Research progress of Mn doped phosphors

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In this review article, Mn applications have been divided into three parts. We provide an overview of recent progress in developing Mn4+ doped red phosphors for promising application in warm white light-emitting diodes. In addition, we summarize reports on Mn2+ co-doped Eu2+ or Ce3+ phosphors that can produce white light with outstanding color rendering index (CRI) and color coordinates via energy transfer processes, and report the application of Mn2+ in green to NIR phosphors.

1 Introduction

The valence electron configuration of transition metal elements is \((n-1)\)d\(^{10}\)ns\(^{2}\), with the common characteristic that the d orbitals are not full. Since the energy of the \((n-1)\)d orbital is similar to that of the ns orbital, and the d electrons can completely or partially participate, there are no significant changes in the chemical properties from left to right in the same cycle. In addition, the transition metal elements can lose s electrons as well as d electrons; hence, the same element has a variety of valence states. Mn\(^{2+}\) and Mn\(^{4+}\) ions are commonly used as light-emitting centers. Nowadays, phosphor-converted white light emitting diodes (pc-WLEDs) based solid-state lighting has received much attention, due to the characteristics of high efficiency, high power density, low applied voltage, reliability, long operation life and environmental friendliness, compared to other traditional lighting technologies, such as incandescent, halogen, xenon and fluorescent lamps. In order to obtain highly efficient WLED devices, several aspects should be taken into account and optimized, including semiconducting components, phosphors and packaging technologies. Phosphors are indispensable components of pc-WLED devices, which can decide the color rendering index, correlated color temperature, rated average life and luminous flux. One method for obtaining these materials is to combine yellow-emitting phosphors with highly efficient (Ga,In)N blue LEDs. For blue chips, Y\(_3\)Al\(_2\)O\(_{12}\):Ce\(^{3+}\) (YAG:Ce) is the commonly used phosphor; however, due to the lack of the red component of YAG:Ce phosphor, this approach has a poor color rendering index and a high correlated color temperature, which cannot fulfill the requirements for general illumination. According to Blasse, Mn\(^{4+}\) absorbs in the entire ultraviolet region and subsequently emits phosphorescence in the deep red (620–680 nm) region, due to transitions from \(^2\)E \(\rightarrow\) \(^4\)A\(_2\), \(^{10}\) which indicates that the phosphor doped Mn\(^{4+}\) can be a good candidate for red fluorescent powder. It has been observed that multi-phase phosphor converted systems commonly lead to the strong reabsorption of the blue light by the red and green phosphors. Moreover, such a device is very complicated and expensive, and it is difficult to control the color balance. To solve the above-mentioned problems, single-emitting component phosphors with UV LEDs chips, which have excellent color rendering indexes and high luminous efficiencies, have been researched. One of the way to generate white light from single-composition phosphors is by co-doping the sensitizer and activator into a crystalline matrix, with the principle of energy transfer from sensitizer to activator, such as Eu\(^{3+}/Mn^{2+}\), Ce\(^{3+}/Mn^{2+}\). Various white light based single-composition phosphors have obtained, such as CaAl\(_2\)Si\(_2\)O\(_6\):Eu\(^{3+}\),Mn\(^{2+}\),\(^{11}\) Ba\(_3\)MgSi\(_2\)O\(_6\):Eu\(^{2+}\),Mn\(^{2+}\),\(^{12}\) KMg\(_4\)\((PO\(_4\))_3:Eu\(^{2+}\),Mn\(^{2+}\),\(^{13}\) BaSi\(_3\)O\(_4\)N\(_2\):Eu\(^{2+}\),Mn\(^{2+}\),\(^{14}\) Ca\(_3\)Gd\(_8\)(SiO\(_4\))\(_6\):Ce\(^{3+}\)/Mn\(^{2+}\),\(^{15}\) and Ca\(_3\)La(GeO\(_4\))\(_6\)PO\(_4\):Ce\(^{3+}\),Mn\(^{2+}\).\(^{16}\) Mn\(^{2+}\) is influenced by the crystal field, and hence it can easily change luminescence from green to deep red and even near-infrared (NIR). Because the near infrared light has the advantages of large penetration depth, less interference, etc., it has become an important topic in the field of optical fiber communication, solid state lasers, light emitting markers and fluorescence immunoassays. The applications of Mn ion doped phosphors are shown in Fig. 1.
2 Application of Mn$^{4+}$

2.1 Introduction to red phosphors

To date, the most practical phosphor is the blue LED chip excited by YAG:Ce$^{3+}$. In order to get white-light emission and a high CRI, improvements must be made by adding a high efficiency red emitting phosphor to the system to solve the problem of the lack of a red component; hence, several red phosphors have been developed. Red fluorescent powder can be divided into three categories: linear emission red phosphor excited by rare earth ions, broadband emission red phosphor excited by rare earth ions, and red phosphor excited by transition metals. With their own characteristics, excellent red phosphors can be obtained by a combination of these.

Linear-emission phosphors excited by the rare ions Eu$^{3+}$ and Sm$^{3+}$ have pure color and high luminous efficiency; however, there is no wide excitation spectrum in the blue and UV regions, which results in a lower luminous efficiency. It is feasible to broaden the excitation spectrum and improve the luminous efficiency, although the effect is not ideal, but this limits the practical application of linear phosphors. Therefore, the application of red emitting Sm$^{3+}$ or Eu$^{3+}$ ions doped phosphors is limited in the present W-LEDs, due to their sharp absorption peaks in the blue and UV region.

In order to create a warm WLED, the development of efficient red phosphors is necessary. The new phosphors are mainly based on rare earth-doped schemes, such as Eu$^{2+}$-doped nitrides, aluminates, oxynitrides and aluminosilicates. Among them, oxynitrides and nitrides have been the subject of maximal studies, due to their excellent performances, e.g., chemical stability and high thermostability; however, these materials must be synthesized at high temperatures and pressures. The strict preparation conditions and high price of raw materials lead to this phosphor being quite expensive. In addition, the reabsorption of the coexistent green emission of WLEDs is not avoided for Eu$^{2+}$-doped red phosphors with blue absorption.

2.2 Luminescence properties of transition metal ions

Transition metal ions have unfilled d orbitals, with the electronic configuration of 3d$^n$ (0 < n < 10). The Hamilton operator is as follows:

$$ H = H_0 + H_{CF} + H_C + H_{SO} $$

where $H_0$ includes the kinetic energy of electrons and the Coulomb energy of the nucleus acting on the electron. $H_{CF}$ is the Coulomb energy between electrons, $H_{SO}$ is the spin–orbit coupling energy and $H_C$ is the energy of the crystal field.

The 3d electron is in the outer layer, which is sensitive to the crystal environment; hence the $H_{CF}$ and the Coulomb energy between electrons ($H_C$) can be compared. The crystal-field spectral term usually has two types of energy representations with weak crystal field and strong crystal field. We always apply the strong crystal field to deal with problems. The first consideration is how the d orbital is affected by the tetrahedral and octahedral crystal fields. The crystal field causes the five degenerate d orbitals to split into the doubly degenerate $e_g$ and triply degenerate $t_{2g}$. The distance between $e_g$ and $t_{2g}$ is denoted as $10Dq$, and $Dq$ is a parameter to characterize the field strength of the octahedral.

