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## Accurate analysis of the size distribution and crystallinity of boron doped Si nanocrystals via Raman and PL spectra

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A narrow size distribution of quantum dots (QDs) is needed for their application in photovoltaics but collection of such information is difficult. This paper demonstrates the application of Raman spectroscopy as a characterisation tool to extract the size distribution and crystalline fraction of Si QD samples fabricated through the sputter-anneal method. Measured Raman spectra of Si QD materials are de-convoluted into four components according to their origins and Raman scattering by Si QD cores is described by a modified one phonon confinement model, while other components are reproduced with Gaussian functions. Through fitting of Raman spectra, Si QD size distributions and Si crystalline fractions are obtained. The results are compared with the values extracted from PL modelling on a series of B doped Si QD samples. The good consistency between the values extracted by these two methods confirms the validity of the Raman model. The result confirms that Si crystallization has been suppressed by B doping as the average Si QD size and Si crystalline fraction are reduced with increased B doping level.

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### Introduction

Silicon quantum dots (Si QDs) embedded in an oxide or nitride matrix are considered to have high energy conversion efficiencies.<sup>1</sup> To realize such efficient photonic devices, a tight control over size and a narrow size distribution are necessary. However, the size distribution of Si QDs embedded in a matrix is very difficult to obtain, considering their non-uniform shapes and remarkable difficulties in TEM imaging.<sup>2,3</sup> Meanwhile, the optical characterization techniques such as PL and Raman spectroscopy are non-destructive and easily-performed but still lack quantitative analysis of the spectra. We proposed a method to accurately estimate the size distribution of such QDs from their photoluminescence (PL) spectra in previous work.<sup>4</sup> The size distribution of nanoparticles from low frequency Raman scattering (LFRS) has been studied<sup>5-7</sup> and in this paper high frequency Raman spectroscopy is introduced as another method to extract such information. Boron doping on Si QDs bilayers is studied by a combination of Raman and PL spectroscopies. Good agreement in term of the deduced Si QD size distribution has been obtained between these two methods.

### Experimental details

Si QD bilayers with varied B doping power are prepared on fused-silica substrates through RF magnetron sputtering system in an argon atmosphere below 1.5 mTorr. Each of the samples consists of 30 bilayers, which are made of alternating SRO and SiO<sub>2</sub> layers with respective thicknesses of 4 nm and 1.8 nm. The volume ratio of Si to SiO<sub>2</sub> is controlled at 66% by adjusting the powers applied to the sputtered targets. In following discussions, samples are denoted as B0, B1, B2, B3, and B4, corresponding to B doping powers of 0 (intrinsic), 12.5, 25, 37.5, and 50 W respectively. Sample annealing is conducted in an 1100 °C N<sub>2</sub> purged environment for 1 hour to facilitate Si QDs nucleation and growth.<sup>2</sup> Annealed samples are characterized with Raman and PL spectroscopies respectively. In Raman measurement, an Ar-ion laser with wavelength of 514 nm is used as incident source without polarization. The diffraction grating of 1800 L mm<sup>-1</sup> is selected, giving an instrumental resolution of 1.7 cm<sup>-1</sup>. The lens magnification (50×) and the exposure time (1 s for 10 times accumulations) adopted are delicately determined after several experiments in order to avoid additional Si crystallization and significant strains induced by a high laser power. The original power of the laser is 25 mW and drops to around 9.9 mW on sample after passing through the optics. The optical penetration depth is about 300 nm for a laser with 514 nm wavelength, covering the total thickness of samples fabricated in this work.<sup>8</sup> In PL measurement, the intensity correction is made using a lamp with a calibrated spectra

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between 250 and 2400 nm. A laser diode with a wavelength of 405 nm is used as the excitation source in room temperature environment, and the PL signal is collected and analysed using a Si CCD camera cooled down to  $-30\text{ }^{\circ}\text{C}$ .

