **14**, 207–217.

50 B. Zhao, G. Shao, B. Fan, W. Guo, Y. Chen and R. Zhang, *Appl. Surf. Sci.*, 2015, **332**, 112–120.

51 C. Chen, Q. Liu, H. Bi, W. You, W. She and R. Che, *Phys. Chem. Chem. Phys.*, 2016, **18**, 26712–26718.

52 X. G. Liu, Z. Q. Qu, D. Y. Geng, Z. Han, J. J. Jiang, W. Liu and Z. D. Zhang, *Carbon*, 2010, **48**, 891–897.

53 H. Lv, G. Ji, W. Liu, H. Zhang and Y. Du, *J. Mater. Chem. C*, 2015, **3**, 10232–10241.

54 H. Lv, H. Zhang, J. Zhao, G. Ji and Y. Du, *Nano Res.*, 2016, **9**, 1813–1822.

55 X. M. Zhang, G. B. Ji, W. Liu, B. Quan, X. H. Liang, C. M. Shang, Y. Cheng and Y. W. Du, *Nanoscale*, 2015, **7**, 12932–12942.

56 G. Z. Wang, Z. Gao, S. W. Tang, C. Q. Chen, F. F. Duan, S. C. Zhao, S. W. Lin, Y. H. Feng, L. Zhou and Q. Y. Qin, *ACS Nano*, 2012, **6**, 11009.

57 X. M. Zhang, G. B. Ji, W. Liu, X. X. Zhang, Q. W. Gao, Y. C. Li and Y. W. Du, *J. Mater. Chem. C*, 2016, **4**, 1860–1870.

58 L. N. Wang, X. L. Jia, Y. F. Li, F. Yang, L. Q. Zhang, L. P. Liu, X. Ren and H. T. Yang, *J. Mater. Chem. A*, 2014, **2**, 14940–14946.

59 Y. C. Yin, X. F. Liu, X. J. Wei, R. H. Yu and J. L. Shui, *ACS Appl. Mater. Interfaces*, 2016, **8**, 34686–34698.

60 H. L. Xing, Z. F. Liu, L. Lin, L. Wang, D. X. Tan, Y. Gan, X. L. Ji and G. C. Xu, *RSC Adv.*, 2016, **6**, 41656–41664.

61 T. K. Zhao, X. G. Ji, W. B. Jin, C. Wang, W. X. Ma, J. J. Gao, A. L. Dang, T. Li, S. M. Shang and Z. F. Zhou, *RSC Adv.*, 2017, **7**, 15903–15910.

62 T. K. Zhao, W. B. Jin, X. L. Ji, J. J. Gao, C. Y. Xiong, A. L. Dang, H. Li, S. M. Shang and Z. F. Zhou, *RSC Adv.*, 2017, **7**, 16196–16203.

63 Y. Cao, Q. Su, R. Che, G. Du and B. Xu, *Synth. Met.*, 2012, **162**, 968–973.

64 X. Sun, J. He, G. Li, J. Tang, T. Wang, Y. Guo and H. Xue, *J. Mater. Chem. C*, 2013, **1**, 765–777.

65 X. Chen, F. Meng, Z. Zhou, X. Tian, L. Shan, S. Zhu, X. Xu, M. Jiang, L. Wang, D. Hui, Y. Wang, J. Lu and J. Gou, *Nanoscale*, 2014, **6**, 8140–8148.