With the single electron orbital wave function of $e_g$ and $t_{2g}$ as basic functions, considering the impact of $H_C$, a new intrinsic function is characterized by $2S^+1I$, where $S$ is the total spin, and there is an irreducible representation of the octahedral group of the intrinsic state. The $2S^+1I$ spectral term is similar to that used in the study of the free ion $2S^+1L$ spectra of rare earth elements. The spin orbit interaction $H_S$ in transition metal ions is weaker than that of rare earth ions, which is usually treated by the perturbation theory and the deviation of the crystal field in the octahedral symmetry. These interactions cause $2S^+1I$ to split into multiple states. The complete theoretical treatment of the transition metal ions eigenstate was given by S. Sugano and Y. Tanabe. Direct coulombic interactions ($H_{Cdir}$) between electrons are characterized by Racah parameters $A$, $B$, $C$

$$ (H_C)_{dir} = \sum_k f_k F^{(k)} $$

where the fitting parameter

$$ F^{(k)} = e^2 \int_0^\infty \int_0^\infty \frac{R_{C}(r_i)R_{C}(r_j)}{r_{ij}^{k+1}} \, dr_i \, dr_j $$

For the $ll'$ configuration, the $k$ in $F^{(k)}$ is an even number from 2 to min (2l, 2l'). For transition metal ions with 3d electrons, the values can be 0, 2, 4. The relationship between Racah parameters and $F^{(k)}$ is as follows:

$$ B = \frac{1}{49} F^{(2)} - \frac{5}{441} F^{(4)} $$

$$ C = \frac{35}{441} F^{(4)} $$

$$ A = F^{(0)} - \frac{1}{9} F^{(4)} $$

Since the $A$ parameter only affects the energy levels, we only consider $B$ and $C$. The ratio of $B/C$ for the all the transition metal ions is almost equal to 4.5, which can be seen as approximately constant. The distance between different spectra of each ion depends on the two parameters, $Dq$ and $B$. In the work of S. Sugano and Y. Tanabe, the Tanabe–Sugano (T–S) (Fig. 2) diagram shows the relationship between the energy levels of metal ions E/B and crystal field strength Dq/B of the d$^2$–d$^6$ configurations in the crystal field. From the T–S graph, we can identify each excitation peak corresponding to energy levels in the excitation spectra.

2.3 The transition of Mn$^{4+}$ in the host

Mn$^{4+}$ doped materials, such as oxides and fluorides, are commonly used as red emitting phosphors. It was discovered that the Mn$^{4+}$ ion can only be stabilized in an octahedral lattice site where it commonly shows strong and broad adsorptions
ranging from 300 to 480 nm and emits red light between 600 and 700 nm with 5 sharp peaks. As the lowest energy state \(^2E(t_{2g})\) hardly changes in different crystal fields, the luminescence properties are very similar from host to host, but intensely depend on the covalence of the Mn\(^{4+}\)-ligand bonding. Moreover, Mn is magnetic and will produce local magnetic ions when introduced into luminescent materials; therefore, Mn may have huge potential application value.

Doped in a solid, the Mn\(^{4+}\) ion has a strong crystal field, due to its highly effective positive charge. Consequently, the spin-forbidden \(^2E_g \rightarrow \;^4A_{2g}\) transition (sharp line) always controls the emission spectrum of Mn\(^{4+}\), while two broad bands corresponding to the \(^4A_{2g} \rightarrow \;^4T_{1g}(^6F)\) and \(^4A_{2g} \rightarrow \;^4T_{2g}\) spin-allowed transitions are observed in the excitation spectra. The third spin-allowed transition (\(^4A_{2g} \rightarrow \;^4T_{1g}(^4P)\)) frequently remains under cover by the host lattice absorption and the charge transfer. The \(^2E_g \rightarrow \;^4A_{2g}\) transition has a wide variation, which is due to the energy of the \(^2E_g\) state in the \(d^3\) electronic configuration being independent of crystal field (Fig. 3).

### 2.4 Spectroscopic parameters of Mn\(^{4+}\) in different hosts

The positions of zero-phonon lines in the \(^2E_g \rightarrow \;^4A_{2g}\) spin-forbidden transition of the Mn\(^{4+}\) ion, 'host cation ligand' bond distance and Racah parameter \(B\) have been summarized numerous crystals. On an average, the distance from Mn\(^{1+}\) to ligand in the fluoride crystals (about 1.67–1.86 Å) is the shortest and the following trends (Fig. 4) are formulated from the graphical representations of the data gathered in Table 1, which is suitable for most compounds. The \(B\) values are greater in fluorides and thus, the emission lines of \(^2E_g \rightarrow \;^4A_{2g}\) are significantly increased in oxides and in materials with the garnet structure (up to 2.11 Å), as shown in Table 1. The energy difference among the \(^4A_{2g} \rightarrow \;^4T_{2}\) and the \(^4A_{2g} \rightarrow \;^4T_{1}\) transitions and the Racah parameter \(B\) can be derived from the expression

\[
\frac{Dq}{B} = \frac{E(4A_{2g} \rightarrow 4T_{1g}) - E(4A_{2g} \rightarrow 4T_{2g})}{Dq} \quad (7)
\]

where \(x\) is defined as

\[
x = \frac{15(x - 8)}{x^2 - 10x} \quad (8)
\]

The Racah parameter \(C\) can be evaluated by the following expression, according to the peak energy of \(^2E_g \rightarrow \;^4A_{2g}\) obtained from the emission spectrum.

\[
\frac{E(2E_g \rightarrow 4A_{1g})}{B} = 3.05 \frac{C}{B} + 7.9 - 1.8 \frac{B}{Dq} \quad (9)
\]

The local crystal-field strength \(Dq\) is determined by the mean peak energy of the \(^2A_2 \rightarrow \;^4T_2\) transition according to the following expression:
The $^{2}E_g$ level is always lower for oxides; however, the value of $B$ is also close to that of fluorides in some special cases (Fig. 4). What causes the condition is that the position of the $^{2}E_g$ level also depends on the value of the second Racah parameter $C$. Such a distinct and explicit separation of fluorides and oxides with respect to the energy position of the $^{2}E_g$ level indicates the potential role of the Mn$^{4+}$ ion as a reliable indicator of covalency/ionicity of a particular host. On the other hand, Mn$^{4+}$ luminescence exhibits sharp emission lines and broad excitation bands in the ultraviolet–blue region, which makes it suitable for excitation by UV and decreases the re-absorption effect when mixing with other yellow and green phosphors. With so many unique features, Mn$^{4+}$ can meet the requirement of an ideal red phosphor for warm LEDs. Moreover, the easy to obtain and cheap manganese raw materials are able to reduce the red phosphors. As a consequence, many researchers have focused on developing the Mn$^{4+}$ activated red phosphors (Table 2).

### 2.5 Examples of Mn$^{4+}$ doped red phosphors

Numerous researchers have focused on the oxide host, due to their outstanding performance such as high stability. Peng et al. have synthesized an efficient red oxide phosphor Sr$_{4}$Al$_{9}$O$_{25}$:Mn$^{4+}$, which is prepared by a traditional solid-state reaction in much milder conditions. The red emission intensity of Sr$_{4}$Al$_{9}$O$_{25}$:Mn$^{4+}$ is greater than that of the reported commercial phosphors, such as (K,Rb)$_{2}$Ge$_{4}$O$_{9}$:Mn$^{4+}$ (ref. 29) and CaAl$_{12}$O$_{19}$:Mn$^{4+}$.