## Results and discussion

Fig. 1 is the high frequency Raman scattering obtained from annealed samples. Main Raman peaks, with different levels of downshifts, are identified near the bulk Si response at  $520\text{ cm}^{-1}$  (labelled) in all samples, indicating the existence of crystalline Si.<sup>9</sup> A one-phonon confinement model was developed by Richter *et al.*<sup>10</sup> accounting for the size confinement effect. Assuming a spherical crystal with diameter  $D$ , the phonon is restricted to the volume of this crystallite, and as a result, the wave function can be expressed as:<sup>10–12</sup>

$$\psi(\mathbf{q}, r) = F_c(r, D)\phi(\mathbf{q}, r) = \psi'(\mathbf{q}, r)u(\mathbf{q}, r) \quad (1)$$

where  $u(\mathbf{q}, r)$  has the periodicity of the lattice and  $F_c(r, D)$  is a phonon weighting/confinement function which can be expanded in Fourier integral as:

$$F_c(r, D) = \int C(\mathbf{q})\exp(i\mathbf{q}\cdot r)d\mathbf{q} \quad (2)$$

Thus

$$C(\mathbf{q}) = \frac{1}{2\pi^3} \int F_c(r, D)\exp(-i\mathbf{q}\cdot r)dr \quad (3)$$

Therefore, the first-order Raman spectrum of a nanocrystal can be calculated as:<sup>13</sup>

$$I(\omega) \propto \int \frac{|C(\mathbf{q})|^2 d\mathbf{q}}{[\omega - \omega'(\mathbf{q})]^2 + \left(\frac{\Gamma_0}{2}\right)^2} \quad (4)$$

where  $\omega'(\mathbf{q})$  is the phonon dispersion relation curve,  $\Gamma_0$  the natural linewidth ( $\Gamma_{\text{Si}} = 3\text{ cm}^{-1}$ ), and  $a$  the lattice constant ( $a_{\text{Si}} = 0.543\text{ nm}$ ) of the material.

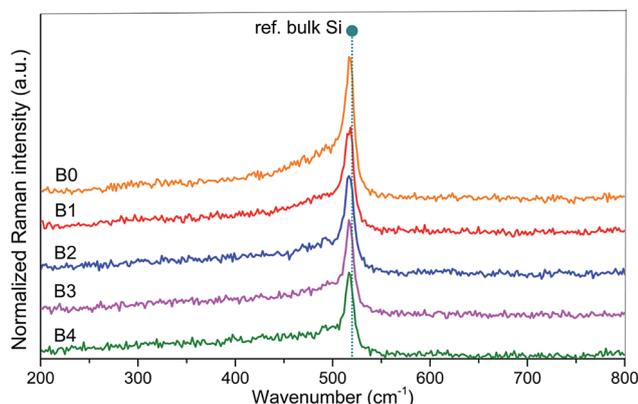


Fig. 1 Raman spectra for B doped samples.

The selection of the phonon confinement function  $F_c(r, D)$  and the integral range are two key parts of in eqn (4). Though the analytical form for  $F_c(r, D)$  can be arbitrary without significant physical reasons, in most case a Gaussian function shown in eqn (5) is adapted to characterize the Raman spectra of Si nanostructures.<sup>10,13</sup>

$$F_c(r, D) = \exp\left(-\frac{\alpha r^2}{D^2}\right) \quad (5)$$

The corresponding Fourier coefficients are given by as:

$$|C(\mathbf{q})|^2 = \frac{D^6}{16(2\pi)^2 \alpha^3} \exp\left(-\frac{q^2 D^2}{2\alpha}\right) \quad (6)$$

In eqn (5),  $\alpha$  is a factor determining the degree of confinement, *i.e.* how quickly the wave function of phonon vanishes towards the boundary of a nanocrystal. In this case, the integral of eqn (5) is extended to the entire first Brillouin zone. However, Faraci *et al.*<sup>12</sup> argued that this approach only observe a single-phonon wave vector, which was not realistic. They suggested a wave packet extended over the size of the nanocrystal; and in their modified model, a confinement function given by a weighted superposition of sinusoidal waves was tested.

$$F_c(r, D) = f(x) = \begin{cases} \sum_n \frac{\sin(k_n r)}{k_n r}, & r \leq \frac{D}{2} \\ 0, & r > \frac{D}{2} \end{cases} \quad (7)$$

where  $k_n = \frac{n\pi}{D}$ ,  $n = 2, 4, 6, \dots, n_{\text{max}} < \frac{2D}{a_{\text{Si}}}$ .