A series of novel red emitting Mg$_{3}$Ga$_{2}$GeO$_{8}$ under near-UV excitation was successfully synthesized by Ding et al. with the traditional high-temperature solid-state reaction. In this Mg$_{3}$Ga$_{2}$GeO$_{8}$ system, there is one tetrahedral site, which is occupied by Ge$^{4+}$/Ga$^{3+}$, and an octahedral site, occupied by Mg$^{2+}$/Ga$^{3+}$. According to radius analysis and luminescence properties of Mg$_{3}$Ga$_{2}$GeO$_{8}$:x%Mn$^{4+}$, we can conclude that Mn$^{4+}$ reliably occupies the Mg$^{2+}$/Ga$^{3+}$ site in this system. When using MnCO$_{3}$ as the raw material, the intensity of the emission is relatively weak, compared to the sample of MnO$_{2}$ as the raw material and the peak shape does not change. In the MnCO$_{3}$, the same Mn$^{2+}$ can just be oxidized to Mn$^{3+}$, rather than Mn$^{4+}$, as shown in Fig. 5. The emission properties of Mg$_{3}$Ga$_{2}$GeO$_{8}$:x%Mn$^{4+}$ are also shown in Fig. 5. It can be easily seen that the emission intensities reach a maximum at $x = 0.5$% and the calculated full-width at half-maximum (FWHM) of the Mn$^{4+}$ emission spectrum is 24 nm, which indicates that the Mg$_{3}$Ga$_{2}$GeO$_{8}$:x%Mn$^{4+}$ phosphor...
possesses high color purity. Furthermore, Mg$_3$Ga$_2$GeO$_8$:0.5\% Mn$^{4+}$ exhibits outstanding quantum efficiency (64.7\%).

The oxide matrix with strong covalent, low thermal vibration and weak polarizability is favourable for red light emission. The Mn$^{4+}$ doped Y$_3$Al$_5$O$_{12}$ phosphor was prepared by a high temperature solid state reaction method. Huang et al. synthesized the Y$_3$Al$_5$O$_{12}$:Mn$^{4+}$ red phosphor and explored its luminescence properties. At 467 nm, this phosphor emits red light.\(^\text{31}\) The main emission peaks are at 643 and 670 nm, all from the $^3E \rightarrow ^1A_2$ transition of Mn$^{4+}$ in Y$_3$Al$_5$O$_{12}$:0.01Mn$^{4+}$. The CIE of the Y$_3$Al$_5$O$_{12}$:0.01Mn$^{4+}$ phosphor is (0.701, 0.299) located in the red area. The excitation spectrum of the emission peak at 670 nm is composed of 400, 450, 468, 474, 482 and 494 nm wide excitation bands.

Novel red-emitting phosphors K$_2$BaGe$_8$O$_{18}$:Mn$^{4+}$ were synthesized via the high temperature solid-state reaction method.\(^\text{32}\) The excitation spectrum of K$_2$BaGe$_8$O$_{18}$:Mn$^{4+}$ monitored at 666 nm shows a wide band with two evident peaks at 319 and 468 nm. Upon 468 nm excitation, the emission spectrum presents a narrow band from 600 to 700 nm (red region) with the maximum emission at around 666 nm, which originates from the $^3E_g \rightarrow ^1A_2$ transition of Mn$^{4+}$ in the GeO$_6$ octahedral environment, as shown in Fig. 6; the CIE coordinates of the phosphors are (0.702, 0.298) and the QY is 32.9% upon 365 nm excitation. The results indicate that this kind of phosphor can be a supplement of red-components in UV and blue light excited w-LEDs.

In the oxide matrix, Mn$^{4+}$ ions replace cations such as Al$^{3+}$, Si$^{4+}$, Zr$^{4+}$, Ti$^{4+}$, Ge$^{4+}$, and occupy the center of the octahedron. When Mn$^{4+}$ replaces ions with different valences, charge compensation is required. Therefore, the emission intensity of the phosphor can be enhanced by performing charge compensation. Mg$^{2+}$ has been very extensively researched as a highly efficient charge compensator. The substitution of the Mg$^{2+}$–Mg$^{2+}$ ion pair to replace the Al$^{3+}$–Al$^{3+}$ ion pair was made under the condition that no additional charge compensation and oxygen atoms are required. Ge$^{4+}$, Ca$^{2+}$, Li$^+$, Na$^+$, K$^+$, Cl$^-$ can also be used as charge compensators to enhance the luminescence properties of red phosphors.\(^\text{33,34}\) Different charge compensators have different effects on the red luminescence properties of Mn$^{4+}$.

The chemical stability of fluorides is appropriate, and the phonon energy is relatively low, so fluorides can be used as the host material for various luminescent rare earth ions. Na$_2$SiF$_6$:Mn$^{4+}$ red fluoride phosphors have been synthesized by the exothermic reduction reaction method, as shown in Fig. 7.\(^\text{36}\) The phosphors have intense blue-light absorption and strong narrow-band red emission. When the Na$_2$SiF$_6$:Mn$^{4+}$ red phosphor was fabricated, a prominent improvement in the CRI of white LEDs was observed. The Na$_2$SiF$_6$:Mn$^{4+}$ phosphors have efficient red luminescence excited by blue light, which is a potential candidate for improving the color reproducibility of white LEDs with simple post-processing, a short reaction time and high yield.

Mn$^{4+}$-activated fluoride phosphors, such as A$_2$BF$_6$ (A = Na, K, Rb, Cs; A$_2$ = Ba, Zn; B = Si, Ge, Ti, Zr, Sn), have attracted...
considerable attention as highly promising red phosphors for warm white light emitting diodes (w-LEDs). These fluoride phosphors have been synthesized via traditional chemical routes with HF-solution. In addition to the possible dangers of hypotoxic HF, the uncontrolled precipitation of fluorides and the extensive processing steps produce large morphological variations, limiting its application. Recently, Song made a prototype w-LED with K3AlF6:Mn4+ as the red light component, via an efficient and water-processable cation-exchange green route. Combined with YAG:Ce3+, the prototype of the sample shows an efficient luminous efficacy beyond 190 lm W−1, along with an excellent color rendering index (CRI, Ra = 84) and a lower correlated color temperature (CCT = 3665 K), as shown in Fig. 8.

Thus far, Mn4+ as an activator doped oxide/fluoride red phosphor has attracted significant attention, due to its special advantages such as relatively simple synthesis methods and low cost. The spectral positions produced by the 2Eg → 4A2g transition can be modified, and the excitation region of the phosphor is located in the ultraviolet and blue band. Therefore, the Mn4+ doped red phosphor can be used as a substitute for rare earth doped red phosphors, and has a bright future in the field of warm white LEDs.

Although a large number of Mn4+ doped red phosphors have been fully studied, there are still some questions worthy of further exploration. For the Mn4+ doped oxide host phosphor, it is desirable to achieve a blue shift of emission from deep red to orange/red and enhanced absorption in the blue wavelength region. One possible solution is doping to increase the absorption in the blue wavelength region for Sr4Al14O25:Mn4+ doped Mg2+; because the neighboring Mn4+–Mn4+ is replaced by Mn4+–Mg2+, the non-radiative decay rate of the 2Eg state decreases and the luminous intensity of the phosphor is enhanced. For fluoride phosphors, how to improve their thermal stability and resistance to humidity needs further study. Surface treatment by external coating can improve the thermal stability and moisture resistance. Finally, the quantum efficiency of materials can be further improved by optimizing synthesis methods and reducing the lattice defects of materials.

3 Application of Mn2+

3.1 Application of Mn2+ in green to NIR phosphors

Most phosphors applied in white LEDs are oxides, sulfides or nitrides doped with rare-earth ions, which play a significant role in controlling the color of the light. Most of these rare-earth ions are very expensive and some chloride and oxide doped rare-earth elements are toxic and harmful, which limit their
development. Consequently, the search for a cheap and nontoxic phosphor has become an important research area and Mn$^{2+}$ doped oxides and nitrides have become hot research topics. With the d electrons within the 3d shell, Mn$^{2+}$ emission can be affected by the crystal field, which results in the emission changing from green to red. When Mn$^{2+}$ is located in the tetrahedral crystal field, it can emit green light. In other words, through changing the crystal field conditions, the spin-forbidden Mn$^{2+}$ forbidden d–d transition ($^4T_1 \rightarrow ^4A_1$) can emit green light.

Rare earth ions (Eu$^{2+}$, Eu$^{3+}$, or Pr$^{3+}$) or transition metal elements (Cr$^{3+}$ or Mn$^{2+}$) can serve as luminescence centers in NIR phosphors. However, for biological applications, these hygroscopic sulphides of Eu$^{2+}$ show a succession of problems in metal-activated oxide compounds, such as Cr$^{3+}$ or Mn$^{2+}$-doped phosphors, which have outstanding chemical stability and suitable chemical instability. Lately the focus has shifted to the transition-emission range in the optical window (from 650 to 1300 nm) and long-lasting persistent luminescence. With the advantages of high signal to noise ratio, deep tissue penetration, and maximum transparency for silicon-based optical fibers and so on, near infrared (NIR) emitting materials have distinct potential for medical diagnostics, laser systems, and optoelectronic devices. On the basis of coordination and the T–S diagram, due to weaker crystal symmetry (angular constraints) and stronger ligand field strength (interionic distance) than that of the four/six-coordinated environments, it is feasible to speculate that the five/seven-coordinated environments of the Mn$^{2+}$ ion are good for NIR emission generation. There is a complete list of known green to NIR persistent phosphors activated by Mn$^{2+}$ in Table 3.

The rare-earth-free narrow-band green-emitting KAlSi$_2$O$_6$:Mn$^{2+}$ phosphor has been synthesized by Ding et al. The excitation and emission spectra are shown in Fig. 9. The five distinct peaks of PL excitation are in good agreement with the Mn$^{2+}$ absorption transitions; specifically, it can match blue GaN chips, which make its application in white LEDs possible. Moreover, a green light emission peak at 513 nm with full-width at half-maximum of 30 nm can be seen from Fig. 9, which is attributed to the spin forbidden d–d transition ($^4T_1 \rightarrow ^4A_1$) of Mn$^{2+}$. The electrons could be relaxed from these excited states to the $^4T_1(G)$ state by a non-radiative relaxation process and then be transferred back to the ground state $^4A_1(S)$, emitting the characteristic green (513 nm) light. The CIE color coordinates of the KAlSi$_2$O$_6$:Mn$^{2+}$ phosphor were calculated to be $x = 0.27$ and $y = 0.64$, which are much closer to the standard green coordinates (0.21, 0.71). The integrated emission intensity of KAS:3% Mn$^{2+}$ decreased to 67.6% (250 °C) of the initial value (25 °C). The chromaticity coordinates of the fabricated white LED are (0.35, 0.36) with CCT of 4775 K. It is demonstrated that the KAlSi$_2$O$_6$:Mn$^{2+}$ phosphor could be simulated to blue-LED for producing efficient white-light.

Moreover, good performance Mn$^{2+}$ doped Na$_2$MgGeO$_3$ phosphors have been reported. Under UV or electron beams, the samples show bright green emission, which is ascribed to the $^4T_1(4G) \rightarrow ^4A_1(6S)$ transition of the Mn$^{2+}$ ions. Based on this transition, the following phosphors were developed:

$$I(t) = I_0 + A \exp(-t/\tau)$$

(11)

where $I(t)$ and $I_0$ represent the emission intensity at time $t$ and $0$, respectively, $A$ is a constant, and $\tau$ is the decay lifetime. For Mn$^{2+}$-heavy doping ($x = 0.10–0.30$), the visible emission exhibited non-exponential decay behavior, probably associated with the prevailing non-radiative transition, which is expressed by the mean decay lifetime ($\tau_m$) estimated by

$$\tau_m = \int_{t_0}^{t_\infty} I(t)/I_{\max} dt$$

(12)

where $I(t)$ is the luminescence intensity at $t$ and $I_{\max} = I(0)$. The non-radiative loss occurring among the Mn$^{2+}$ ions as the Mn$^{2+}$ concentration increases can be demonstrated from the monotonously shortened decay lifetime of the two emission bands.

### Table 3 Known compounds of green to NIR persistent phosphors activated by Mn$^{2+}$

<table>
<thead>
<tr>
<th>Host material</th>
<th>Incorporated site</th>
<th>Emission region (nm)</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>CdSiO$_3$</td>
<td>Cd$^{2+}(v i)$</td>
<td>550–720</td>
<td>50</td>
</tr>
<tr>
<td>MgGeO$_3$</td>
<td>Mg$^{2+}(v i)$</td>
<td>600–720</td>
<td>51</td>
</tr>
<tr>
<td>Li$_2$GeO$_3$</td>
<td>Zn$^{2+}(v i)$</td>
<td>700–900</td>
<td>52</td>
</tr>
<tr>
<td>Ga$_3$(PO$_4$)$_3$</td>
<td>Ca$^{2+}(v i)$</td>
<td>650–750</td>
<td>53</td>
</tr>
<tr>
<td>Ga$_3$Ln(PO$_4$)$_3$</td>
<td>Ca$^{2+}(v i)$</td>
<td>650–750</td>
<td>54</td>
</tr>
<tr>
<td>(Mg$_{1-x}$Zn$_x$)$_3$(PO$_4$)$_2$</td>
<td>Mg$^{2+}/Zn$^{2+}(v i)$</td>
<td>630–800</td>
<td>55</td>
</tr>
<tr>
<td>KMgBO$_3$</td>
<td>Mg$^{2+}(v i)$</td>
<td>600–700</td>
<td>56</td>
</tr>
<tr>
<td>MgAl$_2$O$_4$</td>
<td>Mg$^{2+}(v i)$</td>
<td>500–560</td>
<td>57</td>
</tr>
<tr>
<td>Na$_2$MgGeO$_4$</td>
<td>Mg$^{2+}(v i)$</td>
<td>500–550</td>
<td>58</td>
</tr>
<tr>
<td>KAlSi$_2$O$_6$</td>
<td>Al/Si$^{2+}(v i)$</td>
<td>550–540</td>
<td>59</td>
</tr>
<tr>
<td>Zn$_2$GeO$_4$</td>
<td>Zn$^{2+}(v i)$</td>
<td>500–570</td>
<td>60</td>
</tr>
</tbody>
</table>

$\text{Na}_2\text{Mg}_1-x\text{Si}_2\text{O}_6:x\text{Mn}^{2+}$ green phosphors, with CIE coordinates of (0.163, 0.631) and CCT value of 8643 K, NaAl$_2$O$_4$:Mn$^{2+}$ green phosphors, with CIE chromaticity coordinates of (0.0725, 0.6468), and quantum efficiency of 67.84%, with excitation of 424 nm, Zn$_2$SiO$_4$:Mn$^{2+}$ green phosphors, which exhibit widely tunable CCTs ranging from 3200 to 7000 K and color rendering indices up to 85%, depending on the mixture ratio. Those phosphors with good performance indicators can be candidates for high power peLED applications. In addition, MgAl$_2$O$_4$Mn$^{2+}$7–17 Zn$_2$SiO$_4$:Mn$^{2+}$ green phosphors have also been reported.