Thus, the corresponding Fourier coefficients are given by as:

$$|C_n(\mathbf{q})|^2 = \left| 3 \frac{\sin\left(\frac{qD}{2}\right)}{\pi^3 D^3 q (k_n^2 - q^2)} \right|^2 \quad (8)$$

The form in eqn (7) has components centered at  $r = 0$ , which strongly attenuate to zero at the border of the nanocrystal.<sup>12</sup> It is found that the first component of  $F_c(r, D)$  (*i.e.*  $n = 2$ ) is the most appropriate for simulating the vibrational amplitude of the most Raman-active modes in a Si QD.<sup>12</sup> In this modified approach, the phonon wave vectors  $\mathbf{q}$  in eqn (7) that are contributing to the scattering process are limited to the range  $\left[\frac{n\pi - 1}{D}, \frac{n\pi + 1}{D}\right]$  due to the momentum conservation.<sup>12,14</sup>

This modification can give a good description of the vibrational confinement in Si spheres and a relatively good description of the Raman frequency shifts.<sup>13</sup> Taking the Raman scattering

efficiency  $\rho(\omega) \sim \frac{n(\omega) + 1}{\omega} = \frac{\left(\exp\left(\frac{h\omega}{kT}\right) - 1\right)^{-1} + 1}{\omega}$  into account,<sup>14</sup> the integral first-order Raman spectrum in eqn (4) can be rewritten as:<sup>12,14</sup>



$$I(\omega, D) \propto \rho(\omega) \frac{1}{\pi D^3} \int_{2\pi-1}^{2\pi+1} \frac{4\pi Q^2 \left| 3 \frac{\sin\left(\frac{Q}{2}\right)}{\pi^3 Q(4\pi^2 - Q^2)} \right|^2 \left(\frac{\Gamma}{2}\right)}{\left[\omega - 521 \left(1 - 0.23 \left(\frac{Qa_{\text{Si}}}{2\pi D}\right)^2\right)\right]^2 + \left(\frac{\Gamma}{2}\right)^2} dQ \quad (9)$$

In eqn (9), the size dependency of the integral boundaries is removed with a change of variables. *i.e.*  $Q = qD$ .<sup>14</sup> The common used analytical form of phonon dispersion curve  $\omega'(q)$  for Si are  $\omega'(q) = \sqrt{A + B \cos(qa/2.6)}$  and  $\omega'(q) = 521 \left(1 - 0.23 \left(\frac{qa_{\text{Si}}}{2\pi}\right)^2\right)$ , respectively.<sup>15–17</sup> But Faraci *et al.*<sup>12</sup> has shown that notwithstanding the different analytical form, these two dispersion curves are very close in the common range of interest. Here we adopt the analytical form of  $\omega'(q) = 521 \left(1 - 0.23 \left(\frac{qa_{\text{Si}}}{2\pi}\right)^2\right) = 521 \left(1 - 0.23 \left(\frac{Qa_{\text{Si}}}{2\pi D}\right)^2\right)$  for Si phonon dispersion relation in eqn (9). Fig. 2 demonstrates simulated Raman spectra for different sizes of Si nanocrystals. It is also noticeable that Raman response from Si nanocrystal with diameter of 20 nm is close to the one from bulk Si, and therefore it is safe to state that nanocrystal size larger than this value can be viewed as bulk Si behavior.

Eqn (9) enables experimental analysis of Si QD materials with a monodisperse distribution, *i.e.* uniform size. In reality, however, the size of fabricated Si nanocrystals is found to be spread out, which will make the Raman spectra more complicated. To account for this effect, eqn (9) is required to be projected onto the log-normal size distribution function that is integrated over the possible size range in the following way:<sup>18,19</sup>

$$I(\omega) = \int_{D_{\text{min}}}^{D_{\text{max}}} \Phi(D) I(\omega, D) dD \quad (10)$$

in which the log-normal size distribution with a mean size  $D_0$  and a distribution skewness  $\sigma$  is given as:

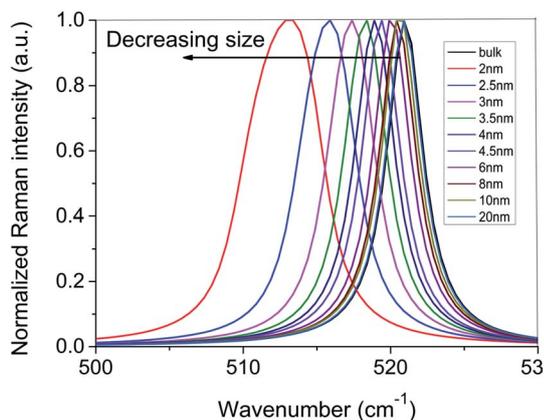


Fig. 2 Simulated Raman spectra (normalized) for varied sizes of Si nanocrystals using eqn (9).