Song et al. have reported the realization of single-band NIR UC emission in a Mn$^{2+}$-doped KZnF$_3$ nanostructure by a heavy doping strategy. The emission spectra of KZn$_{1-m}$Mn$_m$F$_3$ upon 396 nm light excitation. There is a single visible emission located at 585 nm when Mn$^{2+}$ concentration is less than $x = 0.10$, corresponding to the $^4T_{1g}(G) \rightarrow ^4A_{1g}(S)$. An anomalous NIR emission band located at 770 nm appears when $x = 0.10$, while at $x = 0.30$, there is a sharp decrease due to the concentration quenching effect, where visible and NIR emission bands are featured from two different emission centers. The luminescence decay curves (Fig. 10(c) and (d)) were studied in order to obtain more insight into the active centers. The visible emission exhibited single exponential decay at a low concentration ($x = 0.01$ and 0.05), and the decay lifetime was estimated by the following equation:

$$\tau_m = \int_{t_0}^{t_\infty} I(t)/I_{\max} dt$$

(12)
To study the reason for the observed unusual NIR emission in KZn\(_{1-x}\)Mn\(_x\)F\(_3\), the crystal structure of KZnF\(_3\) and the activation process of Mn\(^{2+}\) were further analyzed. KZnF\(_3\) has a typical cubic perovskite structure with a Pm\(_{3}\) space group with its lattice constant \(a = 0.405\) nm, as shown in Fig. 10(a). Zn\(^{2+}\) is substituted by the incorporated Mn\(^{2+}\) ions, leading to the formation of Mn\(^{2+}\)(Zn\(^{2+}\))–Mn\(^{2+}\)(Zn\(^{2+}\)) dimmers, especially at a high doping concentration. The ground state (\(6A_{1g}(S)\)) and the first excited state (\(4T_{1g}(G)\)) are split due to the exchange interaction between Mn\(^{2+}\) and Mn\(^{2+}\) ions, which are mainly governed by the linear combination of two spin states, \(S_i\) and \(S_j\) (i.e., \((S_i + S_j)\), \((S_i - S_j)\)). Therefore, the spin state is \(S = 5, 4, 3, 2, 1, 0\) for the ground state (\((6A_{1g}(S)\)) and \(S = 4, 3, 2, 1\) for the first excited state (\((4T_{1g}(G))\)). The Mn\(^{2+}\)–Mn\(^{2+}\) dimmers can also serve as the chromophoric centers, and the observed NIR luminescence is most likely related to the radiative transition from the excited state (\(S = 1\)) to the ground state with various spin components. There could be a shift in the emission peak caused by the changes in the spin components. The reason for decay lifetime shortening may be that the spin selection rule has changed from spin forbidden (\(S = 5/2 \rightarrow S = 3/2\)) to spin allowed (\(S = 1 \rightarrow S = 1\)), hence, it is reasonable to suggest that the NIR emission in KZn\(_{1-x}\)Mn\(_x\)F\(_3\) originates from the \(6A_{1g}(S)\)\textsuperscript{4+}\textsubscript{G}(G) \(\rightarrow 6A_{1g}(S)\)\textsuperscript{4+}\textsubscript{G}(G) transitions of coupled Mn\(^{2+}\)–Mn\(^{2+}\) dimmers.

Near-infrared (NIR) emitting Li\(_2\)ZnGe\(_3\)O\(_8\) (LZG):\(x\)Mn\(^{2+}\) has been synthesized by a conventional solid state reaction.\(^\text{39}\) The excitation spectrum (PLE) and emission spectrum (PL) of the LZG:\(x\)Mn\(^{2+}\) (\(x = 0.005, 0.01, 0.03, 0.05, 0.07, 0.10, 0.15, 0.20\)) phosphors are presented in Fig. 11(a) and (b), respectively. The emission spectrum shows a broad emission band from 650 to 900 nm, with the main peak at 832 nm. The luminescence properties can be strongly affected by the crystal field within the 3d shell with an increase in Mn\(^{2+}\) concentration. The PL (\(\lambda_{\text{em}} = 832\) nm) intensity rapidly increases, while above this maximum intensity (0.03), a concentration quenching phenomenon occurs, as shown in Fig. 11(d). Essentially, quenching impurities cause concentration quenching, a decline in the emission intensity, as shown in Fig. 11. Non-radiative energy transfer, which include exchange interactions, radiation
reabsorption and multipolar interactions between Mn$^{2+}$ ions within a certain distance, lead to concentration quenching. According to calculations, the concentration quenching of Mn$^{2+}$ ions originates from quadrupole–quadrupole interactions. Besides, the phosphor has a good thermal stability, which expands the scope of its application.

Energy transfer between ions is a rather common and very important physical phenomenon, mainly through physical processes such as collision, energy exchange, and re-absorption of radiation and radiationless transition processes, which cause the ion energy transfer to other ions. Due to the relatively rich energy levels of transition metal ions, especially in the crystals, the energy splitting is affected by the influence of the crystal field, thus becoming more intensive. At this point, there may be the possibility of energy differences between the two levels of the transition metal ions, equal to the difference between other ions. In order to achieve the energy transfer between ions, radiative or radiationless transition processes occur very easily under the action of multipole moments.

Radiationless transitions play a significant role in photo-physics, photobiology and photochemistry, and include internal conversion (IC) (spin-allowed process) and intersystem crossing (ISC) (spin-forbidden process). A simplified Jablonski diagram is presented in Fig. 12, which describes the elementary processes of molecular luminescence.\textsuperscript{68}

Absorption
\[
S_0 + h\nu \xrightarrow{k_s} S_1
\]  \hfill (13)

Fluorescence
\[
S_1 \xrightarrow{k_s} S_0 + h\nu'
\]  \hfill (14)

Internal conversion
\[
S_1 \xrightarrow{k_{ic}} S_0
\]  \hfill (15)

Intersystem crossing
\[
S_1 \xrightarrow{k_{isc}} T_1
\]  \hfill (16)
latter is called the activator.Overlap between the emission spectra of the sensitizer and the excitation spectra of the activator is a prerequisite for the occurrence of resonant transmission. Resonant transmission is a quite significant means of energy transfer in the material activated by rare earth or transition metal elements and organic crystals. In this way, there are two centers S and A, where the center initial state is $S^* + A$, i.e., the S center is in the excited state, the A center is in the ground state. The final state is $S + A$, i.e., the S center is back to the ground state, and the A center is in the excited state. Taking S and A as the dipoles, the transition probability from the initial state to the final state is the $S, A$ resonance energy transfer probability $P_{SA}$:

$$P_{SA} = \left(\frac{R_0}{R}\right)^6 \times \frac{1}{\tau_S}$$

where $R$ is the distance between the center of S and A and $\tau_S$ is the measured lifetime of $S^*$ state, and the $R_0$ is defined as:

$$R_0^6 = \frac{3}{64\pi^2} \times \frac{h^2 e^4}{K} \sigma_A \eta_A \int_0^\infty E^3 \left(\frac{\epsilon_S(E)}{E}\right) dE$$

where $h$ and $e$ are the Planck constant and the velocity of light, respectively, and the $K$ is the dielectric constant of the material. $\sigma_A$ is the total A center absorption cross section and $\eta_A$ is the emission efficiency of the S center. $\epsilon_S(E)$ and $\sigma_A(E)$ are the emission spectrum of S center and absorption spectrum of A centre, respectively (Fig. 13).

### 3.2 The basic properties of Mn$^{2+}$

In the last few years, non-rare earth-based eco-friendly phosphors prepared under milder conditions have received increasing interest. The outer electronic structure of Mn$^{2+}$ is 3d$^5$4s$^2$. On the basis of the Mn$^{2+}$ Tanabe–Sugano diagram, the emission of Mn$^{2+}$ can be easily changed from green to deep red by transforming the crystal field strength. Because the transition is $^4T_1 \rightarrow ^6A_h$, Mn$^{2+}$ has broad band emission in the visible range, depending on the crystal field. Octahedral coordinated Mn$^{2+}$ shows an orange to red emission, while tetrahedral

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**Fig. 11** (a) PLE and (b) PL spectra of the LZG:xMn$^{2+}$ samples ($0 < x < 0.2$). (c) A diagram of the concentration quenching process.$^{52}$

**Fig. 12** Elementary photophysical processes of molecular luminescence.$^{68}$

**Phosphorescence**

$$T_1 \xrightarrow{k_{ts}} S_0 + h\nu$$

**Intersystem crossing**

$$T_1 \xrightarrow{k_{ta}} S_0$$

These elementary processes determine whether a molecule can be a good candidate for a dye molecule for organic solar cells or LED. Besides, there are two other significant elementary processes described in the following:

**Phonon-induced electron transfer**

$$DA \xrightarrow{h\nu} D^*A \rightarrow D^+A^-$$

**Phonon-induced energy transfer**

$$DA \xrightarrow{h\nu} D^*A \rightarrow DA^*$$

All the processes involve two electronic states, i.e., they are non-adiabatic in nature. Vibrational relaxation (VR) exists in molecular luminescence as a non-radiative process, which can happen in either the ground or excited electronic states. The radiative transfer process is also called resonant transmission, if there is a near field force interaction between the two centers and if the energy of light radiation transfer can be one atom to 10 nm, without the help of other nearest neighbor atoms. Dexter first introduced the transfer mechanism to the energy transfer process between the centers of the luminous material, and derived a formula for the probability of resonance energy transfer between the centers. There is an assumption that emission from the ion overlaps another ion, absorbing light, so the light emitting radiation energy is absorbed by another ion, and energy transfer occurs. The former is called the energy sensitizer and the

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coordinated Mn$^{2+}$ shows green emission. The luminescence spectrum of Mn$^{2+}$ consists of a structureless band at peak wavelengths of 490–750 nm, with the half width of 1000–2500 cm$^{-1}$. In general, the crystal field at a tetrahedral site is weaker than that of an octahedral site. When Mn$^{2+}$ is in a weak crystal field, the splitting of the excited energy in orbitals is small, resulting in a higher energy of Mn$^{2+}$ emission, but its energy is low. Generally, when Mn$^{2+}$ is substituted for alkaline-earth metals (Ga$^{2+}$, Sr$^{2+}$, Ba$^{2+}$), the larger lattice distortion is not conducive to the direct excitation of Mn$^{2+}$, which makes it difficult to achieve direct luminescence. Therefore, Mn$^{2+}$ is not suitable for individually serving as an activator with low luminescence efficiency.

Though Mn$^{2+}$ ions emit red light from d $\rightarrow$ d transitions, the 4T$_1$(4G) $\rightarrow$ 6A$_1$(6S) emission in the 3d electron configuration is weak because it is forbidden, which leads to weak absorption in the ultraviolet region and disparate fluorescence properties in different hosts. However, on introducing Eu$^{3+}$ or Ce$^{3+}$ as co-activators, the excitation of Mn$^{2+}$ is covered by the emission of Eu$^{3+}$ or Ce$^{3+}$, which creates the necessary conditions for vibrational transmission, i.e., there is overlap between the emission of the activator and the excitation of the sensitizer. Meanwhile, there may be energy transfer between the sensitizer and activator. On the basis of the resonance energy transfer mechanism, the excitation energy of the sensitizer can induce the activator to emit fluorescence, while the fluorescence intensity decays. White light can be obtained by co-doping Ce$^{3+}$/Mn$^{2+}$ or Eu$^{3+}$/Mn$^{2+}$ under effective resonance-type energy transfer; a schematic diagram of energy transfer is shown in Fig. 14.\(^{72-74}\)

3.3 The co-doping of Ce$^{3+}$/Mn$^{2+}$

The outer electronic structure of Mn$^{2+}$ is 3d$^5$, which presents typical d $\rightarrow$ d transitions when excited. Its photoluminescence range changes from 500 to 700 nm, which depends on the crystal field environment; however the d $\rightarrow$ d transition of Mn$^{2+}$ is forbidden by the selection rules, and the luminescence is very weak since it can only be indirectly excited using sensitized ions or the substrate itself for energy transfer. A large number of

Fig. 13 Schematic diagram of the energy transfer, the dark and light colors indicate the level of energy.

Fig. 14 (a) The schematic energy level diagram for the energy transfer process in the BaMgF$_4$:Ce$^{3+}$, Tb$^{3+}$ phosphors.\(^{75}\) (b) Energy level diagram of KNaSO$_4$:Ce,Dy.\(^{79}\) (c) Energy transfer diagram of Ce$^{5+}$ $\rightarrow$ Mn$^{2+}$.\(^{82}\) (d) Schematic energy-level diagram of Ce$^{5+}$ and Eu$^{2+}$.\(^{85}\)
studies have been conducted on the mechanism of energy transfer, e.g., Ce$^{3+}$ → Tb$^{3+}$, Ce$^{3+}$ → Mn$^{2+}$, Ce$^{3+}$ → Dy$^{3+}$, Ce$^{3+}$ → Eu$^{2+}$, etc., as shown in Table 4, which indicates that the Ce$^{3+}$ ion is a highly efficient ion. The energy level diagram of the energy transfer about Ce is shown in Fig. 14.

The outermost electron configuration of Ce$^{3+}$ is 4f$^1$, which can change to a lower energy level of 5d when absorbing energy. The 5d electronic excited state is different from the 4f electronic state, which is easily shielded by 5s$^2$5p$^6$, resulting in the crystal field having a great influence on it, so the 4f–5d transition of Ce$^{3+}$ in the 200–400 nm spectral region often exhibits strong absorption and excitation spectra.