$$\Phi(D) = \frac{1}{\sigma D \sqrt{2\pi}} \exp\left\{-\frac{\log\left(\frac{D}{D_0}\right)^2}{2\sigma^2}\right\} \quad (11)$$

The effect of distribution skewness, or the degree of size dispersion, is simulated and displayed in Fig. 3. It is identified that as  $\sigma$  value increases, the Raman peak position is slightly moved towards the lower wave number while the asymmetry broadening of the spectra is significantly enhanced. This is because Si nanocrystals with smaller size tend to have higher Raman scattering efficiency according to the expression of  $\rho(\omega)$  and thus their contributions to the final spectra become more obvious. It is also noticeable that as  $\sigma$  is approaching 1/10 of the value of mean size (*i.e.*  $\sigma = 0.4, 0.5$ ), fringes appear in the spectra in the low wavenumber regime. This feature is a consequence of higher Raman scattering efficiency of smaller sized Si QD given in previous discussion.

Apart from the size distribution caused effect, the amorphous Si phase has to be considered in the analysis of the Raman spectrum as well. Si QDs embedded in an amorphous matrix (either amorphous Si or dielectric material) tends to have a core-shell structure, in which the nanoscale Si crystalline core is surrounded by an outer shell.<sup>20,21</sup> The outer shell consists of near-surface fine sized crystalline phases with irregular shapes, which gradually transits to amorphous Si phases (a-Si) before reaching the amorphous matrix.<sup>21–24</sup> The presence of the near-surface crystalline phases manifests itself as a weak intermediate peak near 500–510  $\text{cm}^{-1}$  in Raman spectrum. The a-Si phases are excess Si that is not crystallized after annealing. In terms of Raman spectroscopy, a-Si will introduce broad humps in the low wavenumber tail of the measured spectra, which include: (i) 300  $\text{cm}^{-1}$  phase from the longitudinal acoustic (LA) vibration mode; (ii) 480  $\text{cm}^{-1}$  phase from the transverse optical (TO) vibration mode.<sup>25</sup> The existences of a-Si phases and near-surface states have important influences on both optical and electrical properties of the material. They are closely affected by material

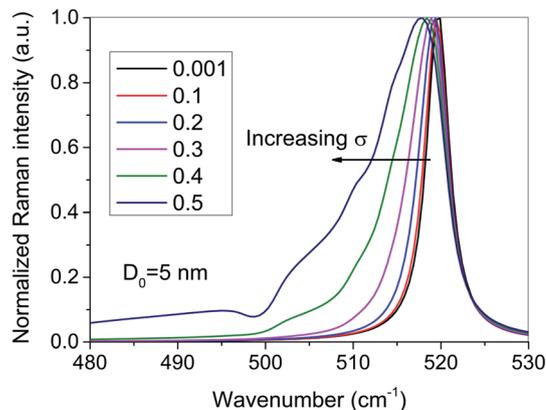


Fig. 3 Simulated Raman spectra (normalized) for Si nanocrystals with mean size  $D_0 = 5$  nm and different distribution skewness  $\sigma$  using eqn (10).



processing parameters, including material stoichiometry, annealing temperature and duration, and doping types and levels. In this paper, these phases will be reproduced with their corresponding Gaussian functions in Raman spectra simulations.

A typical deconvolution of Raman spectrum adapted in this paper is shown in Fig. 4. Under the assumption that the majority of Si QD is hard sphere in shape, the response from c-Si QD cores in the vicinity of  $520\text{ cm}^{-1}$  is fitted with the modified one-phonon confinement model. The intermediate peak from near-surface crystalline states around  $510\text{ cm}^{-1}$  is reproduced using a Gaussian function. As pointed out above, though these states are also crystalline Si components, they are prone to have irregular shapes instead of spheres. This will introduce complicated modifications to the weighing function if the one-phonon confinement model is employed to fit their Raman responses. Therefore, for simplicity, a Gaussian function is applied instead. Apart from c-Si responses, the two a-Si phases near  $300\text{ cm}^{-1}$  and  $480\text{ cm}^{-1}$  are reproduced by their corresponding Gaussian functions, which is a widely adapted method in other researches.<sup>22–24</sup> In the rest parts of this paper, Raman spectrum analysis will all be based on the simulation process described above.