Ca$_{2}$Gd$_{8}$(SiO$_4$)$_6$O$_2$:Ce$^{3+}$,Mn$^{2+}$ has been prepared by Li et al., and the emission spectra are shown in Fig. 15 for different concentrations of Ce$^{3+}$ and Mn$^{2+}$. The phosphors emit blue to white light, and with the increased Mn$^{2+}$ concentration, there was complete conversion to yellow light when excited by UV irradiation, due to the 5d–4f transition of Ce$^{3+}$ and the $^4T_{1g}$ → $^6A_{1g}$ transition of Mn$^{2+}$, respectively. It is noteworthy that the emission peaks of Mn$^{2+}$ have a slightly red shift from 564 nm to 570 nm. The Ce$^{3+}$ ion simultaneously occupies the 4f and 6h sites in the Ca$_2$Gd$_8$(SiO$_4$)$_6$O$_2$ host, and gives different blue emission under different UV excitation. Energy transfer from Ce$^{3+}$ → Mn$^{2+}$ also exists in the phosphor, hence, via precise control of the contents of Mn$^{2+}$ and Ce$^{3+}$, a wide-range white emission can be obtained. On the other hand, changing excitation sources is another way to obtain the color-tunable white emission. It has been shown that for the crystal structure of NaCaBO$_3$, the relative PL intensity of blue to yellow emission increases with the increase of the excitation wavelength from 287 to 330 nm. Moreover, Ca$_3$Gd$_5$(SiO$_4$)$_5$O$_2$:Ce$^{3+}$,Mn$^{2+}$ was also studied, which have potential applications in the field of white EDS. It has been shown that for the crystal structure of NaCaBO$_3$, the relative PL intensity of blue to yellow emission increases with the increase of the excitation wavelength from 287 to 330 nm. Moreover, Ca$_3$Gd$_5$(SiO$_4$)$_5$O$_2$:Ce$^{3+}$,Mn$^{2+}$, has a good CL coefficient (K/S) relation, which was used to calculate the measured reflectance ($R$) for NaCaBO$_3$, to probe the absorption of the host lattice.

$$F(R) = \frac{(1 - R)^2}{2R} = \frac{K}{S} \quad (23)$$

where $R$, $K$, and $S$ represent the reflectance, absorption coefficient, and scattering coefficient, respectively. Hence, the band gap of the NaCaBO$_3$ host was calculated to be approximately 4.48 eV by extrapolation, essentially owing to the transition between the valence band and the conduction band in this host.

From Fig. 16(1) and (2), it is seen that there exists a spectral overlap between the PL spectrum of NaCa$_0.98$BO$_2$:0.01Ce$^{3+}$ and the PLE spectrum of NaCa$_0.98$BO$_2$:0.03Mn$^{2+}$, indicating that a resonance-type energy transfer from the Ce$^{3+}$ to Mn$^{2+}$ ions is expected in the co-doped sample. Due to the energy transfer

### Table 4 The energy transfer about Ce$^{3+}$ in different hosts and doping ions

<table>
<thead>
<tr>
<th>Host</th>
<th>Energy transfer</th>
<th>$\lambda_{ex}$ (nm)</th>
<th>$\lambda_{em}$ (nm)</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>BaMgF$_4$</td>
<td>Ce$^{3+}$ → Tb$^{3+}$</td>
<td>297</td>
<td>488, 545, 584</td>
<td>75</td>
</tr>
<tr>
<td>Ca$_3$(PO$_4$)$_6$S</td>
<td>Ce$^{3+}$ → Tb$^{3+}$</td>
<td>280, 430, 545</td>
<td>410, 550</td>
<td>76</td>
</tr>
<tr>
<td>Sr$_2$MgSiO$_4$</td>
<td>Ce$^{3+}$ → Tb$^{3+}$</td>
<td>228</td>
<td>542</td>
<td>77</td>
</tr>
<tr>
<td>Sr$_2$Gd$_5$(SiO$_4$)$_5$O$_2$</td>
<td>Ce$^{3+}$ → Tb$^{3+}$</td>
<td>348</td>
<td>540, 550</td>
<td>78</td>
</tr>
<tr>
<td>KNaSO$_4$</td>
<td>Ce$^{3+}$ → Dy$^{3+}$</td>
<td>280</td>
<td>330, 483, 556</td>
<td>79</td>
</tr>
<tr>
<td>(Dy$<em>{0.01}$Ce$</em>{0.99-\ldots}$)$_3$Al$_2$O$_12$</td>
<td>Ce$^{3+}$ → Dy$^{3+}$</td>
<td>327, 367</td>
<td>496, 582</td>
<td>80</td>
</tr>
<tr>
<td>Mg$_3$Ca$_4$(PO$_4$)$_6$</td>
<td>Ce$^{3+}$ → Mn$^{2+}$</td>
<td>362, 410, 418</td>
<td>352, 645</td>
<td>81</td>
</tr>
<tr>
<td>K$_2$AEF$_2$O$_7$ (AE = Ca, Sr)</td>
<td>Ce$^{3+}$ → Mn$^{2+}$</td>
<td>320</td>
<td>534, 539</td>
<td>82</td>
</tr>
<tr>
<td>Ca$_3$(PO$_4$)$_6$O</td>
<td>Ce$^{3+}$ → Eu$^{2+}$</td>
<td>350</td>
<td>460, 630</td>
<td>83</td>
</tr>
<tr>
<td>Ca$_3$Y$_2$(SiO$_4$)$_6$</td>
<td>Ce$^{3+}$ → Eu$^{2+}$</td>
<td>356, 280</td>
<td>426, 327</td>
<td>84</td>
</tr>
<tr>
<td>Na$_2$Ca$_4$(PO$_4$)$_6$</td>
<td>Ce$^{3+}$ → Eu$^{2+}$</td>
<td>319</td>
<td>391, 487</td>
<td>85</td>
</tr>
</tbody>
</table>
3.4 The co-doping of Eu\textsuperscript{3+}/Mn\textsuperscript{2+}

Generally speaking, Eu\textsuperscript{3+} presents yellow and blue light, which causes the obtained white light to lack red component, leading to the lower CRI. To increase the CRI of the material, the emission intensity of red light must be improved. Since Mn\textsuperscript{2+} can emit red light, co-doping Eu\textsuperscript{3+} and Mn\textsuperscript{2+} was attempted. For another reason, the emission color of Eu\textsuperscript{3+} singly-doped phosphors can be varied from blue to red in different hosts, nevertheless, the emission color is commonly unchangeable or slightly changeable in certain systems. In order to obtain tuned emission colors and substantial colors, the energy properties from Eu\textsuperscript{3+} to Mn\textsuperscript{2+} are frequently exploited in phosphors.

It has been proved that there is energy transfer between Eu\textsuperscript{3+} and Mn\textsuperscript{2+}, where Eu\textsuperscript{3+} acts as a sensitizer transferring part of its excitation energy into neighboring Mn\textsuperscript{2+} ions. There are many Eu\textsuperscript{3+}/Mn\textsuperscript{2+} co-doped phosphors, such as Sr\textsubscript{2}MgSiO\textsubscript{4}: Eu\textsuperscript{3+}, Mn\textsuperscript{2+},\textsuperscript{90} Sr\textsubscript{3}La(PO\textsubscript{4})\textsubscript{3},\textsuperscript{91} NaScSi\textsubscript{3}O\textsubscript{8},\textsuperscript{92} and NaCaBO\textsubscript{3} (ref. 93) all of which have the properties of spectral adjustment, due to the energy transfer between Eu\textsuperscript{3+} and Mn\textsuperscript{2+}.