Using the modified one-phonon confinement model, the log-normal size distribution function of Si QDs can be extracted by adjusting the mean size  $D_0$  and distribution skewness  $\sigma$ , see inset of Fig. 4. The crystalline fraction  $X_c$  of Si with respect to a-Si can be estimated through eqn (12).<sup>24,26</sup>

$$X_c = \frac{\sum I_c}{\sum I_c + \alpha \sum I_a} \quad (12)$$

in which  $\sum I_c$  and  $\sum I_a$  are sums of integrated intensities for c-Si (*i.e.*  $518\text{ cm}^{-1}$  and  $510\text{ cm}^{-1}$  peaks in Fig. 4), and a-Si (*i.e.*  $300$

$\text{cm}^{-1}$  and  $480\text{ cm}^{-1}$  peaks in Fig. 4) profiles respectively, and  $\alpha = 0.1$  due to the silicon content used.<sup>26,27</sup>

To verify the validity of extracting Si QDs size distribution information from Raman simulation, a model is introduced and used to fit PL spectrum for the same purpose. This PL model correlates the breadth of PL emission with the size distribution of Si QDs, which in return is manifested as distribution of energy bandgaps.<sup>4</sup> Thus the PL intensity as a function of energy  $\hbar\omega$  is expressed in eqn (13):<sup>4</sup>

$$I(\hbar\omega) \propto \int_0^{\hbar\omega} (\hbar\omega - E_g)^2 \cdot (\hbar\omega)^2 \cdot e^{-\frac{\hbar\omega}{kT}} \cdot e^{\frac{\Delta F}{kT}} \cdot f(E_g) dE_g \quad (13)$$

where  $k$  is the Boltzmann constant, and  $E_g$  is energy bandgap.  $\Delta F = E_{Fn} - E_{Fp}$  with  $E_{Fn}$  and  $E_{Fp}$  correspond to quasi-Fermi energies for electrons and holes respectively.  $f(E_g)$  is the probability density function of the energy bandgaps. Again, a log-normal function in eqn (11) is adapted for Si QDs size distribution, then  $f(E_g)$  is given by eqn (14).<sup>4</sup>

$$f(E_g) = \frac{\sqrt{C_{\text{lump}}} \cdot (E_g - E_{g_{\text{bulk}}})^{-\frac{3}{2}}}{2\sigma \sqrt{2\pi} \cdot \frac{C_{\text{lump}}}{E_g - E_{g_{\text{bulk}}}}} \times \exp \left\{ -\frac{\left( \ln \sqrt{\frac{C_{\text{lump}}}{E_g - E_{g_{\text{bulk}}}}} - \ln(D_0) \right)^2}{2\sigma^2} \right\} \quad (14)$$

in which  $E_{g_{\text{bulk}}}$  is the energy bandgap of bulk Si, and  $C_{\text{lump}} = 2\text{ eV nm}^2$  is a coefficient lumping all other parameters that account for quantum confinement effect.<sup>28</sup> Similarly, size distribution of Si QDs can be obtained through simulation of the PL profile. With this alternative approach, it is possible to support conclusions from Raman simulation described previously.

PL spectra of all samples are demonstrated in our previous work.<sup>29</sup> The intensity of PL response initially increases from intrinsic sample (B0) to lightly doped sample (B1), which is followed by a decrease in intensity as the doping power further increases (B2–B4). Simultaneously, the PL peak position is slightly blue-shifted from intrinsic sample to doped samples. Assuming the maxima of the PL profiles are responses from Si QDs, such shifts indicate changes in energy bandgap due to evolution of Si QDs size distributions altered by B doping. Si QDs size distributions extracted through simulations of PL spectra are summarized in Table 1.

According to Fig. 5, extracted values for Si QDs average size from Raman and PL simulations are in good consistency with each other. In general, when the B doping power is increased, the average Si QDs size gets smaller. The crystalline fraction  $X_c$  plotted in the same diagram also demonstrates a reducing trend when B doping level is enhanced. Both of the changes imply the suppression of Si crystallization by B doping. Similar effect has been reported and discussed by Hao *et al.*,<sup>30</sup> in which they attributed this inhibited Si crystallization to the impurity induced local deformations.

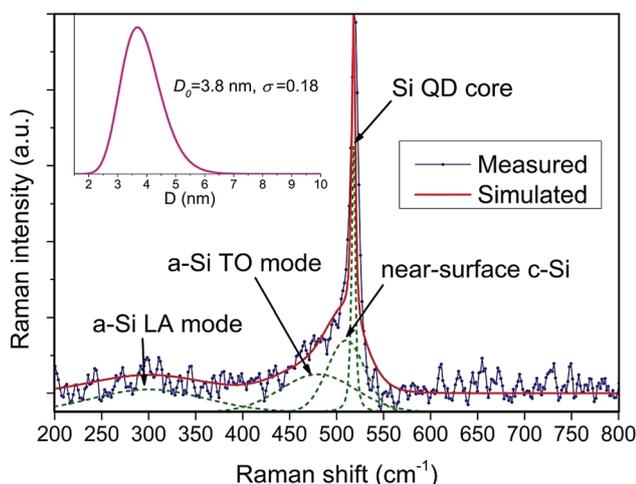
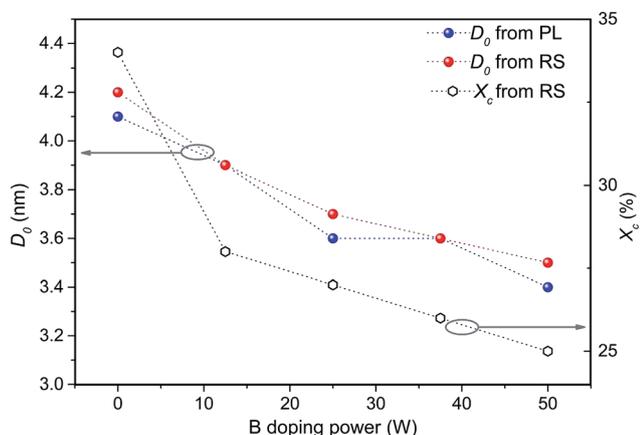


Fig. 4 A typical deconvolution of Raman spectrum: The spectrum consists a c-Si phase that originates from Si QD core in the vicinity of  $520\text{ cm}^{-1}$ , two a-Si phases from LA mode and TO mode near  $300\text{ cm}^{-1}$  and  $480\text{ cm}^{-1}$ , respectively, an intermediate phase due to near-surface c-Si near  $510\text{ cm}^{-1}$ . The inset demonstrates the log-normal size distribution for Si QDs extracted from simulation of the Raman spectrum.



**Table 1** Summary of simulation results of Raman and PL spectroscopies

Samples	$X_c$	Si QDs size distribution from Raman		Si QDs size distribution from PL	
		$D_0$ (nm)	$\sigma$	$D_0$ (nm)	$\sigma$
B0	34%	4.2	0.43	4.1	0.54
B1	28%	3.9	0.40	3.9	0.51
B2	27%	3.7	0.39	3.6	0.51
B3	26%	3.6	0.35	3.6	0.52
B4	25%	3.5	0.35	3.4	0.54

**Fig. 5** Extracted Si QDs average sizes and crystalline fractions from Raman and PL simulations versus B doping power.

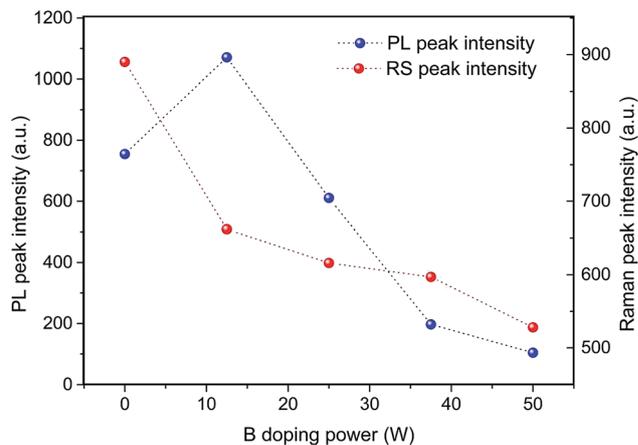
In Table 2,  $\sigma$  values obtained from Raman simulations are lower than those from PL simulations for each sample. This discrepancy can be explained in two aspects. Firstly, the full width half maximum (FWHM) of simulated c-Si peak using the modified one-phonon confinement model are smaller than those of the measured Raman spectra, see Table 2. Though the level of size dispersion (*i.e.*  $\sigma$  value) has little effect on the calculated Raman frequency shift, it has great influence on the width and the shape of the Raman spectra.<sup>31</sup> Therefore, lower values of deduced FWHM indicate possible underestimations of  $\sigma$  from Raman simulations. This is because the Gaussian function for the intermediate peak (near-surface crystalline phases) also contributes to the FWHM of the final simulated Raman profile, however, its size information is not accounted.