Taking Ca\textsubscript{9}Mg(PO\textsubscript{4})\textsubscript{6}F\textsubscript{2}: Eu\textsuperscript{3+}, Mn\textsuperscript{2+} as an example, the excitation and emission spectra of Ca\textsubscript{9}Mg(PO\textsubscript{4})\textsubscript{6}F\textsubscript{2}:0.18Eu\textsuperscript{3+} and Ca\textsubscript{9}Mg(PO\textsubscript{4})\textsubscript{6}F\textsubscript{2}:0.18Mn\textsuperscript{2+} samples are shown in Fig. 15, respectively. It was observed that there is an overlap of Mn\textsuperscript{2+} excitation and Eu\textsuperscript{3+} emission in Fig. 17, which proves the possibility for resonance type energy transfer from Eu\textsuperscript{3+} to Mn\textsuperscript{2+} in the host. Except for the intensity in the Ca\textsubscript{9}Mg(PO\textsubscript{4})\textsubscript{6}F\textsubscript{2}:0.18Eu\textsuperscript{3+},0.18Mn\textsuperscript{2+} phosphor from Fig. 17, it can be further demonstrated by the similar excitation spectra monitored at 545 and 565 nm. Due to the emission spectrum covering the emission band of Eu\textsuperscript{3+} and Mn\textsuperscript{2+}, the tuneable color can be gained via adjusting the Mn\textsuperscript{2+} concentration. Moreover, the lifetime of Eu\textsuperscript{3+} decreases with increasing Mn\textsuperscript{2+}, which provides further confirmation of energy transfer from Eu\textsuperscript{3+} to Mn\textsuperscript{2+} ions. Fig. 17 shows the tunable emission color via changing the ratio of Eu\textsuperscript{3+} and Mn\textsuperscript{2+} doping concentration. The variation of the CIE chromaticity coordinates from blue to yellow with increasing Mn\textsuperscript{2+} doping concentration from 0 to 0.38 is shown in Fig. 17. Eu\textsuperscript{3+}/Mn\textsuperscript{2+}-activated CMPF phosphors can be potentially applied in UV-pumped white LEDs with outstanding properties.\textsuperscript{94}

Liu et al. have synthesized single-phase white light-emitting KCaY(PO\textsubscript{4})\textsubscript{2}:Eu\textsuperscript{3+},Mn\textsuperscript{2+} for light-emitting diode (LED) applications. The absolute quantum efficiency of photoconversion was calculated using the following equation:

\[
A = \frac{L_\text{c}(\lambda) - L_\text{i}(\lambda)}{L_\text{i}(\lambda)} \quad (24)
\]

where \(L_\text{c}(\lambda)\) is the integrated excitation profile when the sample is diffusely illuminated by the integrated sphere’s surface; \(L_\text{i}(\lambda)\) is excited by the incident beam. Furthermore, the quantum efficiency (\(\Phi\)) of KCaY(PO\textsubscript{4})\textsubscript{2}:1%Eu\textsuperscript{3+} phosphors can be calculated by

\[
\Phi = \frac{E_\text{i}(\lambda) - (1 - A)E_\text{c}(\lambda)}{L_\text{i}(\lambda)A} \quad (25)
\]

where \(E_\text{i}(\lambda)\) is the integrated luminescence of the powder upon direct excitation and \(E_\text{c}(\lambda)\) is the integrated. The internal quantum efficiency of KCaY(PO\textsubscript{4})\textsubscript{2}:1%Eu\textsuperscript{3+} was found to be 35.8% and the corresponding external quantum efficiency was 21.4% at the excitation wavelength of 365 nm. Fig. 16 shows the excitation spectra of KCaY(PO\textsubscript{4})\textsubscript{2}:1%Eu\textsuperscript{3+},x%Mn\textsuperscript{2+} phosphors (\(x = 0, 1, 2, 4, 5, 7, 10\)) excited by 365 nm, and the energy transfer efficiency is shown in the inset in Fig. 18. The phosphors generated red and blue emission bands, centered at 652 nm (\(^4\text{T}_1(\text{G}) \rightarrow ^6\text{A}_1(\text{S})\) transition of Mn\textsuperscript{2+}) and 480 nm (4f\textsuperscript{5}5d\textsuperscript{1} \rightarrow 4f\textsuperscript{7} transition of Eu\textsuperscript{3+}). The intensity of the Eu\textsuperscript{3+} at 480 nm decreased as the Mn\textsuperscript{2+} content increased to \(x\). The intensity at 652 nm increased as the Mn\textsuperscript{2+} content increased,
reaching a maximum at $x = 7$ mol%, and then decreased when $x$ exceeded 7 mol%, due to the concentration quenching effect. It was demonstrated that the mechanism of energy transfer is an electric dipole–quadrupole interaction from sensitizer Eu$^{2+}$ to an activator Mn$^{2+}$. Fig. 16 provides a summary of the CIE chromaticity corresponding to the emission spectra excited by 380 nm, which indicates that changing the Mn$^{2+}$ concentration can tune the color from blue through white light and eventually to red in the visible spectral region. The electroluminescence spectrum of white LED lamps fabricated using a NUV 405 nm chip combined with a single-phase white-emitting phosphor KCaY(PO$_4$)$_2$:1% Eu$^{2+}$,4%Mn$^{2+}$ driven by a 350 mA current is shown in Fig. 18. The correlated color temperature is 6507 K and CIE color coordinates are (0.314, 0.329), which indicate that the composition-optimized KCaY(PO$_4$)$_2$:1%Eu$^{2+}$,4%Mn$^{2+}$ phosphor may have promising applications for white-light NUV LEDs.

4 Summary and outlook

Due to the characteristics of the transition metal electronic orbitals, Mn has many valence states, with Mn$^{2+}$ and Mn$^{4+}$ being the most common. As activated ions, they have a wide range of applications. Mn$^{4+}$ doped materials can work as a complement to the commercial phosphor powder “blue chips + Y$_3$Al$_5$O$_{12}$:Ce$^{3+}$ (YAG:Ce)” to make up for the lack of red component, in order to improve the color rendering indexes and the correlated color temperature. They have the advantages of a simple preparation route and low-cost. Mn$^{4+}$ activated oxide and fluoride red phosphors have attracted more and more attention, but low resistance to moisture and unclear structure-related luminescence mechanisms limit their rapid progress in scientific investigation and application of WLED lighting.

Ce$^{3+}$/Mn$^{2+}$, Eu$^{2+}$/Mn$^{2+}$ co-doped materials can produce color tunable phosphors in order to achieve UV excitable single mechanism phosphors to produce white light. In recent years, a single host material emitting white light has made significant progress. In order to further improve its performance, some more studies need to be conducted. Double-doped materials have some deficiencies in the spectral distribution; for example, Eu$^{2+}$/Mn$^{2+}$ or Ce$^{3+}$/Mn$^{2+}$ doped systems in the green emission region are relatively weak, which is certain
to affect the luminescence. Further research is being conducted on how to improve the spectral distribution of dual-doping systems, and how to meet the needs of different white phosphors.

In addition, due to the splitting of the $\text{Mn}^{2+}$ energy level, the emission spectrum can be red shifted to produce deep red and near infrared light. It has been widely used in the fields of optical fiber communication, solid state lasers, light emitting markers and fluorescence immunoassays.

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