**Table 2** Comparison of FWHM of simulated c-Si peaks and measured peaks

Samples	FWHM of simulated c-Si peak ( $\text{cm}^{-1}$ )	FWHM of measured peak ( $\text{cm}^{-1}$ )
B0	10.1	13.5
B1	9.7	13.8
B2	9.1	13.8
B3	8.9	11.3
B4	8.3	11.3

Secondly, in PL simulations, the model employed here takes into consideration of all radiative recombination processes regardless of their actual origins. Since it is not easy to deconvolute emissions of defect states from those of Si QDs, the level of Si QDs size dispersion (*i.e.*  $\sigma$  values) extracted in PL simulations can be overestimated to some extent. Additionally, the aforementioned near-surface states, which are excluded from the one-phonon confinement model in Raman simulations, may cause exciton localization and thus can be PL active.<sup>21,32</sup> As a consequence, the size information these states contain may be included in PL simulations similar to defect-related emissions. Therefore, a  $\sigma$  in between the range of Raman and PL extracted values can be a reasonable good approximation of Si QDs size distribution skewness.

Fig. 6 compares the evolution of relative intensities of Raman and PL peaks as a function of B doping power. To clarify, intensities shown on the graph were average results. According to the experiments, measurement errors for Raman spectra with detected peak intensity around 100 counts were within 6% of the average value; while for PL with peak intensity around 100 counts, the values were less than 10% of the average. Measurement result tends to become more converged as the detected intensity gets higher and thus the error range narrows. Therefore, for the detected intensity levels in both Raman and PL spectra demonstrated here, error bars are not illustrated in Fig. 6 due to their relatively small magnitudes compared to the measurement values. In Raman measurement, the intrinsic sample shows the strongest response. With additions of B dopants, the obtained peak intensity decreases. Under the identical measurement conditions and equipment setups, this tendency, on one hand, is associated with the reduction of Raman scattering volumes due to the decreased Si QDs sizes by increased B doping levels;<sup>33</sup> and on the other hand, is due to the lowered concentration of c-Si vibrational modes caused by decreased Si crystalline fractions according to Fig. 5.<sup>34</sup> The relative intensity of PL peak in Fig. 6 is initially enhanced from sample B0 to B1, which is then reduced when the B doping level further increases (B2–B4). It is verified that the improved PL peak

**Fig. 6** Comparison of PL and Raman peak intensities as a function of B doping power.

intensity in lightly doped sample (B1) is a result of B passivation effect to a-Si dangling bonds.<sup>29</sup> When the B doping level continues to increase, both Auger recombination and impurity induced strains (which will add to non-radiative recombination centers) are enhanced simultaneously.<sup>29</sup> As a combined result, PL emissions are quenched in samples B2–B4.

To conclude, Si QDs average sizes extracted from Raman and PL simulations demonstrate consistent results. It is observed that extracted size distribution skewness from Raman modelling is underestimated, while the one from PL modelling is overestimated. Therefore, a value between them can be considered a good description of size dispersion level. Both of Raman and PL spectroscopies verify that Si crystallization has been suppressed by B doping as average Si QDs size and Si crystalline fraction are reduced with increased B doping level. The decreasing Raman peak intensity with B doping power is caused by reduction in both the scattering volume in smaller Si QDs, and the available vibrational modes. On the other hand, an increase of PL peak intensity is identified when Si QDs are slightly doped, which is benefited from B passivation effect. Further improvement of B doping level leads to reduced PL peak intensities as a result of enhanced Auger process and other non-radiative recombination.

## Conclusion

This paper demonstrates the application of Raman spectroscopy as a characterization tool to extract the size distribution and crystalline fraction of Si QD samples fabricated through sputter-anneal method. Basic principles of Raman spectroscopy and its application in crystalline materials, including both perfect and imperfect crystals, are introduced. Measured Raman spectra of Si QD materials are de-convoluted into four components according to their origins. Raman scattering by Si QD core is described by a modified one-phonon confinement model, while other components are reproduced with Gaussian functions. Through fitting of Raman spectra, Si QDs size distributions and Si crystalline fractions are obtained. The results are compared with the values extracted from PL modelling on a series of B doped Si QD samples. The good consistency between the values extracted by these two methods confirms the validity of the Raman model developed in this paper.

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